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DOI
10.1016/j.physletb.2012.09.016

Publication date
2012

Document Version
Final published version

Published in
Physics Letters B

Citation for published version (APA):

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ATLAS Collaboration

1. Introduction

The search for the Higgs boson [1–3], which in the Standard Model (SM) gives mass to the weak vector bosons and fermions, is one of the most important aspects of the CERN Large Hadron Collider (LHC) physics programme. Direct searches performed at the CERN Large Electron–Positron Collider (LEP) have excluded, at a 95% confidence level (CL), the production of a SM Higgs boson in the range $112 < m_H < 131$ GeV. At the LHC, the combination of ATLAS searches [7], using 1.0–4.9 fb$^{-1}$ of $\sqrt{s} = 7$ TeV data, excluded at the 95% CL the production of a SM Higgs boson in the regions $112.9 < m_H < 115.5$ GeV, $131 < m_H < 238$ GeV and $251 < m_H < 466$ GeV. The CMS combined result [8] using 4.6–4.8 fb$^{-1}$ of $\sqrt{s} = 7$ TeV data excluded the production of a SM Higgs boson at the 95% CL in the region $127 < m_H < 600$ GeV.

The $ZZ \rightarrow \ell^+\ell^-\nu \bar{\nu}$ decay channel offers a substantial branching fraction in combination with a good separation from potential background processes owing to the large transverse momentum, $p_T$, of the electron or muon pair from the leptonic $Z$ boson decay and the large missing transverse momentum from the $Z$ boson decaying to neutrinos. Earlier results in this channel, published in Ref. [9], using 1.0 fb$^{-1}$, were subsequently updated in Ref. [7] with 2.0 fb$^{-1}$ and exclude at the 95% CL the presence of a SM Higgs boson in the range $310 < m_H < 470$ GeV. The most recent CMS analysis [10] in this channel based on 4.6 fb$^{-1}$ excluded at 95% CL $270 < m_H < 440$ GeV.

The data sample considered in the search presented in this Letter was recorded by the ATLAS experiment during the 2011 LHC run at a centre-of-mass energy $\sqrt{s} = 7$ TeV. The integrated luminosity of the data sample, considering only data-taking periods where all relevant detector subsystems were operational, is 4.7 fb$^{-1}$ with an uncertainty of 3.9% [11,12].

2. The ATLAS detector

The ATLAS detector [13] is a multi-purpose particle physics detector with forward–backward symmetric cylindrical geometry. The inner tracking detector (ID) covers $|\eta| < 2.5$ and consists of a silicon pixel detector, a silicon microstrip detector, and a transition radiation tracker. The ID is surrounded by a thin superconducting solenoid providing a 2 T axial magnetic field. A high-granularity lead/liquid-argon (LAr) sampling calorimeter measures the energy and the position of electromagnetic showers with $|\eta| < 3.2$. LAr sampling calorimeters are also used to measure hadronic showers in the end-cap ($1.5 < |\eta| < 3.2$) and forward ($3.1 < |\eta| < 4.9$) regions, while an iron/scintillating-tile calorimeter measures hadronic showers in the central region ($|\eta| < 1.7$). The muon spectrometer surrounds the calorimeters and consists of three large superconducting air–core toroids, each with eight coils, a system of precision tracking chambers ($|\eta| < 2.7$), and fast tracking chambers for triggering ($|\eta| < 2.4$). A three-level trigger system selects events to be recorded for offline analysis.

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3. Data and simulation samples

The \( H \rightarrow ZZ \rightarrow \ell^+\ell^-\nu\bar{\nu} \) signal is modelled using the POWHEG [14,15] event generator, which includes matrix elements for the gluon fusion and vector-boson fusion (VBF) production mechanisms of the Higgs boson up to next-to-leading-order (NLO). POWHEG is interfaced to PYTHIA [16] for the modelling of parton showers. The modelling of final-state radiation is performed with PHOTOS [17], and TAUOLA [18,19] is used for the simulation of \( \tau \) decays. The Higgs boson \( p_T \) spectrum in the gluon fusion process is reweighted to the calculation of Ref. [20], which provides QCD corrections up to NLO and QCD soft-gluon resummations up to next-to-next-to-leading logarithms (NNLL).

