Search for pair production of gluinos decaying via stop and sbottom in events with b-jets and large missing transverse momentum in pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector

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I. INTRODUCTION

Supersymmetry (SUSY) [1–6] is a generalization of space-time symmetries that predicts new bosonic partners to the fermions and new fermionic partners to the bosons of the Standard Model (SM). If R parity is conserved [7], SUSY particles are produced in pairs and the lightest supersymmetric particle (LSP) is stable. The scalar partners of the left- and right-handed quarks, the squarks $\tilde{q}_L$ and $\tilde{q}_R$, can mix to form two mass eigenstates $\tilde{q}_1$ and $\tilde{q}_2$, ordered by increasing mass. SUSY can solve the hierarchy problem [8–11] by preventing “unnatural” fine-tuning in the Higgs sector provided that the superpartners of the top quark (stop, $\tilde{t}_1$ and $\tilde{t}_2$) have masses not too far above the weak scale. Because of the SM weak isospin symmetry, the mass of the left-handed bottom quark scalar partner (sbottom, $\tilde{b}_L$) is tied to the mass of the left-handed top quark scalar partner ($\tilde{t}_L$), and as a consequence the mass of the lightest sbottom $\tilde{b}_1$ is also expected to be close to the weak scale. The fermionic partners of the gluons, the gluinos ($\tilde{g}$), are also constrained by naturalness [12,13] to have a mass around the TeV scale to limit their contributions to the radiative corrections to the stop masses. For these reasons, and because the gluinos are expected to be pair produced with a high cross section at the CERN Large Hadron Collider (LHC), the search for gluino production with decays via stop and sbottom quarks is highly motivated at the LHC.

This paper presents the search for gluino pair production where both gluinos decay either to stops via $\tilde{g} \rightarrow \tilde{t}_1t$ or to sbottoms via $\tilde{g} \rightarrow \tilde{b}_1b$, using a data set of 3.2 fb$^{-1}$ of proton-proton data collected with the ATLAS detector [14] at a center-of-mass energy of $\sqrt{s} = 13$ TeV. Each stop (sbottom) is then assumed to decay to a top (bottom) quark and the LSP: $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^{0}$ ($\tilde{b}_1 \rightarrow b\tilde{\chi}_1^{0}$). The LSP is assumed to be the lightest neutralino $\tilde{\chi}_1^{0}$, the lightest linear superposition of the superpartners of the neutral electroweak and Higgs bosons. The $\tilde{\chi}_1^{0}$ interacts only weakly, resulting in final states with substantial missing transverse momentum of magnitude $E_{T}^{\text{miss}}$. Diagrams of the simplified models [15,16] considered, which are referred to as “Gbb” and “Gtt” in the following, are shown in Figs. 1(a) and 1(b), respectively. The sbottom and stop are assumed to be produced off shell such that the gluinos undergo the three-body decay $\tilde{g} \rightarrow bb\tilde{\chi}_1^{0}$ or $\tilde{g} \rightarrow tt\tilde{\chi}_1^{0}$, and that the only parameters of the simplified models are the gluino and $\tilde{\chi}_1^{0}$ masses.\footnote{Models with on-shell sbottom and stop were studied in Run 1 [17], and the limits on the gluino and the $\tilde{\chi}_1^{0}$ masses were found to be mostly independent of the stop and sbottom masses, except when the stop is very light.}

The Gbb experimental signature consists of four energetic $b$ jets (i.e., jets containing $b$ hadrons) and large $E_{T}^{\text{miss}}$. In order to maintain high signal efficiency, at least three of four required jets must be identified as $b$ jets ($b$ tagged). This requirement is very effective in rejecting $t\bar{t}$ events, which constitute the main background for both the Gbb and Gtt signatures, and which contain only two $b$ jets unless they are produced with additional heavy-flavor jets. The Gtt

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\*Full author list given at the end of the article.
experimental signature also contains four \(b\) jets and \(E_T^{\text{miss}}\), but yields in addition four \(W\) bosons originating from the top quark decays \(t \rightarrow Wb\). Each \(W\) boson can either decay leptonically (\(W \rightarrow \ell \nu\)) or hadronically (\(W \rightarrow q\bar{q}\)). A Gtt event would therefore possess a high jet multiplicity, with as many as 12 jets originating from top quark decays and, potentially, isolated charged leptons. In this paper, pair-produced gluinos decaying via stop and sbottom quarks are searched for using events with high jet multiplicity, of which at least three must be identified as \(b\) jets, large \(E_T^{\text{miss}}\), and either zero leptons (referred to as the Gtt 0-lepton channel) or at least one identified charged lepton\(^2\) (referred to as the Gtt 1-lepton channel). For both the Gbb and Gtt models, several signal regions are designed to cover different ranges of gluino and \(\tilde{\chi}^0_1\) masses. For the Gtt models with a large mass difference (mass splitting) between the gluino and \(\tilde{\chi}^0_1\), the top quarks tend to be highly boosted and their decay products collimated. In the corresponding signal regions, at least one large-radius, trimmed [18] jet, which is reclustered from small-radius jets [19], is required to have a high mass to identify hadronically decaying boosted top quarks.

Pair production of gluinos, with subsequent decays via sbottom quarks, was searched for in ATLAS Run 1 with a similar analysis requiring at least three \(b\)-tagged jets [17]. It excluded gluino masses below 1290 GeV for LSP masses below 400 GeV at 95% confidence level (C.L.). That analysis also searched for gluinos decaying via stop quarks in events with at least three \(b\)-tagged jets and either zero or at least one identified lepton and obtained the best ATLAS limits for the Gtt models with massless and moderately massive LSP [20]. Gluino masses below 1400 GeV were excluded at 95% C.L. for LSP masses below 400 GeV. Pair-produced gluinos with stop-mediated decays have also been searched for by ATLAS in events with high jet multiplicity [21], events with at least one lepton, many jets, and \(E_T^{\text{miss}}\) [22], and events containing pairs of same-sign leptons or three leptons [23], the latter obtaining the best ATLAS limit for Gtt models with compressed mass spectra between the gluino and the LSP in Run 1 [20] and having since been performed in Run 2 [24].

Similar searches performed with the CMS experiment using \(\sqrt{s} = 8\) TeV [25–30] and \(\sqrt{s} = 13\) TeV data [31–34] have produced comparable results to ATLAS searches.

II. ATLAS DETECTOR

The ATLAS detector is a multipurpose particle physics detector with a forward-backward symmetric cylindrical geometry and nearly \(4\pi\) coverage in solid angle.\(^3\) The inner tracking detector (ID) consists of pixel and silicon microstrip detectors covering the pseudorapidity region \(|\eta| < 2.5\), surrounded by a transition radiation tracker, which enhances electron identification in the region \(|\eta| < 2.0\). Before the start of Run 2, the new innermost pixel layer, the Insertable B-Layer (IBL) [35], was inserted at a mean sensor radius of 3.3 cm. The ID is surrounded by a thin superconducting solenoid providing an axial 2 T magnetic field and by a fine-granularity lead/liquid-argon (LAr) electromagnetic calorimeter covering \(|\eta| < 3.2\). A steel/scintillator-tile calorimeter provides coverage for hadronic showers in the central pseudorapidity range \(|\eta| < 1.7\). The end cap and forward regions \((1.5 < |\eta| < 4.9)\) of the hadronic calorimeter are made of LAr active layers with either copper or tungsten as the absorber material. A muon spectrometer with an air-core toroid magnet system surrounds the calorimeters. Three layers of high-precision tracking chambers provide coverage in the range \(|\eta| < 2.7\).

\(^3\)ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point in the center of the detector. The positive \(x\) axis is defined by the direction from the interaction point to the center of the LHC ring, with the positive \(y\) axis pointing upwards, while the beam direction defines the \(z\) axis. Cylindrical coordinates \((r, \phi)\) are used in the transverse plane, \(\phi\) being the azimuthal angle around the \(z\) axis. The pseudorapidity \(\eta\) is defined in terms of the polar angle \(\theta\) by \(\eta = -\ln \tan (\theta/2)\).

