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ATLAS Collaboration

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1. Introduction

The discovery of a Higgs boson at the ATLAS and CMS experiments [1,2] offers a new opportunity to search for physics beyond the Standard Model (SM) by examining the strength and structure of the Higgs boson's interactions with other particles. Thus far, the interactions of the Higgs boson have been probed using the $\kappa$-framework [3], in which the strength of a given coupling is allowed to vary from the SM prediction by a constant value. In this approach, the total rate of a given production and decay channel can differ from the SM prediction, but the kinematic properties of the Higgs boson in each decay channel are unchanged.

An alternative framework for probing physics beyond the SM is the effective field theory (EFT) approach [3–8], whereby the SM Lagrangian is augmented by additional operators of dimension six or higher. Some of these operators produce new tensor structures for the interactions between the Higgs boson and the SM particles, which can modify the shapes of the Higgs boson kinematic distributions as well as the associated jet spectra. The new interactions arise as the low-energy manifestation of new physics that exists at energy scales much larger than the partonic centre-of-mass energies being probed.

In this Letter, the effects of operators that produce anomalous CP-even and CP-odd interactions between the Higgs boson and photons, gluons, W bosons and Z bosons are studied using an EFT-inspired effective Lagrangian. The analysis is performed using a simultaneous fit to five detector-corrected differential cross sections in the $H \rightarrow \gamma\gamma$ decay channel, which were previously published by the ATLAS Collaboration [9]. These are the differential cross sections as functions of the diphoton transverse momentum ($p_T^{\gamma\gamma}$), the number of jets produced in association with the diphoton system ($N_{jets}$), the leading-jet transverse momentum ($p_T^{jet}$), and the invariant mass ($m_{jj}$) and difference in azimuthal angle ($\Delta\phi_{jj}$) of the leading and sub-leading jets in events containing two or more jets. The inclusion of differential information significantly improves the sensitivity to operators that modify the Higgs boson's interactions with $W$ and $Z$ bosons. To perform a simultaneous analysis of these distributions, the statistical correlations between bins of different distributions need to be included in the fit procedure. These correlations are evaluated by analysing the $H \rightarrow \gamma\gamma$ candidate events in the data, and are published as part of this Letter to allow future studies of new physics that produces non-SM kinematic distributions for $H \rightarrow \gamma\gamma$.

2. Higgs effective Lagrangian

The effective Lagrangian used in this analysis is presented in Ref. [8]. In this model, the SM Lagrangian is augmented with the dimension six CP-even operators of the Strongly Interacting Light Higgs formulation [6] and corresponding CP-odd operators. The $H \rightarrow \gamma\gamma$ differential cross sections are mainly sensitive to the
operators that affect the Higgs boson’s interactions with gauge bosons and the relevant terms in the effective Lagrangian can be specified by
\[ \mathcal{L}_{\text{eff}} = \bar{\epsilon}_1 \mathcal{O}_1 + \bar{\epsilon}_2 \mathcal{O}_2 + \bar{\epsilon}_{3W} \mathcal{O}_{3W} + \bar{\epsilon}_{4H} \mathcal{O}_{4H} + \bar{\epsilon}_{4H} \mathcal{O}_{4H}, \]
where \( \bar{\epsilon}_1 \) and \( \bar{\epsilon}_2 \) are ‘Wilson coefficients’ specifying the strength of the new CP-even and CP-odd interactions, respectively, and the dimension-six operators \( \mathcal{O}_i \) are those described in Refs. [8,10]. In the SM, all the Wilson coefficients are equal to zero. The \( \mathcal{O}_1, \mathcal{O}_2 \) operators introduce new interactions between the Higgs boson and two photons. The \( \mathcal{O}_3 \) and \( \mathcal{O}_4 \) operators introduce new interactions between the Higgs boson and two gluons and the analysis presented in this Letter is sensitive to these operators through the gluon fusion production mechanism. The \( \mathcal{O}_{3W} \) and \( \mathcal{O}_{4H} \) operators introduce new \( HWW, HZZ \) and \( HZ\gamma \) interactions. The \( HZZ \) and \( HZ\gamma \) interactions are also impacted by \( \mathcal{O}_{4H} \) and \( \mathcal{O}_{4H} \) and, to a lesser extent, \( \mathcal{O}_2, \mathcal{O}_4 \). The analysis presented in this Letter is sensitive to the \( \mathcal{O}_{3W}, \mathcal{O}_{4H} \) and \( \mathcal{O}_{4H} \) operators through vector-boson fusion and associated production.

