Performance of missing transverse momentum reconstruction in proton-proton collisions at $\sqrt{s} = 7$ TeV with ATLAS

The ATLAS Collaboration*
CERN, 1211 Geneva 23, Switzerland

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Abstract The measurement of missing transverse momentum in the ATLAS detector, described in this paper, makes use of the full event reconstruction and a calibration based on reconstructed physics objects. The performance of the missing transverse momentum reconstruction is evaluated using data collected in $pp$ collisions at a centre-of-mass energy of 7 TeV in 2010. Minimum bias events and events with jets of hadrons are used from data samples corresponding to an integrated luminosity of about 0.3 nb$^{-1}$ and 600 nb$^{-1}$ respectively, together with events containing a $Z$ boson decaying to two leptons (electrons or muons) or a $W$ boson decaying to a lepton (electron or muon) and a neutrino, from a data sample corresponding to an integrated luminosity of about 36 pb$^{-1}$. An estimate of the systematic uncertainty on the missing transverse momentum scale is presented.

1 Introduction

In a collider event the missing transverse momentum is defined as the momentum imbalance in the plane transverse to the beam axis, where momentum conservation is expected. Such an imbalance may signal the presence of unseen particles, such as neutrinos or stable, weakly-interacting supersymmetric (SUSY) particles. The vector momentum imbalance in the transverse plane is obtained from the negative vector sum of the momenta of all particles detected in a $pp$ collision and is denoted as missing transverse momentum, $E_T^{\text{miss}}$. The symbol $E_T^{\text{miss}}$ is used for its magnitude.

A precise measurement of the missing transverse momentum, $E_T^{\text{miss}}$, is essential for physics at the LHC. A large $E_T^{\text{miss}}$ is a key signature for searches for new physics processes such as SUSY and extra dimensions. The measurement of $E_T^{\text{miss}}$ also has a direct impact on the quality of a number of measurements of Standard Model (SM) physics, such as the reconstruction of the top-quark mass in $t\bar{t}$ events. Furthermore, it is crucial in the search for the Higgs boson in the decay channels $H \rightarrow WW$ and $H \rightarrow \tau\tau$, where a good $E_T^{\text{miss}}$ measurement improves the reconstruction of the Higgs boson mass [1].

This paper describes an optimized reconstruction and calibration of $E_T^{\text{miss}}$ developed by the ATLAS Collaboration. The performance achieved represents a significant improvement compared to earlier results [2] presented by ATLAS. The optimal reconstruction of $E_T^{\text{miss}}$ in the ATLAS detector is complex and validation with data, in terms of resolution, scale and tails, is essential. A number of data samples encompassing a variety of event topologies are used. Specifically, the event samples used to assess the quality of the $E_T^{\text{miss}}$ reconstruction are: minimum bias events, events where jets at high transverse momentum are produced via strong interactions described by Quantum Chromodynamics (QCD) and events with leptonically decaying $W$ and $Z$ bosons. This allows the full exploitation of the detector capability in the reconstruction and calibration of different physics objects and optimization of the $E_T^{\text{miss}}$ calculation. Moreover, in events with $W \rightarrow \ell\nu$, where $\ell$ is an electron or muon, the $E_T^{\text{miss}}$ performance can be studied in events where genuine $E_T^{\text{miss}}$ is present due to the neutrino, thus allowing a validation of the $E_T^{\text{miss}}$ scale. In simulated events, the genuine $E_T^{\text{miss}}$, $E_T^{\text{miss, True}}$, is calculated from all generated non-interacting particles in the event and it is also referred to as true $E_T^{\text{miss}}$ in the following.

An important requirement on the measurement of $E_T^{\text{miss}}$ is the minimization of the impact of limited detector coverage, finite detector resolution, the presence of dead regions and different sources of noise that can produce fake $E_T^{\text{miss}}$. The ATLAS calorimeter coverage extends to large pseudorapidities to minimize the impact of high energy particles.

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* e-mail: atlas.publications@cern.ch

1 ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the $z$-axis coinciding with the axis of the beam pipe. The $x$-axis points from the IP to the centre of the LHC ring, and the $y$-axis points upward.
The ATLAS detector [1] is a multipurpose particle physics apparatus with a forward-backward symmetric cylindrical geometry and near 4π coverage in solid angle. The inner tracking detector (ID) covers the pseudorapidity range \(|\eta| < 2.5\), and consists of a silicon pixel detector, a silicon microstrip detector (SCT), and, for \(|\eta| < 2.0\), a transition radiation tracker (TRT). The ID is surrounded by a thin superconducting solenoid providing a 2 T magnetic field. A high-granularity lead/liquid-argon (LAr) sampling electromagnetic calorimeter covers the region \(|\eta| < 3.2\). An iron/scintillator-tile calorimeter provides hadronic coverage in the range \(|\eta| < 1.7\). LAr technology is also used for the hadronic calorimeters in the end-cap region \(1.5 < |\eta| < 3.2\) and for both electromagnetic and hadronic measurements in the forward region up to \(|\eta| < 4.9\). The muon spectrometer surrounds the calorimeters. It consists of three large air-core superconducting toroid systems, precision tracking chambers providing accurate muon tracking out to \(|\eta| = 2.7\), and additional detectors for triggering in the region \(|\eta| < 2.4\).

2 The ATLAS detector

3 Data samples and event selection

During 2010 a large number of proton-proton collisions, at a centre-of-mass energy of 7 TeV, were recorded with stable proton beams as well as nominal magnetic field conditions. Only data with a fully functioning calorimeter, inner detector and muon spectrometer are analysed.

Cuts are applied to clean the data sample against instrumental noise and out-of-time energy deposits in the calorimeter (from cosmic-rays or beam-induced background). Topological clusters reconstructed in the calorimeters (see Sect. 5.1) at the electromagnetic energy (EM) scale are used as the inputs to the jet finder [4]. In this paper the anti-

\[ k_T \]

algorithm [5], with distance parameter \(R = 0.6\), is used for jet reconstruction. The reconstructed jets are required to pass basic jet-quality selection criteria. In particular events are rejected if any jet in the event with transverse momentum \(p_T > 20\) GeV is caused by sporadic noise bursts in the end-cap region, coherent noise in the electromagnetic calorimeter or reconstructed from large out-of-time energy deposits in the calorimeter. These cuts largely suppress the residual sources of fake \(E_T^{\text{miss}}\) due to those instrumental effects which remain after the data-quality requirements.

The 2010 data sets used in this paper correspond to a total integrated luminosity [6, 7] of approximately 600 nb\(^{-1}\) for jet events and to 0.3 nb\(^{-1}\) for minimum bias events. Trigger and selection criteria for these events are described in Sect. 3.1. For the \(Z \to \ell\ell\) and \(W \to \ell\nu\) channels, the samples analysed correspond to an integrated luminosity of approximately 36 pb\(^{-1}\). Trigger and selection criteria, similar to those developed for the W/Z cross-section measurement [8], are applied. These criteria are described in Sects. 3.2 and 3.3.

3.1 Minimum bias and di-jet event selection

For the minimum bias events, only the early period of data taking, with a minimal pile-up contribution, is studied. Selected minimum bias events were triggered by the minimum bias trigger scintillators (MBTS), mounted at each end of the detector in front of the LAr end-cap calorimeter cryostats [9].

Events in the QCD jet sample are required to have passed the first-level calorimeter trigger, which indicates a significant energy deposit in a certain region of the calorimeter, with the most inclusive trigger with a nominal \(p_T\) threshold at 15 GeV. The event sample used in this analysis consists of two subsets of 300 nb\(^{-1}\) each, corresponding to

\[ E_T \]

The EM scale is the basic calorimeter signal scale for the ATLAS calorimeters. It provides the correct scale for energy deposited by electromagnetic showers. It does not correct for the lower energy hadron shower response nor for energy losses in the dead material.

Cylindrical coordinates \((r, \phi)\) are used in the transverse plane, \(\phi\) being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle \(\theta\) as \(\eta = -\ln \tan(\theta/2)\).
two periods with different pile-up and trigger conditions.\(^3\) One subset corresponds to the periods with lower pile-up conditions with, on average, 1 to 1.6 reconstructed vertices per event. The other subset corresponds to the periods with higher pileup conditions, where the peak number of visible inelastic interactions per bunch crossing goes up to 3. In the following, di-jet events are used, selected requiring the presence of exactly two jets with \(p_T > 25\) GeV and \(|\eta| < 4.5\). Jets are calibrated with the local hadronic calibration (see Sect. 5.1).

For each event, at least one good primary vertex is required with a \(z\) displacement from the nominal \(pp\) interaction point less than 200 mm and with at least five associated tracks. After selection, the samples used in the analysis presented here correspond to 14 million minimum bias events and 13 million di-jet events.