\(^2\)The term “lepton” refers exclusively to an electron or a muon in this paper.
while dedicated fast chambers allow triggering in the region $|\eta| < 2.4$. The ATLAS trigger system [36] consists of a hardware-based Level-1 trigger followed by a software-based high level trigger.

**III. DATA AND SIMULATED EVENT SAMPLES**

The data used in this analysis were collected by the ATLAS detector from $pp$ collisions produced by the LHC at a center-of-mass energy of 13 TeV and 25 ns proton bunch spacing. The full data set corresponds to an integrated luminosity of $3.2 \text{ fb}^{-1}$ with an associated uncertainty of $\pm 5\%$, after requiring that all detector subsystems were operational during data recording. The measurement of the integrated luminosity is derived, following a methodology similar to that detailed in Ref. [37], from a calibration of the luminosity scale using a pair of $x$-$y$ beam-separation scans performed in June 2015. Events are required to pass an $E_T^{\text{miss}}$ trigger that is fully efficient for events passing the preselection defined in Sec. V. Each event includes on average 14 additional inelastic $pp$ collisions (“pileup”) in the same bunch crossing.

Simulated event samples are used to model the signal and background processes in this analysis. The signal samples for the Gbb and Gtt processes are generated with up to two additional partons using MADGRAPH5_aMC@NLO samples for the Gbb and Gtt processes are generated with and background processes in this analysis. The signal based high level trigger.

**TABLE I.** List of generators used for the different background processes. Information is given about the pQCD highest-order accuracy used for the normalization of the different samples, the underlying-event tunes and PDF sets considered.

<table>
<thead>
<tr>
<th>Process</th>
<th>Generator + fragmentation/hadronization</th>
<th>Tune</th>
<th>PDF set</th>
<th>Cross-section</th>
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<td>PERUGIA2012</td>
<td>CT10</td>
<td>NNLO + NNLL [61]</td>
</tr>
<tr>
<td>Single top</td>
<td>POWHEG-BOX v2 + PYTHIA-6.428</td>
<td>PERUGIA2012</td>
<td>CT10</td>
<td>NNLO + NNLL [62–64]</td>
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<tr>
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<td>A14</td>
<td>NNPDF2.3 [65]</td>
<td>NLO</td>
</tr>
<tr>
<td>$t\bar{t}h$</td>
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<td>UEEE5</td>
<td>CT10</td>
<td>NLO [66]</td>
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<td>NNLO [67]</td>
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<tr>
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<td>Default</td>
<td>CT10</td>
<td>NNLO [67]</td>
</tr>
</tbody>
</table>
are used to estimate systematic uncertainties on the backgrounds, as described in Sec. VI.

The signal samples are normalized using the best cross sections calculated at NLO in the strong coupling constant, adding the resummation of soft gluon emission at next-to-leading-logarithmic (NLL) accuracy [68–72]. The nominal cross section and the uncertainty are taken from an envelope of cross-section predictions using different PDF sets and factorization and renormalization scales, as described in Ref. [73]. The cross section of gluino pair production in these simplified models is approximately 325 fb for a gluino mass of 1 TeV, falling to 2.8 fb for 1.8 TeV mass gluinos. All background processes are normalized to the best available theoretical calculation for their respective cross sections. The order of this calculation in perturbative QCD (pQCD) for each process is listed in Table I.

**IV. OBJECT RECONSTRUCTION**

Interaction vertices from the proton-proton collisions are reconstructed from at least two tracks with \( p_T > 0.4 \) GeV and are required to be consistent with the beam spot envelope. The primary \( pp \) interaction vertex is identified as the one with the largest sum of squares of the transverse momenta from associated tracks \( (\sum |p_{T,\text{track}}|^2) \) [74].

Basic selection criteria are applied to define candidates for electrons, muons, and jets in the event. An overlap removal procedure is applied to these candidates to prevent double counting. Further requirements are then made to select the final signal leptons and jets from the remaining objects. The details of the object selections and of the overlap removal procedure are given below.

Candidate jets are reconstructed from three-dimensional topological energy clusters [75] in the calorimeter using the anti-\( k_t \) jet algorithm [76] with a radius parameter of 0.4 (small-\( R \) jets). Each topological cluster is calibrated to the electromagnetic scale response prior to jet reconstruction. The reconstructed jets are then calibrated to the particle level by the application of a jet energy scale (JES) derived from simulation and corrections based on 8 TeV data [77,78]. Quality criteria are imposed to reject events that contain at least one jet arising from noncollision sources or detector noise [79]. Further selections are applied to reject jets that originate from pileup interactions [80]. Candidate jets are required to have \( p_T > 20 \) GeV and \( |\eta| < 2.8 \). Signal jets, selected after resolving overlaps with electrons and muons, are required to satisfy the stricter requirement of \( p_T > 30 \) GeV.

A multivariate algorithm using information about the impact parameters of inner detector tracks matched to the jet, the presence of displaced secondary vertices, and the reconstructed flight paths of \( b \) and \( c \) hadrons inside the jet [81–83] is used to tag \( b \) jets. The \( b \) tagging working point with an 85% efficiency, as determined from a simulated sample of \( t\bar{t} \) events, was found to be optimal.

The corresponding rejection factors against jets originating from \( c \) quarks, from \( \tau \) leptons, and from light quarks and gluons in the same sample at this working point are 2.6, 3.8, and 27, respectively.

The candidate small-\( R \) jets are used as inputs for further jet reclustering [19] using the anti-\( k_t \) algorithm with a radius parameter of 1.0. These reclustered jets are then trimmed [18,19] by removing subjets whose \( p_T \) falls below \( f_{\text{cut}} = 5\% \) of the \( p_T \) of the original reclustered jet. The resulting large-\( R \) jets are used to tag high-\( p_T \) boosted top quarks in the event. Selected large-\( R \) jets are required to have \( p_T > 300 \) GeV and to have \( |\eta| < 2.0 \). A large-\( R \) jet is tagged as a top candidate if it has a mass above 100 GeV. The mass of the large-\( R \) jets is computed from the four-momentum sum of its constituent small-\( R \) jets, and the mass of the small-\( R \) jets are computed from the four-momentum sum of the topological clusters that make up the jet, which are assumed to be massless. When it is not explicitly stated otherwise, the term “jets” in this paper refers to small-\( R \) jets.

Electron candidates are reconstructed from energy clusters in the electromagnetic calorimeter and inner detector tracks and are required to satisfy a set of “loose” quality criteria [84–86]. They are also required to have \( |\eta| < 2.47 \). Muon candidates are reconstructed from matching tracks in the inner detector and in the muon spectrometer. They are required to meet “medium” quality criteria, as described in Refs. [87,88] and to have \( |\eta| < 2.5 \). All electron and muon candidates must have \( p_T > 20 \) GeV and survive the overlap removal procedure. Signal leptons are chosen from the candidates with the following isolation requirement—the scalar sum of \( p_T \) of additional inner detector tracks in a cone around the lepton track is required to be \(< 5\% \) of the lepton \( p_T \). The angular separation between the lepton and the \( b \) jet ensuing from a semileptonic top quark decay narrows as the \( p_T \) of the top quark increases. This increased collimation is accounted for by varying the radius of the isolation cone as \( \max(0.2, 10/p_T^\text{lep}) \), where \( p_T^\text{lep} \) is the lepton \( p_T \) expressed in GeV. Signal electrons are further required to meet the “tight” quality criteria, while signal muons are required to satisfy the same medium quality criteria as the muon candidates. Electrons (muons) are matched to the primary vertex by requiring the transverse impact parameter \( d_0 \) to satisfy \( |d_0|/\sigma(d_0) < 5 \) (3), where \( \sigma(d_0) \) is the measured uncertainty in \( d_0 \), and the longitudinal impact parameter \( z_0 \) to satisfy \( |z_0\sin\theta| < 0.5 \) mm. In addition, events containing one or more muon candidates with \( |d_0| > 0.2 \) mm and \( |z_0| > 1 \) mm are rejected to suppress cosmic rays.