\[
H \rightarrow W^+W^- \rightarrow \ell^+\ell^-\nu\bar{\nu}, \quad H \rightarrow ZZ \rightarrow \ell^+\ell^-q\bar{q}, \quad \text{and} \quad H \rightarrow ZZ \rightarrow \ell^+\ell^-\ell^+\ell^- \text{ samples are also simulated using the same generators as for the } H \rightarrow ZZ \rightarrow \ell^+\ell^-\nu\bar{\nu} \text{ samples. These channels, together with the } H \rightarrow ZZ \rightarrow \tau^+\tau^-\nu\bar{\nu} \text{ final state, which is included in the } H \rightarrow ZZ \rightarrow \ell^+\ell^-\nu\bar{\nu} \text{ samples, contribute to the signal yield and are considered as part of the signal. In particular, } H \rightarrow W^+W^- \rightarrow \ell^+\ell^-\nu\bar{\nu} \text{ decays contribute as much as 70% to the signal expectation after the full selection for } m_{T2} = 200 \text{ GeV, decreasing to 13% at } m_{T2} = 300 \text{ GeV, and to less than 4.5% for Higgs boson masses larger than } 400 \text{ GeV. Statistical independence of the analysis with respect to other ATLAS Higgs boson searches [21–23] is ensured through mutually exclusive selection requirements on the dilepton invariant mass, the number of leptons and the missing transverse momentum in the event.}

The Higgs boson production cross sections and decay branching ratios, as well as their uncertainties, are taken from Refs. [24,25]. The cross sections for the gluon fusion process have been calculated to NLO in QCD in Refs. [26–28], and to next-to-next-to-leading-order (NNLO) in Refs. [29–31]. In addition, QCD soft-gluon resummations, calculated in the NNLL approximation [32], are applied for the gluon fusion process. NLO electroweak (EW) corrections are also applied [33,34]. These results are compiled in Refs. [35–37] and assume factorisation between QCD and EW corrections. The cross sections for VBF processes are calculated with full NLO QCD and EW corrections [38–40], and approximate NNLO QCD corrections are included [41]. The uncertainty in the production cross section due to the choice of QCD scale is typically \( \pm 2\% \) for the gluon fusion process, and \( \pm 1\% \) for the VBF process [24]. The uncertainty in the production cross section due to the parton distribution function (PDF) and \( \alpha_S \) is \( \pm 8\% \) for gluon-initiated processes and \( \pm 4\% \) for quark-initiated processes [42–45]. The Higgs boson decay branching ratio [46] to the four-fermion final state is predicted by PROPHET4f [47,48].

The Higgs boson cross sections are calculated with a zero-width approximation. For the Higgs decay, a relativistic Breit–Wigner line shape is applied at the event-generator level. It has been suggested [25,49–51] that effects related to off-shell Higgs boson production and interference with other SM processes may become sizeable for the highest Higgs boson masses (\( m_H > 400 \) GeV) considered in this search. Currently, in the absence of a full lineshape calculation for the production mechanisms as well as a correct account of the interference with SM ZZ production, an estimate of the possible size of such effects is included as a signal normalisation systematic uncertainty, following a parameterisation as a function of \( m_H: 1.5 \times (m_H/\text{TeV})^{2/3} \), for \( m_H \geq 300 \) GeV [25].

Different event generators are chosen to model a range of important background processes. The ALPGEN generator [52] interfaced to HERWIG [53] for parton showers and hadronisation is used to simulate inclusive W/Z boson backgrounds. MCFM [54] is interfaced to HERWIG and JIMMY [55], used for the production of top-pair, single-top and diboson (WW, WZ and ZZ) backgrounds. PYTHIA is used to simulate \( bb \) and \( cc \) samples as well as alternative samples for the inclusive Z boson and ZZ backgrounds. All simulated background samples are scaled to the highest-precision calculations available for the relevant process. An overview of the relevant predictions and their uncertainties is given in Ref. [56].

Generated events are simulated using the ATLAS detector simulation [57] within the GEANT4 framework [58]. Additional \( pp \) interactions in the same and nearby bunch crossings (pile-up) are included in the simulation. The Monte Carlo (MC) samples are reweighted to reproduce the observed distribution of the number of interactions per bunch crossing in the data.

Data used for the search in the muon channel are collected using a single muon trigger with a \( p_T \) threshold of 18 GeV, while in the electron case a logical OR between a single electron trigger with a threshold varying from 20 to 22 GeV and a dielectron trigger with a \( p_T \) threshold of 12 GeV is used.

