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Search for Production of Resonant States in the Photon-Jet Mass Distribution Using pp Collisions at $\sqrt{s} = 7$ TeV Collected by the ATLAS Detector

G. Aad *et al.**

(ATLAS Collaboration)

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This Letter describes a model-independent search for the production of new resonant states in photon + jet events in 2.11 fb^{-1} of proton-proton collisions at $\sqrt{s} = 7$ TeV. We compare the photon + jet mass distribution to a background model derived from data and find consistency with the background-only hypothesis. Given the lack of evidence for a signal, we set 95% credibility level limits on generic Gaussian-shaped signals and on a benchmark excited-quark (q^*) model, excluding 2 TeV Gaussian resonances with cross section times branching fraction times acceptance times efficiency near 5 fb and excluding q^* masses below 2.46 TeV, respectively.

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In the standard model (SM), proton-proton (pp) collisions do not produce photon + jet pairs through a resonance. Direct photon + jet production occurs at tree level via Compton scattering of a quark and a gluon or through quark-antiquark annihilation, the former accounting for the majority of direct photon + jet production at all center-of-mass energies. Events with a high transverse momentum photon and one or more jets can also arise from radiation off final-state quarks or from multijet processes where dijet or higher-order events produce secondary photons during fragmentation of the hard-scatter quarks and gluons [1–4]. The photon + jet invariant mass ($m_{\gamma j}$) [5] distribution resulting from this mixture of processes is smooth and rapidly falling, constituting a promising place to look for resonances.

Despite many possible exotic production mechanisms such as excited quarks [6,7], quirks [8–10], Regge excitations of string theory [11–14], and topological pions [15], the most recent searches for photon + jet resonances were published a decade ago [16–19]. The previous most sensitive search for new phenomena in the photon + jet final state places limits on effective signal cross section, cross section times branching fraction times acceptance times efficiency, of the order of 1 pb and on excited-quark masses up to 460 GeV at the 95% confidence level [16].

This Letter describes a general search for resonant s -channel photon + jet production in 2.11 fb^{-1} of pp collisions at a center-of-mass energy $\sqrt{s} = 7$ TeV with the ATLAS detector. It follows earlier measurements of isolated photon differential cross sections at the Large Hadron Collider (LHC) [20,21]. The entire $m_{\gamma j}$ distribution is fit

to a smooth function to obtain the background to this search. We look for evidence of a narrow resonance, not much wider than the detector mass resolution. This search extends the method used in the search for resonant dijet production [22] to handle the more diverse mixture of processes contributing to the $m_{\gamma j}$ distribution. In the absence of a signal, we use Bayes' theorem to set limits on Gaussian-shaped resonances and on a benchmark excited-quark (q^*) model [6,7].

In the excited-quark model studied here, the LHC could produce single q^* states with vectorlike couplings to the W and Z gauge bosons via the absorption of a gluon by a quark. As in Ref. [7], we define the model by one parameter, the excited-quark mass m_{q^*} , setting the compositeness scale equal to m_{q^*} and SU(3), SU(2), and U(1) coupling multipliers $f_s = f = f^l = 1$. At $m_{q^*} = 2.5$ TeV, this gives branching fractions for $u^* \rightarrow ug$ and $u^* \rightarrow u\gamma$ of 0.85 and 0.02, respectively. The corresponding branching fractions for d^* quarks are 0.85 and 0.005, respectively. We do not make any further assumptions about higher-order corrections or the excited-quark dynamics and neglect scale uncertainties and uncertainties on parton distribution functions in order to provide a convenient benchmark process for theoretical reinterpretation.

We simulate the SM direct photon processes and the q^* model with PYTHIA 6.4.25 [23] using the AMBT1 tune [24], MSTW2007 parton distribution functions [25], and a GEANT4-based detector simulation [26,27]. Supplementary studies of the background shape function are performed with next-to-leading-order JETPHOX 1.3.0 [1,2]. Additional inelastic pp interactions, termed pileup, are included in the event simulation, distributed so as to reproduce the number of collisions per bunch crossing in the data. The mean number of pileup interactions is approximately 6.

A detailed description of the detector is available in Ref. [28]. Photons are detected by a lead-liquid-argon sampling electromagnetic calorimeter (EMC) with an

*Full author list given at the end of the article.

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accordion geometry. In front of the EMC, the inner detector allows an accurate reconstruction of tracks from the primary pp collision point and also from secondary vertices, permitting an efficient reconstruction of photon conversions in the inner detector up to a radius of 80 cm. For $|\eta| < 1.37$ [29], an iron–scintillator-tile calorimeter behind the EMC provides hadronic coverage. The endcap and forward regions, $1.5 < |\eta| < 4.9$, are instrumented with liquid-argon calorimeters for both electromagnetic and hadronic measurements. For further details relevant to photon identification and measurement, see Ref. [20]. For details relevant to jet detection and measurement, see Ref. [30]. Events are collected with a trigger requiring at least one photon candidate with a transverse momentum (p_T) above 80 GeV. The integrated luminosity of the sample is $(2.11 \pm 0.08) \text{ fb}^{-1}$ [31,32].

Events containing at least one photon and at least one jet are selected for analysis. Each event must have a primary vertex with at least five charged-particle tracks with $p_T > 400$ MeV. Multiple vertices can appear when pileup interactions occur for the same bunch crossing. If more than one vertex is found, the primary vertex is taken as the vertex with the highest scalar sum p_T^2 of associated tracks. Photon candidates with $p_T > 85$ GeV and $|\eta| < 1.37$ and jet candidates with $p_T > 30$ GeV and $|\eta| < 2.8$ are used. These objects are identified using criteria that closely follow those applied in the isolated photon cross section measurement [20] and dijet resonance search [22]. Subleading photons or jets are allowed; when more than one photon or jet is found, the highest p_T candidates are selected to constitute the photon + jet pair.

Jets are reconstructed from clusters of calorimeter cells using the anti- k_t clustering algorithm [33] with radius parameter $R = 0.6$. Jet energies are corrected to the hadronic scale [30,34]. Jet candidates are rejected in regions of the calorimeter where the jet energy is not yet measured in an optimal way. Candidates consistent with spurious calorimeter noise or energy spikes are also rejected.

Photon candidates are reconstructed from clusters in the electromagnetic calorimeter and tracking information provided by the inner detector. They satisfy standard ATLAS selection criteria that are designed to reject instrumental backgrounds from hadrons [20]. The photon candidates must meet p_T - and η -dependent requirements on hadronic leakage, shower shapes in the electromagnetic strip layer, and shower shapes in the second sampling layer of the electromagnetic calorimeter. Inner detector tracking information is used to reject electrons and to recover photons converted to e^+e^- pairs. Energy calibrations are applied to photon candidates to account for energy loss in front of the electromagnetic calorimeter and for both lateral and longitudinal leakage. Events are discarded if the leading photon appears in calorimeter cells affected by noise bursts or transient hardware problems.

These photon identification criteria reduce instrumental backgrounds to a negligible level, but much of the substantial background from secondary (jet fragmentation) photons remains. We reduce this background with requirements on nearby calorimeter activity. Associated “isolation” calorimeter energy near the photon candidate is calculated by summing the transverse momentum as measured in electromagnetic and hadronic calorimeter cells inside a cone of radius $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.4$ centered on the photon cluster, but excluding the energy of the photon cluster itself. The isolation energy is corrected on an event-by-event basis for the ambient energy density due to pileup and the underlying event. This isolation energy is required to be less than 7 GeV.

The photon deposits energy in the electromagnetic calorimeter in such a way that it is also reconstructed as a jet. Jets within $\Delta R < 0.2$ of the photon are therefore not considered in this analysis. We require an angular separation $\Delta R(\gamma, \text{jet}) > 0.6$ between the signal photon and all jets with $p_T > 30$ GeV to reduce the background from photons during fragmentation of final-state quarks (fragmentation photons) and to reduce the systematic effects from the leakage of nearby jet showers into the photon isolation energy measurement.

An additional reduction of the fragmentation photon background is achieved by requirements on the photon and jet pseudorapidities. Dijet production rates increase with jet pseudorapidity whereas rates for our assumed s -channel signal would diminish. We restrict the analysis to photons in the barrel calorimeter, $|\eta| < 1.37$, and require $|\eta_\gamma - \eta_j| < 1.4$ between the photon and jet. The former criterion was chosen to avoid kinematic bias of the $m_{\gamma j}$ distribution due to the inclusion of any η range where reconstruction efficiency is lower, such as the barrel-endcap transition region $1.37 < |\eta| < 1.52$. The latter was chosen by optimizing the expected significance using the $|\eta_\gamma - \eta_j|$ distributions found in excited-quark signal simulation and background-dominated control data selected as in the nominal analysis but inverting the photon isolation requirement. This control sample is also used to check the background estimate.

After the above selections, Fig. 1 shows the distribution of the $m_{\gamma j}$ invariant mass in bins equal to the mass resolution. The $m_{\gamma j}$ resolution is about 4% at 600 GeV, improving to 3% at 2 TeV. We determine the combined SM and instrumental background to the search by fitting this distribution to the four-parameter ansatz,

$$f(x \equiv m_{\gamma j}/\sqrt{s}) = p_1(1-x)^{p_2}x^{-p_3-p_4} \ln x. \quad (1)$$

The motivation for this function is discussed in Refs. [16,35–37]. The fit result is also shown in Fig. 1. The bottom panel of the figure shows the statistical significance of the difference between data and the fit in each bin [38]. With a negative log-likelihood test statistic, the p

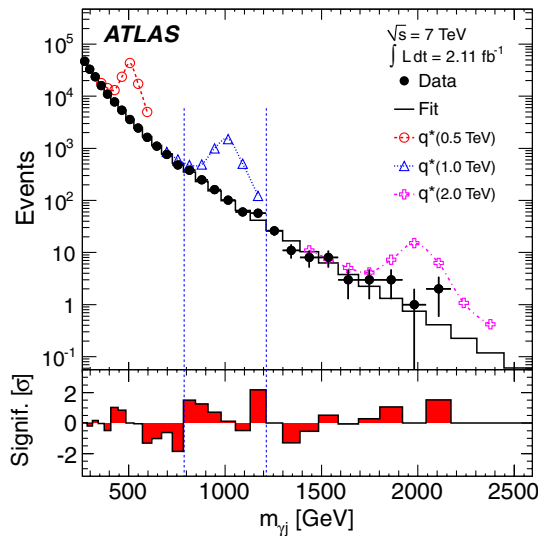


FIG. 1 (color online). Invariant mass of the photon + jet pair for events passing the final selection. Overlaid: The fitted background function integrated over each bin (stepped solid line), the most discrepant region identified by BUMPHUNTER (two dashed vertical lines), and three examples of excited-quark signals, normalized to luminosity, as described in the text. The bottom panel shows the statistical significance of the difference between data and background in each bin.

value is 23%, indicating that the data distribution is compatible with Eq. (1). The functional form also describes the leading-order PYTHIA direct photon prediction for comparable event statistics.

We search for statistical evidence of a resonance in this distribution using the BUMPHUNTER algorithm [39]. The algorithm operates on the binned $m_{\gamma j}$ distribution, comparing the background estimate with the data in mass intervals of varying contiguous bin multiplicities across the entire distribution. For each interval in the scan, it computes the significance of any excess found. The algorithm identifies the interval 784–1212 GeV, indicated by the vertical lines in Fig. 1, as the single most discrepant interval. The significance of the outcome is evaluated using the ensemble of possible outcomes for the significance of any region in the distribution in the background-only hypothesis, obtained by repeating the analysis on pseudodata drawn from the background function. Before including systematic uncertainties, the probability (p value) of observing a background fluctuation at least as significant as the above, including the trials factor, or “look-elsewhere” effect, is 20%. Inclusion of systematic uncertainties renders the p value similarly large.

Lacking evidence of any signal, we exclude two types of photon + jet signals: a generic signal with a Gaussian distribution and arbitrary production cross section, and the excited-quark model. We compute Bayesian limits at 95% credibility level (CL) using a prior probability density that is constant for positive values of the signal production

cross section and zero for unphysical, negative values, as described in Ref. [40]. We consider systematic uncertainties on the expected signal yield due to imperfect knowledge of the detector: the integrated luminosity (3.7%), trigger efficiencies ($< 0.5\%$), and signal photon identification efficiencies (2.0%). The last of these consists of isolation (0.4%), pile-up interactions (0.5%), conversions (1.2%), simulation mismodeling (1.3%), and the extrapolation of the photon identification efficiency to high p_T ($< 0.3\%$). Uncertainties on the photon energy scale (0.5%–1.5%), jet energy scale (2%–4%), and jet energy resolution (5%–15%) contribute through their effects on the signal distribution. These systematic uncertainties are treated as marginalized Gaussian nuisance parameters in the limit calculation.

We also evaluate two systematic uncertainties on the background estimate. To account for the statistical uncertainties on the background fit parameters, we repeatedly fit the background function to pseudodata for each bin drawn from Poisson distributions. The mean of the Poisson distribution for a given bin corresponds to the number of entries actually observed in that bin in the data. We then take the variation in the fit predictions for a given bin, 0.5% of the background at low mass to almost 10% of the background at 2 TeV, as indicative of the systematic uncertainty. We treat this bin-by-bin uncertainty in the limit as fully correlated, using a single nuisance parameter that scales the entire background distribution.

While our function can describe the $m_{\gamma j}$ shape for direct photon production, as modeled in the PYTHIA direct photon + jet simulation, the function need not remain a good description of the full distribution after including nonisolated and fragmentation photon events. For example, the function describes the next-to-leading-order prediction implemented in JETPHOX, which includes the fragmentation photon contributions, for some viable choices of theory parameters but not for others.

The second background systematic uncertainty accounts for any unmodeled features of fragmentation photon events in our isolated photon sample. We fit the background function to the $m_{\gamma j}$ distribution in the control data selected with the inverted isolation requirement, then measure for each $m_{\gamma j}$ bin the magnitude of any deviation from the fit, and assign the ratio of the deviation to the fit expectation as a parametrization bias systematic uncertainty. To extrapolate this uncertainty to large $m_{\gamma j}$ where few control data exist, we fit the tail $m_{\gamma j} > 1$ TeV with a two-degree polynomial.

Figure 2 shows the model-independent limits on the effective cross section, cross section σ times branching fraction B times acceptance A times efficiency ϵ , of a potential signal as a function of the central mass of each signal template. We take the signal line shape to be a Gaussian distribution with one of three widths, $\sigma_G/m_G = 5\%$, 7%, and 10% of the central mass of the

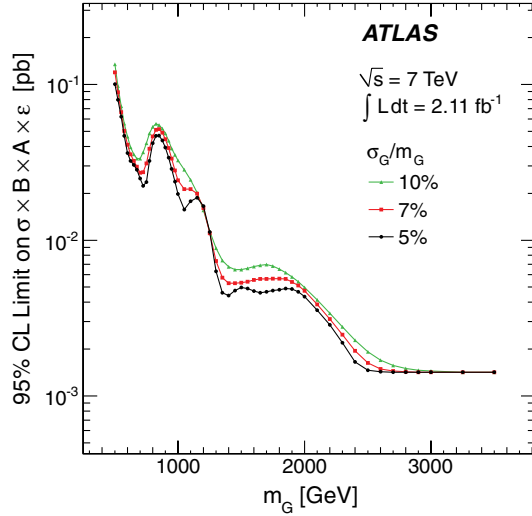


FIG. 2 (color online). The 95% CL upper limits on $\sigma \times B \times A \times \epsilon$ for a hypothetical signal with the Gaussian $m_{\gamma j}$ distribution decaying to a photon and a jet as a function of the signal mass m_G for three values of the relative Gaussian width σ_G/m_G and taking into account systematic uncertainties.

Gaussian. The limit weakens as the width increases and the peak becomes less distinct. For example, for a 1 TeV signal the limit for a width of 10% is 1.6 times the limit for a width of 5%.

The limit on the effective cross section in the excited-quark model is shown in Fig. 3 as a function of the q^* mass. Also shown are $\pm 1\sigma$ and $\pm 2\sigma$ uncertainty bands indicating the underlying distribution of possible limit outcomes in the background-only hypothesis. The solid

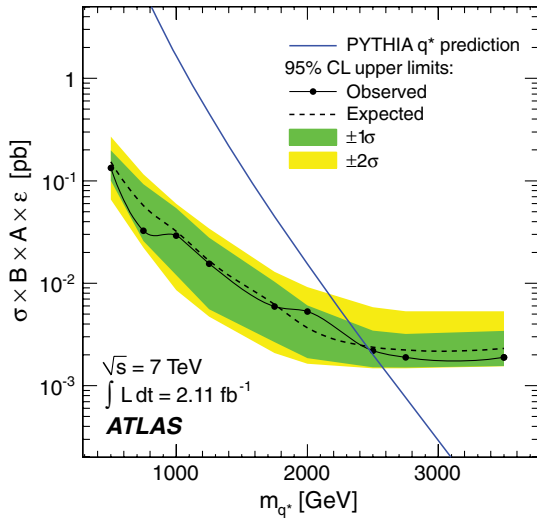


FIG. 3 (color online). The 95% CL upper limits on $\sigma \times B \times A \times \epsilon$ for excited quarks decaying to a photon and a jet as a function of the signal mass m_{q^*} for the PYTHIA q^* prediction and taking into account systematic uncertainties.

line indicates the prediction from the PYTHIA excited-quark implementation. We exclude such excited quarks with masses below 2.46 TeV at 95% CL, complementing the more stringent exclusion below 2.99 TeV on this specific q^* model in the dijet final state [22].

