Search for a light Higgs boson decaying to long-lived weakly interacting particles in proton-proton collisions at \( \sqrt{s} = 7 \) TeV with the ATLAS detector


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Search for a Light Higgs Boson Decaying to Long-Lived Weakly Interacting Particles in Proton-Proton Collisions at $\sqrt{s} = 7$ TeV with the ATLAS Detector

G. Aad et al.*
(ATLAS Collaboration)
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A search for the decay of a light Higgs boson (120–140 GeV) to a pair of weakly interacting, long-lived particles in 1.94 fb$^{-1}$ of proton-proton collisions at $\sqrt{s} = 7$ TeV recorded in 2011 by the ATLAS detector is presented. The search strategy requires that both long-lived particles decay inside the muon spectrometer. No excess of events is observed above the expected background and limits on the Higgs boson production times branching ratio to weakly interacting, long-lived particles are derived as a function of the particle proper decay length.

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A Higgs boson [1–3] below 140 GeV is particularly sensitive to new physics. Many extensions of the standard model (SM) include neutral, weakly coupled particles that can be long lived [4,5] and to which the Higgs boson may decay. These long-lived particles occur in many models, including gauge-mediated extensions of the minimal supersymmetric standard model [6], minimal supersymmetric standard model with $R$-parity violation [7], inelastic dark matter [8], and the hidden valley (HV) scenario [9].

This Letter presents the first ATLAS search for the Higgs boson decay, $h^0 \rightarrow \pi_\nu \pi_\nu$, to two identical neutral particles ($\pi_\nu$) that have a displaced decay to fermion-antifermion pairs. As a benchmark, we take a HV model [9] in which the SM is weakly coupled, by a heavy communicator particle, to a hidden sector that includes a pseudo-scalar, the $\pi_\nu$. Because of the helicity suppression of pseudoscalar decays to low-mass $f \bar{f}$ pairs, the $\pi_\nu$ decays predominantly to heavy fermions, $b\bar{b}$, $c\bar{c}$, and $\tau^+\tau^-$ in the ratio 85:5:8%. The weak coupling between the two sectors leads the $\pi_\nu$ to have a long lifetime. Other, non-HV models with the identical signature, where the $\pi_\nu$ is replaced with another weakly interacting scalar or pseudoscalar particle, are discussed in Refs. [4,10]. Both Tevatron experiments, CDF and D0, performed similar searches for displaced decays in their respective tracking volumes, which limited the proper decay length range they could explore to a few hundred millimeters [11,12].

In many of these beyond-the-SM scenarios, the lifetime of the neutral states is not specified and can have a very large range. The current search covers a range of expected proper decay lengths extending to about 20 m by exploiting the size and layout of the ATLAS muon spectrometer. Consequently the experimental challenge is to develop signature-driven triggers to select displaced decays throughout the ATLAS detector volume [13].

This analysis requires both $\pi_\nu$ decays to occur near the outer radius of the hadronic calorimeter ($r \sim 4$ m) or in the muon spectrometer (MS). Such decays give a $(\eta, \phi)$ cluster of charged and neutral hadrons in the MS. Requiring both $\pi_\nu$’s to have this decay topology improves background rejection. The analysis uses specialized tracking and vertex reconstruction algorithms, described below, to reconstruct vertices in the MS. The analysis strategy takes advantage of the kinematics of the gluon fusion production mechanism and subsequent two-body decay, $h^0 \rightarrow \pi_\nu \pi_\nu$, which results in events with back-to-back $\pi_\nu$’s, by requiring two well-separated vertices $[\Delta R = \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2} > 2]$ [14] in the MS.

The data used in this analysis were collected in the first half of 2011 with the LHC operating at 7 TeV. Applying beam, detector, and data quality requirements resulted in a total integrated luminosity of 1.94 fb$^{-1}$. The integrated luminosity has a relative uncertainty of 3.7% [15,16].