### 3.2 \(Z \rightarrow \ell\ell\) event selection

Candidate \(Z \rightarrow \ell\ell\) events, where \(\ell\) is an electron or a muon, are required to pass an \(e/\gamma\) or muon trigger with a \(p_T\) threshold between 10 and 15 GeV, where the exact trigger selection varies depending on the data period analysed. For each event, at least one good primary vertex, as defined above, is required.

The selection of \(Z \rightarrow \mu\mu\) events requires the presence of exactly two good muons. A good muon is defined to be a muon reconstructed in the muon spectrometer with a matched track in the inner detector with transverse momentum above 20 GeV and \(|\eta| < 2.5\) [10]. Additional requirements on the number of hits used to reconstruct the track in the inner detector are applied. The \(z\) displacement of the muon track from the primary vertex is required to be less than 10 mm. Isolation cuts are applied around the muon track.

The selection of \(Z \rightarrow ee\) events requires the presence of exactly two identified electrons with \(|\eta| < 2.47\), which pass the “medium” identification criteria [8, 11] and have transverse momenta above 20 GeV and \(|\eta| < 1.52\), are not considered for this study. Additional cuts are applied to remove electrons falling into regions where the readout of the calorimeter was not fully operational.

In both the \(Z \rightarrow ee\) and the \(Z \rightarrow \mu\mu\) selections, the two leptons are required to have opposite charge and the reconstructed invariant mass of the di-lepton system, \(m_{\ell\ell}\), is required to be consistent with the \(Z\) mass, \(66 < m_{\ell\ell} < 116\) GeV.

With these selection criteria, about 9000 \(Z \rightarrow ee\) and 13000 \(Z \rightarrow \mu\mu\) events are selected. The estimated background contribution to these samples is less than 2% in both channels [8].

### 3.3 \(W \rightarrow \ell\nu\) event selection

Lepton candidates are selected with lepton identification criteria similar to those used for the \(Z\) analysis. The differences for the selection of \(W \rightarrow e\nu\) events are that the “tight” electron identification criteria [8, 11] are used and an isolation cut is applied on the electron cluster in the calorimeter to reduce contamination from QCD jet background. The event is rejected if it contains more than one reconstructed lepton.

The \(E_T^{\text{miss}}\), calculated as described in Sect. 5, is required to be greater than 25 GeV, and the reconstructed lepton-\(E_T^{\text{miss}}\) transverse mass, \(m_T\), is required to be greater than 50 GeV.

With these selection criteria, about \(8.5 \times 10^4\) \(W \rightarrow e\nu\) and \(1.05 \times 10^5\) \(W \rightarrow \mu\nu\) events are selected. The background contribution to these samples is estimated to be about 5% in both channels [8].

### 4 Monte Carlo simulation samples

Monte Carlo (MC) events are generated using the PYTHIA6 program [12] with the ATLAS minimum bias tune (AMBT1) of the PYTHIA fragmentation and hadronisation parameters [13]. The generated events are processed with the detailed GEANT4 [14] simulation of the ATLAS detector.

The minimum bias MC event samples are generated using non-diffractive as well as single- and double-diffractive processes, where the different components are weighted according to the cross-sections given by the event generators.

The jet MC samples, generated using a 2-to-2 QCD matrix element and subsequent parton shower development, are used for comparison with the two subsets of data taken with different pile-up conditions. In the earlier sample the fraction of events with at least two observed interactions is at most of the order of 8–10%, while in the sample taken later in 2010 this fraction ranges from 10% to more than 50%. These samples are generated in the \(p_T\) range 8–560 GeV, in separated parton \(p_T\) bins to provide a larger statistics also in the high-\(p_T\) bins. Each sample is weighted according to its cross-section.

MC events for the study of SM backgrounds in \(Z \rightarrow \ell\ell\) and \(W \rightarrow \ell\nu\) analyses are also generated using PYTHIA6. The only exceptions are the \(t\bar{t}\) background and the \(W \rightarrow e\nu\) samples used in Sect. 8.2, which are generated with the MC@NLO program [15]. For the study of the total transverse energy of the events, samples produced with PYTHIA8 [16] are used as well.

MC samples were produced with different levels of pile-up in order to reflect the conditions in different data-taking periods.
periods. In particular, two event samples were used for jets: one was simulated with a pile-up model where only pile-up collisions originating from the primary bunch crossing are considered (in-time pile-up) and a second one was simulated with a realistic configuration of the LHC bunch group structure, where pile-up collisions from successive bunch crossings are also included in the simulation. In the case of events containing $Z \rightarrow \ell\ell$ or $W \rightarrow \ell\nu$, MC samples with in-time pile-up configuration are used, because these data correspond to periods where the contribution of out-of-time pileup is small.

The trigger and event selection criteria used for the data are also applied to the MC simulation.

5 $E_T^{\text{miss}}$ reconstruction and calibration

The $E_T^{\text{miss}}$ reconstruction includes contributions from energy deposits in the calorimeters and muons reconstructed in the muon spectrometer. The two $E_T^{\text{miss}}$ components are calculated as:

$$E_T^{\text{miss}} = E_T^{\text{miss,calo}} + E_T^{\text{miss,\mu}}.$$  

(1)

Low-$p_T$ tracks are used to recover low $p_T$ particles which are missed in the calorimeters (see Sect. 5.3.1), and muons reconstructed from the inner detector are used to recover muons in regions not covered by the muon spectrometer (see Sect. 5.2). The two terms in the above equation are referred to as the calorimeter and muon terms, and will be described in more detail in the following sections. The values of $E_T^{\text{miss}}$ and its azimuthal coordinate ($\phi^{\text{miss}}$) are then calculated as:

$$E_T^{\text{miss}} = \sqrt{(E_x^{\text{miss}})^2 + (E_y^{\text{miss}})^2},$$

$$\phi^{\text{miss}} = \arctan(E_y^{\text{miss}}, E_x^{\text{miss}}).$$  

(2)

5.1 Calculation of the $E_T^{\text{miss}}$ calorimeter term

In this paper, the $E_T^{\text{miss}}$ reconstruction uses calorimeter cells calibrated according to the reconstructed physics object to which they are associated. Calorimeter cells are associated with a reconstructed and identified high-$p_T$ parent object in a chosen order: electrons, photons, hadronically decaying $\tau$-leptons, jets and muons. Cells not associated with any such objects are also taken into account in the $E_T^{\text{miss}}$ calculation. Their contribution, named $E_T^{\text{miss,CellOut}}$ hereafter, is important for the $E_T^{\text{miss}}$ resolution [17].

Once the cells are associated with objects as described above, the $E_T^{\text{miss}}$ calorimeter term is calculated as follows (note that the $E_T^{\text{miss,calo,\mu}}$ term is not always added, as explained in Sect. 5.2, and for that reason it is written between parentheses):

$$E_{x(y)}^{\text{miss,calo,\mu}} = E_{x(y)}^{\text{miss,e}} + E_{x(y)}^{\text{miss,\gamma}} + E_{x(y)}^{\text{miss,\tau}} + E_{x(y)}^{\text{miss,jets}} + E_{x(y)}^{\text{miss,softjets}} + (E_{x(y)}^{\text{miss,calo,\mu}}) + E_{x(y)}^{\text{miss,CellOut}},$$  

(3)

where each term is calculated from the negative sum of calibrated cell energies inside the corresponding objects, as:

$$E_{x(y)}^{\text{miss,term}} = -\sum_{i=1}^{N_{\text{term}} \text{cell}} E_i \sin \theta_i \cos \phi_i,$$

$$E_{x(y)}^{\text{miss,term}} = -\sum_{i=1}^{N_{\text{term}} \text{cell}} E_i \sin \theta_i \sin \phi_i,$$

(4)

where $E_i$, $\theta_i$, and $\phi_i$ are the energy, the polar angle and the azimuthal angle, respectively. The summations are over all cells associated with specified objects in the pseudorapidity range $|\eta| < 4.5$.

Because of the high granularity of the calorimeter, it is crucial to suppress noise contributions and to limit the cells used in the $E_T^{\text{miss}}$ sum to those containing a significant signal. This is achieved by using only cells belonging to three-dimensional topological clusters, referred as topoclusters hereafter [18], with the exception of electrons and photons for which a different clustering algorithm is used [11]. The topoclusters are seeded by cells with deposited energy [19] $|E_i| > 4\sigma_{\text{noise}}$, and are built by iteratively adding neighbouring cells with $|E_i| > 2\sigma_{\text{noise}}$ and, finally, by adding all neighbours of the accumulated cells.

