Search for heavy vector-like quarks coupling to light quarks in proton-proton collisions at $\sqrt{s} = 7$ TeV with the ATLAS detector


DOI
10.1016/j.physletb.2012.03.082

Publication date
2012

Document Version
Final published version

Published in
Physics Letters B

Citation for published version (APA):

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Search for heavy vector-like quarks coupling to light quarks in proton–proton collisions at $\sqrt{s} = 7$ TeV with the ATLAS detector

ATLAS Collaboration

1. Introduction

Vector-like quarks (VLQ), defined as quarks for which both chiralities have the same transformation properties under the electroweak group SU(2) × U(1), are predicted by many extensions of the SM, relating to Grand Unification, dynamical electroweak symmetry breaking scenarios or theories with extra dimensions [1–10]. Since the couplings of the light quarks are well constrained, if VLQs exist they are generally expected to only couple sizably to the third generation. However, in certain scenarios, corrections to quark mixings can cancel, relaxing these constraints. The motivation and phenomenology of heavy VLQs coupling to light quarks is discussed for the Tevatron [11], where a baseline model is introduced which considers two degenerate VLQ doublets having hypercharges 1/6 and 7/6 and mixing only with the up quark. This scenario can occur naturally in certain models [12]. Because the doublets are degenerate, cancellations occur which allow VLQ coupling to the first two generations, leading to a potentially strong signal at the Large Hadron Collider (LHC).

Following the notation of more recent work [13] which describes a model-independent approach to VLQ sensitivity at the LHC, a coupling $\kappa_{qQ} = \langle v / m_Q \rangle \kappa_{qQ}$ is defined here, where $q$ stands for any light quark, $Q$ is the VLQ, $m_Q$ is the VLQ mass, $v$ is the Higgs vacuum expectation value and $\kappa_{qQ}$ encodes all the model dependence of the $qVQ$ vertex ($V = W$ or $Z$). Electroweak precision measurements constrain the contribution of heavy quarks to loop diagrams, but under certain conditions, as for the degenerate VLQ doublet model above, mild bounds apply on the dimensionless coupling $\kappa$, allowing it to be as large as $\sim 1$ [13]. The masses of VLQs are not constrained by vacuum stability in the SM [14].

It has been shown that single production provides a favorable process to probe for the existence of these heavy quarks if the coupling to light quarks is large, and that a significant mass reach could be achieved at the LHC with early data [11,13]. Single production of a VLQ occurs via the process $qq' \rightarrow q'Q$ (Fig. 1). A quark produced by this process of gauge boson exchange can have a charge of $5/3, 2/3, -1/3$ or $-4/3$. As a benchmark, we consider theories with only VLQs $U$ of charge $+2/3$ or only with $D$ of charge $-1/3$, without regard to the multiplet structure of the model. The experimental limits obtained on cross section times branching ratio can then be interpreted as limits on the couplings for different VLQ models [13]. The contribution from the $s$-channel diagram is negligible compared to that of the $t$-channel process. Therefore one characteristic of the signal is the presence of a forward jet: after one of the initial state quarks emits the electroweak gauge boson, it will continue in the forward direction with little transverse momentum ($p_T$), while the other quark couples to the $W$ or $Z$ to produce the heavy quark. Because the LHC is a proton–proton collider, the charged current (CC) production of a $D$ quark is expected to have a higher cross section than that of a $U$ quark. Similarly, for the neutral current (NC) process, $U$ quarks are expected to be produced more abundantly. Anti-quark production is suppressed since it involves anti-quarks in the initial state.

Bounds on the mass of new heavy quarks were obtained previously from a search in the pair production process at the Tevatron [15,16] and LHC [17,18]. Limits have also been obtained at the Tevatron [19,20] on single production processes $\sigma(p\bar{p} \rightarrow qQ) \times \text{BR}(Q \rightarrow qW)$, which in the model [11] of degenerate doublets
with $k_{UL} = 1$ and decaying 100% via CC gives a 95% confidence level (C.L.) upper limit exclusion for D quarks with masses up to 690 GeV. Limits at 95% C.L. on $\sigma(p\bar{p} \rightarrow qQ) \times \text{BR}(Q \rightarrow qZ)$ in the same model yield an exclusion of a U quark with $k_{UL} = \sqrt{Z}$ and 100% branching ratio via NC up to a mass of 550 GeV.