Other operators in the full effective Lagrangian of Ref. [8] can also modify Higgs boson interactions. Combinations of some of the CP-even operators have been constrained using global fits to experimental data from LEP and the LHC [8,11,12].

3. Statistical correlations between differential distributions

ATLAS [13] is a multipurpose particle physics detector with cylindrical geometry and nearly 4π coverage in solid angle. The analysis is performed using proton–proton collision data at a centre-of-mass energy \( \sqrt{s} = 8 \text{ TeV} \) and an integrated luminosity of 20.3 fb\(^{-1}\).

The object and event selections used to define the differential distributions are described in detail in Ref. [9]. The statistical correlations between the measured cross sections as a function of different distributions are obtained using a random sampling with replacement method on the detector-level data. This procedure is often referred to as ‘bootstrapping’ [14]. Bootstrapped event samples are constructed from the data by assigning each event a weight pulled from a Poisson distribution with unit mean. The five differential distributions are then reconstructed using the weighted events, and the signal yields in each bin of a differential distribution are determined using an unbinned maximum-likelihood fit of the diphoton invariant mass spectrum (full details of the fit can be found in Ref. [9]). The procedure is repeated 10,000 times with statistically independent weights and the correlation between two bins of different distributions is determined from the scatter graph of the corresponding extracted cross sections. The observed correlations between bins of the measured \( p_T^{\gamma\gamma} \) and \( N_{\text{jets}} \) cross sections are shown in Fig. 1.

The statistical uncertainties on the correlation due to the finite number of bootstrap samples ranges from 0.5% to 1%. The statistical uncertainty on the correlations due to the finite number of events in data is determined to be less than 2% using the statistical overlap and variance of signal and background events in a mass window around the Higgs boson mass. In order to validate this approach, a set of pseudo-experiments was created from input

\[ \text{Fig. 1. Statistical correlations between the measured cross sections in bins of the diphoton transverse momentum and jet multiplicity distributions. The quoted uncertainties refer to the total statistical uncertainty due to the finite number of bootstrapped samples and the finite number of data events.} \]

conditions (with known correlations) chosen to be similar to those in data in terms of purity, kinematics and sample size. For each pseudo-experiment, a value for the correlation is determined using 10,000 bootstrapped samples and compared to the input correlation. No bias due to the bootstrapping is observed in the central value obtained from 500 pseudo-experiments.

As part of this Letter, the correlations computed above are made publicly available in HEPDATA [15], allowing the analysis to be repeated using alternative effective Lagrangians, complete EFT frameworks, or other models with non-SM Higgs boson interactions.

4. Theoretical predictions

The effective Lagrangian has been implemented in FeynRules [10]. Parton-level event samples are produced for specific values of Wilson coefficients by interfacing the universal file output from FeynRules to the MADGRAPH5 [17] event generator. Higgs boson production via gluon fusion is produced with up to two additional partons in the final state using leading-order matrix elements. The 0-, 1- and 2-parton events are merged using the MLM matching scheme [18] and passed through the PYTHIA6 generator [19] to create the fully hadronic final state. Event samples containing a Higgs boson produced either in association with a vector boson or via vector-boson fusion are produced using leading-order matrix elements and passed through the PYTHIA6 generator. For each production mode, the Higgs boson mass is set to 125 GeV [20] and events are generated using the CTEQ6L1 parton distribution function and the AUET2 parameter set [21]. All other Higgs boson production modes are assumed to occur as predicted by the SM.