The various terms in (3) are described in the following:

- $E_{x(y)}^{\text{miss,e}}$, $E_{x(y)}^{\text{miss,\gamma}}$, $E_{x(y)}^{\text{miss,\tau}}$, $E_{x(y)}^{\text{miss,jets}}$, $E_{x(y)}^{\text{miss,softjets}}$ are reconstructed from cells in clusters associated to electrons, photons and $\tau$-jets from hadronically decaying $\tau$-leptons, respectively;
- $E_{x(y)}^{\text{miss,jets}}$ is reconstructed from cells in clusters associated to jets with calibrated $p_T > 20$ GeV;
- $E_{x(y)}^{\text{miss,softjets}}$ is reconstructed from cells in clusters associated to jets with $7$ GeV < $p_T < 20$ GeV;
- $E_{x(y)}^{\text{miss,calo,\mu}}$ is the contribution to $E_T^{\text{miss}}$ originating from the energy lost by muons in the calorimeter (see Sect. 5.2);
- the $E_{x(y)}^{\text{miss,CellOut}}$ term is calculated from the cells in topoclusters which are not included in the reconstructed objects.

All these terms are calibrated independently as described in Sect. 5.3. The final $E_{x(y)}^{\text{miss}}$ is calculated from (1) adding the $E_{x(y)}^{\text{miss,\mu}}$ term, described in Sect. 5.2.

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4This $\eta$ cut is chosen because the MC simulation does not describe data well in the very forward region.

5$\sigma_{\text{noise}}$ is the Gaussian width of the EM cell energy distribution measured in randomly triggered events far from collision bunches.
5.2 Calculation of the $E_T^{\text{miss}}$ muon term

The $E_T^{\text{miss}}$ muon term is calculated from the momenta of muon tracks reconstructed with $|\eta| < 2.7$:

$$E_{x(y)}^{\text{miss},\mu} = - \sum_{\text{muons}} P_{x(y)}^{\mu},$$  \hspace{1cm} (5)

where the summation is over selected muons. In the region $|\eta| < 2.5$, only well-reconstructed muons in the muon spectrometer with a matched track in the inner detector are considered (combined muons). The matching requirement considerably reduces contributions from fake muons (reconstructed muons not corresponding to true muons). These fake muons can sometimes be created from high hit multiplicities in the muon spectrometer in events where some particles from very energetic jets punch through the calorimeter into the muon system.

In order to deal appropriately with the energy deposited by the muon in the calorimeters, $E_{x(y)}^{\text{miss},\text{calo},\mu}$, the muon term is calculated differently for isolated and non-isolated muons, with non-isolated muons defined as those within a distance $\Delta R = \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2} < 0.3$ of a reconstructed jet in the event:

- The $p_T$ of an isolated muon is determined from the combined measurement of the inner detector and muon spectrometer, taking into account the energy deposited in the calorimeters. In this case the energy lost by the muon in the calorimeters ($E_{x(y)}^{\text{miss},\text{calo},\mu}$) is not added to the calorimeter term (see (3)) to avoid double counting of energy.

- For a non-isolated muon, the energy deposited in the calorimeter cannot be resolved from the calorimetric energy depositions of the particles in the jet. The muon spectrometer measurement of the muon momentum after energy loss in the calorimeter is therefore used, so the $E_{x(y)}^{\text{miss},\text{calo},\mu}$ term is added to the calorimeter term (see (3)). Only in cases in which there is a significant mis-match between the spectrometer and the combined measurement, the combined measurement is used and a parameterized estimation of the muon energy loss in the calorimeter [10] is subtracted.

For higher values of pseudorapidity ($2.5 < |\eta| < 2.7$), outside the fiducial volume of the inner detector, there is no matched track requirement and the muon spectrometer $p_T$ alone is used for both isolated and non-isolated muons.

Aside from the loss of muons outside the acceptance of the muon spectrometer ($|\eta| > 2.7$), muons can be lost in other small inactive regions (around $|\eta| = 0$ and $|\eta| \sim 1.2$) of the muon spectrometer. The muons which are reconstructed by segments matched to inner detector tracks extrapolated to the muon spectrometer are used to recover their contributions to $E_T^{\text{miss}}$ in the $|\eta| \sim 1.2$ regions [10].

Although the core of the $E_T^{\text{miss}}$ resolution is not much affected by the muon term, any muons which are not reconstructed, badly measured, or fake, can be a source of fake $E_T^{\text{miss}}$.

5.3 Calibration of $E_T^{\text{miss}}$

The calibration of $E_T^{\text{miss}}$ is performed using the scheme described below, where the cells are calibrated separately according to their parent object:

- The $E_T^{\text{miss},\gamma}$ term is calculated from photons passing the "medium" electron identification requirements, with $p_T > 10$ GeV and calibrated with the default electron calibration [8].

- The $E_T^{\text{miss},\tau}$ term is calculated from photons reconstructed with the "tight" tau identification requirements [11], with $p_T > 10$ GeV at the EM scale. Due to the low photon purity, the default photon calibration is not applied.

- The $E_T^{\text{miss},\tau}$ term is calculated from $\tau$-jets reconstructed with the "tight" $\tau$-identification requirements [19], with $p_T > 10$ GeV, calibrated with the local hadronic calibration (LCW) scheme [20]. The LCW scheme uses properties of clusters to calibrate them individually. It first classifies calorimeter clusters as electromagnetic or hadronic, according to the cluster topology, and then weights each calorimeter cell in clusters according to the cluster energy and the cell energy density. Additional corrections are applied to the cluster energy for the average energy deposited in the non-active material before and between the calorimeters and for unclustered calorimeter energy.

- The $E_T^{\text{miss},\text{softjets}}$ term is calculated from jets (reconstructed using the anti-$k_t$ algorithm with $R = 0.6$) with $7 < p_T < 20$ GeV calibrated with the LCW calibration.

- The $E_T^{\text{miss},\text{jets}}$ term is calculated from jets with $p_T > 20$ GeV calibrated with the LCW calibration and the jet energy scale (JES) factor [21] applied. The JES factor corrects the energy of jets, either at the EM-scale or after cluster calibration, back to particle level. The JES is derived as a function of reconstructed jet $\eta$ and $p_T$ using the generator-level information in MC simulation.

- The $E_T^{\text{miss},\text{CellOut}}$ term is calculated from topoclusters outside reconstructed objects with the LCW calibration and from reconstructed tracks as described in Sect. 5.3.1.

Note that object classification criteria and calibration can be chosen according to specific analysis criteria, if needed.

5.3.1 Calculation of the $E_T^{\text{miss},\text{CellOut}}$ term with a track-cluster matching algorithm

In events with $W$ and $Z$ boson production, the calibration of the $E_T^{\text{miss},\text{CellOut}}$ term is of particular importance because, due to the low particle multiplicity in these events,
this $E_{\text{T}}^\text{miss}$ contribution balances the $W/Z$ boson $p_T$ to a large extent [17]. An energy-flow algorithm is used to improve the calculation of the low-$p_T$ contribution to $E_{\text{T}}^\text{miss}$ ($E_{\text{T}}^\text{miss,CellOut}$). Tracks are added to recover the contribution from low-$p_T$ particles which do not reach the calorimeter or do not seed a topocluster. Furthermore, the track momentum is used instead of the topocluster energy for tracks associated to topoclusters, thus exploiting the better calibration and resolution of tracks at low momentum compared to topoclusters.

Reconstructed tracks with $p_T > 400$ MeV, passing track quality selection criteria such as the number of hits and $\chi^2$ of the track fit, are used for the calculation of the $E_{\text{T}}^\text{miss,CellOut}$ term. All selected tracks are extrapolated to the second layer of the electromagnetic calorimeter and very loose criteria are used for association to reconstructed objects or topoclusters, to avoid double counting. If a track is neither associated to a topocluster nor a reconstructed object, its transverse momentum is added to the calculation of $E_{\text{T}}^\text{miss,CellOut}$.

In the case where the track is associated to a topocluster, its transverse momentum is used for the calculation of the $E_{\text{T}}^\text{miss,CellOut}$ and the topocluster energy is discarded, assuming that the topocluster energy corresponds to the charged particle giving the track. It has to be noticed that there is a strong correlation between the number of particles and topoclusters, so, in general no neutral energy is lost replacing the topocluster by a track, and the neutral topoclusters are kept in most of the cases. If more than one topocluster is associated to a track, only the topocluster with the largest energy is excluded from the $E_{\text{T}}^\text{miss}$ calculation, assuming that this energy corresponds to the track.

6 Study of $E_{\text{T}}^\text{miss}$ performance

In this section the distributions of $E_{\text{T}}^\text{miss}$ in minimum bias, di-jet, $Z \to \ell\ell$ and $W \to \ell\nu$ events from data are compared with the expected distributions from the MC samples. The performance of $E_{\text{T}}^\text{miss}$ in terms of resolution and scale is also derived.

