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Search for massive colored scalars in four-jet final states in $\sqrt{s}=7$ TeV proton–proton collisions with the ATLAS detector

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Abstract A search for pair-produced scalar particles decaying to a four-jet final state is presented. The analysis is performed using an integrated luminosity of 34 pb$^{-1}$ recorded by the ATLAS detector in 2010. No deviation from the Standard Model is observed. For a scalar mass of 100 GeV (190 GeV) the limit on the scalar gluon pair production cross section at 95% confidence level is 1 nb (0.28 nb). When these results are interpreted as mass limits, scalar–gluons (hyperpions) with masses of 100 to 185 GeV (100 to 155 GeV) are excluded at 95% confidence level with the exception of a mass window of width about 5 GeV (15 GeV) around 140 GeV.

At hadron colliders, the search for new phenomena in fully hadronic final states without missing transverse energy or leptons is experimentally challenging because of the large multijet background. Recent studies used the dijet mass spectrum and the dijet angular distribution observed at the LHC [1–3] to search for physics beyond the Standard Model. Some extensions of the Standard Model predict new phenomena in events with higher jet multiplicity. In the six-jet final state CDF [4] and CMS [5] have excluded $R$-parity violating gluinos, the supersymmetric partners of the gluons, with masses from 200 GeV to 280 GeV using a model-independent search. This letter describes a search for pair-produced scalar particles decaying to two jets, leading to a four-jet final state with two jet–jet resonances and no missing transverse energy.

Two scenarios serve as a guideline and motivation for the analysis: the extension of the Standard Model with a new gauge group called “hypercolor” as described in [6–8], and extended supersymmetric models [9–11].

In the hypercolor model, new colored fermions are charged under an additional gauge group (SU(3) hypercolor). In analogy with QCD, new colored fermions are bound into mesons of hypercolor: a color octet vector particle, the coloron, and a color octet pseudoscalar, the hyperpion which decays to two gluons.

In supersymmetric models with Dirac gluinos, a scalar gluon (sgluon) extends the QCD sector, which is made up of the gluon/gluino super-multiplet and an additional gluino/sgluon super-multiplet. As the sgluon has positive $R$-parity [12], light sgluons, i.e. sgluons with masses of the order of 100 GeV, are expected to decay to two gluons with a branching ratio close to 1. Single production of sgluons, loop-induced via supersymmetric particles, is also possible, but these cross sections are several orders of magnitude smaller than those of the studies in Ref. [3]; therefore the previously obtained limit of 1.92 TeV on the color-octet mass does not apply to the models studied in this letter. The pair production cross section does not depend, at leading order, on supersymmetric parameters except the sgluon mass.

In the following, the sgluon pair production will be used as the benchmark process for the production of two heavy objects of equal mass, each decaying into two jets.

ATLAS [13] is a multipurpose detector with nearly $4\pi$ coverage in solid angle. The inner detector, consisting of silicon pixel and microstrip detectors as well as a transition radiation tracker, is immersed in a 2 T solenoidal magnetic field. The finely-segmented, hermetic calorimeter covers $|\eta|<4.9^1$ and provides three-dimensional reconstruction of particle showers. The electromagnetic calorimeter is a lead liquid-argon sampling calorimeter. In the central region it is surrounded by a hadronic calorimeter made of iron and scintillating tiles. For the region $|\eta|>1.7$ the sampling

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1 ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the $z$-axis along the beam pipe. The $x$-axis points from the IP to the centre of the LHC ring, and the $y$ axis points upward. Cylindrical coordinates $(r,\phi)$ are used in the transverse plane, $\phi$ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle $\theta$ as $\eta=-\ln\tan(\theta/2)$.
calorimeter consists of copper or tungsten and liquid argon. The calorimeters are surrounded by the muon spectrometer which consists of three large superconducting toroids, a system of precision tracking chambers, and detectors for triggering.

