Search for resonant WZ production in the WZ →lνl'l' channel in √(s) = 7 TeV pp collisions with the ATLAS detector


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A generic search is presented for a heavy particle decaying to $WZ \to \ell\nu\ell\ell'$ ($\ell, \ell' = e, \mu$) final states. The data were recorded by the ATLAS detector in $\sqrt{s} = 7 \text{ TeV}$ $pp$ collisions at the Large Hadron Collider and correspond to an integrated luminosity of 1.02 fb$^{-1}$. The transverse mass distribution of the selected $WZ$ candidates is found to be consistent with the standard model expectation. Upper limits on the production cross section times branching ratio are derived using two benchmark models predicting a heavy particle decaying to a $WZ$ pair.

III. THE ATLAS DETECTOR

The ATLAS detector [14] is a general-purpose particle detector with an approximately forward-backward symmetric cylindrical geometry, and almost 4$\pi$ coverage in solid angle [15]. The inner tracking detector (ID) covers the pseudorapidity range of $|\eta| < 2.5$ and consists of a silicon pixel detector, a silicon microstrip detector, and a transition radiation tracker. The ID is surrounded by a thin superconducting solenoid providing a 2 T magnetic field, and by a calorimeter system covering an $\eta$ range up to 4.9, which provides three-dimensional reconstruction of particle showers. For $|\eta| < 2.5$, the electromagnetic calorimeter is finely segmented and uses lead and liquid argon (LAr) as active material. The hadronic calorimeter uses steel and scintillating tiles in the barrel region, while the end caps use LAr as the active material and copper as an absorber. The forward calorimeter also uses LAr as an active medium with copper and tungsten as an absorber. The muon spectrometer (MS) is based on one absorber. The muon spectrometer (MS) is based on one absorber.
cathode strip chambers, allow a precise muon momentum measurement up to $|\eta| < 2.7$. Resistive plate and thin-gap chambers provide muon triggering capability up to $|\eta| < 2.4$.

### III. MONTE CARLO SIMULATION

Monte Carlo (MC) simulated samples are used to model signal and background processes. Events are generated at $\sqrt{s} = 7$ TeV, and the detector response simulation [16] is based on the GEANT4 program [17].

The simulation of the signals, both for the EGM $W'$ and the LSTC $\rho T$ production, is based on the LO PYTHIA 18 event generator, with the modified leading-order (LO) [19] parton distribution function (PDF) set MRST2007 LO [20]. By default, PYTHIA also includes $a_T$ production, as discussed below. A mass-dependent $k$ factor is used to rescale the LO PYTHIA prediction to the next-to-next-to-leading-order (NNLO) cross section. The $k$ factor is computed using the ZWPROD program [21] in the approximation of zero width for the resonance; its value decreases with the resonance mass from 1.17 at $m_W = 200$ GeV to 1.08 at $m_W = 1$ TeV.

The LSTC simulated samples correspond to the following set of parameters: number of technicolors $N_{TC} = 4$, charges of up-type and down-type technifermions $Q_U = 1$, $Q_D = 0$, mixing angle between technipions, and electroweak gauge boson longitudinal component $\sin \chi = 1/3$. The $\rho T$ can decay both to $WZ$ and $\pi T W$; if the $\rho T$ and $\pi T$ masses are degenerate, the branching ratio $BR(\rho T \rightarrow WZ)$ is 100%. Two-dimensional exclusion regions are set on the technicolor production in the $(m_{\rho T}, m_{\pi T})$ plane. In addition, for comparison purpose with previous results [13], the relation $m_{\rho T} = m_{\pi T} + m_W$ is used when extracting one-dimensional limits on the $\rho T$ mass, which entails a value of $BR(\rho T \rightarrow WZ) = 98\%$. The axial-vector partner of the $\rho T$, the $a_T$, also decays to $WZ$ and, depending on its mass, contributes to the $WZ$ production cross section. Two scenarios for the value of the mass of the $a_T$ technimeson are considered: $m_{a_T} = 1.1 \times m_{\rho T}$, which is the standard value implemented in PYTHIA, and $m_{a_T} \gg m_{\rho T}$, which is simulated by removing the $a_T$ contribution at the generator level.