The overlap removal procedure between muon and jet candidates is designed to remove those muons that are likely to have originated from the decay of hadrons and to retain the overlapping jet. Jets and muons may also appear in close proximity when the jet results from high-\( p_T \) muon bremsstrahlung, and in such cases the jet should be
removed and the muon retained. Such jets are characterized by having very few matching inner detector tracks. Therefore, if the angular distance \( \Delta R \) between a muon and a jet is within \( \min(0.4, 0.04 + 10 \text{ GeV/}p_T) \) of the axis of a jet,\(^4\) the muon is removed only if the jet has \( \geq 3 \) matching inner detector tracks. If the jet has fewer than three matching tracks, the jet is removed and the muon is kept \([89]\). Overlap removal between electron and jet candidates aims to remove jets that are formed primarily from the showering of a prompt electron and to remove electrons that are produced in the decay chains of hadrons. Since electron showers within the cone of a jet contribute to the measured energy of the jet, any overlap between an electron and the jet must be fully resolved. A \( p_T \)-dependent cone for the purpose of this overlap removal is thus impractical. Consequently, any non-\(b\)-tagged jet whose axis lies \( \Delta R < 0.2 \) from an electron is discarded. If the electron is within \( \Delta R = 0.4 \) of the axis of any jet remaining after this initial overlap removal procedure, the jet is retained and the electron is removed. Finally, electron candidates that lie \( \Delta R < 0.01 \) from muon candidates are removed to suppress contributions from muon bremsstrahlung.

The missing transverse momentum (\( E_T^{\text{miss}} \)) in the event is defined as the magnitude of the negative vector sum transverse momentum (\( \vec{p}_T^{\text{miss}} \)) of all selected and calibrated objects in the event, with an extra term added to account for soft energy that is not associated with any of the selected objects. This soft term is calculated from inner detector tracks matched to the primary vertex to make it more resilient to contamination from pileup interactions \([90,91]\).

Corrections derived from data control samples are applied to simulated events to account for differences between data and simulation in the reconstruction efficiencies, momentum scale, and resolution of leptons \([85–87,92]\) and in the efficiency and false positive rate for identifying \( b \) jets \([82,83]\).

V. EVENT SELECTION

The event selection criteria are defined based on kinematic requirements on the objects defined in Sec. IV and on the following event variables.

Two effective mass variables are used, which would typically have much higher values in pair-produced gluino events than in background events. The \( Gtt \) signal regions employ the inclusive effective mass \( m_{\text{eff}}^{\text{incl}} \),

\[
m_{\text{eff}}^{\text{incl}} = \sum_{i} p_T^{\text{jett}} + \sum_{f} p_T^{\ell} + E_T^{\text{miss}},
\]

where the first and second sums are over the signal jets and leptons, respectively. The signal regions for the \( Gbb \) models, for which four high-\( p_T \) \( b \) jets are expected, are defined using \( m_{\text{eff}}^{b\text{-jets}} \),

\[
m_{\text{eff}}^{b\text{-jets}} = \sum_{i \leq 4} p_T^{\text{jett}} + E_T^{\text{miss}},
\]

where the sum is over the four highest-\( p_T \) (leading) signal jets in the event.

In regions with at least one signal lepton, the transverse mass \( m_T \) of the leading signal lepton (\( \ell \)) and \( E_T^{\text{miss}} \) is used to discriminate between the signal and backgrounds from semileptonic \( t\bar{t} \) and \( W + \) jets events,

\[
m_T = \sqrt{2 p_T^{\ell} E_T^{\text{miss}} (1 - \cos[\Delta \phi(p_T^{\text{miss}}, \ell)])}.
\]

Neglecting resolution effects, \( m_T \) is bounded from above by the \( W \) boson mass for these backgrounds and typically has higher values for \( Gtt \) events. Another useful transverse mass variable is \( m_{\text{eff}}^{b\text{-jets}} \), the minimum transverse mass formed by \( E_T^{\text{miss}} \) and any of the three leading \( b \)-tagged jets in the event,

\[
m_{T,\text{min}}^{b\text{-jets}} = \min_{i \leq 3} \left( \sqrt{2 p_T^{b\text{-jet}} E_T^{\text{miss}} (1 - \cos[\Delta \phi(p_T^{\text{miss}}, b\text{-jet})])} \right).
\]

It is bounded below the top quark mass for semileptonic \( t\bar{t} \) events while peaking at higher values for \( Gbb \) and \( Gtt \) events.

The signal regions require either zero or at least one lepton. The requirement of a signal lepton, with the additional requirements on jets, \( E_T^{\text{miss}} \), and event variables described in Sec. VA, render the multijet background negligible for the \( \geq 1 \)-lepton signal regions. For the 0-lepton signal regions, the minimum azimuthal angle between \( p_T^{\text{miss}} \) and the leading four small-\( R \) jets in the event, \( \Delta \phi_{\text{min}}^{4\text{j}} \), is required to be greater than 0.4,

\[
\Delta \phi_{\text{min}}^{4\text{j}} = \min(|\phi_{\text{jet}1} - \phi_{p_T^{\text{miss}}}|, \ldots, |\phi_{\text{jet}4} - \phi_{p_T^{\text{miss}}}|) > 0.4.
\]

This requirement ensures that the multijet background, which can produce large \( E_T^{\text{miss}} \) if containing poorly measured jets or neutrinos emitted close to the axis of a jet, is also negligible in the 0-lepton signal regions (along with the other requirements on jets, \( E_T^{\text{miss}} \), and event variables described in Sec. VA).

Figure 2 shows the kinematic distributions of \( E_T^{\text{miss}} \), \( m_{\text{eff}}^{\text{incl}} \), \( m_{\text{eff}}^{b\text{-jets}} \), and \( m_T \) for a preselection that requires \( E_T^{\text{miss}} > 200 \text{ GeV} \), at least four signal jets of which at least three must be \( b \) tagged, and \( \Delta \phi_{\text{min}}^{4\text{j}} > 0.4 \).

Figure 3 shows the multiplicity of signal jets, \( b \)-tagged signal jets, top-tagged large-\( R \) jets, and signal leptons in the preselection.
Good agreement between data and simulation is observed. Example signal models with enhanced cross sections are overlaid for comparison.

A. Signal regions

The signal regions are designed by optimizing the expected signal discovery reach for the 2015 data set. They are defined in the leftmost column of Tables II, III, and IV for the Gbb, Gtt 0-lepton, and Gtt 1-lepton channels, respectively, and are discussed below. These tables also contain the definition of the control regions used to normalize the $t\bar{t}$ background, discussed in Sec. VB, and the validation regions used to cross-check the background estimate and which are discussed in Section VC. The following region nomenclature is used in the remainder of the paper. Signal, control, and validation region names start with the prefix “SR,” “CR,” and “VR,” respectively, and with the type of validation region specified for the Gtt validation regions. The name of the region is completed by the type of model targeted and a letter corresponding to the level of mass splitting between the gluino and the LSP. For example the validation region that cross-checks the background prediction. All backgrounds (including $t\bar{t}$) are normalized using the best available theoretical calculation described in Sec. III. The background category “Others” includes $t\bar{t}h$, $tt\bar{t}$, and diboson events. Example signal models with cross sections enhanced by a factor of 100 are overlaid for comparison.