Event samples are produced for different values of a given Wilson coefficient. The particle-level differential cross sections are produced using RIVET [22]. The Professor method [23] is used to interpolate between these samples, for each bin of each distribution, and provides a parameterisation of the effective Lagrangian prediction. The parameterisation function is determined using 11 samples when studying a single Wilson coefficient, whereas

\footnote{The implementation in Ref. [10] involves a redefinition of the gauge boson propagators that results in unphysical amplitudes unless certain physical constants are also redefined. The original implementation did not include the redefinition of these physical constants. However, the impact of redefining the physical constants is found to be less than 1% on the predicted cross sections across the range of Wilson coefficients studied. The relative change in the predicted Higgs boson cross sections as functions of the different Wilson coefficients is also found to agree with that predicted by the Higgs characterisation framework [16], with less than 2% variation across the parameter ranges studied.}
25 samples are used when studying two Wilson coefficients simultaneously. As the Wilson coefficients enter the effective Lagrangian in a linear fashion, second-order polynomials are used to predict the cross sections in each bin. The method was validated by comparing the differential cross sections obtained with the parameterisation function to the predictions obtained with dedicated event samples generated at the specific point in parameter space.

The model implemented in FeynRules fixes the Higgs boson width to be that of the SM, $\Gamma_H = 4.07$ MeV [3]. The cross sections are scaled by $\Gamma_H/\Gamma_H + \Delta\Gamma$, where $\Delta\Gamma$ is the change in partial width due to a specific choice of Wilson coefficient. The change in partial width is determined for each Higgs coupling using the partial-width calculator in MadGraph5 and normalised to reproduce the SM prediction from HDECAY [24].

The leading-order predictions obtained from MadGraph5 are reweighted to account for higher-order QCD and electroweak corrections to the SM process, assuming that these corrections factorise from the new physics effects. The differential cross section as a function of variable $X$ for a specific choice of Wilson coefficient, $c_i$, is given by

$$\frac{d\sigma}{dx} = \sum_{j} \left( \frac{d\sigma_j}{dx} \right)^{\text{ref}} \left( \frac{d\sigma_j}{dx} \right)_{c_i} = 0 \left( \frac{d\sigma_j}{dx} \right)_{c_i=0},$$

where the summation $j$ is over the different Higgs boson production mechanisms, 'MG5' labels the MadGraph5 prediction and 'ref' labels a reference sample for SM Higgs boson production.

The reference sample for Higgs boson production via gluon fusion is simulated using MG5_aMC@NLO [25] with the CT10 parton distribution function [26]. The $H+n$-jets topologies are generated using next-to-leading-order (NLO) matrix elements for each parton multiplicity ($n = 0, 1$ or $2$) and combined using the FxFx merging scheme [27]. The parton-level events are passed through Pythia8 [28] to produce the hadronic final state using the AU2 parameter set [29]. The sample is normalised to the total cross section predicted by a next-to-next-to-leading order plus next-to-next-to-leading-logarithm (NNLO+NNLL) QCD calculation with NLO electroweak corrections applied [3]. The reference sample for Higgs boson production via vector-boson fusion (VBF) is generated at NLO accuracy in QCD using the Powheg Box [30]. The events are generated using the CT10 parton distribution function (PDF) and Pythia8 with the AU2 parameter set. The VBF sample is normalised to an approximate NNLO QCD cross section with NLO electroweak corrections applied [3]. The reference samples for Higgs boson production in association with a vector boson ($VH$, $V = W, Z$) or a top–antitop pair ($t\bar{t}H$) are produced at leading-order accuracy using Pythia8 with the CTEQ6L1 PDF and the 4C parameter set [21]. The ZH and WH samples are normalised to cross sections calculated at NNLO in QCD with NLO electroweak corrections, whereas the tH sample is normalised to a cross section calculated to NLO in QCD [3].

