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Search for extra dimensions using diphoton events in 7 TeV proton–proton collisions with the ATLAS detector

ATLAS Collaboration

A R T I C L E   I N F O

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1. Introduction

The enormous difference between the Planck scale and the electroweak scale is known as the hierarchy problem. A prominent class of new physics models addresses the hierarchy problem through the existence of extra spatial dimensions. In this Letter, we search for evidence of extra dimensions within the context of the models of Arkani-Hamed, Dimopoulos, and Dvali (ADD) [1] and of Randall and Sundrum (RS) [2]. In these models, gravity can propagate in the higher-dimensional bulk, giving rise to a so-called Kaluza–Klein (KK) tower of massive spin-2 graviton excitations (KK gravitons, \(G\)). Due to their couplings to Standard Model (SM) particle–antiparticle pairs, KK gravitons can be investigated in proton–proton (pp) collisions at the Large Hadron Collider (LHC) via a variety of processes, including virtual graviton exchange as well as direct graviton production through gluon–gluon fusion or quark–antiquark annihilation.

The ADD model [1] postulates the existence of \(n\) flat additional spatial dimensions compactified with radius \(R\), in which only gravity propagates. The fundamental Planck scale in the \((4+n)\)-dimensional spacetime, \(M_P\), is related to the apparent scale \(M_{\text{Pl}}\) by Gauss’ law: \(M_P^4 = M_{\text{Pl}}^{n+2} R^n\), where \(M_{\text{Pl}} = M_P/\sqrt{8\pi}\) is the reduced Planck scale. The mass splitting between subsequent KK states is of order \(1/R\). In the ADD model, resolving the hierarchy problem requires typically small values of \(1/R\), giving rise to an almost continuous spectrum of KK graviton states.

While processes involving direct graviton emission depend on \(M_D\), effects involving virtual gravitons depend on the ultraviolet cutoff of the KK spectrum, denoted \(M_S\). The effects of the extra dimensions are typically parametrized by \(\eta_G = F/M_S^2\), where \(\eta_G\) describes the strength of gravity in the presence of the extra dimensions and \(F\) is a dimensionless parameter of order unity reflecting the dependence of virtual KK graviton exchange on the number of extra dimensions. Several theoretical formalisms exist in the literature, using different definitions of \(F\) and, consequently, of \(M_S\):

\[
F = \begin{cases} 
1 & \text{(GRW)} \ [3]; \\
\log\left(\frac{M_{\text{Pl}}^2}{s}\right) & n = 2, \\
\frac{2}{n-2} & n > 2 \ (\text{HLZ}) \ [4]; \\
\pm \frac{2}{\pi} & (\text{Hewett}) \ [5]; 
\end{cases}
\]

(1)

\[
F = \begin{cases} 
1 & \text{(GRW)} \ [3]; \\
\log\left(\frac{M_{\text{Pl}}^2}{s}\right) & n = 2, \\
\frac{2}{n-2} & n > 2 \ (\text{HLZ}) \ [4]; \\
\pm \frac{2}{\pi} & (\text{Hewett}) \ [5]; 
\end{cases}
\]

(2)

\[
F = \pm \frac{2}{\pi} \ (\text{Hewett}) \ [5];
\]

(3)

where \(s\) is the center-of-mass energy of the parton–parton collision. Effects due to ADD graviton exchange would be evidenced by a non-resonant deviation from the SM background expectation. Collider searches for ADD virtual graviton effects have been performed at HERA [6], LEP [7], the Tevatron [8], and the LHC [9,10]. Recent diphoton results from CMS are the most restrictive so far, setting limits on \(M_S\) in the range of 2.3–3.8 TeV [10].