FIG. 2. Distributions of kinematic variables in the preselection region described in the text: (a) $E_T^{miss}$, (b) $m_{T,miss}$, (c) $m_{T,min}$, and (d) $m_T$ (for preselected events with at least one signal lepton). The statistical and experimental systematic uncertainties are included in the uncertainty band, where the systematic uncertainties are defined in Sec. VI. The lower part of each figure shows the ratio of data to the background prediction. All backgrounds (including $t\bar{t}$) are normalized using the best available theoretical calculation described in Sec. III. The background category “Others” includes $t\bar{t}h$, $tt\bar{t}$, and diboson events. Example signal models with cross sections enhanced by a factor of 100 are overlaid for comparison.
FIG. 3. Distributions of the number of (a) signal jets, (b) $b$-tagged jets, (c) top-tagged large-$R$ jets, and (d) signal leptons in the preselection region described in the text. The statistical and experimental systematic uncertainties are included in the uncertainty band, where the systematic uncertainties are defined in Sec. VI. The lower part of each figure shows the ratio of data to the background prediction. All backgrounds (including $t$) are normalized using the best available theoretical calculation described in Sec. III. The background category “Others” includes $i ith, iiiit, and diboson events. Example signal models with cross sections enhanced by a factor of 100 are overlaid for comparison.

and the LSP, respectively. All regions feature stringent cuts on $E_{T}^{miss}$, $m^{4j}$, and the jet transverse momentum $p_{T}^{jet}$.

The experimental signature for the Gtt model is characterized by several high-$p_{T}$ jets of which four are $b$ jets, large $E_{T}^{miss}$, and potentially leptons [Fig. 1(b)]. The Gtt signal regions are classified into regions with a signal lepton veto (0-lepton channel) and regions with at least one signal lepton (1-lepton channel). The Gtt 0-lepton signal regions are defined in the leftmost column of Table III. In all Gtt 0-lepton signal regions at least eight signal jets, $\Delta\phi^{4j} > 0.4$ and $m^{b-jets} > 80$ GeV are required. Three Gtt 0-lepton signal regions are defined to cover Gtt models with decreasing mass splitting between the gluino and the LSP: A ($\gtrsim 1$ TeV), B (between $\approx 200$ GeV and $\approx 1$ TeV), and C ($\lesssim 200$ GeV). In the large and moderate mass splitting scenarios, the top quarks tend to have a large $p_{T}$, and at least one top-tagged large-$R$ jet is required ($N^{top} \geq 1$). The requirements on $E_{T}^{miss}$ and $m_{eff}^{incl}$ decrease with the mass splitting between the gluino and the LSP. However, the required number of $b$-tagged jets $N^{b-jet}$ is tightened to four for the lower mass splitting regions B and C in order to maintain a high background rejection despite the softer signal kinematics.

The Gtt 1-lepton signal regions are defined in the leftmost column of Table IV. Two signal regions A and B are defined to cover Gtt models with decreasing mass difference between the gluino and the LSP. In all signal regions at least one signal lepton, at least six signal jets ($p_{T}^{jet} > 30$ GeV), and $m_{T} > 150$ GeV are required. Region A has tighter requirements on $m_{eff}^{incl}$ ($m_{eff}^{incl} > 1100$ GeV) and the number of top-tagged large-$R$ jets ($N^{top} \geq 1$). Region B has a softer requirement on $m_{eff}^{incl}$ than region A, but it features a tighter cut on $E_{T}^{miss}$ to achieve a satisfactory background rejection without requiring a top-tagged large-$R$ jet.
Criteria common to all regions of the same type

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<th>Variable</th>
<th>Signal region</th>
<th>Control region</th>
<th>Validation region</th>
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<td>= 0</td>
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<tr>
<td>$N_{\text{Signal Lepton}}$</td>
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<td>...</td>
</tr>
<tr>
<td>$\Delta \phi_{\text{min}}^{ij}$</td>
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<td>...</td>
<td>&gt; 0.4</td>
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<tr>
<td>$m_{T,\text{min}}^{b-\text{jets}}$</td>
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<td>...</td>
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<tr>
<td>$m_T$</td>
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<td>&lt;150</td>
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Region A (Large mass splitting)

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</tr>
<tr>
<td>$m_{\text{eff}}^{ij}$</td>
<td>&gt; 1600</td>
<td>&gt; 1200</td>
<td>&lt;1400</td>
</tr>
</tbody>
</table>

Region B (Moderate mass splitting)

<table>
<thead>
<tr>
<th>Variable</th>
<th>Signal region</th>
<th>Control region</th>
<th>Validation region</th>
</tr>
</thead>
<tbody>
<tr>
<td>$p_T^{\text{jet}}$</td>
<td>&gt; 90</td>
<td>&gt; 90</td>
<td>&gt; 90</td>
</tr>
<tr>
<td>$E_T^{\text{miss}}$</td>
<td>&gt; 450</td>
<td>&gt; 300</td>
<td>&gt; 300</td>
</tr>
<tr>
<td>$m_{\text{eff}}^{ij}$</td>
<td>&gt; 1400</td>
<td>&gt; 1000</td>
<td>&lt;1400</td>
</tr>
</tbody>
</table>

Region C (Small mass splitting)

<table>
<thead>
<tr>
<th>Variable</th>
<th>Signal region</th>
<th>Control region</th>
<th>Validation region</th>
</tr>
</thead>
<tbody>
<tr>
<td>$p_T^{\text{jet}}$</td>
<td>&gt; 30</td>
<td>&gt; 30</td>
<td>&gt; 30</td>
</tr>
<tr>
<td>$E_T^{\text{miss}}$</td>
<td>&gt; 500</td>
<td>&gt; 400</td>
<td>&gt; 400</td>
</tr>
<tr>
<td>$m_{\text{eff}}^{ij}$</td>
<td>&gt; 1400</td>
<td>&gt; 1200</td>
<td>&lt;1400</td>
</tr>
</tbody>
</table>

### B. Background estimation and $\bar{t}t$ control regions

The largest background in all signal regions is $\bar{t}t$ produced with additional high-$p_T$ jets. The other relevant backgrounds are $t\bar{t}W$, $t\bar{t}Z$, $t\bar{t}t\bar{t}$, $t\bar{t}h$, single-top, $W +$ jets, $Z +$ jets, and diboson events. All of these smaller backgrounds are estimated with the simulated event samples normalized to the best available theory calculations described in Sec. III. The multijet background is estimated to be negligible in all regions.

For each signal region, the $\bar{t}t$ background is normalized in a dedicated control region. The $\bar{t}t$ normalization factor is

<table>
<thead>
<tr>
<th>Variable</th>
<th>Signal region</th>
<th>Control region</th>
<th>Validation region</th>
</tr>
</thead>
<tbody>
<tr>
<td>$N_{\text{Candidate Lepton}}$</td>
<td>= 0</td>
<td>= 1</td>
<td>= 1</td>
</tr>
<tr>
<td>$N_{\text{Signal Lepton}}$</td>
<td>...</td>
<td>= 1</td>
<td>...</td>
</tr>
<tr>
<td>$\Delta \phi_{\text{min}}^{ij}$</td>
<td>&gt; 0.4</td>
<td>...</td>
<td>&gt; 0.4</td>
</tr>
<tr>
<td>$N_{\text{jet}}$</td>
<td>≥8</td>
<td>≥7</td>
<td>≥7</td>
</tr>
<tr>
<td>$m_{T,\text{min}}^{b-\text{jets}}$</td>
<td>≥80</td>
<td>...</td>
<td>&gt;80</td>
</tr>
<tr>
<td>$m_T$</td>
<td>...</td>
<td>&lt;150</td>
<td>&lt;150</td>
</tr>
</tbody>
</table>

Region A (Large mass splitting)

<table>
<thead>
<tr>
<th>Variable</th>
<th>Signal region</th>
<th>Control region</th>
<th>Validation region</th>
</tr>
</thead>
<tbody>
<tr>
<td>$E_T^{\text{miss}}$</td>
<td>&gt; 400</td>
<td>&gt; 250</td>
<td>&gt; 200</td>
</tr>
<tr>
<td>$m_{\text{eff}}^{\text{incl}}$</td>
<td>&gt; 1700</td>
<td>&gt; 1350</td>
<td>&gt; 1350</td>
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<tr>
<td>$N_{\text{b-tag}}^{\text{top}}$</td>
<td>≥3</td>
<td>≥3</td>
<td>≥3</td>
</tr>
</tbody>
</table>