The ratio of the differential cross sections to the SM predictions for some representative values of the Wilson coefficients are shown in Fig. 2. The impact of the $\xi_g$ and $\xi_t$ coefficients are presented for the gluon fusion production channel and show a large change in the overall cross section normalisation. The $\xi_g$ coefficient also changes the shape of the $\Delta\phi_{jj}$ distribution, which is expected from consideration of the tensor structure of CP-even and CP-odd interactions [31,32]. The impact of the $\xi_{tV}$ and $\xi_{tW}$ coefficients are presented for the VBF + VH production channel and show large shape changes in all of the studied distributions. The

### 5. Limit-setting procedure

Limits on the Wilson coefficients are set by constructing a $\chi^2$ function

$$\chi^2 = \left( \hat{\sigma}_{\text{data}} - \hat{\sigma}_{\text{pred}} \right)^T C^{-1} \left( \hat{\sigma}_{\text{data}} - \hat{\sigma}_{\text{pred}} \right),$$

where $\hat{\sigma}_{\text{data}}$ and $\hat{\sigma}_{\text{pred}}$ are vectors from the measured and predicted cross sections of the five analysed observables, and $C = C_{\text{stat}} + C_{\text{exp}} + C_{\text{pred}}$ is the total covariance matrix defined by the sum of the statistical, experimental and theoretical covariances. The predicted cross section $\hat{\sigma}_{\text{pred}}$ and its associated covariance $C_{\text{pred}}$ are continuous functions of Wilson coefficients. Scans of one or two Wilson coefficients are carried out and the minimum $\chi^2$ value, $\chi^2_{\min}$, is determined. The confidence level (CL) of each scan point can be calculated as

$$1 - \text{CL} = n \int dx \frac{1}{\chi^2}(x; n),$$

with $\chi^2(x; c_i)$ being the $\chi^2$ value evaluated for a given Wilson coefficient $c_i$, and $f(x; m)$ being the $\chi^2$ distribution for $m$ degrees of freedom and $n = 1$ or $\frac{1}{2}$ for two-sided or one-sided limits. The coverage of CL and the effective number of degrees of freedom are determined using ensembles of pseudo-experiments.

The input data vector is compared in Fig. 3 to the SM hypothesis as well as two non-SM hypotheses specified by $\xi_g = 1 \times 10^{-4}$ and $\xi_{tW} = 0.05$, respectively.

The covariance matrix for experimental systematic uncertainties is constructed from all uncertainty sources provided by Ref. [9], which include the jet energy scale and resolution uncertainties, photon energy and resolution uncertainties, and model uncertainties. Identical sources are assumed to be fully correlated across samples using form-factor predictions from VBFNLO [33]. The impact on the $\xi_{tW}$ limits is negligible for $\Delta m > 1$ TeV.

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3 Form factors are sometimes used to regularise the change of the cross section above a momentum scale $\Lambda_{QCD}$. This was investigated by reweighting the VBF + VH

4 For one-dimensional limits on the CP-even (odd) Wilson coefficients, good agreement is found between the asymptotic formula and the pseudo-experiment test statistic with $m = 1$ and $n = 1$ [21]. For the two-dimensional limits on $\xi_g$ versus $\xi_t$ and $\xi_{tW}$, good agreement between pseudo-experiments and asymptotic formula is found for $m = 1$ and $n = 1$. For the two-dimensional limit on $\xi_g$ versus $\xi_t$, good agreement between pseudo-experiments and asymptotic formula is found for $m = 2$ and $n = 1$. 

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**Fig. 2.** Ratio of differential cross sections predicted by specific choices of Wilson coefficient to the differential cross sections predicted by the SM.
the cross correlation have a negligible impact on the results reported here.

The covariance matrix for the theoretical uncertainties is constructed to account for missing higher-order corrections and PDF uncertainties in the SM reference predictions. The uncertainties in the gluon fusion reference samples are: (i) a shape uncertainty, estimated by simultaneously varying the factorisation and renormalisation scales in MG5_aMC@NLO by a factor of 0.5 or 2.0, and (ii) uncertainties on the NNLO+NNLL QCD plus NLO electroweak total cross-section prediction [3], arising from missing higher-order corrections and PDF uncertainties; these uncertainties are assumed to be fully correlated among bins and observables. For VBF, ZH and WH, shape uncertainties are neglected because their impact is expected to be negligible with respect to all other theory uncertainties. Normalisation uncertainties for these processes are taken from Ref. [3].