Signal Monte Carlo (MC) samples were generated using PYTHIA [17,18] to simulate gluon fusion production ($gg \rightarrow h^0$) and the decay of the Higgs boson ($h^0 \rightarrow \pi_\nu \pi_\nu$). Four samples were generated: $m_{h^0} = 120$ and 140 GeV and for each $m_{h^0}$, two $\pi_\nu$ masses of 20 and 40 GeV. The predicted Higgs boson production cross sections [19] are $\sigma(m_{h^0} = 120$ GeV) = $16.6^{+3.3}_{-2.5}$ pb and $\sigma(m_{h^0} = 140$ GeV) = $12.1^{+2.1}_{-1.8}$ pb, and the branching ratio (BR) for $h^0 \rightarrow \pi_\nu \pi_\nu$ is assumed to be 100%. The response of the ATLAS detector was modeled with GEANT4 [20,21].

The effect of multiple $pp$ collisions occurring during the same bunch crossing (pileup) was simulated by superimposing several minimum bias events on the signal event. The MC events were weighted so that the pileup in the simulation agrees with pileup conditions found in data.

ATLAS is a multipurpose detector [22] consisting of an inner tracking detector (ID) surrounded by a superconducting...
solenoid that provides a 2 T field, electromagnetic and hadronic calorimeters and a MS with a toroidal magnetic field. The ID, consisting of silicon pixel and strip detectors and a straw tube tracker, provides precision tracking of charged particles for $|\eta| \leq 2.5$. The calorimeter system covers $|\eta| \leq 4.9$ and has 9.7 interaction lengths at $\eta = 0$. The MS consists of a barrel and two forward spectrometers, each with 16 $\phi$ sectors instrumented with detectors for first level triggering and precision tracking detectors for muon momentum measurement. Each spectrometer has three stations along the muon flight path: inner, middle, and outer. In the barrel, the stations are located at radii of $\sim 4.5$, 7, and 10 m, while in the forward MS, they are located at $|z| \sim 7.5$, 14, and 20 m. This analysis uses muon tracking for $|\eta| \leq 2.4$, where each station is instrumented with two multilayers of precision tracking chambers, monitored drift tubes (MDTs). It also utilizes level 1 [23] (L1) muon triggering in the barrel MS ($|\eta| \leq 1$). The trigger chambers are located in the middle and outer stations. The L1 muon trigger requires hits in the middle station to create a low $p_T$ muon region of interest (RoI) or hits in both the middle and outer stations for a high $p_T$ RoI. The muon RoIs have a spatial extent of $0.2 \times 0.2$ in $\Delta \eta \times \Delta \phi$ and are limited to two RoIs per sector.

A dedicated, signature-driven trigger, the muon RoI cluster trigger [13], was developed to trigger on events with a $\pi^0$ decaying in the MS. It selects events with a cluster of three or more muon RoIs in a $\Delta R = 0.4$ cone in the MS barrel trigger chambers. This trigger configuration implies that one $\pi^0$ must decay in the barrel spectrometer, while the second $\pi^0$ may decay either in the barrel or the forward spectrometer. With this trigger, it is possible to trigger on $\pi^0$ decays at the outer radius of the hadronic calorimeter and in the MS with high efficiency. The backgrounds of punch-through jets [24] and muon bremsstrahlung are suppressed by requiring no calorimeter jets with $E_T \geq 30$ GeV in a cone of $\Delta R = 0.7$ and no ID tracks with $p_T \geq 5$ GeV within a region of $\Delta \eta \times \Delta \phi = 0.2 \times 0.2$ around the RoI cluster center. This isolation criteria result in a negligible loss in the simulated signal while significantly reducing the backgrounds.

As depicted in Fig. 1(a) [25], MC studies show the RoI cluster trigger is $\sim 30\%-50\%$ efficient in the region from 4 to 7 m. The $\pi^0$'s that decay beyond a radius of $\sim 7$ m do not leave hits in the trigger chambers located at $\sim 7$ m, while the $\pi^0$ decays that occur before $r \sim 4$ m are located in the calorimeter and do not produce sufficient activity in the MS to pass the muon RoI cluster trigger. The $m_{\pi^0} = 120$ GeV and $m_{\pi^0} = 40$ GeV sample has a relatively lower efficiency because the $\pi^0$'s have a lower boost and arrive later at the MS. As a result, the trigger signal may be associated with the incorrect bunch crossing, in which case the event is lost.