In conclusion, the photon + jet mass distribution measured in 2.11 fb^{-1} of pp collision data collected at $\sqrt{s} = 7 \text{ TeV}$ by the ATLAS Collaboration has been examined for narrow resonances. The observed distribution extends up to masses of about 2 TeV. It is well described by a smooth function fitted to it and assumed to represent the SM expectation. No evidence for the production of resonances is found. We set limits at 95% CL on the Gaussian line shape and excited-quark signal using Bayesian statistics. The limits on Gaussian resonances, for example, exclude 2 TeV resonances with effective cross sections near 5 fb. We also exclude the excited-quark model in the photon + jet final state for masses up to 2.46 TeV. The limits reported here on the resonant production of new particles in the photon + jet final state are the most stringent limits set to date in this channel.

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G. Aad,⁴⁷ B. Abbott,¹¹⁰ J. Abdallah,¹¹ A. A. Abdelalim,⁴⁸ A. Abdesselam,¹¹⁷ O. Abdinov,¹⁰ B. Abi,¹¹¹ M. Abolins,⁸⁷ H. Abramowicz,¹⁵² H. Abreu,¹¹⁴ E. Acerbi,^{88a,88b} B. S. Acharya,^{163a,163b} L. Adamczyk,³⁷ D. L. Adams,²⁴ T. N. Addy,⁵⁵ J. Adelman,¹⁷⁴ M. Aderholz,⁹⁸ S. Adomeit,⁹⁷ P. Adragna,⁷⁴ T. Adye,¹²⁸ S. Aefsky,²² J. A. Aguilar-Saavedra,^{123b,b} M. Aharrouche,⁸⁰ S. P. Ahlen,²¹ F. Ahles,⁴⁷ A. Ahmad,¹⁴⁷ M. Ahsan,⁴⁰ G. Aielli,^{132a,132b} T. Akdogan,^{18a} T. P. A. Åkesson,⁷⁸ G. Akimoto,¹⁵⁴ A. V. Akimov,⁹³ A. Akiyama,⁶⁶ M. S. Alam,¹ M. A. Alam,⁷⁵ J. Albert,¹⁶⁸ S. Albrand,⁵⁴ M. Aleksa,²⁹ I. N. Aleksandrov,⁶⁴ F. Alessandria,^{88a} C. Alexa,^{25a} G. Alexander,¹⁵² G. Alexandre,⁴⁸ T. Alexopoulos,⁹ M. Alhroob,²⁰ M. Aliev,¹⁵ G. Alimonti,^{88a} J. Alison,¹¹⁹ M. Aliyev,¹⁰ P. P. Allport,⁷² S. E. Allwood-Spiers,⁵² J. Almond,⁸¹ A. Aloisio,^{101a,101b} R. Alon,¹⁷⁰ A. Alonso,⁷⁸ B. Alvarez Gonzalez,⁸⁷ M. G. Alviggi,^{101a,101b} K. Amako,⁶⁵ P. Amaral,²⁹ C. Amelung,²² V. V. Ammosov,¹²⁷ A. Amorim,^{123a,c} G. Amorós,¹⁶⁶ N. Amram,¹⁵² C. Anastopoulos,²⁹ L. S. Ancu,¹⁶ N. Andari,¹¹⁴ T. Andeen,³⁴ C. F. Anders,²⁰ G. Anders,^{57a} K. J. Anderson,³⁰ A. Andreazza,^{88a,88b} V. Andrei,^{57a} M.-L. Andrieux,⁵⁴ X. S. Anduaga,⁶⁹ A. Angerami,³⁴ F. Anghinolfi,²⁹ A. Anisenkov,¹⁰⁶ N. Anjos,^{123a} A. Annovi,⁴⁶ A. Antonaki,⁸ M. Antonelli,⁴⁶ A. Antonov,⁹⁵ J. Antos,^{143b} F. Anulli,^{131a} S. Aoun,⁸² L. Aperio Bella,⁴ R. Apolle,^{117,d} G. Arabidze,⁸⁷ I. Aracena,¹⁴² Y. Arai,⁶⁵ A. T. H. Arce,⁴⁴ J. P. Archambault,²⁸ S. Arfaoui,⁸² J.-F. Arguin,¹⁴ E. Arik,^{18a,a} M. Arik,^{18a} A. J. Armbruster,⁸⁶ O. Arnaez,⁸⁰ C. Arnault,¹¹⁴ A. Artamonov,⁹⁴ G. Artoni,^{131a,131b} D. Arutinov,²⁰ S. Asai,¹⁵⁴ R. Asfandiyarov,¹⁷¹ S. Ask,²⁷ B. Åsman,^{145a,145b} L. Asquith,⁵ K. Assamagan,²⁴ A. Astbury,¹⁶⁸ A. Astvatsatourov,⁵¹ B. Aubert,⁴ E. Auge,¹¹⁴ K. Augsten,¹²⁶ M. Aourousseau,^{144a} G. Avolio,¹⁶² R. Avramidou,⁹ D. Axen,¹⁶⁷ C. Ay,⁵³

G. Azuelos,^{92,e} Y. Azuma,¹⁵⁴ M. A. Baak,²⁹ G. Baccaglioni,^{88a} C. Bacci,^{133a,133b} A. M. Bach,¹⁴ H. Bachacou,¹³⁵ K. Bachas,²⁹ G. Bachy,²⁹ M. Backes,⁴⁸ M. Backhaus,²⁰ E. Badescu,^{25a} P. Bagnaia,^{131a,131b} S. Bahinipati,² Y. Bai,^{32a} D. C. Bailey,¹⁵⁷ T. Bain,¹⁵⁷ J. T. Baines,¹²⁸ O. K. Baker,¹⁷⁴ M. D. Baker,²⁴ S. Baker,⁷⁶ E. Banas,³⁸ P. Banerjee,⁹² Sw. Banerjee,¹⁷¹ D. Banfi,²⁹ A. Bangert,¹⁴⁹ V. Bansal,¹⁶⁸ H. S. Bansil,¹⁷ L. Barak,¹⁷⁰ S. P. Baranov,⁹³ A. Barashkou,⁶⁴ A. Barbaro Galtieri,¹⁴ T. Barber,⁴⁷ E. L. Barberio,⁸⁵ D. Barberis,^{49a,49b} M. Barbero,²⁰ D. Y. Bardin,⁶⁴ T. Barillari,⁹⁸ M. Barisonzi,¹⁷³ T. Barklow,¹⁴² N. Barlow,²⁷ B. M. Barnett,¹²⁸ R. M. Barnett,¹⁴ A. Baroncelli,^{133a} G. Barone,⁴⁸ A. J. Barr,¹¹⁷ F. Barreiro,⁷⁹ J. Barreiro Guimarães da Costa,⁵⁶ P. Barrillon,¹¹⁴ R. Bartoldus,¹⁴² A. E. Barton,⁷⁰ V. Bartsch,¹⁴⁸ R. L. Bates,⁵² L. Batkova,^{143a} J. R. Batley,²⁷ A. Battaglia,¹⁶ M. Battistin,²⁹ F. Bauer,¹³⁵ H. S. Bawa,^{142,f} S. Beale,⁹⁷ B. Beare,¹⁵⁷ T. Beau,⁷⁷ P. H. Beauchemin,¹⁶⁰ R. Beccherle,^{49a} P. Bechtel,²⁰ H. P. Beck,¹⁶ S. Becker,⁹⁷ M. Beckingham,¹³⁷ K. H. Becks,¹⁷³ A. J. Beddall,^{18c} A. Beddall,^{18c} S. Bedikian,¹⁷⁴ V. A. Bednyakov,⁶⁴ C. P. Bee,⁸² M. Begel,²⁴ S. Behar Harpaz,¹⁵¹ P. K. Behera,⁶² M. Beimforde,⁹⁸ C. Belanger-Champagne,⁸⁴ P. J. Bell,⁴⁸ W. H. Bell,⁴⁸ G. Bella,¹⁵² L. Bellagamba,^{19a} F. Bellina,²⁹ M. Bellomo,²⁹ A. Belloni,⁵⁶ O. Beloborodova,^{106,g} K. Belotskiy,⁹⁵ O. Beltramello,²⁹ S. Ben Ami,¹⁵¹ O. Benary,¹⁵² D. Bencheekroun,^{134a} C. Benchouk,⁸² M. Bendel,⁸⁰ N. Benekos,¹⁶⁴ Y. Benhammou,¹⁵² J. A. Benitez Garcia,^{158b} D. P. Benjamin,⁴⁴ M. Benoit,¹¹⁴ J. R. Bensinger,²² K. Benslama,¹²⁹ S. Bentvelsen,¹⁰⁴ D. Berge,²⁹ E. Bergeaas Kuutmann,⁴¹ N. Berger,⁴ F. Berghaus,¹⁶⁸ E. Berglund,¹⁰⁴ J. Beringer,¹⁴ P. Bernat,⁷⁶ R. Bernhard,⁴⁷ C. Bernius,²⁴ T. Berry,⁷⁵ C. Bertella,⁸² A. Bertin,^{19a,19b} F. Bertinelli,²⁹ F. Bertolucci,^{121a,121b} M. I. Besana,^{88a,88b} N. Besson,¹³⁵ S. Bethke,⁹⁸ W. Bhimji,⁴⁵ R. M. Bianchi,²⁹ M. Bianco,^{71a,71b} O. Biebel,⁹⁷ S. P. Bieniek,⁷⁶ K. Bierwagen,⁵³ J. Biesiada,¹⁴ M. Biglietti,^{133a} H. Bilokon,⁴⁶ M. Bindi,^{19a,19b} S. Binet,¹¹⁴ A. Bingul,^{18c} C. Bini,^{131a,131b} C. Biscarat,¹⁷⁶ U. Bitenc,⁴⁷ K. M. Black,²¹ R. E. Blair,⁵ J.-B. Blanchard,¹¹⁴ G. Blanchot,²⁹ T. Blazek,^{143a} C. Blocker,²² J. Blocki,³⁸ A. Blondel,⁴⁸ W. Blum,⁸⁰ U. Blumenschein,⁵³ G. J. Bobbink,¹⁰⁴ V. B. Bobrovnikov,¹⁰⁶ S. S. Bocchetta,⁷⁸ A. Bocci,⁴⁴ C. R. Boddy,¹¹⁷ M. Boehler,⁴¹ J. Boek,¹⁷³ N. Boelaert,³⁵ S. Böser,⁷⁶ J. A. Bogaerts,²⁹ A. Bogdanchikov,¹⁰⁶ A. Bogouch,^{89,a} C. Bohm,^{145a} V. Boisvert,⁷⁵ T. Bold,³⁷ V. Boldea,^{25a} N. M. Bolnet,¹³⁵ M. Bona,⁷⁴ V. G. Bondarenko,⁹⁵ M. Bondioli,¹⁶² M. Boonekamp,¹³⁵ G. Boorman,⁷⁵ C. N. Booth,¹³⁸ S. Bordononi,⁷⁷ C. Borer,¹⁶ A. Borisov,¹²⁷ G. Borissov,⁷⁰ I. Borjanovic,^{12a} S. Borroni,⁸⁶ K. Bos,¹⁰⁴ D. Boscherini,^{19a} M. Bosman,¹¹ H. Boterenbrood,¹⁰⁴ D. Botterill,¹²⁸ J. Bouchami,⁹² J. Boudreau,¹²² E. V. Bouhova-Thacker,⁷⁰ D. Boumediene,³³ C. Bourdarios,¹¹⁴ N. Bousson,⁸² A. Boveia,³⁰ J. Boyd,²⁹ I. R. Boyko,⁶⁴ N. I. Bozhko,¹²⁷ I. Bozovic-Jelisavcic,^{12b} J. Bracinik,¹⁷ A. Braem,²⁹ P. Branchini,^{133a} G. W. Brandenburg,⁵⁶ A. Brandt,⁷ G. Brandt,¹¹⁷ O. Brandt,⁵³ U. Bratzler,¹⁵⁵ B. Brau,⁸³ J. E. Brau,¹¹³ H. M. Braun,¹⁷³ B. Brelier,¹⁵⁷ J. Bremer,²⁹ R. Brenner,¹⁶⁵ S. Bressler,¹⁷⁰ D. Breton,¹¹⁴ D. Britton,⁵² F. M. Brochu,²⁷ I. Brock,²⁰ R. Brock,⁸⁷ T. J. Brodbeck,⁷⁰ E. Brodet,¹⁵² F. Broggi,^{88a} C. Bromberg,⁸⁷ J. Bronner,⁹⁸ G. Brooijmans,³⁴ W. K. Brooks,^{31b} G. Brown,⁸¹ H. Brown,⁷ P. A. Bruckman de Renstrom,³⁸ D. Bruncko,^{143b} R. Bruneliere,⁴⁷ S. Brunet,⁶⁰ A. Bruni,^{19a} G. Bruni,^{19a} M. Bruschi,^{19a} T. Buanes,¹³ Q. Buat,⁵⁴ F. Bucci,⁴⁸ J. Buchanan,¹¹⁷ N. J. Buchanan,² P. Buchholz,¹⁴⁰ R. M. Buckingham,¹¹⁷ A. G. Buckley,⁴⁵ S. I. Buda,^{25a} I. A. Budagov,⁶⁴ B. Budick,¹⁰⁷ V. Büscher,⁸⁰ L. Bugge,¹¹⁶ O. Bulekov,⁹⁵ M. Bunse,⁴² T. Buran,¹¹⁶ H. Burckhart,²⁹ S. Burdin,⁷² T. Burgess,¹³ S. Burke,¹²⁸ E. Busato,³³ P. Bussey,⁵² C. P. Buszello,¹⁶⁵ F. Butin,²⁹ B. Butler,¹⁴² J. M. Butler,²¹ C. M. Buttar,⁵² J. M. Butterworth,⁷⁶ W. Buttinger,²⁷ S. Cabrera Urbán,¹⁶⁶ D. Caforio,^{19a,19b} O. Cakir,^{3a} P. Calafiura,¹⁴ G. Calderini,⁷⁷ P. Calfayan,⁹⁷ R. Calkins,¹⁰⁵ L. P. Caloba,^{23a} R. Caloi,^{131a,131b} D. Calvet,³³ S. Calvet,³³ R. Camacho Toro,³³ P. Camarri,^{132a,132b} M. Cambiaghi,^{118a,118b} D. Cameron,¹¹⁶ L. M. Caminada,¹⁴ S. Campana,²⁹ M. Campanelli,⁷⁶ V. Canale,^{101a,101b} F. Canelli,^{30,h} A. Canepa,^{158a} J. Cantero,⁷⁹ L. Capasso,^{101a,101b} M. D. M. Capeans Garrido,²⁹ I. Caprini,^{25a} M. Caprini,^{25a} D. Capriotti,⁹⁸ M. Capua,^{36a,36b} R. Caputo,⁸⁰ C. Caramarcu,²⁴ R. Cardarelli,^{132a} T. Carli,²⁹ G. Carlino,^{101a} L. Carminati,^{88a,88b} B. Caron,⁸⁴ S. Caron,⁴⁷ G. D. Carrillo Montoya,¹⁷¹ A. A. Carter,⁷⁴ J. R. Carter,²⁷ J. Carvalho,^{123a,i} D. Casadei,¹⁰⁷ M. P. Casado,¹¹ M. Cascella,^{121a,121b} C. Caso,^{49a,49b,a} A. M. Castaneda Hernandez,¹⁷¹ E. Castaneda-Miranda,¹⁷¹ V. Castillo Gimenez,¹⁶⁶ N. F. Castro,^{123a} G. Cataldi,^{71a} F. Cataneo,²⁹ A. Catinaccio,²⁹ J. R. Catmore,²⁹ A. Cattai,²⁹ G. Cattani,^{132a,132b} S. Caughron,⁸⁷ D. Cauz,^{163a,163c} P. Cavalleri,⁷⁷ D. Cavalli,^{88a} M. Cavalli-Sforza,¹¹ V. Cavasinni,^{121a,121b} F. Ceradini,^{133a,133b} A. S. Cerqueira,^{23b} A. Cerri,²⁹ L. Cerrito,⁷⁴ F. Cerutti,⁴⁶ S. A. Cetin,^{18b} F. Cevenini,^{101a,101b} A. Chafaq,^{134a} D. Chakraborty,¹⁰⁵ K. Chan,² B. Chapleau,⁸⁴ J. D. Chapman,²⁷ J. W. Chapman,⁸⁶ E. Chareyre,⁷⁷ D. G. Charlton,¹⁷ V. Chavda,⁸¹ C. A. Chavez Barajas,²⁹ S. Cheatham,⁸⁴ S. Chekanov,⁵ S. V. Chekulaev,^{158a} G. A. Chelkov,⁶⁴ M. A. Chelstowska,¹⁰³ C. Chen,⁶³ H. Chen,²⁴ S. Chen,^{32c} T. Chen,^{32c} X. Chen,¹⁷¹ S. Cheng,^{32a} A. Cheplakov,⁶⁴ V. F. Chepurinov,⁶⁴ R. Cherkaoui El Moursli,^{134e} V. Chernyatin,²⁴ E. Cheu,⁶ S. L. Cheung,¹⁵⁷ L. Chevalier,¹³⁵ G. Chiefari,^{101a,101b} L. Chikovani,^{50a} J. T. Childers,^{57a} A. Chilingarov,⁷⁰ G. Chiodini,^{71a} M. V. Chizhov,⁶⁴ G. Choudalakis,³⁰