The RS model [2] posits the existence of a fifth dimension with “warped” geometry, bounded by two \((3+1)\)-dimensional branes, with the SM fields localized on the so-called TeV brane and gravity originating on the other, dubbed the Planck brane, but capable
of propagating in the bulk. Mass scales on the TeV brane, such as
the Planck mass describing the observed strength of gravity, cor-
respond to mass scales on the Planck brane as given by $M_Q = M_{Pl}e^{-k r_c}$, where $k$ and $r_c$ are the curvature scale and compac-
tification radius of the extra dimension, respectively. The observed
hierarchy of scales can therefore be naturally reproduced in this
model, if $k r_c \approx 12$ [11]. KK gravitons in this model would have
a mass splitting of order 1 TeV and would appear as new reso-
nances. The phenomenology can be described in terms of the mass
of the lightest KK graviton excitation ($m_{\gamma}$) and the dimensionless
coupling to the SM fields, $\frac{k}{M_{Pl}}$. It is theoretically preferred [11]
for $\frac{k}{M_{Pl}}$ to have a value in the range from 0.01 to 0.1. The most
stringent experimental limits on RS gravitons are from the LHC.
For $k/M_{Pl} = 0.1 \sim 1$ fb$^{-1}$ ATLAS results from $G \rightarrow e \mu$ exclude
gravitons below 1.63 TeV [12], assuming leading order (LO) cross
section predictions, and a recent 2.2 fb$^{-1}$ $G \rightarrow \gamma \gamma$ result from
CMS excludes gravitons below 1.84 TeV [10], using next-to-leading
order (NLO) cross section values. These results have surpassed the
limits from searches at the Tevatron [13] and earlier searches at
the LHC [14].

The diphoton final state provides a sensitive channel for this
search due to the clean experimental signature, excellent diphoton
mass resolution, and modest backgrounds, as well as a branching
ratio for graviton decay to diphotons that is twice the value of
that for graviton decay to any individual charged-lepton pair. In
this Letter, we report on a search in the diphoton final state for
evidence of extra dimensions, using a data sample corresponding to
an integrated luminosity of 2.12 fb$^{-1}$ of $\sqrt{s} = 7$ TeV $pp$ colli-
sions, recorded during 2011 with the ATLAS detector at the LHC.
The measurement of the diphoton invariant mass spectrum is in-
terpreted in both the ADD and RS scenarios.

2. The ATLAS detector

The ATLAS detector [15] is a multipurpose particle physics in-
strument with a forward–backward symmetric cylindrical geomet-
try and near 4$\pi$ solid angle coverage.1 Closest to the beamline are
tracking detectors to measure the trajectories of charged particles,
including layers of silicon-based detectors as well as a transition
radiation tracker using straw-tube technology. The tracker is sur-
rounded by a thin solenoid that provides a 2 T magnetic field for
momentum measurements. The solenoid is surrounded by a hermic
calorimeter system, which is particularly important for this analy-
sis. A system of liquid-argon (LAr) sampling calorimeters is di-
vided into a central barrel calorimeter and two endcap calorimeters,
each housed in a separate cryostat. Fine-grained LAr electromagnetic (EM) calorimeters, segmented in three longitudinal
layers, are used to precisely measure the energies of electrons,
positrons and photons for $|\eta| < 3.2$. Most of the EM shower en-
ergy is collected in the second layer, which has a granularity of
$\Delta \eta \times \Delta \phi = 0.025 \times 0.025$. The first layer is segmented into eight strips per middle-layer cell in the $\eta$ direction, extending over
four middle-layer cells in $\phi$, designed to separate photons from
$\pi^0$ mesons. A presampler, covering $|\eta| < 1.81$, is used to correct for
energy lost upstream of the calorimeter. The regions spanning
1.5 < $|\eta|$ < 4.9 are instrumented with LAr calorimetry also for
hadronic measurements, while an iron-scintillator tile calorime-
ter provides hadronic coverage in the range $|\eta| < 1.7$. A muon
spectrometer consisting of three superconducting toroidal magnet
systems, tracking chambers, and detectors for triggering lies out-
side the calorimeter system.

3. Trigger and data selection

The analysis uses data collected between March and September
2011 during stable beam periods of 7 TeV $pp$ collisions. Selected
events had to satisfy a trigger requiring at least two photon can-
didates with transverse energy $E_T > 20$ GeV and satisfying a set
of requirements, referred to as the “loose” photon definition [16],
which includes requirements on the leakage of energy into the
hadronic calorimeter as well as on variables that require the trans-
verse width of the shower, measured in the second EM calorimeter,
layer, be consistent with the narrow width expected for an EM
shower. The loose definition is designed to have high photon effi-
ciency, albeit with reduced background rejection. The trigger was
essentially fully efficient for high mass diphoton events passing the
final selection requirements.