This Letter reports on a search for singly produced VLQs in the ATLAS detector at the LHC. The search is conducted in events with at least two jets and a vector boson, indicated by either two leptons in the final state, and standalone (with at least one lepton) with a track originating from the primary vertex. The trigger requires at least one $p_T$ of the leptons $> 20$ GeV. This includes $K$-factors of 1.22 and 1.25, respectively, to reproduce the inclusive cross sections at next-to-leading order in QCD [29]. MC@NLO [30] is used to simulate $t\bar{t}$ production, giving a cross section of 165 pb. Single top quark events decaying leptonically ($\sigma = 37.5$ pb) are generated with AcerMC [31] combined with parton showering and hadronization by PYTHIA. Diboson backgrounds are simulated with ALPGEN and HERWIG parton shower for the NC channel ($\sigma \times BR = 5.97$ pb), which requires two leptons in the final state, and standalone HERWIG [with a $K$-factor of 1.52 to reproduce the inclusive cross section at next-to-leading order in QCD [32]] for the CC channel ($\sigma \times BR = 69.1$ pb) where a single lepton is required. Multijet backgrounds from QCD processes are derived both from PYTHIA and data samples, described below.

The detector response simulation [33] is based on GEANT4 [34, 35]. The MC samples are generated with superimposed minimum bias events to simulate the conditions that occur in data. In order to improve the modeling of both signal and backgrounds, lepton reconstruction and identification efficiencies, energy scales and resolutions in the MC are corrected to correspond to the values measured in the data.

4. Analysis

The analysis is subdivided into four channels: charged and neutral current, each with either electrons or muons in the final state. Particle definitions and selections are identical in all channels, but signal and control regions for the CC and NC channels are defined independently.

Events are selected in which there is at least one vertex reconstructed with at least three tracks. The vertex with the greatest total transverse momentum, $\sum |p_T|$, of the associated tracks is designated as the primary vertex. The trigger requires at least one cluster in the EM calorimeter with $p_T > 20$ GeV or at least one muon candidate in the MS with a track originating from the primary vertex with $p_T > 18$ GeV. In both cases, the trigger requires a matching ID track.

Electron candidates are required to pass tight quality selection criteria based on the calorimeter shower shape, track quality and track matching with the calorimeter cluster [36]. They must have $p_T > 25$ GeV and lie in the pseudorapidity region $|\eta| < 2.47$, excluding the regions of transition between the central and forward detector sub-elements, 1.37 $< |\eta| < 1.52$. During most run periods

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1. ATLAS uses a right-handed coordinate system with the z-axis along the beam pipe. The x-axis points to the center of the LHC ring, and the y-axis points upward. The pseudorapidity is defined in terms of the polar angle $\theta$ as $\eta = -\ln \tan(\theta/2)$. 

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**Fig. 1.** Vector-like quark production and decay diagrams for s-channel (top) and t-channel (bottom). The thick line indicates the vector-like quark.
of the data set, a region of the EM calorimeter corresponding to about 1% of channels was less efficient than the rest of the detector. An exclusion window around the affected area was defined as $-0.1 < \eta < +1.5$ in pseudorapidity and $-0.9 < \phi < +0.5$ in azimuth. Electrons in this region are removed from data collected during these periods. The same procedure is applied to simulated events corresponding to the fraction of data covered by these run periods. Finally, no more than 4 GeV of transverse energy is allowed outside the core of the electron defined by a cone of size $\Delta R = \sqrt{\Delta \eta^2 + \Delta \phi^2} = 0.2$.

Muon candidates are reconstructed by combining tracks from both the ID and the MS. They are required to pass ID quality requirements [37] and have $p_T > 25$ GeV and $|\eta| < 2.4$. To suppress cosmic rays, muon candidates must have a distance of closest approach to the primary vertex in the longitudinal direction $|z_0| < 5$ mm and in the transverse plane $|d_0| < 0.1$ mm. Isolated muons are selected by requiring that the sum of ID track transverse momentum around the muon track, in a cone of $\Delta R = 0.2$ divided by the $p_T$ of the muon itself be less than 0.1.