S. Chouridou,¹³⁶ I. A. Christidi,⁷⁶ A. Christov,⁴⁷ D. Chromek-Burckhart,²⁹ M. L. Chu,¹⁵⁰ J. Chudoba,¹²⁴
G. Ciapetti,^{131a,131b} K. Ciba,³⁷ A. K. Ciftci,^{3a} R. Ciftci,^{3a} D. Cinca,³³ V. Cindro,⁷³ M. D. Ciobotaru,¹⁶² C. Ciocca,^{19a}
A. Ciocio,¹⁴ M. Cirilli,⁸⁶ M. Citterio,^{88a} M. Ciubancan,^{25a} A. Clark,⁴⁸ P. J. Clark,⁴⁵ W. Cleland,¹²² J. C. Clemens,⁸²
B. Clement,⁵⁴ C. Clement,^{145a,145b} R. W. Clift,¹²⁸ Y. Coadou,⁸² M. Cobal,^{163a,163c} A. Coccaro,^{49a,49b} J. Cochran,⁶³
P. Coe,¹¹⁷ J. G. Cogan,¹⁴² J. Coggeshall,¹⁶⁴ E. Cogneras,¹⁷⁶ J. Colas,⁴ A. P. Colijn,¹⁰⁴ N. J. Collins,¹⁷
C. Collins-Tooth,⁵² J. Collot,⁵⁴ G. Colon,⁸³ P. Conde Muño,^{123a} E. Coniavitis,¹¹⁷ M. C. Conidi,¹¹ M. Consonni,¹⁰³
V. Consorti,⁴⁷ S. Constantinescu,^{25a} C. Conta,^{118a,118b} F. Conventi,^{101a,j} J. Cook,²⁹ M. Cooke,¹⁴ B. D. Cooper,⁷⁶
A. M. Cooper-Sarkar,¹¹⁷ K. Copic,¹⁴ T. Cornelissen,¹⁷³ M. Corradi,^{19a} F. Corriveau,^{84,k} A. Cortes-Gonzalez,¹⁶⁴
G. Cortiana,⁹⁸ G. Costa,^{88a} M. J. Costa,¹⁶⁶ D. Costanzo,¹³⁸ T. Costin,³⁰ D. Côté,²⁹ R. Coura Torres,^{23a}
L. Courneyea,¹⁶⁸ G. Cowan,⁷⁵ C. Cowden,²⁷ B. E. Cox,⁸¹ K. Cranmer,¹⁰⁷ F. Crescioli,^{121a,121b} M. Cristinziani,²⁰
G. Crosetti,^{36a,36b} R. Crupi,^{71a,71b} S. Crépé-Renaudin,⁵⁴ C.-M. Cuciuc,^{25a} C. Cuenca Almenar,¹⁷⁴
T. Cuhadar Donszelmann,¹³⁸ M. Curatolo,⁴⁶ C. J. Curtis,¹⁷ C. Cuthbert,¹⁴⁹ P. Cwetanski,⁶⁰ H. Czirr,¹⁴⁰
Z. Czyczula,¹⁷⁴ S. D'Auria,⁵² M. D'Onofrio,⁷² A. D'Orazio,^{131a,131b} P. V. M. Da Silva,^{23a} C. Da Via,⁸¹
W. Dabrowski,³⁷ T. Dai,⁸⁶ C. Dallapiccola,⁸³ M. Dam,³⁵ M. Dameri,^{49a,49b} D. S. Damiani,¹³⁶ H. O. Danielsson,²⁹
D. Dannheim,⁹⁸ V. Dao,⁴⁸ G. Darbo,^{49a} G. L. Darlea,^{25b} C. Daum,¹⁰⁴ W. Davey,²⁰ T. Davidek,¹²⁵ N. Davidson,⁸⁵
R. Davidson,⁷⁰ E. Davies,^{117,d} M. Davies,⁹² A. R. Davison,⁷⁶ Y. Davygora,^{57a} E. Dawe,¹⁴¹ I. Dawson,¹³⁸
J. W. Dawson,^{5,a} R. K. Daya-Ishmukhametova,²² K. De,⁷ R. de Asmundis,^{101a} S. De Castro,^{19a,19b}
P. E. De Castro Faria Salgado,²⁴ S. De Cecco,⁷⁷ J. de Graat,⁹⁷ N. De Groot,¹⁰³ P. de Jong,¹⁰⁴ C. De La Taille,¹¹⁴
H. De la Torre,⁷⁹ B. De Lotto,^{163a,163c} L. de Mora,⁷⁰ L. De Nooij,¹⁰⁴ D. De Pedis,^{131a} A. De Salvo,^{131a}
U. De Sanctis,^{163a,163c} A. De Santo,¹⁴⁸ J. B. De Vivie De Regie,¹¹⁴ S. Dean,⁷⁶ W. J. Dearnaley,⁷⁰ R. Debye,²⁴
C. Debenedetti,⁴⁵ D. V. Dedovich,⁶⁴ J. Degenhardt,¹¹⁹ M. Dehchar,¹¹⁷ C. Del Papa,^{163a,163c} J. Del Peso,⁷⁹
T. Del Prete,^{121a,121b} T. Delemontex,⁵⁴ M. Deliyergiyev,⁷³ A. Dell'Acqua,²⁹ L. Dell'Asta,²¹ M. Della Pietra,^{101a,j}
D. della Volpe,^{101a,101b} M. Delmastro,⁴ N. Delruelle,²⁹ P. A. Delsart,⁵⁴ C. Deluca,¹⁴⁷ S. Demers,¹⁷⁴ M. Demichev,⁶⁴
B. Demirköz,¹¹¹ J. Deng,¹⁶² S. P. Denisov,¹²⁷ D. Derendarz,³⁸ J. E. Derkaoui,^{134d} F. Derue,⁷⁷ P. Dervan,⁷² K. Desch,²⁰
E. Devetak,¹⁴⁷ P. O. Deviveiros,¹⁰⁴ A. Dewhurst,¹²⁸ B. DeWilde,¹⁴⁷ S. Dhaliwal,¹⁵⁷ R. Dhullipudi,^{24,m}
A. Di Ciaccio,^{132a,132b} L. Di Ciaccio,⁴ A. Di Girolamo,²⁹ B. Di Girolamo,²⁹ S. Di Luise,^{133a,133b} A. Di Mattia,¹⁷¹
B. Di Micco,²⁹ R. Di Nardo,⁴⁶ A. Di Simone,^{132a,132b} R. Di Sipio,^{19a,19b} M. A. Diaz,^{31a} F. Diblen,^{18c} E. B. Diehl,⁸⁶
J. Dietrich,⁴¹ T. A. Dietzsch,^{57a} S. Diglio,⁸⁵ K. Dindar Yagci,³⁹ J. Dingfelder,²⁰ C. Dionisi,^{131a,131b} P. Dita,^{25a}
S. Dita,^{25a} F. Dittus,²⁹ F. Djama,⁸² T. Djobava,^{50b} M. A. B. do Vale,^{23c} A. Do Valle Wemans,^{123a} T. K. O. Doan,⁴
M. Dobbs,⁸⁴ R. Dobinson,^{29,a} D. Dobos,²⁹ E. Dobson,^{29,n} J. Dodd,³⁴ C. Doglioni,¹¹⁷ T. Doherty,⁵² Y. Doi,^{65,a}
J. Dolejsi,¹²⁵ I. Dolenc,⁷³ Z. Dolezal,¹²⁵ B. A. Dolgoshein,^{95,a} T. Dohmae,¹⁵⁴ M. Donadelli,^{23d} M. Donega,¹¹⁹
J. Donini,³³ J. Dopke,²⁹ A. Doria,^{101a} A. Dos Anjos,¹⁷¹ M. Dosil,¹¹ A. Dotti,^{121a,121b} M. T. Dova,⁶⁹ J. D. Dowell,¹⁷
A. D. Doxiadis,¹⁰⁴ A. T. Doyle,⁵² Z. Drasal,¹²⁵ J. Drees,¹⁷³ N. Dressnandt,¹¹⁹ H. Drevermann,²⁹ C. Driouichi,³⁵
M. Dris,⁹ J. Dubbert,⁹⁸ S. Dube,¹⁴ E. Duchovni,¹⁷⁰ G. Duckeck,⁹⁷ A. Dudarev,²⁹ F. Dudziak,⁶³ M. Dührssen,²⁹
I. P. Duerdoth,⁸¹ L. Duflot,¹¹⁴ M.-A. Dufour,⁸⁴ M. Dunford,²⁹ H. Duran Yildiz,^{3a} R. Duxfield,¹³⁸ M. Dwuznik,³⁷
F. Dydak,²⁹ M. Düren,⁵¹ W. L. Ebenstein,⁴⁴ J. Ebke,⁹⁷ S. Eckweiler,⁸⁰ K. Edmonds,⁸⁰ C. A. Edwards,⁷⁵
N. C. Edwards,⁵² W. Ehrenfeld,⁴¹ T. Ehrich,⁹⁸ T. Eifert,²⁹ G. Eigen,¹³ K. Einsweiler,¹⁴ E. Eisenhandler,⁷⁴
T. Ekelof,¹⁶⁵ M. El Kacimi,^{134c} M. Ellert,¹⁶⁵ S. Elles,⁴ F. Ellinghaus,⁸⁰ K. Ellis,⁷⁴ N. Ellis,²⁹ J. Elmsheuser,⁹⁷
M. Elsing,²⁹ D. Emeliyanov,¹²⁸ R. Engelmann,¹⁴⁷ A. Engl,⁹⁷ B. Epp,⁶¹ A. Eppig,⁸⁶ J. Erdmann,⁵³ A. Ereditato,¹⁶
D. Eriksson,^{145a} J. Ernst,¹ M. Ernst,²⁴ J. Ernwein,¹³⁵ D. Errede,¹⁶⁴ S. Errede,¹⁶⁴ E. Ertel,⁸⁰ M. Escalier,¹¹⁴
C. Escobar,¹²² X. Espinal Curull,¹¹ B. Esposito,⁴⁶ F. Etienne,⁸² A. I. Etievre,¹³⁵ E. Etzion,¹⁵² D. Evangelakou,⁵³
H. Evans,⁶⁰ L. Fabbri,^{19a,19b} C. Fabre,²⁹ R. M. Fakhruddinov,¹²⁷ S. Falciano,^{131a} Y. Fang,¹⁷¹ M. Fanti,^{88a,88b}
A. Farbin,⁷ A. Farilla,^{133a} J. Farley,¹⁴⁷ T. Farooque,¹⁵⁷ S. M. Farrington,¹¹⁷ P. Farthouat,²⁹ P. Fassnacht,²⁹
D. Fassouliotis,⁸ B. Fatholahzadeh,¹⁵⁷ A. Favareto,^{88a,88b} L. Fayard,¹¹⁴ S. Fazio,^{36a,36b} R. Febbraro,³³ P. Federic,^{143a}
O. L. Fedin,¹²⁰ W. Fedorko,⁸⁷ M. Fehling-Kaschek,⁴⁷ L. Felgioni,⁸² D. Fellmann,⁵ C. Feng,^{32d} E. J. Feng,³⁰
A. B. Fenyuk,¹²⁷ J. Ferencei,^{143b} J. Ferland,⁹² W. Fernando,¹⁰⁸ S. Ferrag,⁵² J. Ferrando,⁵² V. Ferrara,⁴¹ A. Ferrari,¹⁶⁵
P. Ferrari,¹⁰⁴ R. Ferrari,^{118a} A. Ferrer,¹⁶⁶ M. L. Ferrer,⁴⁶ D. Ferrere,⁴⁸ C. Ferretti,⁸⁶ A. Ferretto Parodi,^{49a,49b}
M. Fiascaris,³⁰ F. Fiedler,⁸⁰ A. Filipčić,⁷³ A. Filippas,⁹ F. Filthaut,¹⁰³ M. Fincke-Keeler,¹⁶⁸ M. C. N. Fiolhais,^{123a,i}
L. Fiorini,¹⁶⁶ A. Firan,³⁹ G. Fischer,⁴¹ P. Fischer,²⁰ M. J. Fisher,¹⁰⁸ M. Flechl,⁴⁷ I. Fleck,¹⁴⁰ J. Fleckner,⁸⁰
P. Fleischmann,¹⁷² S. Fleischmann,¹⁷³ T. Flick,¹⁷³ L. R. Flores Castillo,¹⁷¹ M. J. Flowerdew,⁹⁸ M. Fokitis,⁹
T. Fonseca Martin,¹⁶ D. A. Forbush,¹³⁷ A. Formica,¹³⁵ A. Forti,⁸¹ D. Fortin,^{158a} J. M. Foster,⁸¹ D. Fournier,¹¹⁴