The SM $WZ$ production, which is an irreducible background for this search, is modeled by the MC@NLO event generator [22], which incorporates the next-to-leading-order (NLO) matrix elements into the parton shower by interfacing to the HERWIG program [23]. The underlying event is modeled with JIMMY [24]. Other SM processes that can mimic the same final state include the following: $ZZ \rightarrow \ell \ell \ell \ell$, where one of the leptons is not detected or fails the selection requirements; $Z(\rightarrow \ell \ell) + \gamma$, where the photon is misidentified as an electron; and processes with two identified leptons and jets, namely $Z$ production in association with jets ($Z + \text{jets}$), $t\bar{t}$ and single top processes, where leptons are present from $b$- or $c$-hadron decays or one jet is misidentified as a lepton. SM $ZZ$ events are simulated at LO using HERWIG, and $W/Z + \gamma$ production is modeled with SHERPA [25]. The cross sections for these two processes are corrected to the NLO calculation computed with MCFM [26,27]. The $W/Z + \text{jets}$ process is modeled at LO using ALPGEN [28], and then corrected to the NNLO cross section computed with FEWZ [29]. Single top and $t\bar{t}$ events are simulated at NLO using MC@NLO. The backgrounds due to the $Z + \text{jets}$, $t\bar{t}$ and single top processes (called the “$\ell\ell\ell + \text{jets}$” background in this paper) are estimated using data-driven methods, and the corresponding MC samples mentioned above are used only for cross-checks.

### IV. EVENT SELECTION

The data analyzed are required to have been selected online by a single-lepton ($e$ or $\mu$) trigger with a threshold of 20 GeV on the transverse energy ($E_T$) in the electron case and 18 GeV on the transverse momentum ($p_T$) in the muon case. After applying data quality requirements, the total integrated luminosity of the data set used in this analysis is $1.02 \pm 0.04 \text{fb}^{-1}$ [30,31].

Because of the presence of multiple collisions in a single bunch-crossing, about six on average, each event can have multiple reconstructed primary vertices. The vertex having the largest sum of squared transverse momenta of associated tracks is selected as the primary vertex of the hard collision, and it is used to compute any reconstructed quantity. To reduce the contamination due to cosmic rays, only events where the primary vertex of the hard collision has at least three associated tracks with $p_T > 0.5$ GeV are considered.

Electrons are reconstructed from a combination of an ID track and a calorimeter energy cluster, with $E_T > 25$ GeV and $|\eta| < 1.37$ or $1.52 < |\eta| < 2.47$, avoiding the transition region between the barrel and the end-cap electromagnetic calorimeters. Candidate electrons must satisfy the medium [32] quality definition, which is based on the calorimeter shower shape, track quality, and track matching with the calorimeter cluster. To make sure candidate electrons originate from the primary interaction vertex, they are also required to have a longitudinal impact parameter ($|z_0|$) smaller than 10 mm and a transverse impact parameter ($d_0$) with significance $(|d_0|/\sigma_d)$ smaller than 10, both with respect to the selected primary vertex. In addition, the electron is required to be isolated in the calorimeter such that the sum of the $E_T$ of the clusters around the electron within a cone of $\Delta R = \sqrt{\Delta \eta^2 + \Delta \phi^2} = 0.3$ is less than 4 GeV. Corrections are applied to account for the energy deposition inside the isolation cone due to electron energy leakage and additional pileup collisions.