The benefit of using more than one differential distribution in the analysis is quantified using an ‘Asimov dataset’, which is a representative dataset of the median expected cross-section measurement assuming the SM. For $\varepsilon_g$ and $\varepsilon_y$, the use of a single inclusive distribution ($p_X^{\gamma\gamma}$ or $N_{jets}$) results in the same expected limits as the full five-dimensional fit. For $\varepsilon_g$ and $\varepsilon_y$, the most sensitive variable is found to be $p_X^{\gamma\gamma}$, with a 5% improvement in the expected limits obtained from using the five-dimensional information. For $\varepsilon_{HW}$ and $\varepsilon_{HW}$, the most sensitive variable is $\Delta \phi_{jj}$ and an 18% improvement in the expected limits is obtained from using the five-dimensional fit. In summary, the expected sensitivity for $\varepsilon_g$, $\varepsilon_y$, $\varepsilon_y$ and $\varepsilon_y$ arises mainly from the normalisation of the different production mechanisms, and can be probed using the inclusive distributions that distinguish between the different processes, whereas the $\varepsilon_{HW}$ and $\varepsilon_{HW}$ coefficients benefit more from the full five-dimensional information due to the induced shape changes in the kinematics of the VBF + VH process.

6. Results

The 68% and 95% confidence regions for a two-dimensional scan of $\bar{c}_y$ and $\bar{c}_g$ are shown in Fig. 4, after setting all other Wilson coefficients to zero. These additional interactions can interfere with the corresponding SM interactions. Destructive interference, for example, causes the $H \rightarrow \gamma\gamma$ branching ratio to be zero at $\bar{c}_y \sim 2 \times 10^{-3}$ and the gluon fusion production cross section to be zero at $\bar{c}_g \sim -2.2 \times 10^{-4}$. The impact of these effects is evident in the structure of the obtained limits in the two-dimensional parameter plane.

The 68% and 95% confidence regions for a two-dimensional scan of $\bar{c}_g$ and $\bar{c}_y$ are shown in Fig. 5, after setting all other Wilson coefficients to zero. The $\Delta \phi_{jj}$ distribution is sensitive to the $\bar{c}_g$ parameter through the gluon fusion production mechanism (Figs. 2 and 3) and the limit on $\bar{c}_g$ is improved with the inclusion of this data in the fit. This is evident in Fig. 5 where the limit band is constricted at the largest values of $\bar{c}_g$.

The 68% and 95% confidence regions obtained from scanning $\bar{c}_{HW}$ and $\bar{c}_{HW}$ are shown in Fig. 6, after setting $\bar{c}_{HW} = \bar{c}_{HW}$ and $\bar{c}_{HW} = \bar{c}_{HW}$ to ensure that the partial width for $H \rightarrow Z\gamma$ is unchanged from the SM prediction. As discussed in Section 5, these Wilson coefficients produce large shape changes in all distributions and the obtained limits are strongest when fitting all five distributions simultaneously.

The 95% confidence regions for $\bar{c}_{HW}$ and $\bar{c}_{HW}$ can be translated into the Higgs Characterisation framework [16] and compared to the ATLAS results for non-SM CP-even and CP-odd $H\gamma$ interactions, which were obtained using an angular analysis of the decay

5 Values of $|\bar{c}_{HW} - \bar{c}_{HW}| > 0.033$ lead to a very large decay rate for the $H \rightarrow Z\gamma$ process that is contradicted by the experimental constraints reported by ATLAS [35].
work on the SM, we have used the Higgs boson properties to constrain the parameter space of the model.

Table 1: Observed allowed ranges at 95% CL for the $c_\gamma$, $c_\phi$, and $c_\Delta$ Wilson coefficients and their CP-conjugate partners. Limits on $c_\gamma$, $c_\phi$, and $\bar{c}_\Delta$ are derived with $c_{\phi GW} = 0.90$, whereas limits on $c_{\phi GW}$ and $c_{\gamma GW}$ are derived with $c_{\phi GW} = 0.90$ and $c_{\gamma GW} = c_{\gamma GW}$, respectively. Two allowed regions are observed for $c_\gamma$ and $\bar{c}_\Delta$, with the region between the solutions producing too small $pp \to H \to \gamma\gamma$ cross section due to destructive interference between new interactions and the SM.