The systematic uncertainty of the muon RoI cluster trigger efficiency is evaluated on data using a sample of events containing a punch-through jet. This sample of events is similar to signal events as it contains both low energy photons and charged hadrons in a localized region of the MS. These punch-through jets are selected to be in the barrel calorimeter ($|\eta| \leq 1.4$), have $E_T \geq 20$ GeV, and at least four tracks in the ID, each with $p_T \geq 1$ GeV, and have at least 20 GeV of missing transverse momentum aligned with the jet. To ensure significant activity in the MS, the jet is required to contain at least 300 MDT hits in a cone of $\Delta R = 0.6$, centered around the jet axis [26]. The muon RoI cluster trigger algorithm was run in the vicinity of the punch-through jet for both data and MC events. The distribution of RoIs contained in the cluster for data and MC events, normalized to the number of data events, is shown in Fig. 2. The shapes of the distribution match well between data and MC events. A horizontal line fit to the ratio, as a function of $N_{\text{RoI}}$, yields $1.14 \pm 0.09$, and 14% is taken as the systematic uncertainty. The effects of uncertainties in the jet energy scale (JES) [27], in the initial state radiation (ISR) spectrum [28], and in the amount of pileup were found to be

![ATLAS Simulation](https://example.com/plot.png)
negligible when varying these quantities by their uncertainties.

A specialized tracking and vertex reconstruction algorithm was developed to identify π0’s that decay inside the MS. The decay of a π0 results in a high multiplicity of low p_T particles (1 \( \leq p_T \leq 5 \) GeV) containing \( \sim 10 \) charged particles and \( \sim 5 \) π0’s clustered in a small \( \Delta R \) region of the spectrometer. The π0’s that decay before the last sampling layer of the hadronic calorimeter do not produce a significant number of tracks in the MS. Thus, detectable decay vertices must be located in the region between the outer radius of the hadronic calorimeter and the middle station of the MS. Over a wide range of acceptance in the barrel MS, the total amount of material traversed is roughly 1.3 radiation lengths [22]; therefore, as a consequence of the \( \sim 5 \) π0’s produced in signal events, large electromagnetic showers accompany the \( \sim 10 \) charged particles from π0 decays. The resulting MS environment contains, on average, approximately 800 MDT hits, of which \( \sim 75\% \) are from the electromagnetic showers.

The design of the muon chambers [22] is exploited in order to reconstruct tracks in this busy environment. The separation of the two multilayers inside a single muon chamber provides a powerful tool for track pattern recognition. This separation provides enough of a lever arm to allow, in the barrel, a momentum measurement with acceptable resolution for tracks up to approximately 10 GeV [29]. In the forward spectrometers, the muon chambers are outside the magnetic field region; therefore, it is not possible to measure the track momentum inside of a single chamber. In both cases, the tracklets used in the vertex reconstruction are formed using hits in single muon chambers.

The MS vertex algorithm begins by grouping the tracklets using a simple cone algorithm with \( \Delta R = 0.6 \). In the barrel, the tracklets are extrapolated through the magnetic field, and the vertex position is reconstructed as the point in 

\((r, z)\) that uses the largest number of tracklets to reconstruct a vertex with a \( \chi^2 \) probability greater than 5%. In the forward spectrometer, the reconstructed tracklets do not have a measurement of the momentum; therefore, the vertex is found using a least squares regression that assumes the tracklets are straight lines. Vertices are required to be reconstructed using at least three tracklets, point back to the interaction point (IP) [30] and have | \( \eta \) | \( \leq 2.2 \). After requiring the MS vertex to be separated from ID tracks with \( p_T \geq 5 \) GeV and jets with \( E_T \geq 15 \) GeV by \( \Delta R = 0.4 \) and \( \Delta R = 0.7 \), respectively, the algorithm has an efficiency of \( \sim 40\% \) in signal MC events throughout the barrel region (4 \( \leq r \leq 7.5 \) m) and a resolution of 20 cm in \( z \), 32 cm in \( r \), and 50 mrad in \( \phi \). In the forward spectrometer, the algorithm is \( \sim 40\% \) efficient in the region 8 \( \leq |z| \leq 14 \) m. Figure 1(b) [25] shows the vertex reconstruction efficiency for the barrel reconstruction algorithm in MC signal events that passed the muon RoI cluster trigger.