Jet four-vectors are reconstructed from calorimeter clusters using the anti-$k_t$ algorithm [38] with a radius parameter of 0.4. After correcting for calorimeter non-compensation and inhomogeneities by using $p_T$- and $\eta$-dependent calibration factors [39], jets are required to have $p_T > 25$ GeV and $|\eta| < 4.5$. Events containing jets that fail quality criteria [40] are rejected to ensure an accurate $E_T^{\text{miss}}$ measurement. Furthermore, events containing jets passing through the inefficient region of the EM calorimeter are vetoed. To remove jets originating from other processes within an event, the selected jets are required to have more than 75% of $p_T$-weighted ID tracks associated to the primary vertex. Finally, to avoid counting electrons as jets, any jet candidate within $\Delta R < 0.2$ of a selected electron is removed.

The $E_T^{\text{miss}}$ is calculated as the negative vector of the transverse components of energy deposits in the calorimeters within $|\eta| < 4.5$. For events containing muons, any calorimeter energy deposit from a muon is ignored and the muon energy measured in the MS is used instead [41].

The CC candidates are required to have (i) exactly one electron or muon, (ii) missing transverse momentum $E_T^{\text{miss}} > 50$ GeV, (iii) one jet with $p_T > 50$ GeV and at least one more jet with $p_T > 25$ GeV, (iv) a minimum pseudorapidity separation $|\Delta \eta|$ > 1.0 between the highest-$p_T$ (leading) jet and second or third leading jet, since the presence of a forward jet is expected in signal events, (v) $m_T(E_T^{\text{miss}}) > 40$ GeV, where $m_T(E_T^{\text{miss}}) = \sqrt{2E_T^{\text{miss}}(1 - \cos \Delta \phi_{T, E_T^{\text{miss}}})}$ is the transverse mass of the $W$ candidate, and (vi) an azimuthal angle separation between the lepton and $E_T^{\text{miss}}$ vector $\Delta \phi_{T, E_T^{\text{miss}}} < 2.4$ rad since the $W$ in the signal is expected to be boosted. To reconstruct the mass of the VLQ candidate, the longitudinal momentum $p_z$ of the neutrino is calculated such that the invariant mass of the lepton and $E_T^{\text{miss}}$ equals the mass of the $W$. Of the two solutions, the one which leads to the larger value of $|\Delta \eta|$ between the reconstructed neutrino four-vector and the leading jet is chosen, since the simulation shows it to be the correct solution about 60% of the time. If no real solution is found, the real part of the complex solutions is taken. The system composed of the leading jet and the reconstructed $W$ is taken to be the VLQ candidate.

The NC candidates are required to have exactly two oppositely charged same-flavor leptons with an invariant mass in the range $66 < M(\ell, \ell) < 116$ GeV and a transverse momentum $p_T(\ell, \ell) > 50$ GeV. At least two jets of $p_T > 25$ GeV are required, with the same $|\Delta \eta|$ > 1.0 requirement as described for the CC selection. The invariant mass of the system composed of the two leptons and the leading jet is taken to be the VLQ candidate mass.

### Table 1

<table>
<thead>
<tr>
<th>Process</th>
<th>Electron channel</th>
<th>Muon channel</th>
</tr>
</thead>
<tbody>
<tr>
<td>$W + \text{jets}$</td>
<td>$14.500 \pm 100 \pm 4400$</td>
<td>$16.600 \pm 100 \pm 5000$</td>
</tr>
<tr>
<td>$t\bar{t}$</td>
<td>$236.0 \pm 50 \pm 270$</td>
<td>$2530.0 \pm 50 \pm 290$</td>
</tr>
<tr>
<td>Single top</td>
<td>$700.0 \pm 30 \pm 120$</td>
<td>$740.0 \pm 27 \pm 120$</td>
</tr>
<tr>
<td>Multijet</td>
<td>$670.0 \pm 30 \pm 270$</td>
<td>$340.0 \pm 20 \pm 410$</td>
</tr>
<tr>
<td>$Z + \text{jets}$</td>
<td>$128.1 \pm 11 \pm 90$</td>
<td>$432.1 \pm 21 \pm 170$</td>
</tr>
<tr>
<td>Di boson</td>
<td>$174.0 \pm 13 \pm 53$</td>
<td>$198.0 \pm 14 \pm 62$</td>
</tr>
<tr>
<td>Expected total background</td>
<td>$18.500 \pm 100 \pm 4400$</td>
<td>$20.900 \pm 100 \pm 5100$</td>
</tr>
<tr>
<td>Data</td>
<td>$17.302$</td>
<td>$20.668$</td>
</tr>
<tr>
<td>Expected signal, $D(225 \text{ GeV})$</td>
<td>$23.60 \pm 50 \pm 350$</td>
<td>$2380. \pm 50 \pm 400$</td>
</tr>
<tr>
<td>Expected signal, $D(600 \text{ GeV})$</td>
<td>$133.1 \pm 12 \pm 100$</td>
<td>$133.1 \pm 12 \pm 11$</td>
</tr>
<tr>
<td>Expected signal, $D(1000 \text{ GeV})$</td>
<td>$14 \pm 1 \pm 1$</td>
<td>$14 \pm 1 \pm 1$</td>
</tr>
</tbody>
</table>