Minimum bias, di-jet events and $Z \to \ell\ell$ events are used to investigate the $E_{\text{T}}^\text{miss}$ performance without relying on MC detector simulation. In general, apart from a small contribution from the semi-leptonic decay of heavy-flavour hadrons in jets, no genuine $E_{\text{T}}^\text{miss}$ is expected in these events. Thus most of the $E_{\text{T}}^\text{miss}$ reconstructed in these events is a direct result of imperfections in the reconstruction process or in the detector response.

6.1 $E_{\text{T}}^\text{miss}$ performance in minimum bias and di-jet events

The distributions of $E_{x}^\text{miss}$, $E_{y}^\text{miss}$, $E_{\text{T}}^\text{miss}$ and $\phi^\text{miss}$ for data and MC simulation are shown in Fig. 1 for minimum bias events. The distributions are shown only for events with total transverse energy (see definition at the end of this section) greater than 20 GeV in order to reduce the contamination of fake triggers from the MBTS. Figure 2 shows the distributions of the same variables for the di-jet sample. The di-jet sample corresponding to the periods with higher pileup conditions (see Sect. 3.1) is used. The MC simulation expectations are superimposed, normalized to the number of events in the data.

In di-jet events a reasonable agreement is found between data and simulation for all basic quantities, while there is some disagreement in minimum bias events, attributed to imperfect modelling of soft particle activity in the MC simulation. The better agreement between data and MC simulation in the $\phi^\text{miss}$ distribution for the di-jet sample can be partly explained by the fact that the $E_{\text{T}}^\text{miss}$ is not corrected for the primary vertex position; the primary vertex position in data is better reproduced by the MC simulation for the di-jet sample than in the case of the minimum bias sample.

Events in the tails of the $E_{\text{T}}^\text{miss}$ distributions have been carefully checked, in order to understand the origin of the large measured $E_{\text{T}}^\text{miss}$. The tails are not completely well described by MC simulation, but, both in data and in MC simulation they are in general due to mis-measured jets. In minimum bias events there are more events in the tail in MC simulation and this can be due to the fact that the MC statistics is larger than in data. In di-jet events, there are more events in the tail in data. More MC events would be desirable. In di-jet events there are 19 events with $E_{\text{T}}^\text{miss} > 110$ GeV in the data. The majority of them (13 events) are due to mis-measured jets, where in most of the cases at least one jet points to a transition region between calorimeters. Two events are due to a combination of mis-measured jets with an overlapping muon, and one event is due to a fake high-$p_T$ muon. Finally two events look like good $b\bar{b}$ candidates, and one event has one reconstructed jet and no activity in the other hemisphere.

The events with fake $E_{\text{T}}^\text{miss}$ due to mis-measured jets and jets containing leptonically decays of heavy hadrons can be rejected by a cut based on the azimuthal angle between the jet and $E_{\text{T}}^\text{miss}$, $\Delta\phi(\text{jet}, E_{\text{T}}^\text{miss})$. Since the requirement of event cleaning depends on the physics analysis, the minimal cleaning cut is applied and careful evaluation of tail events is performed in this paper. Analyses that rely on a careful understanding and reduction of the tails of the $E_{\text{T}}^\text{miss}$ distribution (e.g. SUSY searches such as Ref. [3]) have performed more detailed studies to characterize the residual tail in events containing high-$p_T$ jets. These analyses use tighter jet cleaning cuts, track-jet matching, and angular cuts on $\Delta\phi(\text{jet}, E_{\text{T}}^\text{miss})$ to further reduce the fake $E_{\text{T}}^\text{miss}$ tail. In Ref. [3] a fully data-driven method (described in detail in Ref. [17]) was then employed to determine the residual fake $E_{\text{T}}^\text{miss}$ background.
The contributions from jets, soft jets and topoclusters not associated to the reconstructed objects and muons are shown in Fig. 3 for the di-jet events. The data-MC agreement is good for all of the terms contributing to $E_{\text{miss}}^T$. The tails observed in the muon term are mainly due to reconstructed fake muons and to one cosmic-ray muon, which can be rejected by applying a tighter selection for the muons used in the $E_{\text{miss}}^T$ reconstruction, based on $\chi^2$ criteria for the combination, isolation criteria and requirements on the number of hits in muon chambers used for the muon reconstruction.

In the following some distributions are shown for the total transverse energy, $\sum E_T$, which is an important quantity to parameterise and understand the $E_{\text{miss}}^T$ performance. It is defined as:

$$\sum E_T = \sum_{i=1}^{N_{\text{cell}}} E_i \sin \theta_i,$$

(6)

where $E_i$ and $\theta_i$ are the energy and the polar angle, respectively, of calorimeter cells associated to topoclusters within $|\eta| < 4.5$. Cell energies are calibrated according to the scheme described in Sect. 5.3 for $E_{\text{miss}}^T$.

The data distributions of $\sum E_T$ for minimum bias and di-jet events from the subset corresponding to lower pileup conditions (see Sect. 3.1) are compared to MC predictions from two versions of PYTHIA in Fig. 4. The left-hand distributions show comparisons with the ATLAS tune of PYTHIA6. The right-hand distributions show the comparisons with the default tune of PYTHIA8. Due to the limited number of events simulated, the distribution for the di-jet PYTHIA8 MC sample is not smooth, and is zero in the lowest $\sum E_T$ bin populated by data. This is not understood, also if it can be partly explained by the fact that the low $\sum E_T$ region is populated by events from the jet MC sample generated in the lowest parton $p_T$ bin (17–35 GeV), which is the most suppressed by the di-jet selection (a factor about 20 more than other samples) and has a large weight, due to cross-section. Moreover the PYTHIA8 jet MC sample in the 8–17 GeV parton $p_T$ bin is not available. In the case of the minimum bias sample, due to the very limited number of events simulated (about a factor 25 less respect to data), the tails in the PYTHIA8 MC distribution are strongly depleted.

The PYTHIA8 MC [16] version used in this paper has not yet been tuned to the ATLAS data. The current tune [22] uses the CTEQ 6.1 parton distribution functions (PDF).
instead of the MRST LO as used in PYTHIA6, and its diffraction model differs, including higher-\(Q^2\) diffractive processes. The comparison of the mean values and the shapes of the two different MC distributions with data seems to indicate that a better agreement is obtained with the PYTHIA8 but, due to the reduced PYTHIA8 MC statistics, no firm conclusion can be drawn. In the rest of the paper, the PYTHIA6 MC samples with the ATLAS tune are used for comparison with data; this version is used as the baseline for PYTHIA MC samples for 2010 data analyses.

6.2 \(E_T^{\text{miss}}\) performance in \(Z \rightarrow \ell\ell\) events

The absence of genuine \(E_T^{\text{miss}}\) in \(Z \rightarrow \ell\ell\) events, coupled with the clean event signature and the relatively large cross-section, means that it is a good channel to study \(E_T^{\text{miss}}\) performance.

The distributions of \(E_T^{\text{miss}}\) and \(\phi^{\text{miss}}\) for data and MC simulation are shown in Fig. 5 for \(Z \rightarrow ee\) and \(Z \rightarrow \mu\mu\) events. The contributions due to muons are shown for \(Z \rightarrow \mu\mu\) events in Fig. 6. Both the contributions from energy deposited in calorimeter cells associated to muons, taken at the EM scale, and the contributions from reconstructed muons are shown. For \(Z \rightarrow ee\) events, the contributions from electrons, jets, soft jets and topoclusters outside the reconstructed objects are shown separately in Fig. 7. The peak at zero in the distribution of the jet term corresponds to events where there are no jets with \(p_T\) above 20 GeV, and the small values (<20 GeV) in the distribution are due to events with two jets whose transverse momenta balance. The MC simulation expectations, from \(Z \rightarrow \ell\ell\) events and from the dominant SM backgrounds, are superimposed. Each MC sample is weighted with its corresponding cross-section and then the total MC expectation is normalized to the number of events in data. Reasonable agreement between data and MC simulation is observed in all distributions.

Events in the tails of the \(E_T^{\text{miss}}\) distributions in Fig. 5 have been carefully checked. The 22 events with the highest \(E_T^{\text{miss}}\) values, above 60 GeV, have been examined in detail to check whether they are related to cosmic-ray muon backgrounds, fake muons, badly measured jets or jets pointing to dead calorimeter regions. The events in the tails are found to be compatible with either signal candidates, including \(t\bar{t}\), WW and WZ di-boson events, all involving real \(E_T^{\text{miss}}\), or events in which the \(E_T^{\text{miss}}\) vector is close to a jet in the transverse plane. The latter category of events can arise from mis-
measured jets, and be rejected at the analysis level with cuts on $\Delta \phi(\text{jet}, E_T^{\text{miss}})$ (see Sect. 6.1).