The MC description of hadrons and photons in the MS was validated on the same sample of events containing a punch-through jet used to evaluate the trigger performance. The fraction of these jets that produce a MS vertex was compared in data and QCD dijet MC events. Table I shows the fraction of punch-through jets that produce a vertex in data and MC events as a function of the number of MDT hits in a cone of \( \Delta R = 0.6 \) around the jet axis. The data-to-MC ratio is fit to a flat distribution that yields a ratio consistent with unity with a 15% statistical uncertainty, which is taken to be the systematic uncertainty in the vertex reconstruction efficiency. The systematic uncertainties arising from the JES, ISR spectrum, and the amount of pileup were estimated by varying these quantities by their uncertainties and calculating the change in the vertex reconstruction efficiency. The total systematic uncertainty of 16% for the efficiency of reconstructing a vertex is the sum in quadrature of the uncertainties in the efficiency of the isolation criteria due to varying the JES, ISR, and pileup (3%, 3%, and 2%, respectively) and the uncertainty in the comparison of data and MC events (15%).

The final event selection requires two good MS vertices separated by \( \Delta R > 2 \). The background due to events with two jets, both of which punch through the calorimeter, is a negligible contribution to the total background due to the tight isolation criteria applied to each vertex. The background is calculated using a fully data-driven method by

\[
\begin{array}{|c|c|c|}
\hline
\text{Number of MDT hits} & \text{QCD dijet Monte Carlo} & \text{Data} \\
\hline
300 \leq N_{\text{MDT}} < 400 & 10.1 \pm 2.2\% & 9.1 \pm 0.5\% \\
400 \leq N_{\text{MDT}} < 500 & 9.2 \pm 2.8\% & 10.5 \pm 0.7\% \\
500 \leq N_{\text{MDT}} < 600 & 13.1 \pm 5.4\% & 13.0 \pm 0.9\% \\
N_{\text{MDT}} \geq 600 & 16.5 \pm 4.5\% & 16.7 \pm 0.7\% \\
\hline
\end{array}
\]

TABLE I. Fraction of punch-through jets that have a reconstructed vertex in the muon spectrometer for varying numbers of MDT hits for data and QCD Monte Carlo events.
measuring the probability for a random event to contain an MS vertex ($P_{\text{vertex}}$) and the probability of reconstructing a vertex given that the event passed the RoI cluster trigger ($P_{\text{reco}}$). Because $P_{\text{vertex}}$ and $P_{\text{reco}}$ are measured in data, they incorporate backgrounds from cosmic showers, beam halo, and detector noise. The background is calculated as

$$N_{\text{fake}}(2 \text{ MS vertex}) = N(\text{MS vertex, 1 trig})P_{\text{vertex}} + N(\text{MS vertex, 2 trig})P_{\text{reco}}.$$  

$N(\text{MS vertex, 1 trig})$ is the number of events with a single muon RoI cluster trigger object and an isolated MS vertex. $N(\text{MS vertex, 2 trig})$ is the number of events with an isolated vertex and a second RoI cluster trigger object. The first term in the equation is the expected number of background events with one vertex that randomly contain a second vertex. $P_{\text{reco}}$ is the probability to reconstruct a vertex given there was an RoI cluster trigger; thus, the second term in the equation is the expected number of events with two RoI clusters that have two vertices in the MS. $P_{\text{vertex}}$ was measured using zero bias data [31] to be $(9.7 \pm 6.9) \times 10^{-7}$, and $P_{\text{reco}}$ was measured using the events that pass the muon RoI cluster trigger to be $(1.11 \pm 0.01) \times 10^{-2}$. The expected signal would cause, at most, a relative change in $P_{\text{reco}}$ of $\sim 1\%$. $P_{\text{reco}}$ was also measured using a sample of recorded events when there were no collisions. In this sample of noncollision background events, $P_{\text{reco}}$ was measured to be $(7.0 \pm 0.6) \times 10^{-3}$. For calculating the background, the larger value of $P_{\text{reco}}$ $(1.11 \times 10^{-2})$ is taken since it gives a conservative estimate of the background. $N(\text{MS vertex, 1 trig})$ and $N(\text{MS vertex, 2 trig})$ are $15,543$ and $1$, respectively. Therefore, the background is calculated to be $0.03 \pm 0.02$ events.