- A. Foussat,²⁹ A. J. Fowler,⁴⁴ K. Fowler,¹³⁶ H. Fox,⁷⁰ P. Francavilla,^{121a,121b} S. Franchino,^{118a,118b} D. Francis,²⁹ T. Frank,¹⁷⁰ M. Franklin,⁵⁶ S. Franz,²⁹ M. Fraternali,^{118a,118b} S. Fratina,¹¹⁹ S. T. French,²⁷ F. Friedrich,⁴³ R. Froeschl,²⁹ D. Froidevaux,²⁹ J. A. Frost,²⁷ C. Fukunaga,¹⁵⁵ E. Fullana Torregrosa,²⁹ J. Fuster,¹⁶⁶ C. Gabaldon,²⁹ O. Gabizon,¹⁷⁰ T. Gadfort,²⁴ S. Gadomski,⁴⁸ G. Gagliardi,^{49a,49b} P. Gagnon,⁶⁰ C. Galea,⁹⁷ E. J. Gallas,¹¹⁷ V. Gallo,¹⁶ B. J. Gallop,¹²⁸ P. Gallus,¹²⁴ K. K. Gan,¹⁰⁸ Y. S. Gao,^{142,f} V. A. Gapienko,¹²⁷ A. Gaponenko,¹⁴ F. Garbersson,¹⁷⁴ M. Garcia-Sciveres,¹⁴ C. García,¹⁶⁶ J. E. García Navarro,¹⁶⁶ R. W. Gardner,³⁰ N. Garelli,²⁹ H. Garitaonandia,¹⁰⁴ V. Garonne,²⁹ J. Garvey,¹⁷ C. Gatti,⁴⁶ G. Gaudio,^{118a} O. Gaumer,⁴⁸ B. Gaur,¹⁴⁰ L. Gauthier,¹³⁵ I. L. Gavrilenko,⁹³ C. Gay,¹⁶⁷ G. Gaycken,²⁰ J.-C. Gayde,²⁹ E. N. Gazis,⁹ P. Ge,^{32d} C. N. P. Gee,¹²⁸ D. A. A. Geerts,¹⁰⁴ Ch. Geich-Gimbel,²⁰ K. Gellerstedt,^{145a,145b} C. Gemme,^{49a} A. Gemmell,⁵² M. H. Genest,⁹⁷ S. Gentile,^{131a,131b} M. George,⁵³ S. George,⁷⁵ P. Gerlach,¹⁷³ A. Gershon,¹⁵² C. Geweniger,^{57a} H. Ghazlane,^{134b} N. Ghodbane,³³ B. Giacobbe,^{19a} S. Giagu,^{131a,131b} V. Giakoumopoulou,⁸ V. Giangiobbe,¹¹ F. Gianotti,²⁹ B. Gibbard,²⁴ A. Gibson,¹⁵⁷ S. M. Gibson,²⁹ L. M. Gilbert,¹¹⁷ V. Gilewsky,⁹⁰ D. Gillberg,²⁸ A. R. Gillman,¹²⁸ D. M. Gingrich,^{2,e} J. Ginzburg,¹⁵² N. Giokaris,⁸ M. P. Giordani,^{163c} R. Giordano,^{101a,101b} F. M. Giorgi,¹⁵ P. Giovannini,⁹⁸ P. F. Giraud,¹³⁵ D. Giugni,^{88a} M. Giunta,⁹² P. Giusti,^{19a} B. K. Gjelsten,¹¹⁶ L. K. Gladilin,⁹⁶ C. Glasman,⁷⁹ J. Glatzer,⁴⁷ A. Glazov,⁴¹ K. W. Glitza,¹⁷³ G. L. Glonti,⁶⁴ J. R. Goddard,⁷⁴ J. Godfrey,¹⁴¹ J. Godlewski,²⁹ M. Goebel,⁴¹ T. Göpfert,⁴³ C. Goeringer,⁸⁰ C. Gössling,⁴² T. Göttfert,⁹⁸ S. Goldfarb,⁸⁶ T. Golling,¹⁷⁴ S. N. Golovnia,¹²⁷ A. Gomes,^{123a,c} L. S. Gomez Fajardo,⁴¹ R. Gonçalves,⁷⁵ J. Goncalves Pinto Firmino Da Costa,⁴¹ L. Gonella,²⁰ A. Gonidec,²⁹ S. Gonzalez,¹⁷¹ S. González de la Hoz,¹⁶⁶ G. Gonzalez Parra,¹¹ M. L. Gonzalez Silva,²⁶ S. Gonzalez-Sevilla,⁴⁸ J. J. Goodson,¹⁴⁷ L. Goossens,²⁹ P. A. Gorbounov,⁹⁴ H. A. Gordon,²⁴ I. Gorelov,¹⁰² G. Gorfine,¹⁷³ B. Gorini,²⁹ E. Gorini,^{71a,71b} A. Gorišek,⁷³ E. Gornicki,³⁸ S. A. Gorokhov,¹²⁷ V. N. Goryachev,¹²⁷ B. Gosdzik,⁴¹ M. Gosselink,¹⁰⁴ M. I. Gostkin,⁶⁴ I. Gough Eschrich,¹⁶² M. Gouighri,^{134a} D. Goujdami,^{134c} M. P. Goulette,⁴⁸ A. G. Goussiou,¹³⁷ C. Goy,⁴ S. Gozpinar,²² I. Grabowska-Bold,³⁷ P. Grafström,²⁹ K.-J. Grahm,⁴¹ F. Grancagnolo,^{71a} S. Grancagnolo,¹⁵ V. Grassi,¹⁴⁷ V. Gratchev,¹²⁰ N. Grau,³⁴ H. M. Gray,²⁹ J. A. Gray,¹⁴⁷ E. Graziani,^{133a} O. G. Grebenyuk,¹²⁰ T. Greenshaw,⁷² Z. D. Greenwood,^{24,m} K. Gregersen,³⁵ I. M. Gregor,⁴¹ P. Grenier,¹⁴² J. Griffiths,¹³⁷ N. Grigalashvili,⁶⁴ A. A. Grillo,¹³⁶ S. Grinstein,¹¹ Y. V. Grishkevich,⁹⁶ J.-F. Grivaz,¹¹⁴ M. Groh,⁹⁸ E. Gross,¹⁷⁰ J. Grosse-Knetter,⁵³ J. Groth-Jensen,¹⁷⁰ K. Grybel,¹⁴⁰ V. J. Guarino,⁵ D. Guest,¹⁷⁴ C. Guichenev,³³ A. Guida,^{71a,71b} S. Guindon,⁵³ H. Guler,^{84,o} J. Gunther,¹²⁴ B. Guo,¹⁵⁷ J. Guo,³⁴ A. Gupta,³⁰ Y. Gusakov,⁶⁴ V. N. Gushchin,¹²⁷ A. Gutierrez,⁹² P. Gutierrez,¹¹⁰ N. Guttman,¹⁵² O. Gutzwiller,¹⁷¹ C. Guyot,¹³⁵ C. Gwenlan,¹¹⁷ C. B. Gwilliam,⁷² A. Haas,¹⁴² S. Haas,²⁹ C. Haber,¹⁴ H. K. Hadavand,³⁹ D. R. Hadley,¹⁷ P. Haefner,⁹⁸ F. Hahn,²⁹ S. Haider,²⁹ Z. Hajduk,³⁸ H. Hakobyan,¹⁷⁵ D. Hall,¹¹⁷ J. Haller,⁵³ K. Hamacher,¹⁷³ P. Hamal,¹¹² M. Hamer,⁵³ A. Hamilton,^{144b,p} S. Hamilton,¹⁶⁰ H. Han,^{32a} L. Han,^{32b} K. Hanagaki,¹¹⁵ K. Hanawa,¹⁵⁹ M. Hance,¹⁴ C. Handel,⁸⁰ P. Hanke,^{57a} J. R. Hansen,³⁵ J. B. Hansen,³⁵ J. D. Hansen,³⁵ P. H. Hansen,³⁵ P. Hansson,¹⁴² K. Hara,¹⁵⁹ G. A. Hare,¹³⁶ T. Harenberg,¹⁷³ S. Harkusha,⁸⁹ D. Harper,⁸⁶ R. D. Harrington,⁴⁵ O. M. Harris,¹³⁷ K. Harrison,¹⁷ J. Hartert,⁴⁷ F. Hartjes,¹⁰⁴ T. Haruyama,⁶⁵ A. Harvey,⁵⁵ S. Hasegawa,¹⁰⁰ Y. Hasegawa,¹³⁹ S. Hassani,¹³⁵ M. Hatch,²⁹ D. Hauff,⁹⁸ S. Haug,¹⁶ M. Hauschild,²⁹ R. Hauser,⁸⁷ M. Havranek,²⁰ B. M. Hawes,¹¹⁷ C. M. Hawkes,¹⁷ R. J. Hawkins,²⁹ A. D. Hawkins,⁷⁸ D. Hawkins,¹⁶² T. Hayakawa,⁶⁶ T. Hayashi,¹⁵⁹ D. Hayden,⁷⁵ H. S. Hayward,⁷² S. J. Haywood,¹²⁸ E. Hazen,²¹ M. He,^{32d} S. J. Head,¹⁷ V. Hedberg,⁷⁸ L. Heelan,⁷ S. Heim,⁸⁷ B. Heinemann,¹⁴ S. Heisterkamp,³⁵ L. Helary,⁴ C. Heller,⁹⁷ M. Heller,²⁹ S. Hellman,^{145a,145b} D. Hellmich,²⁰ C. Helsen,¹¹ R. C. W. Henderson,⁷⁰ M. Henke,^{57a} A. Henrichs,⁵³ A. M. Henriques Correia,²⁹ S. Henrot-Versille,¹¹⁴ F. Henry-Couannier,⁸² C. Hensel,⁵³ T. Henß,¹⁷³ C. M. Hernandez,⁷ Y. Hernández Jiménez,¹⁶⁶ R. Herrberg,¹⁵ A. D. Hershenhorn,¹⁵¹ G. Herten,⁴⁷ R. Hertenberger,⁹⁷ L. Hervas,²⁹ N. P. Hesse,¹⁰⁴ E. Higón-Rodríguez,¹⁶⁶ D. Hill,^{5,a} J. C. Hill,²⁷ N. Hill,⁵ K. H. Hiller,⁴¹ S. Hillert,²⁰ S. J. Hillier,¹⁷ I. Hinchliffe,¹⁴ E. Hines,¹¹⁹ M. Hirose,¹¹⁵ F. Hirsch,⁴² D. Hirschbuehl,¹⁷³ J. Hobbs,¹⁴⁷ N. Hod,¹⁵² M. C. Hodgkinson,¹³⁸ P. Hodgson,¹³⁸ A. Hoecker,²⁹ M. R. Hoferkamp,¹⁰² J. Hoffman,³⁹ D. Hoffmann,⁸² M. Hohlfeld,⁸⁰ M. Holder,¹⁴⁰ S. O. Holmgren,^{145a} T. Holy,¹²⁶ J. L. Holzbauer,⁸⁷ Y. Homma,⁶⁶ T. M. Hong,¹¹⁹ L. Hooft van Huysduynen,¹⁰⁷ T. Horazdovsky,¹²⁶ C. Horn,¹⁴² S. Horner,⁴⁷ J.-Y. Hostachy,⁵⁴ S. Hou,¹⁵⁰ M. A. Houlden,⁷² A. Hoummada,^{134a} J. Howarth,⁸¹ D. F. Howell,¹¹⁷ I. Hristova,¹⁵ J. Hrivnac,¹¹⁴ I. Hruska,¹²⁴ T. Hryn'ova,⁴ P. J. Hsu,⁸⁰ S.-C. Hsu,¹⁴ G. S. Huang,¹¹⁰ Z. Hubacek,¹²⁶ F. Hubaut,⁸² F. Hügging,²⁰ T. B. Huffman,¹¹⁷ E. W. Hughes,³⁴ G. Hughes,⁷⁰ R. E. Hughes-Jones,⁸¹ M. Huhtinen,²⁹ P. Hurst,⁵⁶ M. Hurwitz,¹⁴ U. Husemann,⁴¹ N. Huseynov,^{64,q} J. Huston,⁸⁷ J. Huth,⁵⁶ G. Iacobucci,⁴⁸ G. Iakovidis,⁹ M. Ibbotson,⁸¹ I. Ibragimov,¹⁴⁰ R. Ichimiya,⁶⁶ L. Iconomidou-Fayard,¹¹⁴ J. Idarraga,¹¹⁴ P. Iengo,^{101a} O. Igonkina,¹⁰⁴ Y. Ikegami,⁶⁵ M. Ikeno,⁶⁵ Y. Ilchenko,³⁹ D. Iliadis,¹⁵³ N. Ilic,¹⁵⁷ D. Imbault,⁷⁷ M. Imori,¹⁵⁴ T. Ince,²⁰ J. Inigo-Golfin,²⁹ P. Ioannou,⁸ M. Iodice,^{133a}

A. Irles Quiles,¹⁶⁶ C. Isaksson,¹⁶⁵ A. Ishikawa,⁶⁶ M. Ishino,⁶⁷ R. Ishmukhametov,³⁹ C. Issever,¹¹⁷ S. Istin,^{18a}
A. V. Ivashin,¹²⁷ W. Iwanski,³⁸ H. Iwasaki,⁶⁵ J. M. Izen,⁴⁰ V. Izzo,^{101a} B. Jackson,¹¹⁹ J. N. Jackson,⁷² P. Jackson,¹⁴²
M. R. Jaekel,²⁹ V. Jain,⁶⁰ K. Jakobs,⁴⁷ S. Jakobsen,³⁵ J. Jakubek,¹²⁶ D. K. Jana,¹¹⁰ E. Jankowski,¹⁵⁷ E. Jansen,⁷⁶
H. Jansen,²⁹ A. Jantsch,⁹⁸ M. Janus,²⁰ G. Jarlskog,⁷⁸ L. Jeanty,⁵⁶ K. Jelen,³⁷ I. Jen-La Plante,³⁰ P. Jenni,²⁹
A. Jeremie,⁴ P. Jež,³⁵ S. Jézéquel,⁴ M. K. Jha,^{19a} H. Ji,¹⁷¹ W. Ji,⁸⁰ J. Jia,¹⁴⁷ Y. Jiang,^{32b} M. Jimenez Belenguer,⁴¹
G. Jin,^{32b} S. Jin,^{32a} O. Jinnouchi,¹⁵⁶ M. D. Joergensen,³⁵ D. Joffe,³⁹ L. G. Johansen,¹³ M. Johansen,^{145a,145b}
K. E. Johansson,^{145a} P. Johansson,¹³⁸ S. Johnert,⁴¹ K. A. Johns,⁶ K. Jon-And,^{145a,145b} G. Jones,⁸¹ R. W. L. Jones,⁷⁰
T. W. Jones,⁷⁶ T. J. Jones,⁷² O. Jonsson,²⁹ C. Joram,²⁹ P. M. Jorge,^{123a} J. Joseph,¹⁴ T. Jovin,^{12b} X. Ju,¹⁷¹ C. A. Jung,⁴²
R. M. Jungst,²⁹ V. Juranek,¹²⁴ P. Jussel,⁶¹ A. Juste Rozas,¹¹ V. V. Kabachenko,¹²⁷ S. Kabana,¹⁶ M. Kaci,¹⁶⁶
A. Kaczmarska,³⁸ P. Kadlecik,³⁵ M. Kado,¹¹⁴ H. Kagan,¹⁰⁸ M. Kagan,⁵⁶ S. Kaiser,⁹⁸ E. Kajomovitz,¹⁵¹ S. Kalinin,¹⁷³
L. V. Kalinovskaya,⁶⁴ S. Kama,³⁹ N. Kanaya,¹⁵⁴ M. Kaneda,²⁹ S. Kaneti,²⁷ T. Kanno,¹⁵⁶ V. A. Kantserov,⁹⁵
J. Kanzaki,⁶⁵ B. Kaplan,¹⁷⁴ A. Kapliy,³⁰ J. Kaplon,²⁹ D. Kar,⁴³ M. Karagounis,²⁰ M. Karagoz,¹¹⁷ M. Karnevskiy,⁴¹
K. Karr,⁵ V. Kartvelishvili,⁷⁰ A. N. Karyukhin,¹²⁷ L. Kashif,¹⁷¹ G. Kasieczka,^{57b} R. D. Kass,¹⁰⁸ A. Kastanas,¹³
M. Kataoka,⁴ Y. Kataoka,¹⁵⁴ E. Katsoufis,⁹ J. Katzy,⁴¹ V. Kaushik,⁶ K. Kawagoe,⁶⁶ T. Kawamoto,¹⁵⁴ G. Kawamura,⁸⁰
M. S. Kayl,¹⁰⁴ V. A. Kazanin,¹⁰⁶ M. Y. Kazarinov,⁶⁴ R. Keeler,¹⁶⁸ R. Kehoe,³⁹ M. Keil,⁵³ G. D. Kekelidze,⁶⁴
J. Kennedy,⁹⁷ C. J. Kenney,¹⁴² M. Kenyon,⁵² O. Kepka,¹²⁴ N. Kerschen,²⁹ B. P. Kerševan,⁷³ S. Kersten,¹⁷³
K. Kessoku,¹⁵⁴ J. Keung,¹⁵⁷ F. Khalil-zada,¹⁰ H. Khandanyan,¹⁶⁴ A. Khanov,¹¹¹ D. Kharchenko,⁶⁴ A. Khodinov,⁹⁵
A. G. Kholodenko,¹²⁷ A. Khomich,^{57a} T. J. Khoo,²⁷ G. Khoriauli,²⁰ A. Khoroshilov,¹⁷³ N. Khovanskiy,⁶⁴
V. Khovanskiy,⁹⁴ E. Khramov,⁶⁴ J. Khubua,^{50b} H. Kim,^{145a,145b} M. S. Kim,² P. C. Kim,¹⁴² S. H. Kim,¹⁵⁹
N. Kimura,¹⁶⁹ O. Kind,¹⁵ B. T. King,⁷² M. King,⁶⁶ R. S. B. King,¹¹⁷ J. Kirk,¹²⁸ L. E. Kirsch,²² A. E. Kiryunin,⁹⁸
T. Kishimoto,⁶⁶ D. Kisieleska,³⁷ T. Kittelmann,¹²² A. M. Kiver,¹²⁷ E. Kladiva,^{143b} J. Klaiber-Lodewigs,⁴²
M. Klein,⁷² U. Klein,⁷² K. Kleinknecht,⁸⁰ M. Klemetti,⁸⁴ A. Klier,¹⁷⁰ P. Klimek,^{145a,145b} A. Klimentov,²⁴
R. Klingenberg,⁴² E. B. Klinkby,³⁵ T. Klioutchnikova,²⁹ P. F. Klok,¹⁰³ S. Klous,¹⁰⁴ E.-E. Kluge,^{57a} T. Kluge,⁷²
P. Kluit,¹⁰⁴ S. Kluth,⁹⁸ N. S. Knecht,¹⁵⁷ E. Kneringer,⁶¹ J. Knobloch,²⁹ E.B.F.G. Knoops,⁸² A. Knue,⁵³ B. R. Ko,⁴⁴
T. Kobayashi,¹⁵⁴ M. Kobel,⁴³ M. Kocian,¹⁴² P. Kodys,¹²⁵ K. Köneke,²⁹ A. C. König,¹⁰³ S. Koenig,⁸⁰ L. Köpke,⁸⁰
F. Koetsveld,¹⁰³ P. Koevesarki,²⁰ T. Koffas,²⁸ E. Koffeman,¹⁰⁴ F. Kohn,⁵³ Z. Kohout,¹²⁶ T. Kohriki,⁶⁵ T. Koi,¹⁴²
T. Kokott,²⁰ G. M. Kolachev,¹⁰⁶ H. Kolanoski,¹⁵ V. Kolesnikov,⁶⁴ I. Koletsou,^{88a} J. Koll,⁸⁷ D. Kollar,²⁹
M. Kollefrath,⁴⁷ S. D. Kolya,⁸¹ A. A. Komar,⁹³ Y. Komori,¹⁵⁴ T. Kondo,⁶⁵ T. Kono,^{41,r} A. I. Kononov,⁴⁷
R. Konoplich,^{107,s} N. Konstantinidis,⁷⁶ A. Kootz,¹⁷³ S. Koperny,³⁷ K. Korcyl,³⁸ K. Kordas,¹⁵³ V. Koreshev,¹²⁷
A. Korn,¹¹⁷ A. Korol,¹⁰⁶ I. Korolkov,¹¹ E. V. Korolkova,¹³⁸ V. A. Korotkov,¹²⁷ O. Kortner,⁹⁸ S. Kortner,⁹⁸
V. V. Kostyukhin,²⁰ M. J. Kotamäki,²⁹ S. Kotov,⁹⁸ V. M. Kotov,⁶⁴ A. Kotwal,⁴⁴ C. Kourkoumelis,⁸ V. Kouskoura,¹⁵³
A. Koutsman,^{158a} R. Kowalewski,¹⁶⁸ T. Z. Kowalski,³⁷ W. Kozanecki,¹³⁵ A. S. Kozhin,¹²⁷ V. Kral,¹²⁶
V. A. Kramarenko,⁹⁶ G. Kramberger,⁷³ M. W. Krasny,⁷⁷ A. Krasznahorkay,¹⁰⁷ J. Kraus,⁸⁷ J. K. Kraus,²⁰ A. Kreisel,¹⁵²
F. Krejci,¹²⁶ J. Kretzschmar,⁷² N. Krieger,⁵³ P. Krieger,¹⁵⁷ K. Kroeninger,⁵³ H. Kroha,⁹⁸ J. Kroll,¹¹⁹ J. Kroseberg,²⁰
J. Krstic,^{12a} U. Kruchonak,⁶⁴ H. Krüger,²⁰ T. Kruker,¹⁶ N. Krumnack,⁶³ Z. V. Krumshcheyn,⁶⁴ A. Kruth,²⁰ T. Kubota,⁸⁵
S. Kuehn,⁴⁷ A. Kugel,^{57c} T. Kuhl,⁴¹ D. Kuhn,⁶¹ V. Kukhtin,⁶⁴ Y. Kulchitsky,⁸⁹ S. Kuleshov,^{31b} C. Kummer,⁹⁷
M. Kuna,⁷⁷ N. Kundu,¹¹⁷ J. Kunkle,¹¹⁹ A. Kupco,¹²⁴ H. Kurashige,⁶⁶ M. Kurata,¹⁵⁹ Y. A. Kurochkin,⁸⁹ V. Kus,¹²⁴
M. Kuze,¹⁵⁶ J. Kvita,¹⁴¹ R. Kwee,¹⁵ A. La Rosa,⁴⁸ L. La Rotonda,^{36a,36b} L. Labarga,⁷⁹ J. Labbe,⁴ S. Lablak,^{134a}
C. Lacasta,¹⁶⁶ F. Lacava,^{131a,131b} H. Lacker,¹⁵ D. Lacour,⁷⁷ V. R. Lacuesta,¹⁶⁶ E. Ladygin,⁶⁴ R. Lafaye,⁴ B. Laforge,⁷⁷
T. Lagouri,⁷⁹ S. Lai,⁴⁷ E. Laisne,⁵⁴ M. Lamanna,²⁹ C. L. Lampen,⁶ W. Lampl,⁶ E. Lancon,¹³⁵ U. Landgraf,⁴⁷
M. P. J. Landon,⁷⁴ H. Landsman,¹⁵¹ J. L. Lane,⁸¹ C. Lange,⁴¹ A. J. Lankford,¹⁶² F. Lanni,²⁴ K. Lantzsch,¹⁷³
S. Laplace,⁷⁷ C. Lapoire,²⁰ J. F. Laporte,¹³⁵ T. Lari,^{88a} A. V. Larionov,¹²⁷ A. Larner,¹¹⁷ C. Lasseur,²⁹ M. Lassnig,²⁹
P. Laurelli,⁴⁶ W. Lavrijsen,¹⁴ P. Laycock,⁷² A. B. Lazarev,⁶⁴ O. Le Dortz,⁷⁷ E. Le Guirriec,⁸² C. Le Maner,¹⁵⁷
E. Le Menedeu,⁹ C. Lebel,⁹² T. LeCompte,⁵ F. Ledroit-Guillon,⁵⁴ H. Lee,¹⁰⁴ J. S. H. Lee,¹¹⁵ S. C. Lee,¹⁵⁰ L. Lee,¹⁷⁴
M. Lefebvre,¹⁶⁸ M. Legendre,¹³⁵ A. Leger,⁴⁸ B. C. LeGeyt,¹¹⁹ F. Legger,⁹⁷ C. Leggett,¹⁴ M. Lehmacher,²⁰
G. Lehmann Miotto,²⁹ X. Lei,⁶ M. A. L. Leite,^{23d} R. Leitner,¹²⁵ D. Lellouch,¹⁷⁰ M. Leltchouk,³⁴ B. Lemmer,⁵³
V. Lendermann,^{57a} K. J. C. Leney,^{144b} T. Lenz,¹⁰⁴ G. Lenzen,¹⁷³ B. Lenzi,²⁹ K. Leonhardt,⁴³ S. Leontsinis,⁹
C. Leroy,⁹² J.-R. Lessard,¹⁶⁸ J. Lesser,^{145a} C. G. Lester,²⁷ A. Leung Fook Cheong,¹⁷¹ J. Levêque,⁴ D. Levin,⁸⁶
L. J. Levinson,¹⁷⁰ M. S. Levitski,¹²⁷ A. Lewis,¹¹⁷ G. H. Lewis,¹⁰⁷ A. M. Leyko,²⁰ M. Leyton,¹⁵ B. Li,⁸² H. Li,^{171,t}
S. Li,^{32b,u} X. Li,⁸⁶ Z. Liang,^{117,v} H. Liao,³³ B. Liberti,^{132a} P. Lichard,²⁹ M. Lichtnecker,⁹⁷ K. Lie,¹⁶⁴ W. Liebig,¹³
R. Lifshitz,¹⁵¹ C. Limbach,²⁰ A. Limosani,⁸⁵ M. Limper,⁶² S. C. Lin,^{150,w} F. Linde,¹⁰⁴ J. T. Linnemann,⁸⁷