Region B (Moderate mass splitting)

<table>
<thead>
<tr>
<th>Variable</th>
<th>Signal region</th>
<th>Control region</th>
<th>Validation region</th>
</tr>
</thead>
<tbody>
<tr>
<td>$E_T^{\text{miss}}$</td>
<td>&gt; 350</td>
<td>&gt; 200</td>
<td>&gt; 200</td>
</tr>
<tr>
<td>$m_{\text{eff}}^{\text{incl}}$</td>
<td>&gt; 1250</td>
<td>&gt; 1000</td>
<td>&gt; 1000</td>
</tr>
<tr>
<td>$N_{\text{b-tag}}^{\text{top}}$</td>
<td>≥4</td>
<td>≥4</td>
<td>≥4</td>
</tr>
</tbody>
</table>

Region C (Small mass splitting)

<table>
<thead>
<tr>
<th>Variable</th>
<th>Signal region</th>
<th>Control region</th>
<th>Validation region</th>
</tr>
</thead>
<tbody>
<tr>
<td>$E_T^{\text{miss}}$</td>
<td>&gt; 350</td>
<td>&gt; 200</td>
<td>&gt; 200</td>
</tr>
<tr>
<td>$m_{\text{eff}}^{\text{incl}}$</td>
<td>&gt; 1250</td>
<td>&gt; 1000</td>
<td>&gt; 1250</td>
</tr>
<tr>
<td>$N_{\text{b-tag}}^{\text{top}}$</td>
<td>≥4</td>
<td>≥4</td>
<td>≥4</td>
</tr>
</tbody>
</table>
required for the total predicted yield to match the data in the control region is used to normalize the $t\bar{t}$ background in the signal region. The control regions are designed to be dominated by $t\bar{t}$ events and to have negligible signal contamination, while being kinematically as close as possible to the corresponding signal region. The latter requirement minimizes the systematic uncertainties associated with extrapolating the normalization factors from the control to the signal regions.

The definitions of the control regions are shown next to the signal regions in Tables II, III, and IV for the Gbb, Gtt 0-lepton, and Gtt 1-lepton channels, respectively. In both the Gbb and Gtt 0-lepton channels, exactly one signal lepton is required. This is motivated by background composition studies using simulated events which show that semileptonic $t\bar{t}$ events, for which the lepton is outside the acceptance or is a hadronically decaying $\tau$ lepton, dominate the $t\bar{t}$ yield in the signal regions. An upper cut on $m_T$ is then applied to ensure orthogonality with the Gtt 1-lepton signal regions and to suppress signal contamination. The jet multiplicity requirement is reduced to seven jets in the Gtt 0-lepton control regions (from eight jets in the signal regions), to accept more events and to obtain a number of jets from top quark decay and parton shower similar to that in the signal region. Approximately 40%–60% of the signal event rates contain a hadronically decaying $\tau$ lepton that is counted as a jet. Orthogonality between Gtt 0-lepton and Gtt 1-lepton control regions is ensured by requiring exactly six jets in the Gtt 1-lepton control regions (as opposed to the requirement of at least six jets in the signal regions). For all Gbb and Gtt 0-lepton control regions, the number of $b$-tagged jets and top-tagged large-$R$ jets is consistent with the signal region. The requirements on $E_T^{\text{miss}}$ and $m_{\text{jet}}$ are, however, relaxed in the control regions to achieve a sufficiently large $t\bar{t}$ yield and small signal contamination ($\lesssim 15\%$). The Gtt 1-lepton control regions are defined by inverting the $m_T$ cut and removing the $m_{T,\text{min}}^{b\text{-jets}}$ requirement. All other requirements are exactly the same as for the signal regions.

### C. Validation regions

Validation regions are defined to cross-check the background prediction in regions that are kinematically close to the signal regions but yet have a small signal contamination. They are designed primarily to cross-check the assumption that the $t\bar{t}$ normalization extracted from the control regions can be accurately extrapolated to the signal regions. Their requirements are shown in the rightmost column(s) of Tables II, III, and IV for the Gbb, Gtt 0-lepton, and Gtt 1-lepton channels, respectively. Their signal contamination is less than approximately 30% for the majority of Gbb and Gtt model points not excluded in Run 1.

One validation region per signal region is defined for the Gbb model. They feature the same requirements as their corresponding signal region except that upper cuts are applied on $m_{T,\text{min}}^{b\text{-jets}}$ and $m_{\text{jet}}$ to reduce signal contamination and ensure orthogonality with the signal regions. In addition, the requirement on $E_T^{\text{miss}}$ is relaxed to obtain a sufficient $t\bar{t}$ yield.

For the Gtt 0-lepton channel, two validation regions per signal region are defined, one requiring exactly one signal lepton (VR1L) and one with a signal lepton veto (VR0L). The regions VR1L have exactly the same criteria as their corresponding control regions except that upper cuts are applied on $m_{T,\text{min}}^{b\text{-jets}}$ and $m_{\text{jet}}$. For such a validation region, the requirement minimizes the systematic uncertainties as associated with extrapolating the normalization factors from the control to the signal regions.

### Table IV. Definitions of the Gtt 1-lepton signal, control, and validation regions. The unit of all kinematic variables is GeV. The jet $p_T$ requirement is also applied to $b$-tagged jets.

<table>
<thead>
<tr>
<th>Criteria common to all regions of the same type</th>
<th>Variable</th>
<th>Signal region</th>
<th>Control region</th>
<th>VR-$m_T$</th>
<th>VR-$m_{T,\text{min}}^{b\text{-jets}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$m_T$</td>
<td>&gt; 150</td>
<td>&lt; 150</td>
<td>&gt; 150</td>
<td>&lt; 150</td>
<td></td>
</tr>
<tr>
<td>$N^{b\text{-tag}}$</td>
<td>$\geq 6$</td>
<td>= 6</td>
<td>$\geq 5$</td>
<td>$\geq 6$</td>
<td></td>
</tr>
<tr>
<td>$N^{b\text{-tag}}$</td>
<td>$\geq 3$</td>
<td>$\geq 3$</td>
<td>= 3</td>
<td>= 3</td>
<td></td>
</tr>
<tr>
<td>Region A (Large mass splitting)</td>
<td>$E_T^{\text{miss}}$</td>
<td>&gt; 200</td>
<td>&gt; 200</td>
<td>&gt; 200</td>
<td>&gt; 200</td>
</tr>
<tr>
<td></td>
<td>$m_{\text{incl}}^{b\text{-tag}}$</td>
<td>&gt; 1100</td>
<td>&gt; 1100</td>
<td>&gt; 600</td>
<td>&gt; 600</td>
</tr>
<tr>
<td></td>
<td>$m_T^{c\text{-jets}}$</td>
<td>&gt; 160</td>
<td>...</td>
<td>&lt; 160</td>
<td>&gt; 140</td>
</tr>
<tr>
<td></td>
<td>$N^{\text{top}}$</td>
<td>$\geq 1$</td>
<td>$\geq 1$</td>
<td>$\geq 1$</td>
<td>$\geq 1$</td>
</tr>
<tr>
<td>Region B (Moderate to small mass splitting)</td>
<td>$E_T^{\text{miss}}$</td>
<td>&gt; 300</td>
<td>&gt; 300</td>
<td>&gt; 200</td>
<td>&gt; 200</td>
</tr>
<tr>
<td></td>
<td>$m_{\text{incl}}^{b\text{-tag}}$</td>
<td>&gt; 900</td>
<td>&gt; 900</td>
<td>&gt; 600</td>
<td>&gt; 600</td>
</tr>
<tr>
<td></td>
<td>$m_T^{c\text{-jets}}$</td>
<td>&gt; 160</td>
<td>...</td>
<td>&lt; 160</td>
<td>&gt; 160</td>
</tr>
</tbody>
</table>
While the theoretical uncertainties in the heavy-flavor fraction of the additional jets in the $t\bar{t} + \text{jets}$ events (i.e., $t\bar{t} + b\bar{b}$ and $t\bar{t} + c\bar{c}$) are large, they affect signal, control, and the 1-lepton validation regions in a similar way, and are thus largely canceled in the semi-data-driven $t\bar{t}$ normalization based on the observed control region yields.