Muon candidates must be reconstructed in both the ID and the MS, and the combined track is required to have $p_T > 25$ GeV and $|\eta| < 2.4$. Good quality is ensured by
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requiring a minimum number of silicon strip and pixel hits associated to the track. To suppress the contribution of muons coming from hadronic jets, the $p_T$ sum of other tracks with $p_T > 1$ GeV, within a cone of $\Delta R = 0.2$ around the muon track, is required to be less than 10% of the muon $p_T$. The muon candidate is required to be compatible with the selected primary vertex, with $|z_0| < 10$ mm and $|d_0|/\sigma_{d_0} < 10$.

The missing transverse momentum $E_T^{\text{miss}}$ is reconstructed, in the range $|\eta| < 4.5$, as the negative vector sum of calorimeter cell transverse energies, calibrated to the electromagnetic scale [33], to which the transverse momenta of identified muons are added.

The $WZ \rightarrow \ell \nu \ell' \ell'$ candidate events are selected by requiring two oppositely charged same-flavor leptons with an invariant mass within 20 GeV of the Z boson mass, plus a third lepton and $E_T^{\text{miss}} > 25$ GeV. The transverse mass of the reconstructed $W$ boson, i.e. $m_T^W = \sqrt{2p_T\ell E_T^{\text{miss}}(1 - \cos \Delta \phi)}$, where $p_T\ell$ is the transverse momentum of the charged lepton and $\Delta \phi$ is the opening angle between the lepton and the $E_T^{\text{miss}}$ direction in the plane transverse to the beam, is required to be greater than 15 GeV to suppress multijet background. Selected events are also required to have exactly three charged leptons to suppress the $ZZ \rightarrow \ell \ell' \ell' \ell'$ background. These selection criteria define the signal region. Four decay channels $eeee$, $e\nu\mu\mu$, $\mu\nu\mu\mu$, and $\mu\mu\mu\mu$ are analyzed separately and then combined. The measurement of the inclusive $pp \rightarrow WZ \rightarrow \ell \nu \ell' \ell'$ cross section has previously been reported by ATLAS [34]. This analysis goes further by using the reconstructed event properties to probe for new phenomena.

After the final selection, the transverse mass of the WZ candidates ($m_T^{WZ}$) is examined for any resonant structure. Here $m_T^{WZ}$ is calculated as $m_T^{WZ} = \sqrt{(E_T^W + E_T^Z)^2 - (p_T^\ell + p_T^\nu)^2 - (p_T^{\ell'} + p_T^{\ell''})^2}$, where $E_T^W$ and $E_T^Z$ are the scalar sums of the transverse energies of the decay products of the Z and W candidates, respectively. The $E_T^{\text{miss}}$ vector is used as the estimator of the transverse momentum of the neutrino arising from the W boson decay.

V. BACKGROUND ESTIMATION

The dominant background for the $WZ$ resonance search comes from SM $WZ$ production. Its contribution is estimated using MC simulation. Simulated events are required to pass the event selection criteria, and the final yield is normalized to the integrated luminosity. Lepton reconstruction and identification efficiencies, energy scale, and resolution in the MC simulation are corrected to the corresponding values measured in the data in order to improve the overall modeling. Other diboson processes such as $ZZ$ and $Z\gamma$ are also estimated using MC simulation.

A data-driven approach is used to estimate the contribution of the $\ell \ell' + \text{jets}$ background in the signal region. It is estimated by selecting a data sample containing two leptons that pass all the quality criteria requested in the lepton selection, and a leptonlike jet, which is defined as a reconstructed object that satisfies all quality criteria requested in the lepton selection, and a leptonlike jet, which is defined as a reconstructed object that satisfies all quality criteria requested in the lepton selection, and a leptonlike jet, which is defined as a reconstructed object that satisfies all quality criteria requested in the lepton selection.