Events were required to have at least one primary collision ver-
tex, with at least three reconstructed tracks. Selected events had
to have at least two photon candidates, each with $E_T > 25$ GeV
and pseudorapidity $|\eta| < 2.37$, with the exclusion of $1.37 <
|\eta| < 1.52$, the transition region between the barrel and endcap
calorimeters. As described in more detail in Ref. [16], photon can-
didates included those classified as unconverted photons, with no
associated track, or photons which converted to electron–positron
pairs, with one or two associated tracks. The two photons were re-
quired to satisfy several quality criteria and to lie outside detector
regions where their energy was not measured in an optimal way.
The two photon candidates each had to satisfy a set of stricter
requirements, referred to as the “tight” photon definition [16], which
included a more stringent selection on the shower width in the
second EM layer and additional requirements on the energy distri-
bution in the first EM calorimeter layer. The tight photon definition
was designed to increase the purity of the photon selection sam-
ples by rejecting most of the remaining jet background, including
jets with a leading neutral hadron (mostly $\pi^0$ mesons) that decay to
a pair of collimated photons.

The isolation transverse energy $E_{Tiso}$ for each photon was calcu-
lated [16] by summing over the cells of both the EM and hadronic
calorimeters that surround the photon candidate within an angular
cone of radius $\Delta R = \sqrt{(\eta - \eta')^2 + (\phi - \phi')^2} < 0.4$, after remov-
ing a central core that contains most of the energy of the photon.
To reduce the jet background further, an isolation requirement was
applied, requiring that each of the two leading photons satisfied
$E_{Tiso} < 5$ GeV. An out-of-core energy correction was applied, to
make $E_{Tiso}$ essentially independent of $E_T$. An ambient energy cor-
rection, based on the measurement of low transverse momentum
jets [17], was also applied, on an event-by-event basis, to remove the
contributions from the underlying event and from “pileup”,
which results from the presence of multiple $pp$ collisions within
the same or nearby bunch crossings.

For events with more than two photon candidates passing all
the selection requirements, the two photons with the highest $E_T$
values were considered. The diphoton invariant mass had to exceed
140 GeV. A total of 6846 events were selected.

4. Monte Carlo simulation studies

Monte Carlo (MC) simulations were performed to study the
detector response for various possible signal models, as well as
to perform some SM background studies. All MC events were
simulated [18] with the ATLAS detector simulation based on
GEANT4 [19] and using ATLAS parameter tunes [20], and were pro-
cessed through the same reconstruction software chain as used

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1 ATLAS uses a right-handed coordinate system with its origin at the nominal in-
teraction point in the centre of the detector and the z-axis along the beam pipe.
Cylindrical coordinates ($r, \phi$) are used in the transverse plane, $\phi$ being the azi-
muthal angle around the beam pipe. The pseudorapidity $\eta$ is defined in terms of
the polar angle $\theta$ by $\eta = -\ln \tan(\theta/2)$. 

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for the data. The MC events were reweighted to mimic the pileup conditions observed in the data.

SM diphoton production was simulated with PYTHIA [21] version 6.424 and MRST2007LOMOD [22] parton distribution functions (PDFs). The PYTHIA events were reweighted as a function of diphoton invariant mass from the differential cross section predicted by the NLO calculation of DIPHOX [23] version 1.3.2. The reweighting factor varied from \( \approx 1.6 \) for a diphoton mass of 140 GeV, decreasing smoothly to unity for large masses. For the DIPHOX calculation, the renormalization scale and the initial and final factorization scales of the model were all set to the diphoton mass. The various scales were varied by a factor of two both up and down, compared to this central value, to evaluate systematic uncertainties. The PDFs were chosen following the recommendations of the PDF4LHC working group [24], with MSTW2008 NLO PDFs [25] used for the NLO predictions, and CTEQ6.6 [26] and MRST2007LOMOD [22] used for systematic comparisons.