### 6.2.1 Measuring $E_T^{\text{miss}}$ response in $Z \rightarrow \ell \ell$ events

From the event topology [17] in events with $Z \rightarrow \ell \ell$ decay one can define an axis in the transverse plane such that the component of $E_T^{\text{miss}}$ along this axis is sensitive to detector resolution and biases. The direction of this axis, $A_Z$, is defined by the reconstructed momenta of the leptons:

$$A_Z = \left( p_T^{\ell^+} + p_T^{\ell^-} \right) / \left| p_T^{\ell^+} + p_T^{\ell^-} \right|,$$

where $p_T^{\ell}$ are the vector transverse momenta of the lepton and anti-lepton. The direction of $A_Z$ thus reconstructs the direction of motion of the $Z$ boson. The perpendicular axis in the transverse plane, $A_{AZ}$, is a unit vector placed at right angles to $A_Z$, with positive direction anticlockwise from the direction of the $Z$ boson.

The mean value of the projection of $E_T^{\text{miss}}$ onto the longitudinal axis, $\langle E_T^{\text{miss}} \cdot A_Z \rangle$, is a measure of the $E_T^{\text{miss}}$ scale, as this axis is sensitive to the balance between the leptons and the hadronic recoil. Figure 8 shows the value of $\langle E_T^{\text{miss}} \cdot A_Z \rangle$ as a function of $p_T^Z$. These mean values are used as a diagnostic to validate the $E_T^{\text{miss}}$ reconstruction algorithms. If the leptons perfectly balanced the hadronic recoil, regardless of the net momentum of the lepton system, then the $E_T^{\text{miss}} \cdot A_Z$ would be zero, independent of $p_T^Z$. Instead, $\langle E_T^{\text{miss}} \cdot A_Z \rangle$ displays a small bias in both the electron and muon channels which is reasonably reproduced by the MC simulation. The observed bias is slightly negative for low values of $p_T^Z$, suggesting either that the $p_T$ of the lepton system is overestimated or that the magnitude of the hadronic recoil is underestimated. The same sign and magnitude of bias is seen in both electron and muon channels, suggesting that the hadronic recoil, here dominated by $E_T^{\text{miss}} \cdot \text{JetOut}$, is the source of this bias.

In Fig. 9 the dependences of $\langle E_T^{\text{miss}} \cdot A_Z \rangle$ on $p_T^Z$ are shown separately for events with $Z \rightarrow \ell \ell$ produced in association with zero jets or with at least one jet, with the jet definition as described in Sect. 3.1. The figure demonstrates that...
there is a negative bias in \( \langle E^\text{miss}_T \cdot A_Z \rangle \) for events with zero jets, which increases with \( p_T^Z \) up to 6 GeV. A similar bias is observed in both electron and muon channels, hence it is interpreted as coming from imperfections in the calibration of the soft hadronic recoil (the \( E^\text{miss}_T, \text{CellOut} \) and the \( E^\text{miss}_T, \text{softjets} \) terms). In events with at least one jet there is a small positive bias in the electron channel at high \( p_T^Z \), which is visible also in the muon channel for \( p_T^Z \) in the region 15–20 GeV.

Figure 10 shows \( \langle E^\text{miss}_T \cdot A_Z \rangle \) for \( Z \rightarrow \ell \ell \) events where there are neither high \( p_T^\ell \) nor soft jets, for two cases of \( E^\text{miss}_T \) reconstruction: calculating the \( E^\text{miss}_T, \text{CellOut} \) term with the track-cluster matching algorithm (see Sect. 5.3.1) or calculating this term from the calorimeter topoclusters only (denoted as \( E^\text{miss}_T, \text{no tracks} \)). The plots show a lower bias for the case with the track-cluster matching algorithm, indicating that it improves the reconstruction of the \( E^\text{miss}_T, \text{CellOut} \) term.

6.3 \( E^\text{miss}_T \) performance in \( W \rightarrow \ell \nu \) events

In this section the \( E^\text{miss}_T \) performance is studied in \( W \rightarrow e\nu \) and \( W \rightarrow \mu\nu \) events. In these events genuine \( E^\text{miss}_T \) is expected due to the presence of the neutrino, therefore the \( E^\text{miss}_T \) scale can be checked.

The distributions of \( E^\text{miss}_T \) and \( \phi^\text{miss} \) in data and in MC simulation are shown in Fig. 11 for \( W \rightarrow e\nu \) and \( W \rightarrow \mu\nu \) events. The contributions due to muons are shown for \( W \rightarrow \mu\nu \) events in Fig. 12. Both, the \( E^\text{miss}_T \) contribution from energy deposited in calorimeter cells associated to muons, taken at the EM scale, and the \( E^\text{miss}_T \) contribution from reconstructed muons are shown. The contributions given by the electrons, jets, soft jets and topoclusters outside reconstructed objects are shown in Fig. 13 for \( W \rightarrow e\nu \) events. The MC expectations are also shown, both from \( W \rightarrow \ell \nu \) events, and from the dominant SM backgrounds. The MC simulation describes all of the quantities well, with the exception that very small data-MC discrepancies are observed in the distribution of the \( E^\text{miss}_T,e \) at low \( E^\text{miss}_T \) values. This can be attributed to the QCD jet background, which would predominantly populate the region of low \( E^\text{miss}_T \) [8], but which is not included in the MC expectation shown.
Fig. 5  Distribution of $E_{\text{miss}}$ (top) and $\phi_{\text{miss}}$ (bottom) as measured in a data sample of $Z \rightarrow e\bar{e}$ (left) and of $Z \rightarrow \mu\bar{\mu}$ (right). The expectation from Monte Carlo simulation is superimposed and normalized to data, after each MC sample is weighted with its corresponding cross-section. The sum of all backgrounds is shown in the lower plots.

Fig. 6  Distribution of $E_{\text{miss}}$ computed with calorimeter cells associated to muons ($E_{\text{miss,calo,\mu}}$) (left) and computed from reconstructed muons ($E_{\text{miss,\mu}}$) (right) for $Z \rightarrow \mu\bar{\mu}$ data. The expectation from Monte Carlo simulation is superimposed and normalized to data, after each MC sample is weighted with its corresponding cross-section.

6.3.1 $E_{\text{miss}}$ linearity in $W \rightarrow \ell\nu$ MC events

The expected $E_{\text{miss}}$ linearity, which is defined as the mean value of the ratio: $(E_{\text{miss}} - E_{\text{miss, \text{True}}}) / E_{\text{miss, \text{True}}}$, is shown as a function of $E_{\text{miss, \text{True}}}$ in Fig. 14 for $W \rightarrow e\nu$ and $W \rightarrow \mu\nu$ MC events. The mean value of this ratio is expected to be zero if the reconstructed $E_{\text{miss}}$ has the correct scale. In Fig. 14, it can be seen that there is a displacement.
Fig. 7  Distribution of $E_{\text{T}}^{\text{miss}}$ computed with cells associated to electrons ($E_{\text{T}}^{\text{miss,e}}$) (top left), jets with $p_T > 20$ GeV ($E_{\text{T}}^{\text{miss,jets}}$) (top right), jets with $7 \text{ GeV} < p_T < 20$ GeV ($E_{\text{T}}^{\text{miss,softjets}}$) (bottom left) and from topoclusters outside reconstructed objects ($E_{\text{T}}^{\text{miss,CellOut}}$) (bottom right) for $Z \rightarrow ee$ data. The expectation from Monte Carlo simulation is superimposed and normalized to data, after each MC sample is weighted with its corresponding cross-section.

Fig. 8  Mean values of $E_{\text{T}}^{\text{miss}} \cdot A_Z$ as a function of $p^Z_T$ in $Z \rightarrow ee$ (left) and $Z \rightarrow \mu\mu$ (right) events.

from zero which varies with the true $E_{\text{T}}^{\text{miss}}$. The bias at low $E_{\text{T}}^{\text{miss, True}}$ values is about 5% and is due to the finite resolution of the $E_{\text{T}}^{\text{miss}}$ measurement. The reconstructed $E_{\text{T}}^{\text{miss}}$ is positive by definition, so the relative difference is positive when the $E_{\text{T}}^{\text{miss, True}}$ is small. The effect extends up to 40 GeV. The bias is in general larger for $W \rightarrow \mu\nu$ events than for $W \rightarrow e\nu$ events. Considering only events with $E_{\text{T}}^{\text{miss, True}} > 40$ GeV, the $E_{\text{T}}^{\text{miss}}$ linearity is better than 1% in $W \rightarrow e\nu$ events, while there is a non-linearity up to about 3% in $W \rightarrow \mu\nu$ events. This may be explained by an
Fig. 9 Mean value of \( E_{\text{miss}} \cdot A_Z \) as a function of \( p_{Z}^T \) requiring either zero jets with \( p_T > 20 \) GeV or at least 1 jet with \( p_T > 20 \) GeV in the event for \( Z \to ee \) (left) and \( Z \to \mu \mu \) (right) events.