ATLAS uses a three-level trigger system. The first level trigger is implemented in hardware, the other two trigger levels are implemented in software. In the analysis particular emphasis was placed on being sensitive to the low mass region in order to exploit the low trigger thresholds of the data recorded in 2010. Therefore, a first level trigger requiring at least four jets with a transverse momentum $p_T$ of at least 5 GeV was used. The trigger efficiency increases as a function of the jet $p_T$ to at least 99% [14] for calibrated jets of $p_T > 55$ GeV. An integrated luminosity of 34 pb$^{-1}$ has been recorded with this trigger, taking into account the prescaling of the trigger at the end of 2010.

The Standard Model (SM) multijet production of four jets with $p_T > 60$ GeV and $|\eta| < 2.8$ has a cross section of approximately 5 nb [15]. Compared to multijet production, even without considering the branching ratios to obtain a final state with four hard jets, other Standard Model backgrounds have much smaller cross sections: $WW$, $t\bar{t}$ with a cross section of 41 pb [16], and $WW$ with a cross section of 171 pb [17] and $WW$ production associated with two jets ($p_T > 20$ GeV, $\eta < 2.8$) with a cross section of 200 pb [18]. Monte Carlo (MC) samples were used to model the SM multijet background. ALPGEN [19] SM multijet samples were generated with the MLM matching scheme [20], interfaced to HERWIG [21] for parton shower and fragmentation and to JIMMY [22] for the simulation of the underlying event. As a cross check, PYTHIA [23] SM dijet samples were compared to the underlying event tune (PDF) [24] and the ALPGEN sample with CTEQ6L1 PDF [25]. For both, the underlying event tune was the ATLAS MC09 tune [26]. The sgluon pair production differential cross section of $41 pb$ [16], $t\bar{t}$ with a cross section of $171 pb$ [17] and $WW$ production associated with two jets ($p_T > 20$ GeV, $\eta < 2.8$) with a cross section of 200 pb [18].

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Fig. 1 Kinematic variables at different stages of the analysis. Data (dots) are compared to the ALPGEN MC sample (solid line). The solid band corresponds to one standard deviation in the jet energy scale. The ratio data/MC is also shown with its statistical uncertainty, which is dominated by the MC statistics. The dashed line corresponds to a sgluon signal of 100 GeV.

(a) The transverse momentum of the 4th jet is shown. (b) The $\Delta R_{jj}$ distribution for the reconstructed sgluon candidate with the highest transverse momentum jet is shown after requiring the transverse momentum to be greater than 55 GeV and pairing the four leading jets into two sgluon candidates. (c) The relative mass difference is shown after the criteria on the $p_T$ and $\Delta R_{jj}$ have been applied. (d) The scattering angle in the 4-jet center-of-mass frame is shown after all other selection criteria have been applied.

<table>
<thead>
<tr>
<th>Cut</th>
<th>Data</th>
<th>ALPGEN MC</th>
<th>sgluon MC</th>
<th>sgluon/ALPGEN</th>
</tr>
</thead>
<tbody>
<tr>
<td>4 jets $p_T &gt; 50$ GeV</td>
<td>568421</td>
<td>568000 ± 8000</td>
<td>27900 ± 800</td>
<td>4.9%</td>
</tr>
<tr>
<td>4 jets $p_T &gt; 55$ GeV</td>
<td>340429</td>
<td>336000 ± 6000</td>
<td>19000 ± 700</td>
<td>5.6%</td>
</tr>
<tr>
<td>$\Delta R_{jj} &lt; 1.6$</td>
<td>56131</td>
<td>55400 ± 1900</td>
<td>4900 ± 350</td>
<td>8.8%</td>
</tr>
<tr>
<td>$</td>
<td>M_1 - M_2</td>
<td>/(M_1 + M_2</td>
<td>&lt; 0.075$</td>
<td>16958</td>
</tr>
<tr>
<td>$</td>
<td>\cos(\theta^*)</td>
<td>&lt; 0.5$</td>
<td>6937</td>
<td>7700 ± 800</td>
</tr>
</tbody>
</table>

Table 1 Cut flow for data, ALPGEN MC sample and sgluon MC ($M_{\text{sgluon}} = 100$ GeV). The ALPGEN MC sample is normalized to the data after the first requirement.
statistical error in the regions B, C and D which feeds into the error on the background prediction.