The overall trigger efficiencies are estimated from MC events after correction of the simulated lepton trigger efficiencies to those observed in data. For signal events passing the full selection criteria, the trigger efficiency is close to 100% in the electron channel and between 94% and 97% in the muon channel. The systematic uncertainties on the trigger efficiency are negligible when compared to the other selection uncertainties.

During the course of the 2011 data-taking, the average number of pile-up interactions increased due to increased beam currents and stronger beam focusing. This changed the average number of interactions per bunch crossing at the start of a fill, from about six in the earlier periods to about 15 in the later periods. The missing transverse momentum resolution is affected by the level of pile-up, resulting in a significant change in the signal-to-background ratio between the earlier and the later periods. To retain the best sensitivity, the search is therefore split between the earlier (2.3 fb\(^{-1}\)) and the later (2.4 fb\(^{-1}\)) periods, hereafter referred to as the “low pile-up data” and the “high pile-up data”, respectively. The selection is unaltered between the periods.

4. Lepton identification and event selection

Electron candidates consist of clusters of energy deposited in the electromagnetic calorimeter that are associated with tracks reconstructed in the ID. The electron candidates must satisfy a set of identification criteria [59] that require the shower profiles to be consistent with those expected for electromagnetic showers and a well-reconstructed ID track pointing to the corresponding cluster. Furthermore, the electron candidates are required to have \( p_T > 20 \) GeV and pseudorapidity \( |\eta| < 2.47 \). The electron transverse momentum is computed from the cluster energy and the track direction at the interaction point.

Muons are identified by reconstructing tracks in the muon spectrometer. These tracks are then extrapolated back to the beam line to find a matching ID track. Details of muon reconstruction and identification can be found in Ref. [60]. Only muons with \( p_T > 20 \) GeV and \( |\eta| < 2.5 \) are considered.

Jets are used in this analysis to reject backgrounds from events with heavy-quark decays or from events with fake missing transverse momentum \( E_T^{miss} \) due to mis-measured jets. For this purpose, jets are reconstructed from clusters of energy deposits in the calorimeters using the anti-\( k_t \) algorithm [61] with a radius parameter \( R = 0.4 \). Jets are calibrated using \( p_T \) and \( \eta \)-dependent correction factors based on MC simulation and validated with data [62]; this calibration corrects for effects of energy from additional proton–proton interactions. Only jets with \( p_T > 25 \) GeV and \( |\eta| < 2.5 \) are considered.

Two conditions are applied to remove leptons associated with jets, such as those originating from semi-leptonic decays of \( b \)-hadrons. Leptons are not considered in the analysis if the scalar
The dominant contribution to the $E_T^{\text{miss}}$ uncertainty comes from both the impact parameter with respect to the primary vertex of tracks associated to the jet and the presence of displaced secondary vertices associated to the jet’s tracks. Jets with a loosened $p_T$ threshold of 20 GeV and $|\eta| < 2.5$ are considered in the $b$-tag veto. The selection applied on the $b$-tagging discriminant achieves an efficiency of about 85% (50%) for identifying $b$-jets ($c$-jets) [64] and a light-jet rejection factor of about ten [65] in inclusive top-pair events.

To exploit the mass-dependent kinematic features of $H \to ZZ \to 4\ell$ production, the search is subdivided into a low Higgs boson mass region ($m_H < 280$ GeV) and a high Higgs boson mass region ($m_H \geq 280$ GeV), where dedicated selection criteria are applied to two important discriminating variables used to reduce the background contributions: $E_T^{\text{miss}}$ and the azimuthal angle between the two leptons, $\Delta \phi (\ell, \ell)$. Events can contribute to one or both search regions depending on whether they satisfy the selection criteria applied on these variables. Figs. 1(b) and 2 show the distributions of these variables after the application of the $m_{\ell\ell}$ window. Since inclusive $Z$ boson production gives rise to a steeply falling $E_T^{\text{miss}}$ distribution, systematic uncertainties on the $E_T^{\text{miss}}$ reconstruction are particularly important in order to estimate this background correctly.

The dominant contribution to the $E_T^{\text{miss}}$ distribution is the contribution from top quark, which is particularly important in order to estimate this background correctly.
larger average number of interactions per bunch crossing, as is illustrated in Fig. 2. In particular, this increases the background from inclusive Z boson production in the high pile-up periods, mostly in the low mass search region.