Data and SM predictions are compared in two dedicated signal-free control regions, selected by requiring the same
VI. SYSTEMATIC UNCERTAINTIES

Different sources of systematic uncertainties have been considered. The first source is related to the lepton trigger, reconstruction, and identification efficiencies. These efficiencies are evaluated with tag-and-probe methods using \( Z \rightarrow \ell \ell, W \rightarrow \ell \nu \), and \( J/\psi \rightarrow \ell \ell \) events [35]. Scale factors are used to correct for differences between data and MC simulation. The lepton trigger efficiency scale factors are compatible with unity, and a systematic uncertainty of 1% is considered. The lepton reconstruction and identification scale factors are close to 1 and have a systematic uncertainty of 1.2% for the electrons and 0.5% for muons [35]. The lepton isolation efficiency uncertainties are estimated to be 2% for electrons and 1% for muons.

The second source of uncertainty is related to the lepton energy, momentum, and \( E_{\text{miss}} \) reconstruction. Additional smearing is applied to the muon \( p_T \) and to the electron cluster energy in the simulation, so that they replicate the \( Z \rightarrow \ell \ell \) invariant mass distributions in data. The uncertainty due to the lepton resolution smearing is of the order of 0.1% [35]. The uncertainty on the \( E_{\text{miss}} \) reconstruction receives contributions from different sources: energy deposits due to additional \( pp \) collisions which are in time and out of time with respect to the bunch-crossing; energy deposits around clusters associated to reconstructed jets and electrons; energy deposits not associated to any reconstructed objects; and muon momentum uncertainties. The total systematic uncertainty on the dominant SM WZ background estimation due to the \( E_{\text{miss}} \) uncertainties lies between 2% and 3%, depending on the channel considered.

The third source of uncertainty is due to the limited knowledge of the theoretical cross sections of SM processes, used both to evaluate WZ, ZZ, and Z\( \gamma \) background contributions, and for subtracting contributions of W and Z leptonic decays from the dijet sample used for the measurement of the correction factor \( f \). An uncertainty of 7% is assigned for the WZ process, 5% for the ZZ process, and 8% for the Z\( \gamma \) process [27], to which the MC statistical uncertainty is added in quadrature.

The fourth source of uncertainty is related to the uncertainty on the \( \ell \ell' + \text{jets} \) background estimation. The systematic uncertainty comes mainly from the uncertainty on \( f \) due to differences in the kinematics and flavor composition of the QCD dijet events with respect to the \( \ell \ell' + \text{jets} \) processes, and differences in event selection criteria for QCD dijet events and WZ candidates. The factor \( f \) is around 0.15 for muons and 0.07 for electrons over the full range of \( p_T \) and \( \eta \), with a relative uncertainty between 5% and 20%. The estimated number of events from the \( \ell \ell' + \text{jets} \) background in the signal region using the data-driven method is \( 6.4 \pm 1.0(\text{stat})^{+1.3}_{-1.2}(\text{syst}) \) events. A MC-based cross-check gives a consistent estimation of \( 4.3 \pm 1.1(\text{syst}) \) events.

The fifth source of uncertainty is related to the estimation of the signal acceptance based on MC simulation. The systematic uncertainty is mainly due to the choice of PDF and is found to be 0.6% when comparing the differences between the predictions of the nominal PDF set MRST2007 LO' and the ones given by MSTW2008 LO [36], using the standard LHAPDF framework [37]. A cross-check has been done using the NNPDF LO' [38], CT09MCs, CT09MC1, and CT09MC2 [39] PDF sets, leading to a compatible uncertainty.

Finally, the luminosity uncertainty is 3.7% [30,31].

VII. RESULTS AND INTERPRETATION

The numbers of events expected and observed after the final selection are reported in Table I. A total of 48 WZ → \( \ell \nu \ell' \ell' \) candidate events are observed in data, to be compared to the SM prediction of \( 45.0 \pm 1.0(\text{stat})^{+4.6}_{-5.2}(\text{syst}) \) events. The expected numbers of events for a \( W' \) with a mass of 750 GeV and a \( p_T \) with a mass of 500 GeV are also reported.