To evaluate the level of multijet background in the CC analysis, a procedure is used based on a fit to the $E_T^{\text{miss}}$ distribution in the range $0 < E_T^{\text{miss}} < 100$ GeV. For this purpose, only selection criteria (i) and (v) above are required. For both the electron and muon modes, template shapes for the non-multijet backgrounds are taken from the MC samples described earlier and summed according to their relative cross sections. The overall normalization of this non-multijet template is left floating. In the electron mode, a sample enriched in objects misidentified as electrons (fakes) is selected from data using medium quality electrons, excluding tight electrons, as defined in [36]. The $E_T^{\text{miss}}$ distribution of this sample serves as the electron multijet template shape. For the muon mode, multijet background is primarily expected to come from heavy-flavor decays. Therefore, the multijet template shape is taken from a PYTHIA sample of $b\bar{b}$ events. For both lepton flavors, a fit to the $E_T^{\text{miss}}$ distribution is performed using the multijet and non-multijet templates to determine the normalization of the multijet component. The modeling of the multijet background was tested in a control region defined by the range of $10 < E_T^{\text{miss}} < 30$ GeV. The modeling of the kinematic variable distributions, and in particular of the VLQ candidate mass, was found to agree with data within statistics.

In the case of the NC selection, the multijet background is estimated from data-driven studies to be negligible. In the electron channel the selection for electron candidates is changed to require the medium criteria, excluding tight electrons, to obtain a fake di-electron template shape in $m_{e\mu}$, which is then scaled to make the total background expectation match the data with the same selection. In the muon channel the isolation requirement is inverted, and the $m_{\mu\mu}$ template scaled to the data in the same way.

With the above selections, the observed event yields and corresponding predictions are given in Tables 1 and 2. From these yields, no significant excess is observed in the data that can be accounted for by a VLQ signal. Fig. 2 shows the invariant mass distributions of the reconstructed VLQ candidate in the signal regions for both channels. These distributions are used in a binned likelihood fit to extract signal yields and production cross section upper limits. A slight shape discrepancy between data and MC is apparent in Fig. 2. Before extracting an upper limit, a correction is applied to the MC background shape, as described in Section 6.

### 5. Systematic uncertainties

Systematic uncertainties on the simulation of the signal arise from uncertainties in PDFs and the factorization and renormalization scales. In order to estimate the uncertainty due to the parton...
Table 2
Expected and observed event counts in the kinematically allowed VLQ mass range after the final selection in the NC channel with an integrated luminosity of 1.04 fb$^{-1}$. Uncertainties are statistical and systematic, respectively. The signal predictions assume a coupling $\kappa_{VLQ} = 1$.