The VR0L regions have similar requirements on their corresponding signal regions except that the requirements on $E_T^{\text{miss}}$, $m^{\text{recoil}}_T$, and the number of $b$-tagged jets are loosened to achieve sufficient event yields. Furthermore, the criterion $m^{b\text{-jets}}_{T,\text{min}} < 80 \text{ GeV}$ is applied to all VR0L regions to ensure orthogonality with the signal regions. The regions VR0L test the extrapolation of the $t\bar{t}$ normalization from a 1-lepton to a 0-lepton region. Simulation studies show that the VR0L regions have a composition of semileptonic $t\bar{t}$ events (in particular of hadronically decaying $\tau$ leptons) similar to that in the signal regions, while the control and VR1L regions are by construction dominated by semileptonic $t\bar{t}$ events with a muon or an electron. Two requirements are different between Gtt 1-lepton control regions and their corresponding signal regions: the requirement on $m^{b\text{-jets}}_{T,\text{min}}$ (absent in the control regions) and the requirement on $m_T$ (inverted in the control regions). Therefore, two validation regions per signal region are defined for the Gtt 1-lepton channel, VR-$m_T$ and VR-$m^{b\text{-jets}}_{T,\text{min}}$, which respectively test, one at a time, the extrapolations over $m_T$ and $m^{b\text{-jets}}_{T,\text{min}}$. Exactly three $b$-tagged jets are required for all 1-lepton validation regions to limit the signal contamination and to be close to the signal regions. For the VR-$m_T$ regions, the same requirement $m_T > 150 \text{ GeV}$ as in the signal region is applied but the criterion on $m^{b\text{-jets}}_{T,\text{min}}$ is inverted. Other requirements are relaxed to achieve sufficiently large background yields and small signal contamination. For the VR-$m^{b\text{-jets}}_{T,\text{min}}$ regions, the signal region requirement on $m^{b\text{-jets}}_{T,\text{min}}$ is applied (slightly loosened to 140 GeV instead of 160 GeV in region A) and the criterion on $m_T$ is inverted. Again, other requirements are generally relaxed. Simulation studies show that $t\bar{t}$ dilepton events dominate in the signal regions, in particular due to the requirement on $m_T$, while semileptonic $t\bar{t}$ events dominate in the control regions. This extrapolation is cross-checked by the VR-$m_T$ regions, which have a $t\bar{t}$ dileptonic fraction similar to that in the signal regions.

VI. SYSTEMATIC UNCERTAINTIES

The largest sources of detector-related systematic uncertainties in this analysis relate to the JES, jet energy resolution (JER), and the $b$-tagging efficiencies and mistagging rates. The JES uncertainties are obtained by extrapolating the uncertainties derived from $\sqrt{s} = 8 \text{ TeV}$ data and simulations to $\sqrt{s} = 13 \text{ TeV}$ [77]. The uncertainties in the energy scale of the small-$R$ jets are propagated to the reclustered large-$R$ jets, which use them as inputs. The JES uncertainties are especially important in the Gtt signal regions, since these regions require high jet multiplicities. The impact of these uncertainties on the expected background yields in these regions is between 10% and 25%. Uncertainties in the JER are similarly derived from dijet asymmetry measurements in Run 1 data and extrapolated to $\sqrt{s} = 13 \text{ TeV}$. The impact of the JER uncertainties on the background yields are in the range of 1%–10%.

Uncertainties in the measured $b$-tagging efficiencies and mistagging rates are the subleading sources of experimental uncertainties in the Gtt 1-lepton signal regions and the leading source in the Gtt 0-lepton and Gbb regions. Uncertainties measured in $\sqrt{s} = 8 \text{ TeV}$ data are extrapolated to $\sqrt{s} = 13 \text{ TeV}$, with the addition of the new IBL system in Run 2 taken into account. Uncertainties for jet $p_T$ above 300 GeV are estimated using simulated events. The impact of the $b$-tagging uncertainties on the expected background yields in the Gbb and Gtt 0-lepton signal regions is around 22%–30%, and around 15% in the Gtt 1-lepton signal regions.

The uncertainties associated with lepton reconstruction and energy measurements have very small impact on the final results. All lepton and jet measurement uncertainties are propagated to the calculation of $E_T^{\text{miss}}$, and additional uncertainties are included in the scale and resolution of the soft term. The overall impact of the $E_T^{\text{miss}}$ soft term uncertainties on the expected background yields is 5% or less.

Uncertainties in the modeling of the $t\bar{t}$ background are evaluated using additional samples varied by each systematic uncertainty. Hadronization and parton showering uncertainties are estimated using a sample generated with POWHEG and showered by HERWIG++ v2.7.1 [49] with the UEESE5 underlying-event tune [60]. Systematic uncertainties in the modeling of initial- and final-state radiation are explored with two alternative settings of POWHEG, both of which are showered by PYTHIA v6.428 as for the nominal sample. The first of these uses the PERUGIA2012 radHi tune and has the renormalization and factorization scales set to twice the nominal value, resulting in more radiation in the final state. It also has $h_{\text{damp}}$ set to $2m_{\text{top}}$. The second sample, using the PERUGIA2012 radLo tune, has $h_{\text{damp}} = m_{\text{top}}$, and the renormalization and factorization scales are set to half of their nominal values, resulting in less radiation in the event. In each case, the uncertainty is taken as the deviation in the expected yield of $t\bar{t}$ background with respect to the nominal sample. The uncertainty due to the choice of generator is estimated by comparing the expected yields obtained using a $t\bar{t}$ sample generated with MADGRAPH5_aMC@NLO, and one that is generated with POWHEG. Both of these samples are showered with HERWIG++ v2.7.1. Finally, a 30% uncertainty is assigned to the cross section of $t\bar{t}$ events with additional heavy-flavor
jets in the final state, in accordance with the results of the ATLAS measurement of this cross section at $\sqrt{s} = 8$ TeV [93]. Uncertainties in single-top and $W/Z + \text{jets}$ background processes are similarly estimated by comparisons between the nominal sample and samples with different generators, showering models, and radiation tunes. An additional 5% uncertainty is included in the cross section of single-top processes [94]. A 50% constant uncertainty is assigned to each of the remaining small backgrounds. The impact of the modeling uncertainties for the smaller backgrounds on these yields is consistently below 10% in all signal regions. The uncertainties in the cross sections of signal processes are determined from an envelope of different cross-section predictions, as described in Sec. III.

The cumulative impact of the systematic uncertainties listed above on the background yields ranges between 23% and 63%, depending on the signal region. The typical impact on the signal yields is in the range 10%–30%.

VII. RESULTS

The SM background expectation is determined separately in each signal region with a profile likelihood fit [95], referred to as a background-only fit. The fit uses as a constraint the observed event yield in the associated control region to adjust the $t\bar{t}$ normalization, assuming that a signal does not contribute to this yield, and applies that normalization factor to the number $t\bar{t}$ events predicted by simulation in the signal region. The numbers of observed and predicted events in each control region are described by Poisson probability density functions. The systematic uncertainties in the expected values are included in the fit as nuisance parameters. They are constrained by Gaussian distributions with widths corresponding to the sizes of the uncertainties and are treated as correlated, when appropriate, between the various regions. The product of the various probability density functions forms the likelihood, which the fit maximizes by adjusting the $t\bar{t}$ normalization and the nuisance parameters within their constraints. The inputs to the fit for each signal region are the number of events observed in its associated control region and the number of events predicted by simulation in each region for all background processes.