The overall acceptance times trigger, reconstruction, and selection efficiencies \( (A \times \epsilon) \) for EGM \( W' \rightarrow W \rightarrow \ell \nu \ell' \ell' \) and the LSTC \( \rho_T \rightarrow W \rightarrow \ell \nu \ell' \ell' \) events as implemented in PYTHIA is shown in Table II for various \( WZ \) resonance masses. The value of \( A \times \epsilon \) is 6.2% for \( m_W = 200 \text{ GeV} \) and increases to 20.5% for \( m_W = 1 \text{ TeV} \). The corresponding \( A \times \epsilon \) for the LSTC \( \rho_T \) is found to be slightly lower than that of the EGM \( W' \) due to the fact that the PYTHIA implementation of the \( \rho_T \rightarrow WZ \) process does not account for the polarizations of vector bosons in their decay. A massive \( W' \) boson is expected to decay predominantly to longitudinally polarized W and Z bosons, as is the \( \rho_T \) technimeson. While the production and decay with spin correlations is fully implemented in PYTHIA for \( W' \), spin correlation information is not considered in the decay of the W and Z bosons in the \( \rho_T \) case; hence, they each decay isotropically in their respective rest frames. This leads to a softer lepton \( p_T \) spectrum and consequently lower \( A \times \epsilon \). The interpretation of the data in terms of \( \rho_T \) production is performed in two different manners: the first uses the PYTHIA implementation of \( \rho_T \) production and decay, and the second assumes that \( A \times \epsilon \) for the \( \rho_T \) is equal to that of the \( W' \).
The transverse mass distribution of the WZ candidates is presented in Fig. 2 for data and background expectations together with possible contributions from $W'$ and $p_T$ using PYTHIA. The $\ell\ell' +$ jets and $Z\gamma$ background contributions to the $m^{WZ}_T$ distribution are extrapolated using exponential functions to extend over the full $m^{WZ}_T$ signal region. The transverse mass distribution is used to build a log-likelihood ratio test statistic [40], which allows the compatibility of the data with the presence of a signal in addition to the background to be assessed, in a modified frequentist approach [41]. Confidence levels for the signal plus background hypothesis, $\text{CL}_{s+b}$, and background-only hypothesis, $\text{CL}_b$, are computed by integrating the log-likelihood ratio distributions obtained from simulated pseudoeperiments using Poisson statistics. The confidence level for the signal hypothesis $\text{CL}_s$, defined as the ratio $\text{CL}_{s+b}/\text{CL}_b$, is used to determine the exclusion limits.

The probability that the background fluctuations give rise to an excess at least as large as that observed in data has been computed as $p$-value $= 1 - \text{CL}_b$ and is reported in Table III for the signal hypothesis of a $W'$ particle with mass from 200 GeV to 1 TeV. Since no statistically significant excess is observed for any value of the $W'$ mass, limits are derived on the production cross section times branching ratio $|\sigma \times \text{BR}(W' \rightarrow WZ)|$ for a $W'$ decaying to $WZ$, already corrected for the $A \times \epsilon$ of the leptonic decay $WZ \rightarrow \ell\ell'\ell'\ell'$. The 95% CL limit on $\sigma \times \text{BR}(W' \rightarrow WZ)$ is defined as the value giving $\text{CL}_s = 0.05$. The upper limit on $\sigma \times \text{BR}(W' \rightarrow WZ)$ for $pp \rightarrow W' \rightarrow WZ$ as a function of the $W'$ mass is shown in Fig. 3(a), and the values are reported in Table III. Simulation of $W'$ bosons is performed for $m_{W'}$ between 200 GeV and 1 TeV with a 150 to 250 GeV mass spacing, and an interpolation procedure provides $m^{WZ}_T$.