<table>
<thead>
<tr>
<th>Process</th>
<th>Electron channel</th>
<th>Muon channel</th>
</tr>
</thead>
<tbody>
<tr>
<td>$Z + \text{jets}$</td>
<td>3250 ± 60 ± 430</td>
<td>5350 ± 70 ± 700</td>
</tr>
<tr>
<td>$t\bar{t}$</td>
<td>58 ± 8 ± 3</td>
<td>90 ± 9 ± 5</td>
</tr>
<tr>
<td>Diboson</td>
<td>38 ± 6 ± 4</td>
<td>58 ± 8 ± 4</td>
</tr>
<tr>
<td>Expected total background</td>
<td>3350 ± 60 ± 430</td>
<td>5500 ± 70 ± 700</td>
</tr>
<tr>
<td>Data</td>
<td>3105</td>
<td>5070</td>
</tr>
<tr>
<td>Expected signal, $U(225 \text{ GeV})$</td>
<td>192 ± 14 ± 9</td>
<td>339 ± 18 ± 19</td>
</tr>
<tr>
<td>Expected signal, $U(600 \text{ GeV})$</td>
<td>15 ± 3.9 ± 0.6</td>
<td>23 ± 4.8 ± 0.7</td>
</tr>
<tr>
<td>Expected signal, $U(1000 \text{ GeV})$</td>
<td>1.9 ± 1.4 ± 0.1</td>
<td>2.7 ± 1.6 ± 0.1</td>
</tr>
</tbody>
</table>

For signal and background events, the jet-energy-scale uncertainty is calculated by shifting the $p_T$ of all jets up and down by factors that vary as a function of $p_T$ and $\eta$. The factors range from 4.6% for jets with $p_T = 20$ GeV to 2.5% for jets with $p_T$ above 60 GeV [39]. This procedure results in an uncertainty of about 20% on the background normalization, and about 5% on the signal efficiency. The jet-energy-resolution uncertainty is calculated by smearing the $p_T$ of each jet depending on the jet $p_T$ and $\eta$, typically by around 10%. This source of uncertainty is found to impact both the background normalization and signal efficiency by about 1%. The lepton-energy-scale uncertainty is evaluated and found to be much less than 1% for both signal and background. The effect of the previously mentioned EM calorimeter inefficiency is also found to be much less than 1%. Uncertainties also arise from the trigger, identification, and reconstruction efficiency corrections applied to the MC simulation. They affect the signal efficiency uncertainty by 1–2% depending on the mass. The rate uncertainty from MC statistics after event selection is 3–5%. Finally, the uncertainty on the luminosity is 3.7% [21]. None of the systematics studied have been found to significantly affect the shape of the VLQ candidate mass distribution.

6. Results

To determine signal yields, a binned maximum likelihood fit is performed using template histograms of the VLQ candidate mass distribution. The fit is performed separately for each signal mass. The electron and muon final states are fitted simultaneously. The overall signal and background normalizations are left floating in the fit. Systematic uncertainties on the template normalizations are incorporated as Gaussian-distributed nuisance parameters, as are the signal efficiency systematics used in determining the cross section limits. Signal template shapes are taken from MC, while background templates are as shown in Fig. 2, with an additional correction described next.

A heavy VLQ signal would appear as a peak on top of a smooth background in the VLQ candidate invariant mass distribution. It is therefore important to have a good estimate of the background shape in the region around a signal mass hypothesis. The fit procedure described above makes use of the full range of mass, but the normalization is dominated by the lower mass region where the number of events is higher. A small shape difference between Monte Carlo and data can therefore yield a systematic bias in the fit at high mass. For that reason a correction is applied to the background model for each signal mass. It is obtained from linear fits to the reconstructed invariant mass of the ratio of data/MC after the full event selection, excluding bins in the range [−200, +100] GeV around each signal mass tested. The asymmetric choice in the excluded mass is motivated by the fact that the expected signal has a low mass tail. The uncertainty in the slope is taken as a systematic shape uncertainty. It was verified that no significant difference to the fit results arose from choosing a narrower excluded mass window, or even no exclusion at all.
expected limits, using a Gaussian prior for all nuisance parameters and including the shape uncertainty from the linear correction.

The 95% C.L. exclusion limits on $\sigma(pp \to Qq) \times BR(Q \to Vq)$ as a function of the VLQ mass, based on the CLs method [43], are shown in Fig. 3. Taking the intersection of the observed (expected) cross section limits with the central value of the theoretical cross section, masses below 900 GeV (840 GeV) are excluded for the CC channel and 760 GeV (820 GeV) for the NC channel, assuming a coupling $\kappa_{UD}^2 = 1$ and a 100% branching ratio for VLQs to decay to a vector boson and a jet. Within the $\pm 1\sigma$ theoretical uncertainties, the observed CC mass limit ranges from 870–920 GeV. The corresponding range for the NC limit is 730–770 GeV. Limits for each mass tested are given in Table 3. The fourth and fifth columns show an interpretation of the cross section limits in terms of limits on the couplings $k_{UD}$ and $k_{UH}$, in each case assuming only $D$ production or only $U$ production, respectively, and 100% branching fraction to a vector boson and jet.