E. Lipeles,¹¹⁹ L. Lipinsky,¹²⁴ A. Lipniacka,¹³ T. M. Liss,¹⁶⁴ D. Lissauer,²⁴ A. Lister,⁴⁸ A. M. Litke,¹³⁶ C. Liu,²⁸ D. Liu,¹⁵⁰ H. Liu,⁸⁶ J. B. Liu,⁸⁶ M. Liu,^{32b} S. Liu,² Y. Liu,^{32b} M. Livan,^{118a,118b} S. S. A. Livermore,¹¹⁷ A. Lleres,⁵⁴ J. Llorente Merino,⁷⁹ S. L. Lloyd,⁷⁴ E. Lobodzinska,⁴¹ P. Loch,⁶ W. S. Lockman,¹³⁶ T. Loddenkoetter,²⁰ F. K. Loebinger,⁸¹ A. Loginov,¹⁷⁴ C. W. Loh,¹⁶⁷ T. Lohse,¹⁵ K. Lohwasser,⁴⁷ M. Lokajicek,¹²⁴ J. Loken,¹¹⁷ V. P. Lombardo,⁴ R. E. Long,⁷⁰ L. Lopes,^{123a,c} D. Lopez Mateos,⁵⁶ J. Lorenz,⁹⁷ M. Losada,¹⁶¹ P. Loscutoff,¹⁴ F. Lo Sterzo,^{131a,131b} M. J. Losty,^{158a} X. Lou,⁴⁰ A. Lounis,¹¹⁴ K. F. Loureiro,¹⁶¹ J. Love,²¹ P. A. Love,⁷⁰ A. J. Lowe,^{142,f} F. Lu,^{32a} H. J. Lubatti,¹³⁷ C. Luci,^{131a,131b} A. Lucotte,⁵⁴ A. Ludwig,⁴³ D. Ludwig,⁴¹ I. Ludwig,⁴⁷ J. Ludwig,⁴⁷ F. Luehring,⁶⁰ G. Luijckx,¹⁰⁴ D. Lumb,⁴⁷ L. Luminari,^{131a} E. Lund,¹¹⁶ B. Lund-Jensen,¹⁴⁶ B. Lundberg,⁷⁸ J. Lundberg,^{145a,145b} J. Lundquist,³⁵ M. Lungwitz,⁸⁰ G. Lutz,⁹⁸ D. Lynn,²⁴ J. Lys,¹⁴ E. Lytken,⁷⁸ H. Ma,²⁴ L. L. Ma,¹⁷¹ J. A. Macana Goia,⁹² G. Maccarrone,⁴⁶ A. Macchiolo,⁹⁸ B. Maček,⁷³ J. Machado Miguens,^{123a} R. Mackeprang,³⁵ R. J. Madaras,¹⁴ W. F. Mader,⁴³ R. Maenner,^{57c} T. Maeno,²⁴ P. Mättig,¹⁷³ S. Mättig,⁴¹ L. Magnoni,²⁹ E. Magradze,⁵³ Y. Mahalalel,¹⁵² K. Mahboubi,⁴⁷ G. Mahout,¹⁷ C. Maiani,^{131a,131b} C. Maidantchik,^{23a} A. Maio,^{123a,c} S. Majewski,²⁴ Y. Makida,⁶⁵ N. Makovec,¹¹⁴ P. Mal,¹³⁵ B. Malaescu,²⁹ Pa. Malecki,³⁸ P. Malecki,³⁸ V. P. Maleev,¹²⁰ F. Malek,⁵⁴ U. Mallik,⁶² D. Malon,⁵ C. Malone,¹⁴² S. Maltezos,⁹ V. Malyshev,¹⁰⁶ S. Malyukov,²⁹ R. Mameghani,⁹⁷ J. Mamuzic,^{12b} A. Manabe,⁶⁵ L. Mandelli,^{88a} I. Mandić,⁷³ R. Mandrysch,¹⁵ J. Maneira,^{123a} P. S. Mangedard,⁸⁷ I. D. Manjavidze,⁶⁴ A. Mann,⁵³ P. M. Manning,¹³⁶ A. Manousakis-Katsikakis,⁸ B. Mansoulie,¹³⁵ A. Manz,⁹⁸ A. Mapelli,²⁹ L. Mapelli,²⁹ L. March,⁷⁹ J. F. Marchand,²⁸ F. Marchese,^{132a,132b} G. Marchiori,⁷⁷ M. Marcisovsky,¹²⁴ A. Marin,^{21,a} C. P. Marino,¹⁶⁸ F. Marroquim,^{23a} R. Marshall,⁸¹ Z. Marshall,²⁹ F. K. Martens,¹⁵⁷ S. Marti-Garcia,¹⁶⁶ A. J. Martin,¹⁷⁴ B. Martin,²⁹ B. Martin,⁸⁷ F. F. Martin,¹¹⁹ J. P. Martin,⁹² Ph. Martin,⁵⁴ T. A. Martin,¹⁷ V. J. Martin,⁴⁵ B. Martin dit Latour,⁴⁸ S. Martin-Haugh,¹⁴⁸ M. Martinez,¹¹ V. Martinez Outschoorn,⁵⁶ A. C. Martyniuk,¹⁶⁸ M. Marx,⁸¹ F. Marzano,^{131a} A. Marzin,¹¹⁰ L. Masetti,⁸⁰ T. Mashimo,¹⁵⁴ R. Mashinistov,⁹³ J. Masik,⁸¹ A. L. Maslennikov,¹⁰⁶ I. Massa,^{19a,19b} G. Massaro,¹⁰⁴ N. Massol,⁴ P. Mastrandrea,^{131a,131b} A. Mastroberardino,^{36a,36b} T. Masubuchi,¹⁵⁴ M. Mathes,²⁰ P. Matricon,¹¹⁴ H. Matsumoto,¹⁵⁴ H. Matsunaga,¹⁵⁴ T. Matsushita,⁶⁶ C. Mattravers,^{117,d} J. M. Maugain,²⁹ J. Maurer,⁸² S. J. Maxfield,⁷² D. A. Maximov,^{106,g} E. N. May,⁵ A. Mayne,¹³⁸ R. Mazini,¹⁵⁰ M. Mazur,²⁰ M. Mazzanti,^{88a} E. Mazzoni,^{121a,121b} S. P. Mc Kee,⁸⁶ A. McCarn,¹⁶⁴ R. L. McCarthy,¹⁴⁷ T. G. McCarthy,²⁸ N. A. McCubbin,¹²⁸ K. W. McFarlane,⁵⁵ J. A. McFayden,¹³⁸ H. McGlone,⁵² G. Mchedlidze,^{50b} R. A. McLaren,²⁹ T. McLaughlan,¹⁷ S. J. McMahan,¹²⁸ R. A. McPherson,^{168,k} A. Meade,⁸³ J. Mechnich,¹⁰⁴ M. Mechtel,¹⁷³ M. Medinnis,⁴¹ R. Meera-Lebbai,¹¹⁰ T. Meguro,¹¹⁵ R. Mehdiyev,⁹² S. Mehlhase,³⁵ A. Mehta,⁷² K. Meier,^{57a} B. Meirose,⁷⁸ C. Melachrinou,³⁰ B. R. Mellado Garcia,¹⁷¹ L. Mendoza Navas,¹⁶¹ Z. Meng,^{150,t} A. Mengarelli,^{19a,19b} S. Menke,⁹⁸ C. Menot,²⁹ E. Meoni,¹¹ K. M. Mercurio,⁵⁶ P. Mermod,⁴⁸ L. Merola,^{101a,101b} C. Meroni,^{88a} F. S. Merritt,³⁰ A. Messina,²⁹ J. Metcalfe,¹⁰² A. S. Mete,⁶³ C. Meyer,⁸⁰ C. Meyer,³⁰ J.-P. Meyer,¹³⁵ J. Meyer,¹⁷² J. Meyer,⁵³ T. C. Meyer,²⁹ W. T. Meyer,⁶³ J. Miao,^{32d} S. Michal,²⁹ L. Micu,^{25a} R. P. Middleton,¹²⁸ S. Migas,⁷² L. Mijović,⁴¹ G. Mikenberg,¹⁷⁰ M. Mikestikova,¹²⁴ M. Mikuž,⁷³ D. W. Miller,³⁰ R. J. Miller,⁸⁷ W. J. Mills,¹⁶⁷ C. Mills,⁵⁶ A. Milov,¹⁷⁰ D. A. Milstead,^{145a,145b} D. Milstein,¹⁷⁰ A. A. Minaenko,¹²⁷ M. Miñano Moya,¹⁶⁶ I. A. Minashvili,⁶⁴ A. I. Mincer,¹⁰⁷ B. Mindur,³⁷ M. Mineev,⁶⁴ Y. Ming,¹⁷¹ L. M. Mir,¹¹ G. Mirabelli,^{131a} L. Miralles Verge,¹¹ A. Misiejuk,⁷⁵ J. Mitrevski,¹³⁶ G. Y. Mitrofanov,¹²⁷ V. A. Mitsou,¹⁶⁶ S. Mitsui,⁶⁵ P. S. Miyagawa,¹³⁸ K. Miyazaki,⁶⁶ J. U. Mjörnmark,⁷⁸ T. Moa,^{145a,145b} P. Mockett,¹³⁷ S. Moed,⁵⁶ V. Moeller,²⁷ K. Mönig,⁴¹ N. Möser,²⁰ S. Mohapatra,¹⁴⁷ W. Mohr,⁴⁷ S. Mohrdieck-Möck,⁹⁸ A. M. Moiseev,^{127,a} R. Moles-Valls,¹⁶⁶ J. Molina-Perez,²⁹ J. Monk,⁷⁶ E. Monnier,⁸² S. Montesano,^{88a,88b} F. Monticelli,⁶⁹ S. Monzani,^{19a,19b} R. W. Moore,² G. F. Moorhead,⁸⁵ C. Mora Herrera,⁴⁸ A. Moraes,⁵² N. Morange,¹³⁵ J. Morel,⁵³ G. Morello,^{36a,36b} D. Moreno,⁸⁰ M. Moreno Llácer,¹⁶⁶ P. Morettini,^{49a} M. Morii,⁵⁶ J. Morin,⁷⁴ A. K. Morley,²⁹ G. Mornacchi,²⁹ S. V. Morozov,⁹⁵ J. D. Morris,⁷⁴ L. Morvaj,¹⁰⁰ H. G. Moser,⁹⁸ M. Mosidze,^{50b} J. Moss,¹⁰⁸ R. Mount,¹⁴² E. Mountricha,^{9,x} S. V. Mouraviev,⁹³ E. J. W. Moyse,⁸³ M. Mudrinic,^{12b} F. Mueller,^{57a} J. Mueller,¹²² K. Mueller,²⁰ T. A. Müller,⁹⁷ T. Mueller,⁸⁰ D. Muenstermann,²⁹ A. Muir,¹⁶⁷ Y. Munwes,¹⁵² W. J. Murray,¹²⁸ I. Mussche,¹⁰⁴ E. Musto,^{101a,101b} A. G. Myagkov,¹²⁷ M. Myska,¹²⁴ J. Nadal,¹¹ K. Nagai,¹⁵⁹ K. Nagano,⁶⁵ Y. Nagasaka,⁵⁹ M. Nagel,⁹⁸ A. M. Nairz,²⁹ Y. Nakahama,²⁹ K. Nakamura,¹⁵⁴ T. Nakamura,¹⁵⁴ I. Nakano,¹⁰⁹ G. Nanava,²⁰ A. Napier,¹⁶⁰ M. Nash,^{76,d} N. R. Nation,²¹ T. Nattermann,²⁰ T. Naumann,⁴¹ G. Navarro,¹⁶¹ H. A. Neal,⁸⁶ E. Nebot,⁷⁹ P. Yu. Nechaeva,⁹³ A. Negri,^{118a,118b} G. Negri,²⁹ S. Nektarijevic,⁴⁸ A. Nelson,¹⁶² S. Nelson,¹⁴² T. K. Nelson,¹⁴² S. Nemecek,¹²⁴ P. Nemethy,¹⁰⁷ A. A. Nepomuceno,^{23a} M. Nessi,^{29,y} M. S. Neubauer,¹⁶⁴ A. Neusiedl,⁸⁰ R. M. Neves,¹⁰⁷ P. Nevski,²⁴ P. R. Newman,¹⁷ V. Nguyen Thi Hong,¹³⁵ R. B. Nickerson,¹¹⁷ R. Nicolaidou,¹³⁵ L. Nicolas,¹³⁸ B. Nicquevert,²⁹ F. Niedercorn,¹¹⁴ J. Nielsen,¹³⁶ T. Niinikoski,²⁹ N. Nikiforou,³⁴