Fig. 10 Mean value of \( E_{\text{miss}} \cdot A_Z \) as a function of \( p_{Z}^T \) in \( Z \to ee \) (left) and \( Z \to \mu \mu \) (right) for events with no jets with \( p_T > 7 \) GeV. The default \( E_{\text{miss}} \) is compared with \( E_{\text{miss}} \) calculated in the same way with the exception that the track-cluster matching algorithm is not used for the calculation of \( E_{\text{miss,CaloOut}} \).

underestimation of the \( E_{\text{miss,calo,\mu}} \) term, in which too few calorimeter cells are associated to the reconstructed muon.

6.4 \( E_{\text{miss}} \) resolution

A more quantitative evaluation of the \( E_{\text{miss}} \) performance can be obtained from a study of the \( (E_{\text{miss}}^x, E_{\text{miss}}^y) \) resolutions as a function of \( \sum E_{T} \). In \( Z \to \ell \ell \) events, as well as in minimum bias and QCD jet events, no genuine \( E_{\text{miss}}^\ell \) is expected, so the resolution of the two \( E_{\text{miss}} \) components is measured directly from reconstructed quantities, assuming that the true values of \( E_{x} \) and \( E_{y} \) are equal to zero. The resolution is estimated from the width of the combined distribution of \( E_{x} \) and \( E_{y} \) (denoted \( (E_{x}, E_{y}) \) distribution) in bins of \( \sum E_{T} \). The core of the distribution is fitted, for each \( \sum E_{T} \) bin, with a Gaussian over twice the expected resolution obtained from previous studies [17] and the fitted width, \( \sigma \), is examined as a function of \( \sum E_{T} \). The \( E_{\text{miss}} \) resolution follows an approximately stochastic behaviour as a function of \( \sum E_{T} \), which can be described with the function \( \sigma = k \cdot \sqrt{\sum E_{T}} \), but deviations from this simple law are expected in the low \( \sum E_{T} \) region due to noise and in the very large \( \sum E_{T} \) region due to the constant term.

Figure 15 (left) shows the resolution from data at \( \sqrt{s} = 7 \) TeV for \( Z \to \ell \ell \) events, minimum bias and di-jet events as a function of the total transverse energy in the event, obtained by summing the \( p_T \) of muons and the \( \sum E_{T} \) in calorimeters, calculated as described in Sect. 6.1. If the resolution is shown as a function of the \( \sum E_{T} \) in calorimeters, a difference between \( Z \to ee \) and \( Z \to \mu \mu \) events is observed due to the fact that \( \sum E_{T} \) includes electron momenta in \( Z \to ee \) events while muon momenta are not included in \( Z \to \mu \mu \) events.

The resolution of the two \( E_{\text{miss}} \) components is fitted with the simple function given above. The fits are acceptable and are of similar quality for all different channels studied. This allows to use the parameter \( k \) as an estimator for the resolution and to compare it in various physics channels in
Fig. 11 Distribution of $E_{\text{miss}}^T$ (top) and $\phi_{\text{miss}}$ (bottom) as measured in a data sample of $W \rightarrow e\nu$ (left) and $W \rightarrow \mu\nu$ (right) events. The expectation from Monte Carlo simulation is superimposed and normalized to data, after each MC sample is weighted with its corresponding cross-section. The sum of all backgrounds is shown in the lower plots.

Fig. 12 Distribution of $E_{\text{miss}}^T$ computed with cells from muons ($E_{\text{miss,calo}}^T$, left) and reconstructed muons ($E_{\text{miss,Muon}}^T$, right) for $W \rightarrow \mu\nu$ data. The expectation from Monte Carlo simulation is superimposed and normalized to data, after each MC sample is weighted with its corresponding cross-section.

There is a reasonable agreement in the $E_{\text{T}}^{\text{miss}}$ resolution in the different physics channels, as can be seen from the fit parameters $k$ reported in the figure. The $k$ parameter has fit values ranging from 0.42 GeV$^{1/2}$ for $Z \rightarrow \ell\ell$ events to 0.51 GeV$^{1/2}$ for di-jet events. The $E_{\text{T}}^{\text{miss}}$ resolution is better in $Z \rightarrow \ell\ell$ events because the lepton momenta are measured with better precision than jets.
Fig. 13 Distribution of $E_{\text{miss}}$ computed with cells associated to electrons ($E_{\text{miss},e}$) (top left), jets with $p_T > 20$ GeV ($E_{\text{miss,jets}}$) (top right), jets with $p_T < 20$ GeV ($E_{\text{miss,softjets}}$) (bottom left) and from topoclusters outside reconstructed objects ($E_{\text{miss,CellOut}}$) (bottom right) for data. The expectation from Monte Carlo simulation is superimposed and normalized to data, after each MC sample is weighted with its corresponding cross-section.

In Fig. 14 the $E_T^{\text{miss}}$ linearity in $W \rightarrow e\nu$ and $W \rightarrow \mu\nu$ MC events as a function of $E_T^{\text{miss, True}}$ is shown. In $W$ events the resolution of the two $E_T^{\text{miss}}$ components is estimated from the width of $(E_T^{\text{miss}} - E_T^{\text{miss, True}})^2$ in bins of $\sum E_T$, fitted with a Gaussian as explained above. There is a reasonable agreement in the $E_T^{\text{miss}}$ resolution in the different MC channels studied with the fitted value of $k$ ranging from $0.42 \text{ GeV}^{1/2}$ for $Z \rightarrow \ell \ell$ events to $0.50 \text{ GeV}^{1/2}$ for di-jet events. As observed for data, the $E_T^{\text{miss}}$ resolution is better in $Z \rightarrow \ell \ell$ events and slightly better in $W \rightarrow \ell\nu$ events, due to the presence of the leptons which are more precisely measured.

The resolution in MC minimum bias events is slightly worse than in data. This is probably due to imperfections of the modelling of soft particle activity in MC simulation, while there is a good data-MC agreement in the resolution for other channels.

7 Evaluation of the systematic uncertainty on the $E_T^{\text{miss}}$ scale

For any analysis using $E_T^{\text{miss}}$, it is necessary to be able to evaluate the systematic uncertainty on the $E_T^{\text{miss}}$ scale. The
$E_{\text{miss}}^{\text{T}}$, as defined in Sect. 5.3, is the sum of several terms corresponding to different types of reconstructed objects. The uncertainty on each individual term can be evaluated given the knowledge of the reconstructed objects [8, 23] that are used to build it and this uncertainty can be propagated to $E_{\text{miss}}^{\text{T}}$. The overall systematic uncertainty on the $E_{\text{miss}}^{\text{T}}$ scale is then calculated by combining the uncertainties on each term.

The relative impact of the uncertainty of the constituent terms on $E_{\text{miss}}^{\text{T}}$ differs from one analysis to another depending on the final state being studied. In particular, in events containing $W$ and $Z$ bosons decaying to leptons, uncertainties on the scale and resolution in the measurements of the charged leptons, together with uncertainties on the jet energy scale, need to be propagated to the systematic uncertainty estimate of $E_{\text{miss}}^{\text{T}}$. Another significant contribution to the $E_{\text{miss}}^{\text{T}}$ scale uncertainty in $W$ and $Z$ boson final states comes from the contribution of topoclusters outside reconstructed objects and from soft jets. In the next three subsections, two complementary methods for the evaluation of the systematic uncertainty on the $E_{\text{miss}}^{\text{T},\text{CellOut}}$ and the $E_{\text{miss}}^{\text{T},\text{softjets}}$ terms are described. Finally the overall $E_{\text{miss}}^{\text{T}}$ uncertainty for $W \to \ell \nu$ events is calculated.

### Table 1

<table>
<thead>
<tr>
<th>Variation</th>
<th>Description</th>
</tr>
</thead>
<tbody>
<tr>
<td>Dead Material</td>
<td>5% increase in the inner detector material</td>
</tr>
<tr>
<td></td>
<td>0.1 $X_0$ in front of the cryostat of the EM barrel calorimeter</td>
</tr>
<tr>
<td></td>
<td>0.05 $X_0$ between presampler and EM barrel calorimeter</td>
</tr>
<tr>
<td></td>
<td>0.1 $X_0$ in the cryostat after the EM barrel calorimeter</td>
</tr>
<tr>
<td></td>
<td>density of material in barrel-endcap transition of the EM calorimeter $\times 1.5$</td>
</tr>
<tr>
<td>FTFP_BERT</td>
<td>An alternative shower model for hadronic interaction in GEANT4</td>
</tr>
<tr>
<td>QGSP</td>
<td>An alternative shower model for hadronic interaction in GEANT4</td>
</tr>
<tr>
<td>PYTHIA Perugia 2010 tune</td>
<td>An alternative setting of the PYTHIA parameters with increased final state radiation and more soft particles</td>
</tr>
</tbody>
</table>

7.1 Evaluation of the systematic uncertainty on the $E_{\text{miss}}^{\text{T},\text{CellOut}}$ scale using Monte Carlo simulation

There are several possible sources of systematic uncertainty in the calculation of $E_{\text{miss}}^{\text{T},\text{CellOut}}$. These sources include inaccuracies in the description of the detector material, the choice of shower model and the model for the underlying event in the simulation. The systematic uncertainty due to each of these sources is estimated with dedicated MC simulations. The MC jet samples, generated with PYTHIA, are those used to assess the systematic uncertainty on the jet energy scale [21]. Table 1 lists the simulation samples considered, referred to in the following as “variations” with respect to the nominal sample.