SHERPA [27] version 1.2.3 was used with MRST2007LOMOD [22] PDFs to simulate the various ADD scenarios for a variety of \( M_5 \) values. Due to the interference between the SM and gravity-mediated contributions, it is necessary to simulate events according to the full differential cross section as a function of the diphoton mass. A generator-level cut was applied to restrict the signal simulation to diphoton masses above 200 GeV. The ADD MC samples were used to determine the signal acceptance (\( A \)) and selection efficiency (\( \epsilon \)). The acceptance, defined as the percentage of diphoton signal events with the two highest \( E_\gamma \) photons passing the applied \( \eta_\gamma \) and \( \eta_\gamma \) cuts, varied somewhat for the various ADD implementations and fell from typical values of \( \approx 20\% \) for \( M_5 = 1.5 \) TeV down to \( \approx 15\% \) for \( M_5 = 3 \) TeV, due mostly to the variations in the \( \eta_\gamma \) distributions. The selection efficiency, for events within the detector acceptance, was found to be \( \approx 70\% \).

RS model MC signal samples were produced using the implementation of the RS model in PYTHIA [21] version 6.424, which is fully specified by providing the values of \( m_G \) and \( k/M_{Pl} \). MC signal samples were produced for a range of \( m_G \) and \( k/M_{Pl} \) values, using the MRST2007LOMOD [22] PDFs. The products of \( A \times \epsilon \) for the RS signal models were in the range \((53–80\%),\) slowly rising with increasing graviton mass. The reconstructed shape of the graviton resonance was modeled by convolving the graviton Breit–Wigner lineshape with a double-sided Crystal Ball (CB) function to describe the detector response. The natural width of the Breit–Wigner lineshape with a double-sided Crystal Ball (CB) function, which includes a Gaussian core to model the detector resolution, was found to be \( \approx 0.1 \) from \( 8 \) GeV up to \( \approx 30 \) GeV for \( m_G \) values from \( 800 \) GeV to \( 2200 \) GeV, respectively. The parameters of the CB function, which includes a Gaussian core to model the detector resolution matched to exponential functions on both sides to model the modest non-Gaussian tails, were determined by fitting to the reconstructed MC signals. The fitted values of \( \sigma \) of the Gaussian core approached a value of \( \approx 1 \% \) for high \( m_G \) values, as expected given the current value of the constant term in the EM calorimeter energy resolution, and were found to be independent of \( k/M_{Pl} \). The EM energy resolution has been verified in data using \( Z \rightarrow ee \) decays [28], and MC used to describe the modest differences between the response to photons versus electrons. The fitted values of the CB parameters varied smoothly with \( m_G \). Fitting this mass dependence provided a signal parametrization that was used to describe signals with any values of \( m_G \) and \( k/M_{Pl} \).

5. Background evaluation

The largest background for this analysis is the irreducible background due to SM \( \gamma \gamma \) production. The shape of the diphoton invariant mass spectrum from this background was estimated using MC, reweighting the PYTHIA samples to the differential cross section predictions of DIPHOX.

Another significant background component is the reducible background that includes events in which one or both of the reconstructed photon candidates result from a different physics object being misidentified as a photon. This background is dominated by \( \gamma + j \) events, and \( jj \) events, with one or two jets faking photons, respectively. Backgrounds with electrons faking photons, such as the Drell–Yan production of electron–positron pairs as well as \( W/Z + \gamma \) and \( t\bar{t} \) processes, were verified using MC to be small after the event selection and were neglected. Several background-enriched control samples were defined in order to determine the shape of the reducible background using data-driven techniques. In all control samples, the two photon candidates were required to pass the same isolation cut as for the signal selection, since removing the isolation requirement was seen to modify the diphoton mass spectrum. The first control sample contained those events where one of the photon candidates passed the tight requirement applied for the signal selection. However, the other photon candidate was required to fail the tight photon identification definition, but to pass the loose requirement; the latter restriction was applied to avoid any trigger bias, as the trigger required two loose photons. This sample is enriched in \( \gamma + j \) events, where the photon passed the tight requirement and a jet passed the loose one, and also in \( jj \) events where both photon candidates were due to jets. A second control sample, dominated by \( jj \) events, was similarly defined, but both photon candidates were required to fail the tight photon identification while passing the loose definition.