- A. Nikiforov,¹⁵ V. Nikolaenko,¹²⁷ K. Nikolaev,⁶⁴ I. Nikolic-Audit,⁷⁷ K. Nikolics,⁴⁸ K. Nikolopoulos,²⁴ H. Nilsen,⁴⁷ P. Nilsson,⁷ Y. Ninomiya,¹⁵⁴ A. Nisati,^{131a} T. Nishiyama,⁶⁶ R. Nisius,⁹⁸ L. Nodulman,⁵ M. Nomachi,¹¹⁵ I. Nomidis,¹⁵³ M. Nordberg,²⁹ B. Nordkvist,^{145a,145b} P. R. Norton,¹²⁸ J. Novakova,¹²⁵ M. Nozaki,⁶⁵ L. Nozka,¹¹² I. M. Nugent,^{158a} A.-E. Nuncio-Quiroz,²⁰ G. Nunes Hanninger,⁸⁵ T. Nunnemann,⁹⁷ E. Nurse,⁷⁶ T. Nyman,²⁹ B. J. O'Brien,⁴⁵ S. W. O'Neale,^{17a} D. C. O'Neil,¹⁴¹ V. O'Shea,⁵² L. B. Oakes,⁹⁷ F. G. Oakham,^{28e} H. Oberlack,⁹⁸ J. Ocariz,⁷⁷ A. Ochi,⁶⁶ S. Oda,¹⁵⁴ S. Odaka,⁶⁵ J. Odier,⁸² H. Ogren,⁶⁰ A. Oh,⁸¹ S. H. Oh,⁴⁴ C. C. Ohm,^{145a,145b} T. Ohshima,¹⁰⁰ H. Ohshita,¹³⁹ T. Ohsugi,⁵⁸ S. Okada,⁶⁶ H. Okawa,¹⁶² Y. Okumura,¹⁰⁰ T. Okuyama,¹⁵⁴ A. Olariu,^{25a} M. Olcese,^{49a} A. G. Olchevski,⁶⁴ M. Oliveira,^{123a,i} D. Oliveira Damazio,²⁴ E. Oliver Garcia,¹⁶⁶ D. Olivito,¹¹⁹ A. Olszewski,³⁸ J. Olszowska,³⁸ C. Omachi,⁶⁶ A. Onofre,^{123az} P. U. E. Onyisi,³⁰ C. J. Oram,^{158a} M. J. Oreglia,³⁰ Y. Oren,¹⁵² D. Orestano,^{133a,133b} I. Orlov,¹⁰⁶ C. Oropeza Barrera,⁵² R. S. Orr,¹⁵⁷ B. Osculati,^{49a,49b} R. Ospanov,¹¹⁹ C. Osuna,¹¹ G. Otero y Garzon,²⁶ J. P. Ottersbach,¹⁰⁴ M. Ouchrif,^{134d} F. Ould-Saada,¹¹⁶ A. Ouraou,¹³⁵ Q. Ouyang,^{32a} A. Ovcharova,¹⁴ M. Owen,⁸¹ S. Owen,¹³⁸ V. E. Ozcan,^{18a} N. Ozturk,⁷ A. Pacheco Pages,¹¹ C. Padilla Aranda,¹¹ S. Pagan Griso,¹⁴ E. Paganis,¹³⁸ F. Paige,²⁴ P. Pais,⁸³ K. Pajchel,¹¹⁶ G. Palacino,^{158b} C. P. Paleari,⁶ S. Palestini,²⁹ D. Pallin,³³ A. Palma,^{123a} J. D. Palmer,¹⁷ Y. B. Pan,¹⁷¹ E. Panagiotopoulou,⁹ B. Panes,^{31a} N. Panikashvili,⁸⁶ S. Panitkin,²⁴ D. Pantea,^{25a} M. Panuskova,¹²⁴ V. Paolone,¹²² A. Papadelis,^{145a} Th.D. Papadopoulou,⁹ A. Paramonov,⁵ W. Park,^{24,aa} M. A. Parker,²⁷ F. Parodi,^{49a,49b} J. A. Parsons,³⁴ U. Parzefall,⁴⁷ E. Pasqualucci,^{131a} S. Passaggio,^{49a} A. Passeri,^{133a} F. Pastore,^{133a,133b} Fr. Pastore,⁷⁵ G. Pásztor,^{48,bb} S. Patariaia,¹⁷³ N. Patel,¹⁴⁹ J. R. Pater,⁸¹ S. Patricelli,^{101a,101b} T. Pauly,²⁹ M. Pecsý,^{143a} M. I. Pedraza Morales,¹⁷¹ S. V. Peleganchuk,¹⁰⁶ H. Peng,^{32b} R. Pengo,²⁹ A. Penson,³⁴ J. Penwell,⁶⁰ M. Perantoni,^{23a} K. Perez,^{34,cc} T. Perez Cavalcanti,⁴¹ E. Perez Codina,¹¹ M. T. Pérez García-Estañ,¹⁶⁶ V. Perez Reale,³⁴ L. Perini,^{88a,88b} H. Pernegger,²⁹ R. Perrino,^{71a} P. Perrodo,⁴ S. Persebe,^{3a} A. Perus,¹¹⁴ V. D. Peshekhonov,⁶⁴ B. A. Petersen,²⁹ J. Petersen,²⁹ T. C. Petersen,³⁵ E. Petit,⁴ A. Petridis,¹⁵³ C. Petridou,¹⁵³ E. Petrollo,^{131a} F. Petrucci,^{133a,133b} D. Petschull,⁴¹ M. Petteni,¹⁴¹ R. Pezoa,^{31b} A. Phan,⁸⁵ P. W. Phillips,¹²⁸ G. Piacquadio,²⁹ E. Piccaro,⁷⁴ M. Piccinini,^{19a,19b} S. M. Piec,⁴¹ R. Piegai,²⁶ D. T. Pignotti,¹⁰⁸ J. E. Pilcher,³⁰ A. D. Pilkington,⁸¹ J. Pina,^{123a,c} M. Pinamonti,^{163a,163c} A. Pinder,¹¹⁷ J. L. Pinfold,² J. Ping,^{32c} B. Pinto,^{123a,c} O. Pirotte,²⁹ C. Pizio,^{88a,88b} M. Plamondon,¹⁶⁸ M.-A. Pleier,²⁴ A. V. Pleskach,¹²⁷ A. Poblaguev,²⁴ S. Poddar,^{57a} F. Podlyski,³³ L. Poggioli,¹¹⁴ T. Poghosyan,²⁰ M. Pohl,⁴⁸ F. Polci,⁵⁴ G. Polesello,^{118a} A. Policicchio,^{36a,36b} A. Polini,^{19a} J. Poll,⁷⁴ V. Polychronakos,²⁴ D. M. Pomarede,¹³⁵ D. Pomeroy,²² K. Pommès,²⁹ L. Pontecorvo,^{131a} B. G. Pope,⁸⁷ G. A. Popeneciu,^{25a} D. S. Popovic,^{12a} A. Poppleton,²⁹ X. Portell Bueso,²⁹ C. Posch,²¹ G. E. Pospelov,⁹⁸ S. Pospisil,¹²⁶ I. N. Potrap,⁹⁸ C. J. Potter,¹⁴⁸ C. T. Potter,¹¹³ G. Poulard,²⁹ J. Poveda,¹⁷¹ R. Prabhu,⁷⁶ P. Pralavorio,⁸² A. Pranko,¹⁴ S. Prasad,⁵⁶ R. Pravahan,⁷ S. Prell,⁶³ K. Pretzl,¹⁶ L. Pribyl,²⁹ D. Price,⁶⁰ J. Price,⁷² L. E. Price,⁵ M. J. Price,²⁹ D. Prieur,¹²² M. Primavera,^{71a} K. Prokofiev,¹⁰⁷ F. Prokoshin,^{31b} S. Protopopescu,²⁴ J. Proudfoot,⁵ X. Prudent,⁴³ M. Przybycien,³⁷ H. Przysiezniak,⁴ S. Psoroulas,²⁰ E. Ptacek,¹¹³ E. Pueschel,⁸³ J. Purdham,⁸⁶ M. Purohit,^{24,aa} P. Puzo,¹¹⁴ Y. Pylypchenko,⁶² J. Qian,⁸⁶ Z. Qian,⁸² Z. Qin,⁴¹ A. Quadt,⁵³ D. R. Quarrie,¹⁴ W. B. Quayle,¹⁷¹ F. Quinonez,^{31a} M. Raas,¹⁰³ V. Radescu,^{57b} B. Radics,²⁰ T. Rador,^{18a} F. Ragusa,^{88a,88b} G. Rahal,¹⁷⁶ A. M. Rahimi,¹⁰⁸ D. Rahm,²⁴ S. Rajagopalan,²⁴ M. Rammensee,⁴⁷ M. Rammes,¹⁴⁰ A. S. Randle-Conde,³⁹ K. Randrianarivony,²⁸ P. N. Ratoff,⁷⁰ F. Rauscher,⁹⁷ M. Raymond,²⁹ A. L. Read,¹¹⁶ D. M. Rebuzzi,^{118a,118b} A. Redelbach,¹⁷² G. Redlinger,²⁴ R. Reece,¹¹⁹ K. Reeves,⁴⁰ A. Reichold,¹⁰⁴ E. Reinherz-Aronis,¹⁵² A. Reinsch,¹¹³ I. Reisinger,⁴² D. Reljic,^{12a} C. Rembser,²⁹ Z. L. Ren,¹⁵⁰ A. Renaud,¹¹⁴ P. Renkel,³⁹ M. Rescigno,^{131a} S. Resconi,^{88a} B. Resende,¹³⁵ P. Reznicek,⁹⁷ R. Rezvani,¹⁵⁷ A. Richards,⁷⁶ R. Richter,⁹⁸ E. Richter-Was,⁴ M. Ridel,⁷⁷ M. Rijpstra,¹⁰⁴ M. Rijssenbeek,¹⁴⁷ A. Rimoldi,^{118a,118b} L. Rinaldi,^{19a} R. R. Rios,³⁹ I. Riu,¹¹ G. Rivoltella,^{88a,88b} F. Rizatdinova,¹¹¹ E. Rizvi,⁷⁴ S. H. Robertson,^{84,k} A. Robichaud-Veronneau,¹¹⁷ D. Robinson,²⁷ J. E. M. Robinson,⁷⁶ M. Robinson,¹¹³ A. Robson,⁵² J. G. Rocha de Lima,¹⁰⁵ C. Roda,^{121a,121b} D. Roda Dos Santos,²⁹ D. Rodriguez,¹⁶¹ Y. Rodriguez Garcia,¹⁶¹ A. Roe,⁵³ S. Roe,²⁹ O. Røhne,¹¹⁶ V. Rojo,¹ S. Rolli,¹⁶⁰ A. Romaniouk,⁹⁵ M. Romano,^{19a,19b} V. M. Romanov,⁶⁴ G. Romeo,²⁶ L. Roos,⁷⁷ E. Ros,¹⁶⁶ S. Rosati,^{131a} K. Rosbach,⁴⁸ A. Rose,¹⁴⁸ M. Rose,⁷⁵ G. A. Rosenbaum,¹⁵⁷ E. I. Rosenberg,⁶³ P. L. Rosendahl,¹³ O. Rosenthal,¹⁴⁰ L. Rossetlet,⁴⁸ V. Rossetti,¹¹ E. Rossi,^{131a,131b} L. P. Rossi,^{49a} M. Rotaru,^{25a} I. Roth,¹⁷⁰ J. Rothberg,¹³⁷ D. Rousseau,¹¹⁴ C. R. Royon,¹³⁵ A. Rozanov,⁸² Y. Rozen,¹⁵¹ X. Ruan,^{114,dd} I. Rubinskiy,⁴¹ B. Ruckert,⁹⁷ N. Ruckstuhl,¹⁰⁴ V. I. Rud,⁹⁶ C. Rudolph,⁴³ G. Rudolph,⁶¹ F. Rühr,⁶ F. Ruggieri,^{133a,133b} A. Ruiz-Martinez,⁶³ V. Rumiantsev,^{90,a} L. Rumyantsev,⁶⁴ K. Runge,⁴⁷ Z. Rurikova,⁴⁷ N. A. Rusakovich,⁶⁴ D. R. Rust,⁶⁰ J. P. Rutherford,⁶ C. Ruwiedel,¹⁴ P. Ruzicka,¹²⁴ Y. F. Ryabov,¹²⁰ V. Ryadovikov,¹²⁷ P. Ryan,⁸⁷ M. Rybar,¹²⁵ G. Rybkin,¹¹⁴ N. C. Ryder,¹¹⁷ S. Rzaeva,¹⁰ A. F. Saavedra,¹⁴⁹ I. Sadeh,¹⁵² H.F.-W. Sadrozinski,¹³⁶

- R. Sadykov,⁶⁴ F. Safai Tehrani,^{131a} H. Sakamoto,¹⁵⁴ G. Salamanna,⁷⁴ A. Salamon,^{132a} M. Saleem,¹¹⁰ D. Salihagic,⁹⁸
 A. Salnikov,¹⁴² J. Salt,¹⁶⁶ B. M. Salvachua Ferrando,⁵ D. Salvatore,^{36a,36b} F. Salvatore,¹⁴⁸ A. Salvucci,¹⁰³
 A. Salzburger,²⁹ D. Sampsonidis,¹⁵³ B. H. Samset,¹¹⁶ A. Sanchez,^{101a,101b} H. Sandaker,¹³ H. G. Sander,⁸⁰
 M. P. Sanders,⁹⁷ M. Sandhoff,¹⁷³ T. Sandoval,²⁷ C. Sandoval,¹⁶¹ R. Sandstroem,⁹⁸ S. Sandvoss,¹⁷³ D. P. C. Sankey,¹²⁸
 A. Sansoni,⁴⁶ C. Santamarina Rios,⁸⁴ C. Santoni,³³ R. Santonico,^{132a,132b} H. Santos,^{123a} J. G. Saraiva,^{123a}
 T. Sarangi,¹⁷¹ E. Sarkisyan-Grinbaum,⁷ F. Sarri,^{121a,121b} G. Sartisohn,¹⁷³ O. Sasaki,⁶⁵ N. Sasao,⁶⁷
 I. Satsounkevitch,⁸⁹ G. Sauvage,⁴ E. Sauvan,⁴ J. B. Sauvan,¹¹⁴ P. Savard,^{157,e} V. Savinov,¹²² D. O. Savu,²⁹
 L. Sawyer,^{24,m} D. H. Saxon,⁵² L. P. Says,³³ C. Sbarra,^{19a} A. Sbrizzi,^{19a,19b} O. Scallon,⁹² D. A. Scannicchio,¹⁶²
 M. Scarcella,¹⁴⁹ J. Schaarschmidt,¹¹⁴ P. Schacht,⁹⁸ U. Schäfer,⁸⁰ S. Schaepe,²⁰ S. Schaezel,^{57b} A. C. Schaffer,¹¹⁴
 D. Schaile,⁹⁷ R. D. Schamberger,¹⁴⁷ A. G. Schamov,¹⁰⁶ V. Scharf,^{57a} V. A. Schegelsky,¹²⁰ D. Scheirich,⁸⁶
 M. Schernau,¹⁶² M. I. Scherzer,³⁴ C. Schiavi,^{49a,49b} J. Schieck,⁹⁷ M. Schioppa,^{36a,36b} S. Schlenker,²⁹ J. L. Schlereth,⁵
 E. Schmidt,⁴⁷ K. Schmieden,²⁰ C. Schmitt,⁸⁰ S. Schmitt,^{57b} M. Schmitz,²⁰ A. Schöning,^{57b} M. Schott,²⁹
 D. Schouten,^{158a} J. Schovancova,¹²⁴ M. Schram,⁸⁴ C. Schroeder,⁸⁰ N. Schroer,^{57c} S. Schuh,²⁹ G. Schuler,²⁹
 J. Schultes,¹⁷³ H.-C. Schultz-Coulon,^{57a} H. Schulz,¹⁵ J. W. Schumacher,²⁰ M. Schumacher,⁴⁷ B. A. Schumm,¹³⁶
 Ph. Schune,¹³⁵ C. Schwanenberger,⁸¹ A. Schwartzman,¹⁴² Ph. Schwemling,⁷⁷ R. Schwienhorst,⁸⁷ R. Schwierz,⁴³
 J. Schwindling,¹³⁵ T. Schwindt,²⁰ M. Schwoerer,⁴ W. G. Scott,¹²⁸ J. Searcy,¹¹³ G. Sedov,⁴¹ E. Sedykh,¹²⁰ E. Segura,¹¹
 S. C. Seidel,¹⁰² A. Seiden,¹³⁶ F. Seifert,⁴³ J. M. Seixas,^{23a} G. Sekhniaidze,^{101a} D. M. Seliverstov,¹²⁰ B. Sellden,^{145a}
 G. Sellers,⁷² M. Seman,^{143b} N. Semprini-Cesari,^{19a,19b} C. Serfon,⁹⁷ L. Serin,¹¹⁴ R. Seuster,⁹⁸ H. Severini,¹¹⁰
 M. E. Sevir,⁸⁵ A. Sfyrta,²⁹ E. Shabalina,⁵³ M. Shamim,¹¹³ L. Y. Shan,^{32a} J. T. Shank,²¹ Q. T. Shao,⁸⁵ M. Shapiro,¹⁴
 P. B. Shatalov,⁹⁴ L. Shaver,⁶ K. Shaw,^{163a,163c} D. Sherman,¹⁷⁴ P. Sherwood,⁷⁶ A. Shibata,¹⁰⁷ H. Shichi,¹⁰⁰
 S. Shimizu,²⁹ M. Shimojima,⁹⁹ T. Shin,⁵⁵ M. Shiyakova,⁶⁴ A. Shmeleva,⁹³ M. J. Shochet,³⁰ D. Short,¹¹⁷
 S. Shrestha,⁶³ M. A. Shupe,⁶ P. Sicho,¹²⁴ A. Sidoti,^{131a} F. Siegert,⁴⁷ Dj. Sijacki,^{12a} O. Silbert,¹⁷⁰ J. Silva,^{123a,c}
 Y. Silver,¹⁵² D. Silverstein,¹⁴² S. B. Silverstein,^{145a} V. Simak,¹²⁶ O. Simard,¹³⁵ Lj. Simic,^{12a} S. Simion,¹¹⁴
 B. Simmons,⁷⁶ M. Simonyan,³⁵ P. Sinervo,¹⁵⁷ N. B. Sinev,¹¹³ V. Sipica,¹⁴⁰ G. Siragusa,¹⁷² A. Sircar,²⁴
 A. N. Sisakyan,⁶⁴ S. Yu. Sivoklov,⁹⁶ J. Sjölin,^{145a,145b} T. B. Sjørnsen,¹³ L. A. Skinnari,¹⁴ H. P. Skottowe,⁵⁶
 K. Skovpen,¹⁰⁶ P. Skubic,¹¹⁰ N. Skvorodnev,²² M. Slater,¹⁷ T. Slavicek,¹²⁶ K. Sliwa,¹⁶⁰ J. Sloper,²⁹ V. Smakhtin,¹⁷⁰
 S. Yu. Smirnov,⁹⁵ L. N. Smirnova,⁹⁶ O. Smirnova,⁷⁸ B. C. Smith,⁵⁶ D. Smith,¹⁴² K. M. Smith,⁵² M. Smizanska,⁷⁰
 K. Smolek,¹²⁶ A. A. Snesarev,⁹³ S. W. Snow,⁸¹ J. Snow,¹¹⁰ J. Snuverink,¹⁰⁴ S. Snyder,²⁴ M. Soares,^{123a} R. Sobie,^{168,k}
 J. Sodomka,¹²⁶ A. Soffer,¹⁵² C. A. Solans,¹⁶⁶ M. Solar,¹²⁶ J. Solc,¹²⁶ E. Soldatov,⁹⁵ U. Soldevila,¹⁶⁶
 E. Solfaroli Camillocci,^{131a,131b} A. A. Solodkov,¹²⁷ O. V. Solovyanov,¹²⁷ N. Soni,² V. Sopko,¹²⁶ B. Sopko,¹²⁶
 M. Sosebee,⁷ R. Soualah,^{163a,163c} A. Soukharev,¹⁰⁶ S. Spagnolo,^{71a,71b} F. Spanò,⁷⁵ R. Spighi,^{19a} G. Spigo,²⁹
 F. Spila,^{131a,131b} R. Spiwoks,²⁹ M. Spousta,¹²⁵ T. Spreitzer,¹⁵⁷ B. Spurlock,⁷ R. D. St. Denis,⁵² T. Stahl,¹⁴⁰
 J. Stahlman,¹¹⁹ R. Stamen,^{57a} E. Stanecka,³⁸ R. W. Stanek,⁵ C. Stanescu,^{133a} S. Stapnes,¹¹⁶ E. A. Starchenko,¹²⁷
 J. Stark,⁵⁴ P. Staroba,¹²⁴ P. Starovoitov,⁹⁰ A. Staude,⁹⁷ P. Stavina,^{143a} G. Stavropoulos,¹⁴ G. Steele,⁵² P. Steinbach,⁴³
 P. Steinberg,²⁴ I. Stekl,¹²⁶ B. Stelzer,¹⁴¹ H. J. Stelzer,⁸⁷ O. Stelzer-Chilton,^{158a} H. Stenzel,⁵¹ S. Stern,⁹⁸
 K. Stevenson,⁷⁴ G. A. Stewart,²⁹ J. A. Stillings,²⁰ M. C. Stockton,²⁹ K. Stoerig,⁴⁷ G. Stoicea,^{25a} S. Stonjek,⁹⁸
 P. Strachota,¹²⁵ A. R. Stradling,⁷ A. Straessner,⁴³ J. Strandberg,¹⁴⁶ S. Strandberg,^{145a,145b} A. Strandlie,¹¹⁶
 M. Strang,¹⁰⁸ E. Strauss,¹⁴² M. Strauss,¹¹⁰ P. Strizenece,^{143b} R. Ströhmer,¹⁷² D. M. Strom,¹¹³ J. A. Strong,^{75,a}
 R. Stroynowski,³⁹ J. Strube,¹²⁸ B. Stugu,¹³ I. Stumer,^{24,a} J. Stupak,¹⁴⁷ P. Sturm,¹⁷³ N. A. Styles,⁴¹ D. A. Soh,^{150,v}
 D. Su,¹⁴² H. S. Subramania,² A. Succurro,¹¹ Y. Sugaya,¹¹⁵ T. Sugimoto,¹⁰⁰ C. Suhr,¹⁰⁵ K. Suita,⁶⁶ M. Suk,¹²⁵
 V. V. Sulin,⁹³ S. Sultansoy,^{3d} T. Sumida,⁶⁷ X. Sun,⁵⁴ J. E. Sundermann,⁴⁷ K. Suruliz,¹³⁸ S. Sushkov,¹¹
 G. Susinno,^{36a,36b} M. R. Sutton,¹⁴⁸ Y. Suzuki,⁶⁵ Y. Suzuki,⁶⁶ M. Svatos,¹²⁴ Yu. M. Sviridov,¹²⁷ S. Swedish,¹⁶⁷
 I. Sykora,^{143a} T. Sykora,¹²⁵ B. Szeless,²⁹ J. Sánchez,¹⁶⁶ D. Ta,¹⁰⁴ K. Tackmann,⁴¹ A. Taffard,¹⁶² R. Tafirout,^{158a}
 N. Taiblum,¹⁵² Y. Takahashi,¹⁰⁰ H. Takai,²⁴ R. Takashima,⁶⁸ H. Takeda,⁶⁶ T. Takeshita,¹³⁹ M. Talby,⁸²
 A. Talyshev,^{106,g} M. C. Tamsett,²⁴ J. Tanaka,¹⁵⁴ R. Tanaka,¹¹⁴ S. Tanaka,¹³⁰ S. Tanaka,⁶⁵ Y. Tanaka,⁹⁹ K. Tani,⁶⁶
 N. Tannoury,⁸² G. P. Tappern,²⁹ S. Tapprogge,⁸⁰ D. Tardif,¹⁵⁷ S. Tarem,¹⁵¹ F. Tarrade,²⁸ G. F. Tartarelli,^{88a} P. Tas,¹²⁵
 M. Tasevsky,¹²⁴ E. Tassi,^{36a,36b} M. Tatarkhanov,¹⁴ Y. Tayalati,^{134d} C. Taylor,⁷⁶ F. E. Taylor,⁹¹ G. N. Taylor,⁸⁵
 W. Taylor,^{158b} M. Teinturier,¹¹⁴ M. Teixeira Dias Castanheira,⁷⁴ P. Teixeira-Dias,⁷⁵ K. K. Temming,⁴⁷ H. Ten Kate,²⁹
 P. K. Teng,¹⁵⁰ S. Terada,⁶⁵ K. Terashi,¹⁵⁴ J. Terron,⁷⁹ M. Testa,⁴⁶ R. J. Teuscher,^{157,k} J. Thadome,¹⁷³ J. Therhaag,²⁰
 T. Theveneaux-Pelzer,⁷⁷ M. Thioye,¹⁷⁴ S. Thoma,⁴⁷ J. P. Thomas,¹⁷ E. N. Thompson,³⁴ P. D. Thompson,¹⁷
 P. D. Thompson,¹⁵⁷ A. S. Thompson,⁵² E. Thomson,¹¹⁹ M. Thomson,²⁷ R. P. Thun,⁸⁶ F. Tian,³⁴ M. J. Tibbetts,¹⁴