A stronger limit in the CC channel may be obtained by repeating the CC analysis, requiring a negatively charged lepton because the SM background from $W^- + jets$ is lower than for $W^+ + jets$. The upper limits on $\sigma(pp \to D^{+}q) \times BR(D^{-} \to W^-u)$ are given in the sixth column of Table 3.

7. Conclusion

A search for single production of vector-like quarks coupling to light generations has been presented. No evidence is found for such quarks above the expected background in either the CC or NC channel. Upper limits on the production cross section times branching ratio to a vector boson and a jet were determined at 95% confidence level. Assuming couplings $k_{UD}^2$ and $k_{UH}^2$, in each case assuming only $D$ production or only $U$ production, respectively, and 100% branching fraction to a vector boson and jet.

Since no significant excess of data over the background prediction is observed in either channel, limits as function of the VLQ mass are obtained based on the likelihood fits. Pseudoexperiments are generated by sampling the likelihood function to compute the expected limits, using a Gaussian prior for all nuisance parameters and including the shape uncertainty from the linear correction.

Table 3
Observed upper limits at 95% confidence level on the cross section times branching ratio $\sigma(pp \to Qq) \times BR(Q \to Vq)$ as a function of mass and the corresponding upper limit on a model-independent heavy-to-light quark coupling. The final column shows the limit on the CC process after selecting negatively charged leptons.

<table>
<thead>
<tr>
<th>Mass (GeV)</th>
<th>CC $\sigma \times BR$ [pb]</th>
<th>NC $\sigma \times BR$ [pb]</th>
<th>$k_{UD}^2$</th>
<th>$k_{UH}^2$</th>
<th>CC− $\sigma \times BR$ [pb]</th>
</tr>
</thead>
<tbody>
<tr>
<td>225</td>
<td>15</td>
<td>18</td>
<td>0.075</td>
<td>0.21</td>
<td>12</td>
</tr>
<tr>
<td>300</td>
<td>17</td>
<td>11</td>
<td>0.24</td>
<td>0.31</td>
<td>5.6</td>
</tr>
<tr>
<td>400</td>
<td>5.3</td>
<td>2.4</td>
<td>0.21</td>
<td>0.19</td>
<td>3.8</td>
</tr>
<tr>
<td>500</td>
<td>2.1</td>
<td>1.4</td>
<td>0.19</td>
<td>0.26</td>
<td>1.1</td>
</tr>
<tr>
<td>600</td>
<td>1.9</td>
<td>1.5</td>
<td>0.37</td>
<td>0.56</td>
<td>1.9</td>
</tr>
<tr>
<td>700</td>
<td>2.2</td>
<td>1.0</td>
<td>0.86</td>
<td>0.75</td>
<td>2.2</td>
</tr>
<tr>
<td>800</td>
<td>0.93</td>
<td>1.0</td>
<td>0.66</td>
<td>1.33</td>
<td>0.97</td>
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<tr>
<td>900</td>
<td>0.80</td>
<td>0.9</td>
<td>1.0</td>
<td>2.1</td>
<td>0.70</td>
</tr>
<tr>
<td>1000</td>
<td>0.91</td>
<td>1.1</td>
<td>1.9</td>
<td>4.0</td>
<td>0.50</td>
</tr>
</tbody>
</table>

Fig. 3. Upper limits at 95% confidence level on the cross section times branching ratio $\sigma(pp \to Qq) \times BR(Q \to Vq)$ as a function of mass. The leading-order (LO) theoretical cross section assumes $k_{UD}^2 = 1$ and $k_{UH}^2 = 1$ on the top and bottom, respectively. The width of the dark band around it corresponds to the theoretical uncertainty described in the text. The expected cross section upper limit is determined by the median result of pseudoexperiments, and is shown with its $1\sigma$ and $2\sigma$ uncertainties, respectively.