The estimated background yields, the observed number of data events, and the predicted signal yield predicted by PYTHIA for a $W'$ boson with a mass of 750 GeV and a $p_T$ technimeson with a mass of 500 GeV are shown after applying all signal selection cuts, for each of the four channels considered and for their combination. For the $p_T$ production, the relation $m_{\gamma} = 1.1 \times m_{p_T}$ is used. Where one error is quoted, it includes all sources of systematic uncertainty. Where two errors are given, the first comes from the limited statistics of the data and the second includes systematic uncertainties.

<table>
<thead>
<tr>
<th></th>
<th>$ee$</th>
<th>$\mu ee$</th>
<th>$\mu \mu \mu$</th>
<th>$\mu \nu \mu$</th>
<th>Combined</th>
</tr>
</thead>
<tbody>
<tr>
<td>$WZ$</td>
<td>6.2</td>
<td>7.6</td>
<td>9.2</td>
<td>11.6</td>
<td>34.6</td>
</tr>
<tr>
<td>$ZZ$</td>
<td>0.25</td>
<td>0.48</td>
<td>0.37</td>
<td>0.63</td>
<td>1.7</td>
</tr>
<tr>
<td>$Z\gamma$</td>
<td>1.3</td>
<td>-</td>
<td>1.0</td>
<td>-</td>
<td>2.3</td>
</tr>
<tr>
<td>$\ell\ell' +$ jets</td>
<td>1.1</td>
<td>1.3</td>
<td>3.0</td>
<td>1.0</td>
<td>2.3</td>
</tr>
<tr>
<td>Overall backgrounds</td>
<td>8.9</td>
<td>9.4</td>
<td>13.6</td>
<td>13.6</td>
<td>6.4</td>
</tr>
<tr>
<td>Data</td>
<td>9</td>
<td>7</td>
<td>16</td>
<td>16</td>
<td>48</td>
</tr>
</tbody>
</table>
TABLE III. Expected and observed limits on the $\sigma \times BR(W^' \to WZ)$ (pb) for $W'$ production decaying to $WZ$, as a function of the $W'$ mass. The $p$ values are also reported.

<table>
<thead>
<tr>
<th>$W'$ mass (GeV)</th>
<th>Expected</th>
<th>Observed</th>
<th>$p$ value</th>
</tr>
</thead>
<tbody>
<tr>
<td>200</td>
<td>7.31</td>
<td>7.62</td>
<td>0.43</td>
</tr>
<tr>
<td>250</td>
<td>5.26</td>
<td>6.55</td>
<td>0.34</td>
</tr>
<tr>
<td>300</td>
<td>2.74</td>
<td>3.38</td>
<td>0.28</td>
</tr>
<tr>
<td>350</td>
<td>1.72</td>
<td>2.06</td>
<td>0.25</td>
</tr>
<tr>
<td>400</td>
<td>1.18</td>
<td>1.48</td>
<td>0.25</td>
</tr>
<tr>
<td>450</td>
<td>0.92</td>
<td>1.07</td>
<td>0.23</td>
</tr>
<tr>
<td>500</td>
<td>0.76</td>
<td>0.93</td>
<td>0.21</td>
</tr>
<tr>
<td>550</td>
<td>0.61</td>
<td>0.79</td>
<td>0.19</td>
</tr>
<tr>
<td>600</td>
<td>0.54</td>
<td>0.63</td>
<td>0.26</td>
</tr>
<tr>
<td>650</td>
<td>0.51</td>
<td>0.56</td>
<td>0.33</td>
</tr>
<tr>
<td>700</td>
<td>0.48</td>
<td>0.53</td>
<td>0.34</td>
</tr>
<tr>
<td>750</td>
<td>0.49</td>
<td>0.52</td>
<td>0.34</td>
</tr>
<tr>
<td>800</td>
<td>0.45</td>
<td>0.50</td>
<td>0.37</td>
</tr>
<tr>
<td>850</td>
<td>0.46</td>
<td>0.47</td>
<td>0.38</td>
</tr>
<tr>
<td>900</td>
<td>0.50</td>
<td>0.50</td>
<td>0.39</td>
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<tr>
<td>950</td>
<td>0.44</td>
<td>0.44</td>
<td>0.40</td>
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<tr>
<td>1000</td>
<td>0.48</td>
<td>0.46</td>
<td>0.35</td>
</tr>
</tbody>
</table>