The obtained limits on $c_{\phi GW}$ and $c_{\gamma GW}$ are also not excluded by current signal strength measurements. For example, the signal strength in the $H \to ZZ^*$ and $H \to WW^*$ channels is predicted to be approximately 1 for $c_{\phi GW} = 0.1$, which is consistent with the dedicated measurements.

The 95% confidence regions for a one-dimensional scan of the Wilson coefficients are given in Table 1.

7. Summary

The strength and structure of the Higgs boson's interactions with other particles have been investigated using an effective Lagrangian. Limits are placed on anomalous CP-even and CP-odd interactions between the Higgs boson and photons, gluons, $W$-bosons and $Z$-bosons, using a fit to five differential cross sections previously measured by ATLAS in the $H \to \gamma\gamma$ decay channel at $\sqrt{s} = 8$ TeV [9]. No significant deviations from the SM predictions are observed. To allow a simultaneous fit to all distributions, the statistical correlations between these distributions have been determined by re-analysing the candidate $H \to \gamma\gamma$ events in the proton–proton collision data. These correlations are made publicly [15] available to allow for future analysis of theories with non-SM Higgs boson interactions.

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References

ATLAS Collaboration


1 Department of Physics, University of Adelaide, Adelaide, Australia
2 Physics Department, SUNY Albany, Albany, NY, United States
3 Department of Physics, University of Alberta, Edmonton, AB, Canada
4 (a) Department of Physics, Ankara University, Ankara; (b) Istanbul Aydin University, Istanbul; (c) Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey
5 LAPP, CNRS/IN2P3 and Université Savoie Mont Blanc, Annecy-le-Vieux, France
6 High Energy Physics Division, Argonne National Laboratory, Argonne, IL, United States
7 Department of Physics, University of Arizona, Tucson, AZ, United States
8 Department of Physics, The University of Texas at Arlington, Arlington, TX, United States
9 Physics Department, University of Athens, Athens, Greece
10 Physics Department, National Technical University of Athens, Zografou, Greece
11 Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
12 Institut de Física d’Altes Energies and Departamento de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain
13 Institute of Physics, University of Belgrade, Belgrade, Serbia
14 Department for Physics and Technology, University of Bergen, Bergen, Norway
15 Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, CA, United States
16 Department of Physics, Humboldt University, Berlin, Germany
17 Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
18 School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
19 (a) Department of Physics, Bogazici University, Istanbul; (b) Department of Physics Engineering, Gaziantep University, Gaziantep; (c) Department of Physics, Dogaş University, Istanbul, Turkey
20 (a) INFN Sezione di Bologna; (b) Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy
21 Physikalisches Institut, Universität von Bonn, Bonn, Germany
22 Department of Physics, Boston University, Boston, MA, United States
23 Department of Physics, Brandeis University, Waltham, MA, United States
24 (a) Universidade Federal do Rio de Janeiro COPPE/UFRJ, Rio de Janeiro; (b) Electrical Circuits Department, Federal University of Juiz de Fora (UFJE), Juiz de Fora; (c) Federal University of Sao Joao del Rei (UFJF), Sao Joao del Rei; (d) Instituto de Física, Universidade de Sao Paulo, Sao Paulo, Brazil
25 Physics Department, Brookhaven National Laboratory, Upton, NY, United States
26 (a) Transilvania University of Brasov, Brasov; (b) National Institute of Physics and Nuclear Engineering, Bucharest; (c) National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj Napoca; (d) University Politehnica Bucharest, Bucharest; (e) West University in Timisoara, Timisoara, Romania
27 Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
28 Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
29 Department of Physics, Carleton University, Ottawa, ON, Canada
30 CERN, Geneva, Switzerland
31 Enrico Fermi Institute, University of Chicago, Chicago, IL, United States
32 (a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; (b) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
33 (a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; (b) Department of Modern Physics, University of Science and Technology of China, Anhui; (c) Department of Physics, Nanjing University, Jiangsu; (d) School of Physics, Shandong University, Shandong; (e) Department of Physics and Astronomy, Shanghai Key Laboratory for Particle Physics and Cosmology, Shanghai Jiao Tong University, Shanghai; (f) Physics Department, Tsinghua University, Beijing, 100084, China
34 Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France
35 Nevis Laboratory, Columbia University, Irvington, NY, United States
36 Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