- T. Tic,¹²⁴ V. O. Tikhomirov,⁹³ Y. A. Tikhonov,^{106,g} S. Timoshenko,⁹⁵ P. Tipton,¹⁷⁴ F. J. Tique Aires Viegas,²⁹ S. Tisserant,⁸² B. Toczek,³⁷ T. Todorov,⁴ S. Todorova-Nova,¹⁶⁰ B. Toggerson,¹⁶² J. Tojo,⁶⁵ S. Tokár,^{143a} K. Tokunaga,⁶⁶ K. Tokushuku,⁶⁵ K. Tollefson,⁸⁷ M. Tomoto,¹⁰⁰ L. Tompkins,³⁰ K. Toms,¹⁰² G. Tong,^{32a} A. Tonoyan,¹³ C. Topfel,¹⁶ N. D. Topilin,⁶⁴ I. Torchiani,²⁹ E. Torrence,¹¹³ H. Torres,⁷⁷ E. Torró Pastor,¹⁶⁶ J. Toth,^{82,bb} F. Touchard,⁸² D. R. Tovey,¹³⁸ T. Trefzger,¹⁷² L. Tremblet,²⁹ A. Tricoli,²⁹ I. M. Trigger,^{158a} S. Trincaz-Duvoid,⁷⁷ T. N. Trinh,⁷⁷ M. F. Tripana,⁶⁹ W. Trischuk,¹⁵⁷ A. Trivedi,^{24,aa} B. Trocmé,⁵⁴ C. Troncon,^{88a} M. Trotter-McDonald,¹⁴¹ M. Trzebinski,³⁸ A. Trzuppek,³⁸ C. Tsarouchas,²⁹ J. C.-L. Tseng,¹¹⁷ M. Tsiakiris,¹⁰⁴ P. V. Tsiareshka,⁸⁹ D. Tsionou,^{4,ee} G. Tsipolitis,⁹ V. Tsiskaridze,⁴⁷ E. G. Tskhadadze,^{50a} I. I. Tsukerman,⁹⁴ V. Tsulaia,¹⁴ J.-W. Tsung,²⁰ S. Tsuno,⁶⁵ D. Tsybychev,¹⁴⁷ A. Tua,¹³⁸ A. Tudorache,^{25a} V. Tudorache,^{25a} J. M. Tuggle,³⁰ M. Turala,³⁸ D. Turecek,¹²⁶ I. Turk Cakir,^{3e} E. Turlay,¹⁰⁴ R. Turra,^{88a,88b} P. M. Tuts,³⁴ A. Tykhonov,⁷³ M. Tylmad,^{145a,145b} M. Tyndel,¹²⁸ G. Tzanakos,⁸ K. Uchida,²⁰ I. Ueda,¹⁵⁴ R. Ueno,²⁸ M. Ugland,¹³ M. Uhlenbrock,²⁰ M. Uhrmacher,⁵³ F. Ukegawa,¹⁵⁹ G. Unal,²⁹ D. G. Underwood,⁵ A. Undrus,²⁴ G. Unel,¹⁶² Y. Unno,⁶⁵ D. Urbaniec,³⁴ G. Usai,⁷ M. Uslenghi,^{118a,118b} L. Vacavant,⁸² V. Vacek,¹²⁶ B. Vachon,⁸⁴ S. Vahsen,¹⁴ J. Valenta,¹²⁴ P. Valente,^{131a} S. Valentinetti,^{19a,19b} S. Valkar,¹²⁵ E. Valladolid Gallego,¹⁶⁶ S. Vallecorsa,¹⁵¹ J. A. Valls Ferrer,¹⁶⁶ H. van der Graaf,¹⁰⁴ E. van der Kraaij,¹⁰⁴ R. Van Der Leeuw,¹⁰⁴ E. van der Poel,¹⁰⁴ D. van der Ster,²⁹ N. van Eldik,⁸³ P. van Gemmeren,⁵ Z. van Kesteren,¹⁰⁴ I. van Vulpen,¹⁰⁴ M. Vanadia,⁹⁸ W. Vandelli,²⁹ G. Vandoni,²⁹ A. Vaniachine,⁵ P. Vankov,⁴¹ F. Vannucci,⁷⁷ F. Varela Rodriguez,²⁹ R. Vari,^{131a} E. W. Varnes,⁶ D. Varouchas,¹⁴ A. Vartapetian,⁷ K. E. Varvell,¹⁴⁹ V. I. Vassilakopoulos,⁵⁵ F. Vazeille,³³ G. Vegni,^{88a,88b} J. J. Veillet,¹¹⁴ C. Vellidis,⁸ F. Veloso,^{123a} R. Veness,²⁹ S. Veneziano,^{131a} A. Ventura,^{71a,71b} D. Ventura,¹³⁷ M. Venturi,⁴⁷ N. Venturi,¹⁵⁷ V. Vercesi,^{118a} M. Verducci,¹³⁷ W. Verkerke,¹⁰⁴ J. C. Vermeulen,¹⁰⁴ A. Vest,⁴³ M. C. Vetterli,^{141,e} I. Vichou,¹⁶⁴ T. Vickey,^{144b,ff} O. E. Vickey Boeriu,^{144b} G. H. A. Viehhauser,¹¹⁷ S. Viel,¹⁶⁷ M. Villa,^{19a,19b} M. Villaplana Perez,¹⁶⁶ E. Vilucchi,⁴⁶ M. G. Vincker,²⁸ E. Vinek,²⁹ V. B. Vinogradov,⁶⁴ M. Virchaux,^{135,a} J. Virzi,¹⁴ O. Vitells,¹⁷⁰ M. Viti,⁴¹ I. Vivarelli,⁴⁷ F. Vives Vaque,² S. Vlachos,⁹ D. Vladoiu,⁹⁷ M. Vlasak,¹²⁶ N. Vlasov,²⁰ A. Vogel,²⁰ P. Vokac,¹²⁶ G. Volpi,⁴⁶ M. Volpi,⁸⁵ G. Volpini,^{88a} H. von der Schmitt,⁹⁸ J. von Loeben,⁹⁸ H. von Radziewski,⁴⁷ E. von Toerne,²⁰ V. Vorobel,¹²⁵ A. P. Vorobiev,¹²⁷ V. Vorwerk,¹¹ M. Vos,¹⁶⁶ R. Voss,²⁹ T. T. Voss,¹⁷³ J. H. Vossebeld,⁷² N. Vranjes,^{12a} M. Vranjes Milosavljevic,¹⁰⁴ V. Vrba,¹²⁴ M. Vreeswijk,¹⁰⁴ T. Vu Anh,⁸⁰ R. Vuillemet,²⁹ I. Vukotic,¹¹⁴ W. Wagner,¹⁷³ P. Wagner,¹¹⁹ H. Wahlen,¹⁷³ J. Wakabayashi,¹⁰⁰ J. Walbersloh,⁴² S. Walch,⁸⁶ J. Walder,⁷⁰ R. Walker,⁹⁷ W. Walkowiak,¹⁴⁰ R. Wall,¹⁷⁴ P. Waller,⁷² C. Wang,⁴⁴ H. Wang,¹⁷¹ H. Wang,^{32b,gg} J. Wang,¹⁵⁰ J. Wang,⁵⁴ J. C. Wang,¹³⁷ R. Wang,¹⁰² S. M. Wang,¹⁵⁰ A. Warburton,⁸⁴ C. P. Ward,²⁷ M. Warsinsky,⁴⁷ P. M. Watkins,¹⁷ A. T. Watson,¹⁷ I. J. Watson,¹⁴⁹ M. F. Watson,¹⁷ G. Watts,¹³⁷ S. Watts,⁸¹ A. T. Waugh,¹⁴⁹ B. M. Waugh,⁷⁶ M. Weber,¹²⁸ M. S. Weber,¹⁶ P. Weber,⁵³ A. R. Weidberg,¹¹⁷ P. Weigell,⁹⁸ J. Weingarten,⁵³ C. Weiser,⁴⁷ H. Wellenstein,²² P. S. Wells,²⁹ M. Wen,⁴⁶ T. Wenaus,²⁴ S. Wendler,¹²² Z. Weng,^{150,v} T. Wengler,²⁹ S. Wenig,²⁹ N. Wermes,²⁰ M. Werner,⁴⁷ P. Werner,²⁹ M. Werth,¹⁶² M. Wessels,^{57a} C. Weydert,⁵⁴ K. Whalen,²⁸ S. J. Wheeler-Ellis,¹⁶² S. P. Whitaker,²¹ A. White,⁷ M. J. White,⁸⁵ S. R. Whitehead,¹¹⁷ D. Whiteson,¹⁶² D. Whittington,⁶⁰ F. Wicek,¹¹⁴ D. Wicke,¹⁷³ F. J. Wickens,¹²⁸ W. Wiedenmann,¹⁷¹ M. Wielers,¹²⁸ P. Wienemann,²⁰ C. Wiglesworth,⁷⁴ L. A. M. Wiik-Fuchs,⁴⁷ P. A. Wijeratne,⁷⁶ A. Wildauer,¹⁶⁶ M. A. Wildt,^{41,r} I. Wilhelm,¹²⁵ H. G. Wilkens,²⁹ J. Z. Will,⁹⁷ E. Williams,³⁴ H. H. Williams,¹¹⁹ W. Willis,³⁴ S. Willocq,⁸³ J. A. Wilson,¹⁷ M. G. Wilson,¹⁴² A. Wilson,⁸⁶ I. Wingerter-Seez,⁴ S. Winkelmann,⁴⁷ F. Winklmeier,²⁹ M. Wittgen,¹⁴² M. W. Wolter,³⁸ H. Wolters,^{123a,i} W. C. Wong,⁴⁰ G. Wooden,⁸⁶ B. K. Wosiek,³⁸ J. Wotschack,²⁹ M. J. Woudstra,⁸³ K. W. Wozniak,³⁸ K. Wraight,⁵² C. Wright,⁵² M. Wright,⁵² B. Wrona,⁷² S. L. Wu,¹⁷¹ X. Wu,⁴⁸ Y. Wu,^{32b,hh} E. Wulf,³⁴ R. Wunstorf,⁴² B. M. Wynne,⁴⁵ S. Xella,³⁵ M. Xiao,¹³⁵ S. Xie,⁴⁷ Y. Xie,^{32a} C. Xu,^{32b,x} D. Xu,¹³⁸ G. Xu,^{32a} B. Yabsley,¹⁴⁹ S. Yacoob,^{144b} M. Yamada,⁶⁵ H. Yamaguchi,¹⁵⁴ A. Yamamoto,⁶⁵ K. Yamamoto,⁶³ S. Yamamoto,¹⁵⁴ T. Yamamura,¹⁵⁴ T. Yamanaka,¹⁵⁴ J. Yamaoka,⁴⁴ T. Yamazaki,¹⁵⁴ Y. Yamazaki,⁶⁶ Z. Yan,²¹ H. Yang,⁸⁶ U. K. Yang,⁸¹ Y. Yang,⁶⁰ Y. Yang,^{32a} Z. Yang,^{145a,145b} S. Yanush,⁹⁰ Y. Yao,¹⁴ Y. Yasu,⁶⁵ G. V. Ybeles Smit,¹²⁹ J. Ye,³⁹ S. Ye,²⁴ M. Yilmaz,^{3c} R. Yoosofmiya,¹²² K. Yorita,¹⁶⁹ R. Yoshida,⁵ C. Young,¹⁴² S. Youssef,²¹ D. Yu,²⁴ J. Yu,⁷ J. Yu,¹¹¹ L. Yuan,^{32a,ii} A. Yurkewicz,¹⁰⁵ B. Zabinski,³⁸ V. G. Zaets,¹²⁷ R. Zaidan,⁶² A. M. Zaitsev,¹²⁷ Z. Zajacova,²⁹ L. Zanello,^{131a,131b} P. Zarzhitsky,³⁹ A. Zaytsev,¹⁰⁶ C. Zeitnitz,¹⁷³ M. Zeller,¹⁷⁴ M. Zeman,¹²⁴ A. Zemla,³⁸ C. Zendler,²⁰ O. Zenin,¹²⁷ T. Ženiš,^{143a} Z. Zinonos,^{121a,121b} S. Zenz,¹⁴ D. Zerwas,¹¹⁴ G. Zevi della Porta,⁵⁶ Z. Zhan,^{32d} D. Zhang,^{32b,gg} H. Zhang,⁸⁷ J. Zhang,⁵ X. Zhang,^{32d} Z. Zhang,¹¹⁴ L. Zhao,¹⁰⁷ T. Zhao,¹³⁷ Z. Zhao,^{32b} A. Zhemchugov,⁶⁴ S. Zheng,^{32a} J. Zhong,¹¹⁷ B. Zhou,⁸⁶ N. Zhou,¹⁶² Y. Zhou,¹⁵⁰ C. G. Zhu,^{32d} H. Zhu,⁴¹ J. Zhu,⁸⁶ Y. Zhu,^{32b} X. Zhuang,⁹⁷ V. Zhuravlov,⁹⁸ D. Zieminska,⁶⁰ R. Zimmermann,²⁰ S. Zimmermann,²⁰ S. Zimmermann,⁴⁷