For the average mass distribution, the shape of the background in the signal region is modeled by that in the control region B, parameterized by a fit to

\[ f(x) = (x - p_1)p_2 \cdot e^{-x \cdot p_3 - x^2 \cdot p_4} \]

where \( p_1, p_2, p_3, p_4 \) are the free parameters and \( x \) is the reconstructed average mass. The parameters \( p_1 \) and \( p_2 \) describe the rising edge of the distribution; whereas, the \( p_3 \) and \( p_4 \) parameters model its tail. In region B the selection criteria on the scattering angle is inverted with respect to the signal region, but not the one on the relative mass difference, leading to the best description of the background shape in the signal region.

The normalization of the background in the signal region is derived from the number of events in the control regions: \( N_{\text{extrapolation}}^A = N_B \cdot N_C / N_D \). This is referred to as the ABCD method. The effect of the correlation between the two variables used to define the different regions is neglected since the correlation is less than 0.2%. Performing a closure test on the ALPGEN and PYTHIA MC samples, the number of events in the signal region agreed with the prediction derived from the three other regions within the statistical error.

In Fig. 2 the result obtained with the data is shown in the signal region (region A) for sgluon masses of 100, 140, 160 and 190 GeV. The data in region A is compared to the data in the control region B, and to the fit in region B, where the data and the fit are each normalized using the ABCD method.
Table 3 Comparison of data in signal region with background prediction. The first column shows the $p_T$ requirement applied on the 4th leading jet in $p_T$, the second column the observed number of events in the signal region. The third column shows the prediction of the ABCD method. Only the statistical uncertainty is indicated. The fourth column is the $\chi^2/\text{NDF}(A, B)$ between the shapes of the reconstructed average mass distribution in regions A and B. The last column shows $\chi^2/\text{NDF}(B)$ for the fit to the background region

<table>
<thead>
<tr>
<th>$p_T^{\text{min}}$ (4th jet) [GeV]</th>
<th>Data</th>
<th>ABCD prediction</th>
<th>$\chi^2/\text{NDF}(A, B)$</th>
<th>$\chi^2/\text{NDF}(B)$</th>
</tr>
</thead>
<tbody>
<tr>
<td>49</td>
<td>11732</td>
<td>11410 ± 150</td>
<td>1.31</td>
<td>0.77</td>
</tr>
<tr>
<td>55</td>
<td>6937</td>
<td>6740 ± 120</td>
<td>1.02</td>
<td>1.05</td>
</tr>
<tr>
<td>60</td>
<td>4098</td>
<td>3980 ± 90</td>
<td>0.85</td>
<td>1.09</td>
</tr>
<tr>
<td>66</td>
<td>2532</td>
<td>2460 ± 70</td>
<td>1.04</td>
<td>0.87</td>
</tr>
<tr>
<td>71</td>
<td>1590</td>
<td>1580 ± 60</td>
<td>1.18</td>
<td>0.98</td>
</tr>
<tr>
<td>77</td>
<td>1069</td>
<td>1030 ± 50</td>
<td>1.39</td>
<td>0.61</td>
</tr>
<tr>
<td>82</td>
<td>701</td>
<td>720 ± 40</td>
<td>1.59</td>
<td>1.04</td>
</tr>
<tr>
<td>88</td>
<td>480</td>
<td>517 ± 34</td>
<td>1.32</td>
<td>1.00</td>
</tr>
<tr>
<td>93</td>
<td>322</td>
<td>364 ± 29</td>
<td>0.94</td>
<td>1.22</td>
</tr>
<tr>
<td>99</td>
<td>318</td>
<td>266 ± 25</td>
<td>1.08</td>
<td>1.22</td>
</tr>
<tr>
<td>104</td>
<td>162</td>
<td>187 ± 21</td>
<td>1.05</td>
<td>1.13</td>
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<tr>
<td>110</td>
<td>116</td>
<td>151 ± 19</td>
<td>1.42</td>
<td>1.44</td>
</tr>
</tbody>
</table>