37 (a) INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati; (b) Dipartimento di Fisica, Università della Calabria, Rende, Italy
38 (a) AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow; (b) Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland
39 Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland
40 Physics Department, Southern Methodist University, Dallas, TX, United States
41 Physics Department, University of Texas at Dallas, Richardson, TX, United States
42 DESY, Hamburg and Zeuthen, Germany
43 Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
44 Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
45 Department of Physics, Duke University, Durham, NC, United States
46 SUPA – School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
47 INFN Laboratori Nazionali di Frascati, Frascati, Italy
48 Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany
49 Section de Physique, Université de Genève, Genève, Switzerland
50 INFN Sezione di Genova; (a) Dipartimento di Fisica, Università di Genova, Genova, Italy
51 (a) E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi; (b) High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
52 II Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
53 SUPA – School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
54 II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany
55 Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS/IN2P3, Grenoble, France
56 Department of Physics, Hampton University, Hampton, VA, United States
57 Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, MA, United States
58 (a) Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; (b) Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany
59 Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
60 (a) Department of Physics, The Chinese University of Hong Kong, Shatin, N.T., Hong Kong; (b) Department of Physics, The University of Hong Kong, Hong Kong; (c) Department of Physics, The Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China
61 Department of Physics, Indiana University, Bloomington, IN, United States
62 Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
63 University of Iowa, Iowa City, IA, United States
64 Department of Physics and Astronomy, Iowa State University, Ames, IA, United States
65 Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia
66 KER, High Energy Accelerator Research Organization, Tsukuba, Japan
67 Graduate School of Science, Kobe University, Kobe, Japan
68 Faculty of Science, Kyoto University, Kyoto, Japan
Also at Department of Physics, King's College London, London, United Kingdom.
b Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.
c Also at Novosibirsk State University, Novosibirsk, Russia.
da Also at TRUMF, Vancouver, BC, Canada.
b Also at Department of Physics, California State University, Fresno, CA, United States of America.
c Also at Department of Physics, University of Fribourg, Fribourg, Switzerland.
d Also at Departamento de Fisica e Astronomia, Faculdade de Ciencias, Universidade do Porto, Portugal.
e Also at Tomsk State University, Tomsk, Russia.
f Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France.
g Also at Universita di Napoli Parthenope, Napoli, Italy.
h Also at Institute of Particle Physics (IPP), Canada.
i Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom.
j Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia.
k Also at Louisiana Tech University, Ruston, LA, United States of America.
l Also at Instituto Catalana de Recerca i Studis Avançats, ICREA, Barcelona, Spain.
m Also at Graduate School of Science, Osaka University, Osaka, Japan.
 Also at Department of Physics, National Tsing Hua University, Taiwan.
 Also at Department of Physics, The University of Texas at Austin, Austin, TX, United States of America.
 Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia.
i Also at CERN, Geneva, Switzerland.
j Also at Georgian Technical University (GTU), Tbilisi, Georgia.
k Also at Manhattan College, New York, NY, United States of America.
l Also at Hellenic Open University, Patras, Greece.
m Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.
 Also at LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France.
n Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan.
o Also at School of Physics, Shandong University, Shandong, China.
p Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia.
q Also at Section de Physique, Université de Genève, Geneva, Switzerland.
r Also at International School for Advanced Studies (SISSA), Trieste, Italy.
s Also at Department of Physics and Astronomy, University of South Carolina, Columbia, SC, United States of America.
t Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China.
u Also at Faculty of Physics, M.V. Lomonosov Moscow State University, Moscow, Russia.
w Also at National Research Nuclear University MEPhI, Moscow, Russia.
x Also at Department of Physics, Stanford University, Stanford, CA, United States of America.
y Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.
z Also at University of Malaya, Department of Physics, Kuala Lumpur, Malaysia.