The diphoton invariant mass distributions were compared for these control samples. To check for any kinematic bias, the control sample with one tight and one loose photon candidate was further divided, with the \( \gamma j (j\gamma) \) subsample being defined as the case with the tight photon being the photon candidate with the highest (second highest) transverse energy. The diphoton invariant mass distributions of all three control subsamples were found to be consistent with each other, within statistical uncertainties. The sum of the control samples was used to provide the best estimate of the reducible background shape. Variations among the subsamples were taken into account as a source of systematic uncertainty in the reducible background prediction.

The data control samples have relatively few events in the high diphoton mass signal region. It was therefore necessary to extrapolate the reducible background shape to higher masses, which was done by fitting with a smooth function of the form \( f(x) = p_1 \times x^{p_2 + p_3} \log x \), where \( x = m_{\gamma\gamma} \) and \( p_i \) are the fit parameters. This functional form has been used in previous ATLAS resonance searches [12,29], and describes well the shape of the control data samples.

The total background, calculated as the sum of the irreducible and reducible components, was normalized to the number of data events in a low mass control region with diphoton masses between 140 and 400 GeV, in which possible ADD and RS signals have been excluded by previous searches. The fraction of the total background in this region that is due to the irreducible background is defined as the purity of the sample. The purity (\( p \)) was determined by three complementary methods. The most precise measurement resulted from a method previously used in Refs. [30,31] that examines the \( E^{\text{iso}}_T \) values of the two photon candidates. Templates for the \( E^{\text{iso}}_T \) distributions of true photons and of fake photons from jets were both determined from the data. The shape for fake photons was found using a sample of photon candidates that failed at least one of a subset of several of the selection requirements used for the tight photon definition. The shape for photons was found from the tight photon sample, after subtracting the fake photon shape normalized to match the number of candidates with large
values (greater than 10 GeV) of $E_T^{\gamma\gamma}$. In addition, for $jj$ events, due to the observed significant ($\approx 20\%$) correlation between the $E_T^{jj}$ values of the two photon candidates, a two-dimensional template was formed using events in which both photon candidates failed the tight identification. An extended maximum likelihood fit to the two-dimensional distribution formed from the $E_T^{jj}$ values of the two photon candidates was performed in order to extract the contributions from $\gamma\gamma$, $jj$, $j\gamma$, and $jj$ events. The fit was performed using the photon and fake photon $E_T^{\gamma\gamma}$ templates, as well as the two-dimensional $jj$ template. The resultant value of the purity in the low mass control region was $p = 71^{+5}_{-9}\%$. The uncertainty was determined by varying the subset of tight selection criteria failed by fake photon candidates, and then repeating the purity determination. Cross checks using either the DIPHOX prediction for the absolute normalization of the irreducible component, or fitting the shapes of the irreducible and reducible backgrounds to the data in the low mass control region, yielded consistent, but less precise results. The result from the isolation method was therefore used as the best estimate of the purity, and the total SM background prediction was set equal to the sum of the irreducible and reducible components, weighted appropriately by this purity value and normalized to data in the low mass control region.

6. Systematic uncertainties

Systematic uncertainties in the DIPHOX prediction for the shape of the irreducible background were obtained by varying the scales of the model and the PDFs, while keeping the overall normalization fixed in the low mass control region in which the total background prediction was normalized to the data. The resultant systematic uncertainties range from a few percent at low masses, up to $\approx 15\%$ for diphoton masses of $\approx 2$ TeV. Systematic uncertainties in the reducible background shape were obtained by comparing the results of the extrapolation fit for the various control data subsamples, in each case maintaining the overall normalization to the data in the low mass control region. The resultant uncertainties increase from $\approx 5\%$ for low masses to $\approx 100\%$ at a mass of $\approx 2$ TeV.

The systematic uncertainty on the shape of the total background was obtained by adding in quadrature the uncertainties on the shapes of the irreducible and reducible background components, weighted appropriately to account for the purity. In addition, there is a contribution, which is roughly constant with a value of $\approx 10\%$ for diphoton masses above 800 GeV, introduced by varying the purity value within its uncertainty. An additional overall uncertainty of $\approx 2\%$ was included due to the finite statistics of the data sample in the low mass control region.