The expected signals in regions A and B, normalized to the nominal sgluon cross section, are also shown. The prediction of the background from the ABCD prediction, the $\chi^2$ per degree of freedom (NDF) between the shapes of the distributions in region A and B ($\chi^2/\text{NDF}(A, B)$), as well as the $\chi^2/\text{NDF}(B)$ in the background region for the fit of the background function. No significant deviation is observed between the data-driven background prediction and the data. Therefore limits are set on the excluded cross section using a profile likelihood ratio with the CL$_s$ approach [30]. The shapes of the average mass distribution for signal and background in region A are parametrised and used in the likelihood. The signal shape is modeled with a Gaussian distribution and the background shape with the parametrisation of the ABCD method. The signal contamination in the control regions is taken into account in the likelihood. A Gaussian distribution is used to simulate the signal contamination in region B; whereas, in region C and D no assumption is made on the signal shape since only the number of events is used in these two regions.

The different sources of systematic uncertainty and their effect are summarized in Table 4. The uncertainty on the integrated luminosity is 3.4% [31]. The trigger efficiency is estimated in minimum bias data to be 99 ± 1%. The signal acceptance and contamination are taken from the full simulation Monte Carlo samples with a statistical uncertainty of 5% (in region A) by fitting the efficiencies as a function of the sgluon mass. The jet energy scale uncertainty is propagated to the signal [29], affecting the selection efficiency by 15%. A second effect of the JES uncertainty on the signal is a ±2% shift of the signal mass peak position. The impact of the jet energy resolution uncertainty on the signal mass peak width is 10%. The impact of the choice of the PDF for the signal generation was estimated to be less than 2%. Finally a systematic error, reflecting the statistics available to check the prediction of the ABCD method in the absence of new physics, is assigned to the background prediction. Gaussian nuisance parameters are implemented in the likelihood corresponding to the errors taking into account the correlations, e.g. the error on the luminosity is common to the ABCD regions. The contamination of the regions B, C and D by the signal is also taken into account in the likelihood.

For each tested mass, the observed and expected median CL$_s$ are determined as a function of the signal cross section. The analysis is performed for masses from 100 to 200 GeV in steps of 10 GeV. The resulting excluded cross section, shown in Fig. 3, is 1 nb at 100 GeV and 280 pb at 190 GeV. Converting this result into a mass limit for a
branching ratio of 1 to gluon pairs, using a leading order cross section [9] with CTEQ6L1 [25], sgluons with masses from 100 GeV to 185 GeV are excluded at 95% confidence level with the exception of a mass window of about 5 GeV around 140 GeV. The sgluon cross section used was checked at $\sqrt{s} = 14$ TeV with Ref. [9] and was found to agree at the percent level. The centrality of the hyperpions compared to the sgluons increases due to the additional contribution of the s-channel coloron exchange. This property should increase the selection efficiency due to the presence of the requirement on the scattering angle. However, the hyperpion cross section was scaled down from the sgluon cross section according to Ref. [7], which makes the limits less stringent. Hyperpions with masses of 100 GeV to 155 GeV are excluded with the exception of a mass window of 15 GeV around 140 GeV.

To conclude, four-jet events have been analyzed by the ATLAS experiment, searching for the pair production of a new scalar particle decaying to two jets. The data in the signal region is in good agreement with the data-driven background estimation. No evidence for new phenomena was found. Cross section limits as a function of the mass of the scalar particle have been determined. Interpreting the cross section limit, sgluons (hyperpions) with masses from 100 GeV to 185 GeV (155 GeV) are excluded at 95% CL. A mass window of about 5 GeV (15 GeV) around 140 GeV remains unexcluded for the sgluons (hyperpions).

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References

The ATLAS Collaboration

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