No events in the data sample pass the selection requiring two isolated, back-to-back vertices in the muon spectrometer. Since no significant excess over the background prediction is found, exclusion limits for $\sigma_{h^0} \times \text{BR}(h^0 \rightarrow \pi^0 \pi^0)$ are set by rejecting the signal hypothesis at the 95% confidence level (CL) using the CLs procedure [32]. Figure 3 shows the 95% CL upper limit on $\sigma_{h^0} \times \text{BR}(h^0 \rightarrow \pi^0 \pi^0)/\sigma_{\text{SM}}$ as a function of the $\pi^0$ proper decay length ($c\tau$) in multiples of the SM Higgs boson cross section, $\sigma_{\text{SM}}$. As expected, the Higgs boson and $\pi^0$ mass combinations with the largest boosts leading to larger $\beta y c \tau$ have the smallest exclusion limits.

In 1.94 fb$^{-1}$ of $pp$ collision data at a center-of-mass energy of $7$ TeV, there is no evidence of an excess of events containing two isolated, back-to-back vertices in the ATLAS muon spectrometer. Using the model of a light Higgs boson decaying to weakly interacting, long-lived pseudoscalars, limits have been placed on the pseudoscalar proper decay length. Table II shows the broad range of $\pi^0$ proper decay lengths that have been excluded at the 95% CL, assuming 100% branching ratio for $h^0 \rightarrow \pi^0 \pi^0$. These limits also apply to models in which the Higgs boson decays to a pair of weakly interacting scalars that, in turn, decay to heavy quark pairs.

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<th>$m_{h^0}$ (GeV)</th>
<th>$m_{\pi^0}$ (GeV)</th>
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<td>$0.50 &lt; c\tau &lt; 20.65$ m</td>
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<tr>
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<td>$1.10 &lt; c\tau &lt; 26.75$ m</td>
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![ATLAS](https://example.com/atlas.png)

FIG. 3 (color online). Observed 95% upper limits on the process $h^0 \rightarrow \pi^0 \pi^0$, vs the $\pi^0$ proper decay length, expressed as a multiple of the SM cross section for Higgs boson production. Exclusion limits assume 100% branching ratio for the Higgs boson decaying to $\pi^0$s.
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[12] The ATLAS Collaboration uses a right-handed coordinate system with its origin at the nominal IP in the center of the detector and the z axis coinciding with the beam pipe. The x axis points from the IP to the center of the LHC ring, and the y axis points upward. Cylindrical coordinates (r, φ) are used in the transverse plane, with φ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle θ as η = −ln tan(θ/2).

(ATLAS Collaboration)

1 University at Albany, Albany, New York, USA
2 Department of Physics, University of Alberta, Edmonton, Alberta, Canada
3 Department of Physics, Ankara University, Ankara, Turkey
4 Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey
5 Turkish Atomic Energy Authority, Ankara, Turkey
6 LAPP, CNRS/IN2P3 and Université de Savoie, Annecy-le-Vieux, France
7 High Energy Physics Division, Argonne National Laboratory, Argonne, Illinois, USA
8 Department of Physics, University of Arizona, Tucson, Arizona, USA
9 Department of Physics, The University of Texas at Arlington, Arlington, Texas, USA
10 Physics Department, National Technical University of Athens, Zografou, Greece
11 Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
12 Institut de Fisica d’Altes Energies and Departament de Fisica de la Universitat Autonoma de Barcelona and ICREA, Barcelona, Spain
13 Institute of Physics, University of Belgrade, Belgrade, Serbia
14 Vinca Institute of Nuclear Sciences, University of Belgrade, Belgrade, Serbia
15 Department for Physics and Technology, University of Bergen, Bergen, Norway
16 Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, California, USA
17 Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
18 School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
19 Department of Physics, Bogazici University, Istanbul, Turkey
20 Division of Physics, Dogus University, Istanbul, Turkey
21 Department of Physics, Gaziantep University, Gaziantep, Turkey
22 Department of Physics, Istanbul Technical University, Istanbul, Turkey
23 INFN Sezione di Bologna, Italy
24 Dipartimento di Fisica, Università di Bologna, Bologna, Italy
25 Physikalisches Institut, University of Bonn, Bonn, Germany
26 Department of Physics, Boston University, Boston, Massachusetts, USA
27 Department of Physics, Brandeis University, Waltham, Massachusetts, USA
28 Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil
29 Federal University of Juiz de Fora (UFJF), Juiz de Fora, Brazil
30 Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei, Brazil
31 Instituto de Fisica, Universidade de Sao Paulo, Sao Paulo, Brazil
32 Physics Department, Brookhaven National Laboratory, Upton, New York, USA
33 National Institute of Physics and Nuclear Engineering, Bucharest, Romania
34 University Politehnica Bucharest, Bucharest, Romania
35 West University in Timisoara, Timisoara, Romania
36 Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
37 Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
38 Department of Physics, Carleton University, Ottawa, Ontario, Canada
39 CERN, Geneva, Switzerland
40 Enrico Fermi Institute, University of Chicago, Chicago, Illinois, USA