The total background systematic uncertainty starts at $\approx 2\%$ for $m_{\gamma\gamma} = 140$ GeV, rises to $\approx 15\%$ by 700 GeV and then increases slowly up to almost 20% for the highest $m_{\gamma\gamma}$ values, above 2 TeV.

Systematic uncertainties on the signal yields were evaluated separately for the ADD and RS models. Since the differences were small, for simplicity the higher value was taken and applied to both models. The systematic uncertainties considered for the signal yield include the 3.7% uncertainty on the integrated luminosity [32], and a 1% uncertainty to account for the limited signal MC statistics. A value of 1% for the uncertainty on the bunch crossing identification (BCID) efficiency accounts for the ability of the Level 1 trigger hardware to pick the correct BCID when signal pulse saturation occurs in the trigger digitization. In addition, a value of 2% was applied for the uncertainty on the efficiency of the diphoton trigger. An uncertainty of 2.5% was applied due to the influence of pileup on the signal efficiency. Finally, a value of 4.3% was taken to account for the uncertainty in the selection and identification of the pair of photons, including uncertainties due to the photon isolation cut, the description of the detector material, the tight photon identification requirements, and extrapolation to the high photon $E_T$ values typical of the signal models. Uncertainties due to the current knowledge of the EM energy scale and resolution were verified to have a negligible impact. Adding all effects in quadrature, the total systematic uncertainty on the signal yields was 6.7%.

Uncertainties in the theoretical signal cross sections due to PDFs and due to the NLO approximation were considered. The uncertainties due to PDFs range from $\approx 10$–15% for ADD models and from $\approx 5$–10% for RS models. The authors of Refs. [33,34] have privately updated their calculations of the NLO signal cross sections for 14 TeV, and provided k-factors to the LHC experiments to scale from LO to NLO cross section values for the case of 7 TeV pp collisions. The NLO k-factor values, evaluated in our case for $|\eta| < 2.5$, have some modest dependence on the diphoton mass as well as on $M_S$ for the ADD model, and on the $k/\kappa_{\gamma\gamma}$ value for the RS model. However, the variations are within the theoretical uncertainty. For simplicity, therefore, constant values of 1.70 and 1.75 were assumed for the ADD and RS models, respectively, and an uncertainty in the k-factor value of $\pm 0.1$ was assigned to account for the variations.

7. Results and interpretation

Fig. 1 shows the observed invariant mass distribution of diphoton events, with the predicted SM background superimposed as well as ADD and RS signals for certain choices of the model parameters. The reducible background component is shown separately, in addition to the total background expectation, which sums the reducible and irreducible contributions. The shaded bands around each contribution indicate the corresponding uncertainty. The bottom plot of Fig. 1 shows the statistical significance, measured in standard deviations and based on Poisson distributions, of the difference between the data and the expected background in each bin. The significance was calculated and displayed as detailed in Ref. [35], and plotted as positive (negative) where there was an excess (deficit) in the data in a given bin. Table 1 lists, in bins of diphoton mass, the expected numbers of events for the irreducible and reducible background components, as well as for the total background, and also the numbers of observed data events. Both Fig. 1 and Table 1 demonstrate that there is agreement between the observed mass distribution and the expectation from the SM backgrounds over the entire diphoton mass range; no evidence is seen for either resonant or non-resonant deviations which would indicate the presence of a signal due to new physics. An analysis using the BUMPHunter [36] tool found that the probability, given the background-only hypothesis, of observing discrepancies at least as large as observed in the data was 0.28, indicating quantitatively the good agreement between the data and the expected SM background.