Figure 4 shows the results of the background-only fit to the control regions, extrapolated to the validation regions. The number of events predicted by the background-only fit is compared to the data in the upper panel. The pull, defined by the difference between the observed number of events ($n_{\text{obs}}$) and the predicted background yield ($n_{\text{pred}}$) divided by the total uncertainty ($\sigma_{\text{tot}}$), is shown for each region in the lower panel. No evidence of significant background mis-modeling is observed in the validation regions. There is a certain tendency for the predicted background to be above the data, in particular for the Gtt-0L validation regions, but the results in the validation regions of a given channel are not independent. The validation and control regions of different mass splittings can overlap, with the overlap fraction ranging from approximately 30% to 70% for Gtt-0L. Furthermore, the uncertainties in the predicted

![Figure 4](attachment:image.png)

**FIG. 4.** Results of the likelihood fit extrapolated to the validation regions. The $t\bar{t}$ normalization is obtained from the fit to the control regions. The upper panel shows the observed number of events and the predicted background yield. The background category “Others” includes $t\bar{t}h$, $t\bar{t}t\bar{t}$, and diboson events. The lower panel shows the pulls in each validation region.
yield are dominated by the same (correlated) systematic uncertainties. The background is dominated by \( \bar{t}t \) events in all Gbb and Gtt signal regions. The subdominant contributions in the Gbb and Gtt 0-lepton signal regions are \( Z(\nu\bar{\nu}) + \text{jets} \) and \( W(\nu\bar{\nu}) + \text{jets} \) events, where for \( W + \text{jets} \) events the lepton is a nonidentified electron or muon or is a hadronically decaying \( \tau \) lepton. In the Gtt 1-lepton signal regions, the subdominant backgrounds are single-top, \( t\bar{t}W \) and \( t\bar{t}Z \).

Figure 6 shows the \( E_T^{\text{miss}} \) distributions in data and simulated samples for SR-Gbb-B, SR-Gtt-0L-C, and SR-Gtt-1L-A, after relaxing the \( E_T^{\text{miss}} \) threshold to 200 GeV.

### TABLE V. Results of the likelihood fit extrapolated to the Gbb signal regions. The uncertainties shown include all systematic uncertainties. The data in the signal regions are not included in the fit. The row “MC-only background prediction” provides the total background prediction when the \( \bar{t}t \) normalization is obtained from a theoretical calculation \[61\]. The \( \bar{t}t \) normalization factor \( \mu_\bar{t} \) obtained from the corresponding \( \bar{t}t \) control region is also provided. The background category “Others” includes \( t\bar{t}h, t\bar{t}A, \) and diboson events. Expected yields for two example Gbb models are also shown.

<table>
<thead>
<tr>
<th></th>
<th>SR-Gbb-A</th>
<th>SR-Gbb-B</th>
<th>SR-Gbb-C</th>
</tr>
</thead>
<tbody>
<tr>
<td>Observed events</td>
<td>0</td>
<td>0</td>
<td>5</td>
</tr>
<tr>
<td>Fitted background events</td>
<td>1.3 ± 0.4</td>
<td>1.5 ± 0.6</td>
<td>7.6 ± 1.7</td>
</tr>
<tr>
<td>( \bar{t}t )</td>
<td>0.63 ± 0.30</td>
<td>0.9 ± 0.5</td>
<td>4.3 ± 1.5</td>
</tr>
<tr>
<td>( Z + \text{jets} )</td>
<td>0.23 ± 0.08</td>
<td>0.23 ± 0.09</td>
<td>1.2 ± 0.5</td>
</tr>
<tr>
<td>( W + \text{jets} )</td>
<td>0.17 ± 0.06</td>
<td>0.13 ± 0.05</td>
<td>0.82 ± 0.28</td>
</tr>
<tr>
<td>Single-top</td>
<td>0.25 ± 0.14</td>
<td>0.15 ± 0.14</td>
<td>0.65 ± 0.33</td>
</tr>
<tr>
<td>( t\bar{t}W/Z )</td>
<td>&lt;0.1</td>
<td>&lt;0.1</td>
<td>0.22 ± 0.12</td>
</tr>
<tr>
<td>Others</td>
<td>&lt;0.1</td>
<td>&lt;0.1</td>
<td>0.39 ± 0.22</td>
</tr>
<tr>
<td>MC-only background prediction</td>
<td>1.7</td>
<td>1.5</td>
<td>6.7</td>
</tr>
<tr>
<td>( \mu_\bar{t} )</td>
<td>0.64 ± 0.33</td>
<td>1.0 ± 0.4</td>
<td>1.2 ± 0.4</td>
</tr>
<tr>
<td>Gbb ( (m_\bar{g} = 1700 \text{ GeV}, m_{\chi_1^0} = 200 \text{ GeV}) )</td>
<td>3.8</td>
<td>3.5</td>
<td>4.2</td>
</tr>
<tr>
<td>Gbb ( (m_\bar{g} = 1400 \text{ GeV}, m_{\chi_1^0} = 800 \text{ GeV}) )</td>
<td>5.3</td>
<td>7.2</td>
<td>10.5</td>
</tr>
</tbody>
</table>

### TABLE VI. Results of the likelihood fit extrapolated to the Gtt 0-lepton signal regions. The uncertainties shown include all systematic uncertainties. The data in the signal regions are not included in the fit. The row “MC-only background prediction” provides the total background prediction when the \( \bar{t}t \) normalization is obtained from a theoretical calculation \[61\]. The \( \bar{t}t \) normalization factor \( \mu_\bar{t} \) obtained from the corresponding \( \bar{t}t \) control region is also provided. The category “Others” includes \( t\bar{t}h, t\bar{t}A, \) and diboson events. Expected yields for two example Gtt models are also shown.

<table>
<thead>
<tr>
<th></th>
<th>SR-Gtt-0L-A</th>
<th>SR-Gtt-0L-B</th>
<th>SR-Gtt-0L-C</th>
</tr>
</thead>
<tbody>
<tr>
<td>Observed events</td>
<td>1</td>
<td>1</td>
<td>1</td>
</tr>
<tr>
<td>Fitted background events</td>
<td>2.1 ± 0.5</td>
<td>2.9 ± 1.8</td>
<td>3.4 ± 1.8</td>
</tr>
<tr>
<td>( \bar{t}t )</td>
<td>1.4 ± 0.4</td>
<td>2.4 ± 1.7</td>
<td>2.6 ± 1.8</td>
</tr>
<tr>
<td>( Z + \text{jets} )</td>
<td>0.22 ± 0.09</td>
<td>0.11 ± 0.06</td>
<td>0.14 ± 0.07</td>
</tr>
<tr>
<td>( W + \text{jets} )</td>
<td>0.19 ± 0.08</td>
<td>0.14 ± 0.06</td>
<td>0.18 ± 0.08</td>
</tr>
<tr>
<td>Single-top</td>
<td>0.17 ± 0.17</td>
<td>0.14 ± 0.13</td>
<td>0.17 ± 0.15</td>
</tr>
<tr>
<td>( t\bar{t}W/Z )</td>
<td>&lt;0.1</td>
<td>&lt;0.1</td>
<td>0.10 ± 0.07</td>
</tr>
<tr>
<td>Others</td>
<td>&lt;0.1</td>
<td>&lt;0.1</td>
<td>0.20 ± 0.17</td>
</tr>
<tr>
<td>MC-only background prediction</td>
<td>1.8</td>
<td>1.9</td>
<td>2.5</td>
</tr>
<tr>
<td>( \mu_\bar{t} )</td>
<td>1.3 ± 0.4</td>
<td>1.8 ± 0.8</td>
<td>1.5 ± 0.7</td>
</tr>
<tr>
<td>Gtt ( (m_\bar{g} = 1600 \text{ GeV}, m_{\chi_1^0} = 200 \text{ GeV}) )</td>
<td>3.8</td>
<td>2.8</td>
<td>2.9</td>
</tr>
<tr>
<td>Gtt ( (m_\bar{g} = 1400 \text{ GeV}, m_{\chi_1^0} = 800 \text{ GeV}) )</td>
<td>2.0</td>
<td>3.7</td>
<td>4.1</td>
</tr>
</tbody>
</table>
TABLE VII. Results of the likelihood fit extrapolated to the Gtt 1-lepton signal regions. The uncertainties shown include all systematic uncertainties. The data in the signal regions are not included in the fit. The row “MC-only background prediction” provides the total background prediction when the $\bar{t}\bar{t}$ normalization is obtained from a theoretical calculation [61]. The $\bar{t}\bar{t}$ normalization factor $\mu_{\bar{t}\bar{t}}$ obtained from the corresponding $\bar{t}\bar{t}$ control region is also provided. The category “Others” includes $\bar{t}h$, $\bar{t}\bar{t}i$, and diboson events. Expected yields for two example Gtt models are also shown.