Fig. 3. Upper limits at 95% confidence level on the cross section times branching ratio $\sigma(pp \to Qq) \times BR(Q \to Vq)$ as a function of mass. The leading-order (LO) theoretical cross section assumes $k_{UD}^2 = 1$ and $k_{UH}^2 = 1$ on the top and bottom, respectively. The width of the dark band around it corresponds to the theoretical uncertainty described in the text. The expected cross section upper limit is determined by the median result of pseudoexperiments, and is shown with its $1\sigma$ and $2\sigma$ uncertainties, respectively.

Acknowledgements

We thank A. Atre, M. Carena, T. Han, and J. Santiago for the MadGraph code used to produce the signal MC samples.

We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently.

We acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWF, Austria; ANAS, Azerbaijan; SSTC, Belarus; CNPq and FAPESP, Brazil; NSERC, NRC and CFI, Canada; CONICYT, Chile; CAS, MOST and NSFC, China; COLCIENCIAS, Colombia; MSMT CR, MPO CR and VSC CR, Czech Republic; DNRF, DNSRC and Lundbeck Foundation, Denmark; ARTEMIS, European Union; IN2P3-CNRS, CEA-DSM/IRFU, France; GNAS, Georgia; BMBF,
References

150 School of Physics, University of Sydney, Sydney, Australia
151 Institute of Physics, Academia Sinica, Taipei, Taiwan
152 Department of Physics, Technion, Israel Inst. of Technology, Haifa, Israel
153 Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
154 Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
155 International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan
156 Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
157 Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
158 Department of Physics, University of Toronto, Toronto, ON, Canada
159 (\textsuperscript{160}) TRIUMF, Vancouver, BC, \textsuperscript{161} Department of Physics and Astronomy, York University, Toronto, ON, Canada
159 (\textsuperscript{160}) Department of Pure and Applied Sciences, University of Tsukuba, 1-1-1 Tenmacho, Tsukuba, Ibaraki 305-8571, Japan
160 Science and Technology Center, Tufts University, Medford, MA, United States
161 Centre de Investigaciones, Universidad Antonio Narino, Bogota, Colombia
162 Department of Physics and Astronomy, University of California Irvine, Irvine, CA, United States
163 INFN Gruppo Collegato di Udine; ICTP, Trieste; Departamento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy
164 Department of Physics, University of Illinois, Urbana, IL, United States
165 Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden
166 Department of Physics, University of Wisconsin, Madison, WI, United States
167 Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel
168 Department of Physics, University of Wisconsin, Madison, WI, United States
169 Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany
170 Department of Physics, Yale University, New Haven, CT, United States
171 Yerevan Physics Institute, Yerevan, Armenia
172 Department of Physics, Yale University, New Haven, CT, United States
173 Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany
174 Department of Physics, University of Texas, Austin, TX, United States
175 Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
176 Domaine scientifique de la Doua, Centre de Calcul CNRS/IN2P3, Villeurbanne Cedex, France
177 Department of Physics, Bergische Universität Wuppertal, Wuppertal, Germany
178 Department of Physics, Technion, Israel Inst. of Technology, Haifa, Israel
179 School of Physics, University of Sydney, Sydney, Australia
181 Domaine scientifique de la Doua, Centre de Calcul CNRS/IN2P3, Villeurbanne Cedex, France

\textsuperscript{a} Also at Laboratorio de Instrumentacnio e Fisica Experimental de Particulas – LIP, Lisboa, Portugal.
\textsuperscript{b} Also at Faculdade de Ciencias and CPNUL, Universidade de Lisboa, Lisboa, Portugal.
\textsuperscript{c} Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom.
\textsuperscript{d} Also at TRIUMF, Vancouver, BC, Canada.
\textsuperscript{e} Also at School of Physics, Shandong University, Shandong, China.
\textsuperscript{f} Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France.
\textsuperscript{g} Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China.
\textsuperscript{h} Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan.
\textsuperscript{i} Also at DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l'Univers), CEA Saclay (Commissariat a l'Energie Atomique), Gif-sur-Yvette, France.
\textsuperscript{j} Also at Section de Physique, Université de Genève, Geneva, Switzerland.
\textsuperscript{k} Also at Department of Physics, University of California, Santa Barbara, CA, United States.
\textsuperscript{l} Also at Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France.
\textsuperscript{m} Deceased.