Given the absence of evidence for a signal, 95% CL upper limits were determined on the ADD and RS signal cross sections, using a Bayesian approach [37] with a flat prior on the signal cross section. The systematic uncertainties were incorporated as Gaussian-distributed nuisance parameters and integrated over. To set limits on the ADD model, the number of observed events with diphoton invariant mass in a high mass signal region was compared with the expected total SM background. To optimise the expected limit, the ADD signal search region was chosen as $m_{\gamma\gamma} > 1.1$ TeV. There are 2 observed events in this signal region, with a background expectation of 1.33 $\pm$ 0.26 events,
between the two background components, from the uncertainty on the purity. How-
background components do not include the contribution, which is anti-correlated
statistical and systematic uncertainties. The errors on the irreducible and reducible
malized to the corresponding number of observed events. The errors include both
400 GeV, corresponds to the control region in which the total background was nor-
events, in different diphoton mass bins. The first row, with masses from 140 to
Fig. 1. The observed invariant mass distribution of diphoton events, superimposed
with the predicted SM background and expected signals for ADD and RS models
with certain choices of parameters. The bin width is constant in log(m_{\gamma\gamma}). The bin-
by-bin significance of the difference between data and background is shown in the

Table 1
The expected numbers of events for the irreducible and reducible background com-
ponents and for the total background, as well as the numbers of observed data events,
in different diphoton mass bins. The first row, with masses from 140 to 400 GeV, corresponds to the control region in which the total background was normal-
ized to the corresponding number of observed events. The errors include both
statistical and systematic uncertainties. The errors on the irreducible and reducible
background components do not include the contribution, which is anti-correlated
between the two background components, from the uncertainty on the purity. How-
ever, this contribution is included in the errors listed for the total background.

<table>
<thead>
<tr>
<th>Mass range (GeV)</th>
<th>Background expectation</th>
<th>Observed events</th>
</tr>
</thead>
<tbody>
<tr>
<td>[140, 400)</td>
<td>4738 ± 180</td>
<td>6674</td>
</tr>
<tr>
<td>[400, 500)</td>
<td>90.0 ± 8.5</td>
<td>109.9 ± 9.2</td>
</tr>
<tr>
<td>[500, 600)</td>
<td>31.7 ± 4.0</td>
<td>37.0 ± 4.2</td>
</tr>
<tr>
<td>[600, 700)</td>
<td>13.7 ± 2.3</td>
<td>15.7 ± 2.4</td>
</tr>
<tr>
<td>[700, 800)</td>
<td>6.2 ± 1.2</td>
<td>6.9 ± 1.3</td>
</tr>
<tr>
<td>[800, 900)</td>
<td>3.1 ± 0.4</td>
<td>3.4 ± 0.5</td>
</tr>
<tr>
<td>[900, 1000)</td>
<td>1.6 ± 0.2</td>
<td>1.8 ± 0.3</td>
</tr>
<tr>
<td>[1000, 1100)</td>
<td>1.0 ± 0.2</td>
<td>1.0 ± 0.2</td>
</tr>
<tr>
<td>[1100, 1200)</td>
<td>0.50 ± 0.09</td>
<td>0.54 ± 0.11</td>
</tr>
<tr>
<td>[1200, 1300)</td>
<td>0.29 ± 0.07</td>
<td>0.31 ± 0.07</td>
</tr>
<tr>
<td>[1300, 1400)</td>
<td>0.14 ± 0.04</td>
<td>0.15 ± 0.04</td>
</tr>
<tr>
<td>[1400, 1500)</td>
<td>0.13 ± 0.04</td>
<td>0.14 ± 0.04</td>
</tr>
<tr>
<td>&gt; 1500</td>
<td>0.18 ± 0.09</td>
<td>0.19 ± 0.09</td>
</tr>
</tbody>
</table>

where the uncertainty includes both statistical and systematic er-
ors. The observed (expected) 95% CL upper limit is 2.49 (1.94) fb for
the product of the cross section due to new physics mul-
tipled by the acceptance and efficiency. The cross section re-
result can be translated into limits on \eta_G and, subsequently, on
the parameter M_S of the ADD model. As summarized in Ta-
ble 2, assuming a k-factor of 1.70, the 95% CL lower limits on M_S
range between 2.27 and 3.53 TeV, depending on the num-
ber of extra dimensions assumed and the ADD model imple-
mentation. LO results are also included in Table 2, for re-
ference.