<table>
<thead>
<tr>
<th></th>
<th>SR-Gtt-1L-A</th>
<th>SR-Gtt-1L-B</th>
</tr>
</thead>
<tbody>
<tr>
<td>Observed events</td>
<td>2</td>
<td>0</td>
</tr>
<tr>
<td>Fitted background events</td>
<td>1.2 ± 0.6</td>
<td>1.2 ± 0.8</td>
</tr>
<tr>
<td>$\bar{t}\bar{t}$</td>
<td>0.8 ± 0.6</td>
<td>0.8 ± 0.7</td>
</tr>
<tr>
<td>$Z +$ jets</td>
<td>...</td>
<td>&lt;0.1</td>
</tr>
<tr>
<td>$W +$ jets</td>
<td>&lt;0.1</td>
<td>&lt;0.1</td>
</tr>
<tr>
<td>Single-top</td>
<td>0.18 ± 0.14</td>
<td>0.14 ± 0.12</td>
</tr>
<tr>
<td>$\bar{t}\tilde{W}/Z$</td>
<td>0.14 ± 0.08</td>
<td>0.15 ± 0.09</td>
</tr>
<tr>
<td>Others</td>
<td>&lt;0.1</td>
<td>&lt;0.1</td>
</tr>
<tr>
<td>MC-only background prediction</td>
<td>1.3</td>
<td>1.2</td>
</tr>
<tr>
<td>$\mu_{\bar{t}\bar{t}}$</td>
<td>0.86 ± 0.28</td>
<td>1.0 ± 0.4</td>
</tr>
<tr>
<td>Gtt ($m_\tilde{g} = 1600$ GeV, $m_\tilde{\chi}_1^0 = 200$ GeV)</td>
<td>3.4</td>
<td>3.3</td>
</tr>
<tr>
<td>Gtt ($m_\tilde{g} = 1400$ GeV, $m_\tilde{\chi}_1^0 = 800$ GeV)</td>
<td>3.8</td>
<td>4.6</td>
</tr>
</tbody>
</table>

VIII. INTERPRETATION

Since no significant excess over the expected background from SM processes is observed, the data are used to derive one-sided upper limits at 95% C.L. Model-independent limits on the number of beyond-the-SM (BSM) events for each signal region are derived with pseudoexperiments using the CL_s prescription [96]. They can be translated into upper limits on the visible BSM cross section ($\sigma_{\text{vis}}$), where $\sigma_{\text{vis}}$ is defined as the product of acceptance, reconstruction efficiency, and production cross section. The results are given in Table VIII, where the observed ($S_{\text{obs}}^{95}$) and expected ($S_{\text{exp}}^{95}$) 95% C.L. upper limits on the number of BSM events are also provided.

The measurement is used to place exclusion limits on gluino and LSP masses in the Gbb and Gtt simplified models. The results are obtained using the CL_s prescription in the asymptotic approximation [97]. The signal contamination in the control regions and the experimental
systematic uncertainties in the signal are taken into account for this calculation. For the Gbb models, the results are obtained from the Gbb signal region with the best expected sensitivity at each point of the parameter space of each model. For the Gtt models, the 0- and 1-lepton channels both contribute to the sensitivity, and they are combined in a simultaneous fit to enhance the sensitivity of the analysis. This is performed by considering all possible permutations between the three Gtt 0-lepton and the two Gtt 1-lepton signal regions for each point of the parameter space, and the best expected combination is used. The 95% C.L. observed and expected exclusion limits for the Gbb and Gtt models are shown in the LSP and gluino mass plane in Figs. 7(a) and 7(b), respectively. The $\sigma_{\text{SUSY}}$ theory lines around the observed limits are obtained by changing the SUSY cross section by 1 standard deviation ($\pm 1\sigma$), as described in Sec. III. The yellow band around the expected limit shows the $\pm 1\sigma$ uncertainty, including all statistical and systematic uncertainties except the theoretical uncertainties in the SUSY cross section. It has been checked that the observed exclusion limits obtained from pseudoexperiments differ by less than 25 GeV from the asymptotic approximation in gluino or LSP mass in the combined limits in Fig. 7, although the difference can be up to 50 GeV when using
For the Gbb models, gluinos with masses below 1.78 TeV are excluded at the 95% C.L. for LSP masses below 800 GeV. At high gluino masses, the exclusion limits are driven by the SR-Gbb-A and SR-Gbb-B signal regions. The best exclusion limit on the LSP mass is approximately 1.0 TeV, which is reached for a gluino mass of approximately 1.6 TeV. The exclusion limit is dominated by SR-Gbb-C for high LSP masses. For the Gtt models, gluinos up to 1.8 TeV are excluded for massless LSP. For LSP masses below 700 GeV, gluino masses below 1.76 TeV are excluded. For large gluino masses, the exclusion limits are driven by the combination of SR-Gtt-IL-B and SR-Gtt-OL-A. The LSP exclusion extends up to approximately 975 GeV, corresponding to a gluino mass of approximately 1.5 TeV–1.6 TeV. The best exclusion limits are obtained by the combination of SR-Gtt-IL-B and SR-Gtt-OL-C for high LSP masses. The ATLAS exclusion limits obtained with the full $\sqrt{s} = 8$ TeV data set are also shown in Fig. 7. The results largely improve on the $\sqrt{s} = 8$ TeV limits despite the lower integrated luminosity. The exclusion limit on the gluino mass is extended by approximately 500 GeV and 400 GeV for the Gbb and Gtt models for massless LSP, respectively. This improvement is primarily attributable to the increased center-of-mass energy of the LHC. The addition of the IBL pixel layer in Run 2, which improves the capability to tag $b$ jets [35], also particularly benefits this analysis that employs a data set requiring at least three $b$-tagged jets. The sensitivity of the data analysis is also improved with respect to the $\sqrt{s} = 8$ TeV analysis [17] by using top-tagged large-$R$ jets, lepton isolation adapted to a busy environment, and the $m_{T,\text{min}}$ variable.

### IX. CONCLUSION

A search for pair-produced gluinos decaying via sbottom or stop is presented. LHC proton-proton collision data from the full 2015 data-taking period were analyzed, corresponding to an integrated luminosity of 3.2 fb$^{-1}$ collected at $\sqrt{s} = 13$ TeV by the ATLAS detector. Several signal regions are designed for different scenarios of gluino and LSP masses. They require several high-$p_T$ jets, of which at least three must be $b$ tagged, large $E_T^{\text{miss}}$, and, either zero or at least one charged lepton. For the gluino models with stop-mediated decays in which there is a large mass difference between the gluino and the LSP, large-$R$ jets identified as originating from highly boosted top quarks are employed. The background is dominated by $t\bar{t}$ + jets, which is normalized in dedicated control regions. No excess is found above the predicted background of each signal region. Model-independent limits are set on the $\sqrt{s} = 8$ TeV visible cross section for new physics processes. Exclusion limits are set on gluino and LSP masses in the simplified gluino models with stop-mediated and sbottom-mediated decays. For LSP masses below approximately 700 GeV, gluino masses of less than 1.78 TeV and 1.76 TeV are excluded at the 95% C.L. for the gluino models with sbottom-mediated and stop-mediated decays, respectively. These results significantly extend the exclusion limits obtained with the $\sqrt{s} = 8$ TeV data set.

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