The estimate of the uncertainty on $E_{\text{miss}}^{\text{T},\text{CellOut}}$ for a variation $i$ is determined by calculating the percentage difference between the mean value of this term for the nominal sample, labelled $\mu_0$, and that for the variation sample, labelled $\mu_i$. This approach assumes that the variations affect the total scale and none of the variations introduces a shape dependence in the $E_{\text{miss}}^{\text{T},\text{CellOut}}$ term, as verified in Ref. [24]. In order to cross-check for a possible dependence
on the event total transverse energy, the relative difference \( R_i = (\mu_i - \mu_0)/\mu_0 \) between different variations is computed in bins of \( \sum E_T \) for the jet samples. No significant dependence of \( R_i \) on \( \sum E_T \) is observed. A cross-check on the topology dependence is done using \( W \to \ell \nu \) samples simulated by introducing the variations \( i \). Table 2 shows the \( R_i \) values as computed in both the QCD jet samples and the \( W \to \ell \nu \) samples. The results are consistent, showing that the estimated uncertainty does not have a large dependence on the event topology.

A symmetric systematic uncertainty on the \( E_T^{\text{CellOut}} \) scale is obtained by summing in quadrature the estimated uncertainties averaged between simulated jet and \( W \) events. The total estimated uncertainty on the \( E_T^{\text{CellOut}} \) term is 2.6%.

### 7.2 Evaluation of the systematic uncertainty on the \( E_T^{\text{CellOut}} \) scale from the topocluster energy scale uncertainty

The uncertainty on the scale of the \( E_T^{\text{CellOut}} \) term, which is built from topoclusters with a correction based on tracks (see Sect. 5.3.1), can also be calculated from the topocluster energy scale uncertainties. These uncertainties can be estimated from comparisons between data and MC simulation using the \( E_T^{\text{CellOut}} \) response from single tracks, measured by summing the energies of all calorimeter clusters around a single isolated track [25]. The effects of these uncertainties on the \( E_T^{\text{CellOut}} \) term can be evaluated by varying the energy scale of topoclusters that contribute to the \( E_T^{\text{CellOut}} \) term in \( W \to e\nu \) MC samples, as was done in Ref. [8].

The shift in the topocluster energy scale is applied by multiplying the topocluster energy by the function:

\[
1 \pm a \times (1 + b/p_T),
\]

with \( a = 3(10)\% \) for \(|\eta| < (>) 3.2 \) and \( b = 1.2 \) GeV.

The \( a \) parameter in (8) addresses the uncertainty on the cluster energy scale, obtained by comparing the ratio of the cluster energy and the measured track momentum, \( E/p \), in data and MC simulation [25]. The value in the forward region, where tracks cannot be used to validate the energy scale, is estimated from the transverse momentum balance of one jet in the central region and one jet in the forward region in events with only two jets at high transverse momenta.

The \( b \) parameter in (8) addresses the possible change in the clustering efficiency and scale in a non-isolated environment. To go from the response for single isolated particles to the cluster energy scale, possible effects from the noise thresholds in the configuration with nearby particles are taken into account.

Because of threshold effects, more energy is clustered for nearby particles than for isolated ones. In an hypothetic worst case scenario, the environment is so busy that the clustering algorithm is forced to cluster all the deposited energy, with no bias due to the noise thresholds. Therefore, the maximal size of the noise threshold effect can be evaluated by comparing the ratio \( E_{\text{Cell}}/p \) of the total energy \( E_{\text{Cell}} \) deposited into all cells around an isolated track to the track momentum, to the ratio \( E/p \) of the clustered energy \( E \) to the track momentum, in data and MC simulation.

The fractional \( E_T^{\text{CellOut}} \) uncertainty is evaluated from:

\[
(\Delta E_T^{\text{CellOut}+} + \Delta E_T^{\text{CellOut}-})/(2 \times E_T^{\text{CellOut}}),
\]

where

\[
\Delta E_T^{\text{CellOut}+} = |E_T^{\text{CellOut}+} - E_T^{\text{CellOut}}|,
\]

\[
\Delta E_T^{\text{CellOut}-} = |E_T^{\text{CellOut}-} - E_T^{\text{CellOut}}|,
\]

with \( E_T^{\text{CellOut}+} \) and \( E_T^{\text{CellOut}-} \) obtained by shifting the topocluster energies up and down, respectively, using (8). The value of the fractional \( E_T^{\text{CellOut}} \) uncertainty is found to be approximately 13%, decreasing slightly with increasing \( \sum E_T^{\text{Cell}} \). This uncertainty is much larger than the uncertainty due to the detector description estimated from the first three lines of Table 2. The main reason is that the values of \( a \) and \( b \) which enter into (8) are conservative, to include the effects described above. In particular the cluster energy uncertainty in the forward region is conservatively estimated, since the uncertainty cannot be evaluated using tracks. Moreover, the procedure does not take into account the fact that when the clusters are shifted up in \( p_T \), some of them can form jets above threshold and they are therefore included in the soft jet term in \( E_T^{\text{CellOut}} \). These clusters should be removed from the \( E_T^{\text{CellOut}} \), they are in fact kept and this increases the uncertainty. It should also be noted that in the calculation of \( E_T^{\text{CellOut}} \) the track momentum is

---

### Table 2

Systematic uncertainties \( (R_i) \) on \( E_T^{\text{CellOut}} \) associated with variations in the dead material (all the variations listed in Table 1 are applied at the same time), in the calorimeter shower modelling (FTFP_BERT, QGSP) and in the event generator settings (PYTHIA Perugia 2010 tune)

<table>
<thead>
<tr>
<th>Variation</th>
<th>jet events</th>
<th>( W ) production</th>
</tr>
</thead>
<tbody>
<tr>
<td>Dead Material</td>
<td>((-0.5 \pm 0.1)%)</td>
<td>((-0.6 \pm 0.2)%)</td>
</tr>
<tr>
<td>FTFP_BERT</td>
<td>((0.1 \pm 0.4)%)</td>
<td>((0.5 \pm 0.2)%)</td>
</tr>
<tr>
<td>QGSP</td>
<td>((-1.6 \pm 0.4)%)</td>
<td>((-2.2 \pm 0.2)%)</td>
</tr>
<tr>
<td>PYTHIA Perugia 2010 tune</td>
<td>((-1.7 \pm 0.1)%)</td>
<td>((-1.5 \pm 0.2)%)</td>
</tr>
</tbody>
</table>

---

\(^6\)In this uncertainty evaluation using MC simulation, the uncertainty on the absolute electromagnetic energy scale in the calorimeters should also be taken into account. For the bulk of the LAr barrel electromagnetic calorimeter a 1.5% uncertainty is found on the cell energy measurement, increasing to 5% for the presampler and 3% for the tile calorimeter [25].
used instead of the topocluster energy when there is a track-
topocluster matching (see Sect. 5.3.1). This would result in a
reduced uncertainty due to the more precise measurement of
the track momentum, which is not taken into account here.
Further study is expected to provide a reduction in this un-
certainty in future, by considering the described effects in
detail.

To give an estimate of the $E^\text{miss,CellOut}$ systematic un-
certainty, the calorimeter contribution can be taken from
Sect. 7.2, and the uncertainty from the event generator set-
tings from Sect. 7.1 (PYTHIA Perugia 2010 tune). This
results in a total systematic uncertainty on the scale of
$E^\text{miss,CellOut}$ of about 13%, which slightly decreases when
$\sum E_T^{\text{CellOut}}$ increases.

7.3 Evaluation of the systematic uncertainty on the $E_T^{\text{miss,softjets}}$ scale

The same procedure described in the previous sections is
used to assess the systematic uncertainty on the $E_T^{\text{miss}}$ term
calculated from soft jets (see Sect. 5.1).