To determine the limits on the RS model, the observed invar-
ant mass distribution was compared to templates of the expected
backgrounds and varying amounts of signal for various graviton
masses and k/M_{Pl} values. A likelihood function was defined as
the product of the Poisson probabilities over all mass bins in the

search region, defined as m_{\gamma\gamma} > 500 GeV, where the Poisson
probability in each bin was evaluated for the observed number of data
events given the expectation from the template. The total signal
acceptance as a function of mass was propagated into the expec-
tation. The theory uncertainties were not included in the limit
calculation, but are indicated by showing the theory prediction as
a band with a width equal to the combined theory uncertainty
when plotting the results. The resultant limits are summarized in
Table 3. Using a constant k-factor value of 1.75, the 95% CL lower
limits from the diphoton channel are m_G > 0.79 (1.85) TeV for
k/M_{Pl} = 0.01 (0.1).

The RS model results can be combined with the previously pub-
lished ATLAS results [12] from the dilepton final state, where, as-
suming LO cross sections and k/M_{Pl} = 0.1, RS gravitons with
masses below 1.51 (1.45) TeV were excluded at 95% CL using data
samples of 1.08 (1.21) fb^{-1} to search for G \rightarrow ee (G \rightarrow \mu\mu). To en-
sure their statistical independence, the selection cuts of the dipo-
ton analysis included a veto of any events which were also selected
by the 1.08 fb^{-1} G \rightarrow ee analysis. In performing the combination,
correlations were considered between the systematic uncertainties
in the \gamma\gamma and ee channels. In the ee analysis [12], the background
prediction was normalized such that the expected and observed
numbers of events in the region of the Z peak agreed, eliminating
the dependence of the ee result on the measured integrated lumi-
nosity. Therefore, the \gamma\gamma and ee signal predictions were treated
as uncorrelated, since there should be no correlation in the lumi-
nosity and efficiency uncertainties. The systematic uncertainty on
the QCD dijet background was treated as being correlated; how-
ever, this background was quite small so the effect was minor. The
PDF and scale uncertainties were treated as correlated across all
three channels, and affect the irreducible background in the \gamma\gamma
channel as well as the Drell–Yan background in the

Table 2
95% CL limits on the value of M_S (in TeV) for various implementations of the ADD
model, using both LO (k-factor = 1) and NLO (k-factor = 1.70) theory cross section
calculations.

<table>
<thead>
<tr>
<th>k-Factor value</th>
<th>Channel(s) used</th>
<th>95% CL limit [TeV]</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>GRW, Hewett, HLZ</td>
<td></td>
</tr>
<tr>
<td></td>
<td>Pos</td>
<td>Neg</td>
</tr>
<tr>
<td>1</td>
<td>2.73</td>
<td>2.44</td>
</tr>
<tr>
<td>1.70</td>
<td>2.97</td>
<td>2.66</td>
</tr>
</tbody>
</table>

Table 3
95% CL lower limits on the mass (GeV) of the lightest RS graviton, for various values
of k/M_{Pl}. The results are shown for the diphoton channel alone and for the combi-
nation of the diphoton channel with the dilepton results of Ref. [12], using both LO
(k-factor = 1) and NLO (k-factor = 1.75) theory cross section calculations.

<table>
<thead>
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<tr>
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<td></td>
<td>Pos</td>
<td>Neg</td>
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<tr>
<td>1</td>
<td>0.01</td>
<td>0.03</td>
</tr>
<tr>
<td>1.75</td>
<td>0.79</td>
<td>1.0</td>
</tr>
</tbody>
</table>

1400 > 1300 > 1200 > 1100 > 1000 > 900 > 800 > 700 > 600 > 500 > 400 > 300 > 200 > 100 > 50.
8. Summary

Using a dataset corresponding to 2.12 fb⁻¹, an analysis of the diphoton final state was used to set 95% CL lower limits of between 2.27 and 3.53 TeV on the parameter \(\frac{M_g}{M_\text{Pl}}\) of the ADD large extra dimension scenario, depending on the number of extra dimensions and the theoretical formalism used. The diphoton results also exclude at 95% CL RS graviton masses below 0.79 (1.85) TeV for the dimensionless RS coupling \(k/\sqrt{M_\text{Pl}} = 0.01 (0.1)\). Combining with the previous ATLAS dilepton analyses further tightens these limits to exclude at 95% CL RS graviton masses below 0.80 (1.95) TeV for \(k/\sqrt{M_\text{Pl}} = 0.01 (0.1)\).

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