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<td>Belarus, Minsk, Republic of</td>
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<tr>
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<td>Dipartimento di Scienze Fisiche, Università di Napoli, Napoli, Italy</td>
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<td>Department of Physics and Astronomy, University of New Mexico, Albuquerque,</td>
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<td>Institute for Mathematics, Astrophysics and Particle Physics, Radboud</td>
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<td>Amsterdam, Amsterdam, Netherlands</td>
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</tr>
<tr>
<td>Department of Physics, Northern Illinois University, DeKalb, Illinois, USA</td>
<td></td>
</tr>
<tr>
<td>Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia</td>
<td></td>
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<tr>
<td>Department of Physics, New York University, New York, New York, USA</td>
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<tr>
<td>Homer L. Dodge Department of Physics and Astronomy, University of</td>
<td>Oklahoma, USA</td>
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<td>Oklahoma, USA</td>
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<tr>
<td>Department of Physics, Oklahoma State University, Stillwater, Oklahoma,</td>
<td>USA</td>
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<td>Palacký University, RCPTM, Olomouc, Czech Republic</td>
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<td>Center for High Energy Physics, University of Oregon, Eugene, Oregon, USA</td>
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<td>LAE, Univ. Paris-Sud and CNRS/IN2P3, Orsay, France</td>
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<tr>
<td>Graduate School of Science, Osaka University, Osaka, Japan</td>
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<td>Department of Physics, University of Oslo, Oslo, Norway</td>
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<td>Department of Physics, Oxford University, Oxford, United Kingdom</td>
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<td>INFN Sezione di Pavia, Italy</td>
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<td>Dipartimento di Fisica, Università di Pavia, Pavia, Italy</td>
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<td>Department of Physics, University of Pennsylvania, Philadelphia,</td>
<td>Pennsylvania, USA</td>
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<td>Petersburg Nuclear Physics Institute, Gatchina, Russia</td>
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<td>INFN Sezione di Pisa, Italy</td>
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<td>Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy</td>
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<tr>
<td>Department of Physics and Astronomy, University of Pittsburgh,</td>
<td>Pittsburgh, Pennsylvania, USA</td>
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<tr>
<td>Laboratorio de Instrumentación e Física Experimental de Partículas-LIP,</td>
<td>Lisboa, Portugal</td>
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<tr>
<td>Departamento de Física Teorica y del Cosmos and CAFPE, Universidad de</td>
<td>Granada, Spain</td>
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<tr>
<td>Institute of Physics, Academy of Sciences of the Czech Republic, Praha,</td>
<td>Czech Republic</td>
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<tr>
<td>Charles University in Prague, Praha, Czech Republic</td>
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<tr>
<td>Czech Technical University in Prague, Praha, Czech Republic</td>
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<tr>
<td>State Research Center Institute for High Energy Physics, Protvino, Russia</td>
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<tr>
<td>Particle Physics Department, Rutherford Appleton Laboratory, Didcot,</td>
<td>United Kingdom</td>
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<td>United Kingdom</td>
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<tr>
<td>Physics Department, University of Regina, Regina, Saskatchewan, Canada</td>
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<tr>
<td>INFN Sezione di Roma I, Italy</td>
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<tr>
<td>Dipartimento di Fisica, Università La Sapienza, Roma, Italy</td>
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<td>INFN Sezione di Roma Tor Vergata, Italy</td>
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<td>Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy</td>
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<td>INFN Sezione di Roma Tre, Italy</td>
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<td>Dipartimento di Fisica, Università Roma Tre, Roma, Italy</td>
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</tbody>
</table>
Department of Physics and Astronomy, University of Sheikhfi, Sheikhfi, United Kingdom