shape templates with a 50 GeV spacing. The $m_T^{WZ}$ shapes from the fully simulated signal samples have been fitted with a Crystal Ball function using ROOFIT [42]. The obtained Crystal Ball parameters are fitted as a function of the $W'$ mass, and the functional value for these parameters is then used to build the $m_T^{WZ}$ templates for the intermediate $W'$ mass points. The observed (expected) exclusion limit on the $W'$ mass is found to be 760 (776) GeV.

The observed (expected) limits on $\sigma \times BR(p_T \rightarrow WZ)$ for the $p_T$ technimeson are shown in Fig. 3(b) assuming $m_{p_T} = 1.1 m_{p_T}$ and unpolarized $W$ and $Z$ decays. This corresponds to an observed (expected) limit on the $p_T$ mass of 467 (506) GeV. A limit on the $p_T$ mass of 456 (482) GeV is obtained if $m_{p_T} \gg m_{p_T}$. Assuming $A \times \epsilon$ for the $p_T$ signal to be equal to that of the $W'$ signal, which is estimated by accounting for predominantly longitudinal $W$ and $Z$ polarization, the observed (expected) limit on the $p_T$ mass is 483 (553) GeV for $m_{p_T} = 1.1 m_{p_T}$ and 469 (507) GeV for $m_{p_T} \gg m_{p_T}$. Table IV summarizes these limits, which all assume the relation $m_{p_T} = m_{p_T} + m_W$.

Figure 4 shows the 95% CL expected and observed excluded regions in the $(m_{p_T}, m_{WZ})$ plane for $m_{p_T} = 1.1 m_{p_T}$ and $m_{p_T} \gg m_{p_T}$, respectively. Results are shown under the two assumptions on $A \times \epsilon$ for the $p_T$ signal.

VIII. CONCLUSION

A search for resonant production of a pair of $WZ$ bosons with three charged leptons in the final state has been performed using 1.02 $fb^{-1}$ of data collected with the ATLAS detector in $pp$ collisions at $\sqrt{s} = 7$ TeV at the Large Hadron Collider. No significant excess of events is observed, and upper limits are derived on the production cross section times branching ratio of new physics using the transverse mass of the $WZ$ system. EGM $W'$ bosons with masses up to 760 GeV are excluded at 95% CL. Using

![FIG. 3 (color online). The observed and expected limits on $\sigma \times BR(W' \rightarrow WZ)$ for $W' \rightarrow WZ$ (a) and $pp \rightarrow p_T, a_T \rightarrow WZ$ (b). The theoretical prediction is shown with a systematic uncertainty of 5% due to the choice of PDF and is estimated by comparing the differences between the predictions of the nominal PDF set MRST2007 LO' and the ones given by MSTW2008 LO PDF using the LHAPDF framework. The green and yellow bands represent, respectively, the 1$\sigma$ and 2$\sigma$ uncertainties on the expected limit.](image-url)
FIG. 4 (color online). The 95% CL expected and observed excluded mass regions in the \((m_{\rho_T},m_{\pi_T})\) plane for \(m_{\pi_T} = 1.1 m_{\rho_T}\) (a) and \(m_{\pi_T} \gg m_{\rho_T}\) (b), above the curves. Two different assumptions about the \(p_T\) signal \(A \times \epsilon\) are used: assuming a \(p_T\) signal where \(A \times \epsilon\) is equal to that of the \(W\) signal, and assuming a \(p_T\) signal where \(A \times \epsilon\) is obtained through its implementation in PYTHIA.

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