M. Ziolkowski,¹⁴⁰ R. Zitoun,⁴ L. Živković,³⁴ V. V. Zmouchko,^{127,a} G. Zobernig,¹⁷¹ A. Zoccoli,^{19a,19b}
Y. Zolnierowski,⁴ A. Zsenei,²⁹ M. zur Nedden,¹⁵ V. Zutshi,¹⁰⁵ and L. Zwalinski²⁹

(ATLAS Collaboration)

- ¹University at Albany, Albany, New York, USA
²Department of Physics, University of Alberta, Edmonton AB, Canada
^{3a}Department of Physics, Ankara University, Ankara, Turkey
^{3b}Department of Physics, Dumlupinar University, Kutahya, Turkey
^{3c}Department of Physics, Gazi University, Ankara, Turkey
^{3d}Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey
^{3e}Turkish Atomic Energy Authority, Ankara, Turkey
⁴LAPP, CNRS/IN2P3 and Université de Savoie, Annecy-le-Vieux, France
⁵High Energy Physics Division, Argonne National Laboratory, Argonne, Illinois, USA
⁶Department of Physics, University of Arizona, Tucson, Arizona, USA
⁷Department of Physics, The University of Texas at Arlington, Arlington, Texas, USA
⁸Physics Department, University of Athens, Athens, Greece
⁹Physics Department, National Technical University of Athens, Zografou, Greece
¹⁰Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
¹¹Institut de Física d'Altes Energies and Departament de Física de la Universitat Autònoma de Barcelona and ICREA, Barcelona, Spain
^{12a}Institute of Physics, University of Belgrade, Belgrade, Serbia
^{12b}Vinca Institute of Nuclear Sciences, University of Belgrade, Belgrade, Serbia
¹³Department for Physics and Technology, University of Bergen, Bergen, Norway
¹⁴Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, California, USA
¹⁵Department of Physics, Humboldt University, Berlin, Germany
¹⁶Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
¹⁷School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
^{18a}Department of Physics, Bogazici University, Istanbul, Turkey
^{18b}Division of Physics, Dogus University, Istanbul, Turkey
^{18c}Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey
^{18d}Department of Physics, Istanbul Technical University, Istanbul, Turkey
^{19a}INFN Sezione di Bologna, Italy
^{19b}Dipartimento di Fisica, Università di Bologna, Bologna, Italy
²⁰Physikalisches Institut, University of Bonn, Bonn, Germany
²¹Department of Physics, Boston University, Boston, Massachusetts, USA
²²Department of Physics, Brandeis University, Waltham, Massachusetts, USA
^{23a}Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil
^{23b}Federal University of Juiz de Fora (UFJF), Juiz de Fora, Brazil
^{23c}Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei, Brazil
^{23d}Instituto de Física, Universidade de Sao Paulo, Sao Paulo, Brazil
²⁴Physics Department, Brookhaven National Laboratory, Upton, New York, USA
^{25a}National Institute of Physics and Nuclear Engineering, Bucharest, Romania
^{25b}University Politehnica Bucharest, Bucharest, Romania
^{25c}West University in Timisoara, Timisoara, Romania
²⁶Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
²⁷Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
²⁸Department of Physics, Carleton University, Ottawa ON, Canada
²⁹CERN, Geneva, Switzerland
³⁰Enrico Fermi Institute, University of Chicago, Chicago, Illinois, USA
^{31a}Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile
^{31b}Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
^{32a}Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China
^{32b}Department of Modern Physics, University of Science and Technology of China, Anhui, China
^{32c}Department of Physics, Nanjing University, Jiangsu, China
^{32d}School of Physics, Shandong University, Shandong, China
³³Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Aubiere Cedex, France
³⁴Nevis Laboratory, Columbia University, Irvington, New York, USA
³⁵Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark

- ^{36a}INFN Gruppo Collegato di Cosenza, Italy
- ^{36b}Dipartimento di Fisica, Università della Calabria, Arcavata di Rende, Italy
- ³⁷AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland
- ³⁸The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland
- ³⁹Physics Department, Southern Methodist University, Dallas, Texas, USA
- ⁴⁰Physics Department, University of Texas at Dallas, Richardson, Texas, USA
- ⁴¹DESY, Hamburg and Zeuthen, Germany
- ⁴²Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
- ⁴³Institut für Kern- und Teilchenphysik, Technical University Dresden, Dresden, Germany
- ⁴⁴Department of Physics, Duke University, Durham, North Carolina, USA
- ⁴⁵SUPA - School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
- ⁴⁶INFN Laboratori Nazionali di Frascati, Frascati, Italy
- ⁴⁷Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg i.Br., Germany
- ⁴⁸Section de Physique, Université de Genève, Geneva, Switzerland
- ^{49a}INFN Sezione di Genova, Italy
- ^{49b}Dipartimento di Fisica, Università di Genova, Genova, Italy
- ^{50a}E. Andronikashvili Institute of Physics, Tbilisi State University, Tbilisi, Georgia
- ^{50b}High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
- ⁵¹II Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
- ⁵²SUPA - School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
- ⁵³II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany
- ⁵⁴Laboratoire de Physique Subatomique et de Cosmologie, Université Joseph Fourier and CNRS/IN2P3 and Institut National Polytechnique de Grenoble, Grenoble, France
- ⁵⁵Department of Physics, Hampton University, Hampton, Virginia, USA
- ⁵⁶Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, Massachusetts, USA
- ^{57a}Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
- ^{57b}Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
- ^{57c}ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany
- ⁵⁸Faculty of Science, Hiroshima University, Hiroshima, Japan
- ⁵⁹Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
- ⁶⁰Department of Physics, Indiana University, Bloomington, Indiana, USA
- ⁶¹Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
- ⁶²University of Iowa, Iowa City, Iowa, USA
- ⁶³Department of Physics and Astronomy, Iowa State University, Ames, Iowa, USA
- ⁶⁴Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia
- ⁶⁵KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
- ⁶⁶Graduate School of Science, Kobe University, Kobe, Japan
- ⁶⁷Faculty of Science, Kyoto University, Kyoto, Japan
- ⁶⁸Kyoto University of Education, Kyoto, Japan
- ⁶⁹Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
- ⁷⁰Physics Department, Lancaster University, Lancaster, United Kingdom
- ^{71a}INFN Sezione di Lecce, Italy
- ^{71b}Dipartimento di Fisica, Università del Salento, Lecce, Italy
- ⁷²Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom
- ⁷³Department of Physics, Jožef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia
- ⁷⁴School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom
- ⁷⁵Department of Physics, Royal Holloway University of London, Surrey, United Kingdom
- ⁷⁶Department of Physics and Astronomy, University College London, London, United Kingdom
- ⁷⁷Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France
- ⁷⁸Fysiska institutionen, Lunds Universitet, Lund, Sweden
- ⁷⁹Departamento de Física Teórica C-15, Universidad Autónoma de Madrid, Madrid, Spain
- ⁸⁰Institut für Physik, Universität Mainz, Mainz, Germany
- ⁸¹School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
- ⁸²CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
- ⁸³Department of Physics, University of Massachusetts, Amherst, Massachusetts, USA
- ⁸⁴Department of Physics, McGill University, Montreal QC, Canada
- ⁸⁵School of Physics, University of Melbourne, Victoria, Australia
- ⁸⁶Department of Physics, The University of Michigan, Ann Arbor, Michigan, USA
- ⁸⁷Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan, USA
- ^{88a}INFN Sezione di Milano, Italy
- ^{88b}Dipartimento di Fisica, Università di Milano, Milano, Italy

- ⁸⁹*B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus*
- ⁹⁰*National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Republic of Belarus*
- ⁹¹*Department of Physics, Massachusetts Institute of Technology, Cambridge, Massachusetts, USA*
- ⁹²*Group of Particle Physics, University of Montreal, Montreal QC, Canada*
- ⁹³*P.N. Lebedev Institute of Physics, Academy of Sciences, Moscow, Russia*
- ⁹⁴*Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia*
- ⁹⁵*Moscow Engineering and Physics Institute (MEPhI), Moscow, Russia*
- ⁹⁶*Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia*
- ⁹⁷*Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany*
- ⁹⁸*Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany*
- ⁹⁹*Nagasaki Institute of Applied Science, Nagasaki, Japan*
- ¹⁰⁰*Graduate School of Science, Nagoya University, Nagoya, Japan*
- ^{101a}*INFN Sezione di Napoli, Italy*
- ^{101b}*Dipartimento di Scienze Fisiche, Università di Napoli, Napoli, Italy*
- ¹⁰²*Department of Physics and Astronomy, University of New Mexico, Albuquerque, New Mexico, USA*
- ¹⁰³*Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands*
- ¹⁰⁴*Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands*
- ¹⁰⁵*Department of Physics, Northern Illinois University, DeKalb, Illinois, USA*
- ¹⁰⁶*Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia*
- ¹⁰⁷*Department of Physics, New York University, New York, New York, USA*
- ¹⁰⁸*Ohio State University, Columbus, Ohio, USA*
- ¹⁰⁹*Faculty of Science, Okayama University, Okayama, Japan*
- ¹¹⁰*Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, Oklahoma, USA*
- ¹¹¹*Department of Physics, Oklahoma State University, Stillwater, Oklahoma, USA*
- ¹¹²*Palacký University, RCPTM, Olomouc, Czech Republic*
- ¹¹³*Center for High Energy Physics, University of Oregon, Eugene, Oregon, USA*
- ¹¹⁴*LAL, Univ. Paris-Sud and CNRS/IN2P3, Orsay, France*
- ¹¹⁵*Graduate School of Science, Osaka University, Osaka, Japan*
- ¹¹⁶*Department of Physics, University of Oslo, Oslo, Norway*
- ¹¹⁷*Department of Physics, Oxford University, Oxford, United Kingdom*
- ^{118a}*INFN Sezione di Pavia, Italy*
- ^{118b}*Dipartimento di Fisica, Università di Pavia, Pavia, Italy*
- ¹¹⁹*Department of Physics, University of Pennsylvania, Philadelphia, Pennsylvania, USA*
- ¹²⁰*Petersburg Nuclear Physics Institute, Gatchina, Russia*
- ^{121a}*INFN Sezione di Pisa, Italy*
- ^{121b}*Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy*
- ¹²²*Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, Pennsylvania, USA*
- ^{123a}*Laboratorio de Instrumentacao e Fisica Experimental de Particulas - LIP, Lisboa, Portugal*
- ^{123b}*Departamento de Fisica Teorica y del Cosmos and CAFPE, Universidad de Granada, Granada, Spain*
- ¹²⁴*Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic*
- ¹²⁵*Faculty of Mathematics and Physics, Charles University in Prague, Praha, Czech Republic*
- ¹²⁶*Czech Technical University in Prague, Praha, Czech Republic*
- ¹²⁷*State Research Center Institute for High Energy Physics, Protvino, Russia*
- ¹²⁸*Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom*
- ¹²⁹*Physics Department, University of Regina, Regina SK, Canada*
- ¹³⁰*Ritsumeikan University, Kusatsu, Shiga, Japan*
- ^{131a}*INFN Sezione di Roma I, Italy*
- ^{131b}*Dipartimento di Fisica, Università La Sapienza, Roma, Italy*
- ^{132a}*INFN Sezione di Roma Tor Vergata, Italy*
- ^{132b}*Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy*
- ^{133a}*INFN Sezione di Roma Tre, Italy*
- ^{133b}*Dipartimento di Fisica, Università Roma Tre, Roma, Italy*
- ^{134a}*Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca, Morocco*
- ^{134b}*Centre National de l'Energie des Sciences Techniques Nucleaires, Rabat, Morocco*
- ^{134c}*Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech, Morocco*
- ^{134d}*Faculté des Sciences, Université Mohamed Premier and LTPM, Oujda, Morocco*
- ^{134e}*Faculté des Sciences, Université Mohammed V- Agdal, Rabat, Morocco*
- ¹³⁵*DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l'Univers), CEA Saclay (Commissariat a l'Energie Atomique), Gif-sur-Yvette, France*
- ¹³⁶*Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, California, USA*

- ¹³⁷*Department of Physics, University of Washington, Seattle, Washington, USA*
- ¹³⁸*Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom*
- ¹³⁹*Department of Physics, Shinshu University, Nagano, Japan*
- ¹⁴⁰*Fachbereich Physik, Universität Siegen, Siegen, Germany*
- ¹⁴¹*Department of Physics, Simon Fraser University, Burnaby BC, Canada*
- ¹⁴²*SLAC National Accelerator Laboratory, Stanford, California, USA*
- ^{143a}*Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava, Slovak Republic*
- ^{143b}*Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic*
- ^{144a}*Department of Physics, University of Johannesburg, Johannesburg, South Africa*
- ^{144b}*School of Physics, University of the Witwatersrand, Johannesburg, South Africa*
- ^{145a}*Department of Physics, Stockholm University, Sweden*
- ^{145b}*The Oskar Klein Centre, Stockholm, Sweden*
- ¹⁴⁶*Physics Department, Royal Institute of Technology, Stockholm, Sweden*
- ¹⁴⁷*Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook, New York, USA*
- ¹⁴⁸*Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom*
- ¹⁴⁹*School of Physics, University of Sydney, Sydney, Australia*
- ¹⁵⁰*Institute of Physics, Academia Sinica, Taipei, Taiwan*
- ¹⁵¹*Department of Physics, Technion: Israel Inst. of Technology, Haifa, Israel*
- ¹⁵²*Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel*
- ¹⁵³*Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece*
- ¹⁵⁴*International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan*
- ¹⁵⁵*Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan*
- ¹⁵⁶*Department of Physics, Tokyo Institute of Technology, Tokyo, Japan*
- ¹⁵⁷*Department of Physics, University of Toronto, Toronto ON, Canada*
- ^{158a}*TRIUMF, Vancouver BC, Canada*
- ^{158b}*Department of Physics and Astronomy, York University, Toronto ON, Canada*
- ¹⁵⁹*Institute of Pure and Applied Sciences, University of Tsukuba, 1-1-1 Tennodai, Tsukuba, Ibaraki 305-8571, Japan*
- ¹⁶⁰*Science and Technology Center, Tufts University, Medford, Massachusetts, USA*
- ¹⁶¹*Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia*
- ¹⁶²*Department of Physics and Astronomy, University of California Irvine, Irvine, California, USA*
- ^{163a}*INFN Gruppo Collegato di Udine, Italy*
- ^{163b}*ICTP, Trieste, Italy*
- ^{163c}*Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy*
- ¹⁶⁴*Department of Physics, University of Illinois, Urbana, Illinois, USA*
- ¹⁶⁵*Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden*
- ¹⁶⁶*Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain*
- ¹⁶⁷*Department of Physics, University of British Columbia, Vancouver BC, Canada*
- ¹⁶⁸*Department of Physics and Astronomy, University of Victoria, Victoria BC, Canada*
- ¹⁶⁹*Waseda University, Tokyo, Japan*
- ¹⁷⁰*Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel*
- ¹⁷¹*Department of Physics, University of Wisconsin, Madison, Wisconsin, USA*
- ¹⁷²*Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany*
- ¹⁷³*Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany*
- ¹⁷⁴*Department of Physics, Yale University, New Haven, Connecticut, USA*
- ¹⁷⁵*Yerevan Physics Institute, Yerevan, Armenia*
- ¹⁷⁶*Domaine Scientifique de la Doua, Centre de Calcul CNRS/IN2P3, Villeurbanne Cedex, France*

^aDeceased.

^bAlso at Laboratorio de Instrumentacao e Fisica Experimental de Particulas - LIP, Lisboa, Portugal.

^cAlso at Faculdade de Ciencias and CFNUL, Universidade de Lisboa, Lisboa, Portugal.

^dAlso at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom.

^eAlso at TRIUMF, Vancouver BC, Canada.

^fAlso at Department of Physics, California State University, Fresno, CA, USA.

^gAlso at Novosibirsk State University, Novosibirsk, Russia.

^hAlso at Fermilab, Batavia, IL, USA.

ⁱAlso at Department of Physics, University of Coimbra, Coimbra, Portugal.

^jAlso at Università di Napoli Parthenope, Napoli, Italy.

^kAlso at Institute of Particle Physics (IPP), Canada.

- ^lAlso at Department of Physics, Middle East Technical University, Ankara, Turkey.
- ^mAlso at LA Tech University, Ruston, LA, USA.
- ⁿAlso at Department of Physics and Astronomy, University College London, London, United Kingdom.
- ^oAlso at Group of Particle Physics, University of Montreal, Montreal QC, Canada.
- ^pAlso at Department of Physics, University of Cape Town, Cape Town, South Africa.
- ^qAlso at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.
- ^rAlso at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany.
- ^sAlso at Manhattan College, New York, NY, USA.
- ^tAlso at School of Physics, Shandong University, Shandong, China.
- ^uAlso at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France.
- ^vAlso at School of Physics and Engineering, Sun Yat-sen University, Guanzhou, China.
- ^wAlso at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan.
- ^xAlso at DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l'Univers), CEA Saclay (Commissariat à l'Energie Atomique), Gif-sur-Yvette, France.
- ^yAlso at Section de Physique, Université de Genève, Geneva, Switzerland.
- ^zAlso at Departamento de Física, Universidade de Minho, Braga, Portugal.
- ^{aa}Also at Department of Physics and Astronomy, University of South Carolina, Columbia, SC, USA
- ^{bb}Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.
- ^{cc}Also at CA Institute of Technology, Pasadena CA, USA.
- ^{dd}Also at Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China.
- ^{ee}Also at Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom.
- ^{ff}Also at Department of Physics, Oxford University, Oxford, United Kingdom.
- ^{gg}Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.
- ^{hh}Also at Department of Physics, The University of Michigan, Ann Arbor, MI, USA.
- ⁱⁱAlso at Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France.