Fig. 2 shows that, at high \( E_{\text{miss}} \), the data and the combined background expectation agree within systematic uncertainties. In the low \( m_H \) region, events are required to satisfy \( E_{\text{miss}} > 66 \text{ GeV} \), whilst in the high \( m_H \) region the requirement is \( E_{\text{miss}} > 82 \text{ GeV} \). These selection criteria reduce significantly the backgrounds from processes with no or modest genuine missing transverse momentum originating from unobserved neutrinos. In the low \( m_H \) region a significant fraction of signal events is also removed by this criterion.

The boost of the Z bosons originating from a Higgs boson decay increases with \( m_H \), thus reducing the expected opening angle between the leptons. In the low (high) \( m_H \) region an upper bound, \( \Delta\phi(\ell, \ell) < 2.64 \) (2.25), is therefore applied. In the low \( m_H \) region a lower bound, \( \Delta\phi(\ell, \ell) > 1 \), is also applied, to reduce backgrounds from processes with no or modest genuine missing transverse momentum originating from unobserved neutrinos. In the low \( m_H \) region a significant fraction of signal events is also removed by this criterion.

Finally, to further reduce the background from events with fake \( E_T^{\text{miss}} \) due to mis-measured jets, events are rejected if the azimuthal angle between the missing transverse momentum vector and the direction of the \( Z \to \ell\ell \) boson candidate is \( \Delta\phi(\vec{p}_T^{\text{miss}}, \vec{p}_T^\ell) < 1 \).

The efficiency of the event selection is very similar in the electron and muon channels, ranging from 3.1% for \( m_H = 200 \text{ GeV} \) to about 43% for \( m_H = 600 \text{ GeV} \).

5. Background normalisation and control regions

Standard Model pair-production of Z bosons has a final state identical to the signal, and is therefore expected to survive most of the applied selection criteria and form a continuum in the transverse mass distribution (defined in Section 7). The normalisation for this background is obtained from a calculation including next-to-leading-order terms [66] with an additional 6% term to account for missing quark-box diagrams \((gg \to ZZ)\) [67]. An 11% normalisation uncertainty is assigned to this background. This combines the theoretical uncertainty for Z pair production estimated in Ref. [25] with an additional modelling uncertainty related to the used Monte Carlo model, estimated from comparison to other models. The ZZ background is taken from the MC simulation. A systematic uncertainty to account for shape uncertainties is derived using Pythia as an alternative event generator. WW and WZ backgrounds are normalised in a similar way. For the WZ background the normalisation is verified using a control sample in which the presence of exactly three leptons is required, where the minimum \( p_T \) for the third lepton is 10 GeV. Fig. 3(a) shows the \( E_{\text{miss}} \) distribution in this control region, which is dominated by WZ background for \( E_{\text{miss}} > 40 \text{ GeV} \) and well modelled by the MC simulations, in the high \( E_{\text{miss}} \) region.

The background from top-quark events is taken from the MC prediction. This prediction is verified to agree with data in two
The background from inclusive $Z$ boson production is taken from the MC prediction and verified in a control region populated with events rejected by the $\Delta\phi(p_\text{T}^{\text{miss}}, \vec{p}_\text{T}^{\text{jet}})$ selection criterion after the $E_\text{T}^{\text{miss}}$ requirement. The full $\Delta\phi(p_\text{T}^{\text{miss}}, \vec{p}_\text{T}^{\text{jet}})$ distribution is shown in Fig. 4. At low $\Delta\phi(p_\text{T}^{\text{miss}}, \vec{p}_\text{T}^{\text{jet}})$, where the inclusive $Z$ boson background dominates, a small discrepancy is observed between the data and the expected backgrounds. This discrepancy is within the systematic uncertainties applied on the inclusive $Z$ boson background.

The signal efficiencies and overall background expectations are similar in the electron and muon channels; therefore only combined results are presented. The numbers of candidate $H \rightarrow ZZ \rightarrow e^+e^-\mu^+\mu^-$ events selected in data and the expected yields from signal and background processes are shown in Table 1.

### 6. Systematic uncertainties

The systematic uncertainties applied in this search include experimental uncertainties related to the selection and calibration of electrons, muons and jets, which are also all explicitly propagated to the $E_\text{T}^{\text{miss}}$ calculation. Uncertainties on the $b(c)$-tagging efficiency as well as the light-jet mis-tagging rate are also applied.