Department of Physics, Shanghai University, Shanghái, China

Department of Physics, Simon Fraser University, Burnaby, British Columbia, Canada

DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat à l’Énergie Atomique), Gif-sur-Yvette, France

Department of Nuclear Physics, Technion: Israel Institute of Technology, Haifa, Israel

Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece

International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo, Japan

Graduate School of Science and Technology, University of Tokyo, Tokyo, Japan

Department of Physics, University of Tokyo, Tokyo, Japan

The Oskar Klein Centre, Stockholm, Sweden

Physics Department, Royal Institute of Technology, Stockholm, Sweden

Department of Physics & Astronomy, Stony Brook University, Stony Brook, New York, USA

Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom

School of Physics, University of the Witwatersrand, Johannesburg, South Africa

School of Physics, University of the Witwatersrand, Johannesburg, South Africa

Department of Physics, Stockholm University, Sweden

The Oskar Klein Centre, Stockholm, Sweden

Physics Department, Royal Institute of Technology, Stockholm, Sweden

Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic

Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava, Slovak Republic

Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece

International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo, Japan

Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan

Department of Physics, Tokyo Institute of Technology, Tokyo, Japan

Department of Physics, University of Toronto, Toronto ON, Canada

TRIUMF, Vancouver, British Columbia, Canada

Department of Physics and Astronomy, University of British Columbia, Vancouver, British Columbia, Canada

Department of Physics and Astronomy, University of California Irvine, Irvine, California, USA

INFN Gruppo Collegato di Udine, Italy

ICTP, Trieste, Italy

Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy

Department of Physics, University of Illinois, Urbana, Illinois, USA

Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden

Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain

Department of Physics, University of British Columbia, Vancouver, British Columbia, Canada

Department of Physics and Astronomy, University of Victoria, Victoria, British Columbia, Canada

Waseda University, Tokyo, Japan

Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel

Department of Physics, University of Wisconsin, Madison Wisconsin, USA

Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany

Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany

Department of Physics, Yale University, New Haven, Connecticut, USA

Yerevan Physics Institute, Yerevan, Armenia

Domaine scientifique de la Doua, Centre de Calcul CNRS/IN2P3, Villeurbanne Cedex, France

Deceased.

Also at Laboratorio de Instrumentacao e Física Experimental de Partículas-LIP, Lisboa, Portugal.

Also at Faculdade de Ciencias and CFNUL, Universidade de Lisboa, Lisboa, Portugal.

Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom.
eAlso at TRIUMF, Vancouver BC, Canada.
fAlso at Department of Physics, California State University, Fresno CA, USA.
gAlso at Novosibirsk State University, Novosibirsk, Russia.
hAlso at Fermilab, Batavia IL, USA.
iAlso at Department of Physics, University of Coimbra, Coimbra, Portugal.
jAlso at Università di Napoli Parthenope, Napoli, Italy.
kAlso at Institute of Particle Physics (IPP), Canada.
lAlso at Department of Physics, Middle East Technical University, Ankara, Turkey.
mAlso at Louisiana Tech University, Ruston LA, USA.

Also at Department of Physics and Astronomy, University College London, London, United Kingdom.
Also at Group of Particle Physics, University of Montreal, Montreal QC, Canada.
Also at Department of Physics, University of Cape Town, Cape Town, South Africa.
Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.
Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany.

Also at Manhattan College, New York NY, USA.
Also at School of Physics, Shandong University, Shandong, China.
Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France.
Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China.
Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan.
Also at DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat a l’Energie Atomique), Gif-sur-Yvette, France.
Also at Section de Physique, Université de Genève, Geneva, Switzerland.
Also at Departamento de Física, Universidade de Minho, Braga, Portugal.

Also at Department of Physics and Astronomy, University of South Carolina, Columbia SC, USA.
Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.
Also at California Institute of Technology, Pasadena CA, USA.
Also at Institute of Physics, Jagiellonian University, Krakow, Poland.
Also at Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China.
Also at Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom.
Also at Department of Physics, Oxford University, Oxford, United Kingdom.

Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.
Also at Department of Physics, The University of Michigan, Ann Arbor MI, USA.
Also at Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France.