Normalisation uncertainties for the signal (gluon fusion $t^+t^-$ and $VBF$ $gg$) [24] and for the diboson backgrounds (11%, see Section 5) are obtained from theory; uncertainties for the inclusive $Z$ boson production (2.5%), top quark production (9%), inclusive $W$ boson production (10%), and for multijet production in the electron channel (50%) are estimated from data. These normalisation uncertainties are not the full systematic uncertainties applied...
Fig. 5. The transverse mass distribution of $H \rightarrow ZZ \rightarrow \ell^+\ell^-\nu\bar{\nu}$ candidates. Figures (a) and (b) refer to the low $m_H$ search region and figures (c) to (f) to the high $m_H$ search region. The top, centre and bottom figures show the expected signal for a Higgs boson with a mass of 200 GeV, 400 GeV and 600 GeV, respectively. Figures on the left correspond to the low pile-up data and figures on the right to the high pile-up data. The hashed area represents the combined systematic uncertainties on the total expected background.

on these backgrounds. Where a normalisation uncertainty is obtained from data, additional detector related systematic uncertainties are applied if they are likely to affect the background in the control region and in the signal region differently. Where a normalisation error is taken from theory, full detector related systematic uncertainties are applied.

For the signal, an additional uncertainty is applied to account for the possible effect of theoretical uncertainties on the acceptance. This uncertainty is obtained from signal samples containing variations in the modelling of initial- and final-state radiation, the renormalisation and factorisation scales, which are varied to half and twice their nominal values, as well as variations in the underlying event tune, which was changed from the default AUE12B [68] to PERUGIA2011 [69], both based on LHC data. An overall uncertainty of 8.4% (3.4%) is assigned in the low (high) mass search regions, obtained by the addition in quadrature of the largest deviations observed for each variation.

The luminosity uncertainty is 3.7% and 4.1% for the low and high pile-up data, respectively, based on the calibration described in Refs. [11,12]. Where appropriate, systematic uncertainties are treated as correlated between the signal and the different background expectations. The total systematic uncertainty on the signal and on each of the background contributions can be seen in Table 1. In most cases the assigned normalisation uncertainties dominate the total systematic uncertainty, except for the inclusive $Z$ boson background, for which the jet energy scale and resolution uncertainties dominate, and the top-quark background for which the $b$-tagging uncertainty dominates. In the low $m_H$ search using the high pile-up data, the uncertainty on the inclusive $Z$ boson background uncertainty dominates the overall uncertainty.
Results of a search for a heavy SM Higgs boson with a mass in the range 200 < \( m_H < 600 \) GeV decaying to \( ZZ \rightarrow \ell^+\ell^-\nu\bar{\nu} \) have been presented. These results are based on a data sample corresponding to an integrated luminosity of 4.7 fb\(^{-1}\) and \( \sqrt{s} = 7 \) TeV, recorded with the ATLAS detector at the LHC. No evidence for a signal is observed and cross section limits are placed over the mass range considered in this search, excluding the production of a SM Higgs boson in the region 319 < \( m_H < 558 \) GeV at the 95% CL.

Acknowledgements

We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently.

We acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWF, Austria; ANAS, Azerbaijan; SSTC, Belarus; CNPq and FAPESP, Brazil; NSERC, NRC and CFI, Canada; CERN; CONICYT, Chile; CAS, MOST and NSFC, China; COLCIENCIAS, Colombia; MSMT CR, MPO CR and VSC CR, Czech Republic; DNRF, DNSRC and Lundbeck Foundation, Denmark; EPLANET and ERC, European Union; IN2P3-CNRS, CEA-DSM/IRFU, France; GNAS, Georgia; BMBF, DFG, HGF, MPG and AvH Foundation, Germany; GSRT, Greece; ISF, MINERVA, GIF, DIP and Benoziyo Center, Israel; INFN, Italy; MEXT and JSPS, Japan; CNRST, Morocco; FOM and NWO, Netherlands; RCN, Norway; MNISW, Poland; GRICES and FCT, Portugal; MERSYS (MECTS), Romania; MES of Russia and ROSATOM, Russian Federation; JINR; MSTD, Serbia; MSSR, Slovakia; ARRS and MVZT, Slovenia; DST/NRF, South Africa; MICINN, Spain; SRC and-wallenberg Foundation, Sweden; SER, SNSF and Cantons of Bern and Geneva, Switzerland; NSC, Taiwan; TAEK, Turkey; STFC, the Royal Society and Leverhulme Trust, United Kingdom; DOE and NSF, United States of America.

The crucial computing support from all WLCG partners is acknowledged gratefully, in particular from CERN and the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CC-IN2P3 (France), KIT/GridKA (Germany), INFN-CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (UK) and BNL (USA) and in the Tier-2 facilities worldwide.

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