Using the MC approach described in Sect. 7.1, it is found
that the uncertainty on $E_T^{\text{miss,softjets}}$ does not exhibit a large
dependence on the event $\sum E_T$, as was also found for the un-
certainty on the $E_T^{\text{miss,CellOut}}$ scale. The results are consistent
between the QCD jet samples and the $W$ samples, as can be
seen from Table 3 which gives the systematic uncertainties
$R_i$ as computed in jet samples and in $W \to \ell\nu$ samples.
A total, symmetric, systematic uncertainty of about 3.3%
on the $E_T^{\text{miss,softjets}}$ term is obtained by combining the results
in Table 3, as was done in Sect. 7.1. With the same data-
driven approach utilising the uncertainty on the topocluster
energy scale described in Sect. 7.2, the systematic uncer-
tainty on $E_T^{\text{miss,softjets}}$ is evaluated to be about 10%.

As for $E_T^{\text{miss,CellOut}}$, the uncertainty on the $E_T^{\text{miss,softjets}}$
scale found by shifting the topocluster energies is larger
than the uncertainty estimated from MC simulation. To give
an estimate of the systematic uncertainty on $E_T^{\text{miss,softjets}}$, the
contribution from the calorimeter response can be taken
from the data-driven evaluation and the contribution from the
event generator settings from Table 3. This results in an
overall systematic uncertainty of about 10% on $E_T^{\text{miss,softjets}}$
slightly increasing as $\sum E_T$ increases.

7.4 Evaluation of the overall systematic uncertainty
on the $E_T^{\text{miss}}$ scale in $W \to e\nu$ and $W \to \mu\nu$ events

Using as inputs the systematic uncertainties on the differ-
ent reconstructed objects [8, 21] and on $E_T^{\text{miss,CellOut}}$ and
$E_T^{\text{miss,softjets}}$ evaluated in the previous sections, the over-
al $E_T^{\text{miss}}$ systematic uncertainty in $W \to e\nu$ and $W \to \mu\nu$ events is estimated. The same method can be applied
to any final state event topology. Figure 16 shows, for both
$W \to e\nu$ and $W \to \mu\nu$ events, the systematic uncertainties on each of the terms $E_T^{\text{miss,e}} (E_T^{\text{miss,\mu}})$, $E_T^{\text{miss,jets}}$
$E_T^{\text{miss,softjets}}$ and $E_T^{\text{miss,CellOut}}$ as a function of their individu-
all contribution to $\sum E_T$ labelled $\sum E_T^{\text{term}}$. All the uncer-
tainties are calculated with the formulae in (9) and (10). In
the same figure the uncertainty on $E_T^{\text{miss}}$ due to the uncer-
tainties on the different terms is also shown as a function
of the total $\sum E_T$, together with the overall uncertainty on
$E_T^{\text{miss}}$, obtained by combining the partial terms. The uncer-
tainties on $E_T^{\text{miss,softjets}}$ and $E_T^{\text{miss,CellOut}}$ are considered to be
fully correlated. In $W \to e\nu$ and $W \to \mu\nu$ events, se-
lected as described in Sect. 3.3, the overall uncertainty on
the $E_T^{\text{miss}}$ scale increases with $\sum E_T$ from ~1% to ~7%. It
is estimated to be, on average, about 2.6% for both channels.
The $E_T^{\text{miss}}$ scale uncertainty depends on the event topol-
ogy because the contribution of a given $E_T^{\text{miss}}$ term can vary
for different final states.

8 Determination of the $E_T^{\text{miss}}$ scale from $W \to \ell\nu$ events

The determination of the absolute $E_T^{\text{miss}}$ scale is im-
portant in a range of analyses involving $E_T^{\text{miss}}$ measurements,
ranging from precision measurements to searches for new
physics.

In this section two complementary methods to determine
the absolute scale of $E_T^{\text{miss}}$ using $W \to \ell\nu$ events are de-
dscribed. The first method uses a fit to the distribution of
the transverse mass, $m_T$, of the lepton-$E_T^{\text{miss}}$ system, and
is sensitive both to the scale and the resolution of $E_T^{\text{miss}}$.
The second method uses the interdependence of the neu-
trino and lepton momenta in the $W \to e\nu$ channel, and the
$E_T^{\text{miss}}$ scale is determined as a function of the reconstructed
electron transverse momentum. Both methods allow checks
on the agreement between data and MC simulation for the
$E_T^{\text{miss}}$ scale.

8.1 Reconstructed transverse mass method

The method described in this section uses the shape of the
$m_T$ distribution and is sensitive to both the $E_T^{\text{miss}}$ resolu-
tion and scale. The lepton transverse momentum, $p_T^{\ell}$, and
Fig. 16 Fractional systematic uncertainty (calculated as in (9) and (10)) on different $E_T^{\text{miss}}$ terms as a function of respective $\sum E_T^{\text{term}}$ (left) and contributions of different term uncertainties on $E_T^{\text{miss}}$ uncertainty as a function of $\sum E_T$ (right) in MC $W \rightarrow e\nu$ events (top) and $W \rightarrow \mu\nu$ events (bottom). The overall systematic uncertainty on the $E_T^{\text{miss}}$ scale, obtained combining the various contributions is shown in the right plots (filled circles). The uncertainties on $E_T^{\text{miss}}$, softjets and $E_T^{\text{miss}, \text{CellOut}}$ are considered to be fully correlated

\[
m_T = \sqrt{2p_T^{\ell}E_T^{\text{miss}}(1 - \cos \phi)},
\]

where $\phi$ is the azimuthal angle between the lepton momentum and $E_T^{\text{miss}}$ directions. The true $m_T$ is reconstructed from the simulation under the hypothesis that $E_T^{\text{miss}}$ is entirely due to the neutrino momentum, $p_T^{\nu}$. Template histograms of the $m_T$ distributions are generated by convoluting the true transverse mass distribution with a Gaussian function:

\[
x^{(y)} \rightarrow \alpha E_T^{\text{miss}, \text{True}} \ast \text{Gauss}(0, k \cdot \sqrt{\Sigma E_T}),
\]

where the parameters $\alpha$ and $k$ are the $E_T^{\text{miss}}$ scale and resolution respectively.

The $\alpha$ and $k$ parameters are determined through a fit of the $m_T$ distributions to data using a linear combination of signal and background $m_T$ distributions obtained from simulation. All the backgrounds, with the exception of the jet background, are evaluated from the same MC samples used in Sect. 6.3 and the normalization is fixed according to their cross-sections. The shape of the jet background is also evaluated from MC simulation and its normalization is obtained from the fit, in addition to $\alpha$ and $k$.

To select $W \rightarrow \mu\nu$ events, the same criteria as described in Sect. 3.3 are used, with the exception that no cut on $E_T^{\text{miss}}$ is applied and a looser cut, $m_T > 30$ GeV, is applied in order that the background normalization can be fitted. The $\alpha$ and $k$ parameters obtained from the fit are shown in Table 4, together with the numbers of events for the signal and backgrounds and the $\chi^2$/ndof of the fit. In the table, instead of the values of $\alpha$, the values of $\alpha - 1 = (E_T^{\text{miss}, \text{True}} - E_T^{\text{miss}, \text{True}})/E_T^{\text{miss}, \text{True}}$ are reported, in order to compare with the result in Sects. 6.3.1 and 8.2. The results for the $\alpha$ and $k$ parameters using the $m_T$ distribution of the simulated signal are also shown in Table 4, and they are in good agreement with the results from data. The result of the fit to data and MC simulation is shown in Fig. 17.

To select $W \rightarrow e\nu$ events, the selection described in Sect. 3.3 is used with the addition of tighter cuts. A cut $E_T^{\text{miss}} > 36$ GeV is applied to exclude the region where the $E_T^{\text{miss}}$ response is not linear (see Fig. 14). A cut $m_T > 40$
The results obtained with this method are compatible, at
the few percent level, with the results shown in Fig. 14 and
Fig. 15, which were derived using only simulation. From
those figures, for the $W \rightarrow \mu \nu$ channel $\alpha - 1$ has values
up to 3% and the resolution is $0.47 \sqrt{\sum E_T}$; for the $W \rightarrow e\nu$
channel $\alpha - 1$ is close to zero for high $E_T^{\text{miss}}$ values and
the resolution is $0.47 \sqrt{\sum E_T}$.

The uncertainty due to background subtraction is already
included in the uncertainty reported in Table 4. The system-
atic uncertainty on $\alpha - 1$ is determined to be about 1% for
each channel, by checking the stability of the results us-
ing different cuts on $E_T^{\text{miss}}$ and using a different generator,
MC@NLO. In summary, with this method the $E_T^{\text{miss}}$
absolute scale is determined from $W \rightarrow \ell \nu$ events, in a data
sample corresponding to an integrated luminosity of about
36 pb$^{-1}$, with an uncertainty (adding the uncertainties re-
ported in Table 4 with the systematic uncertainty) of about
1.5% and about 2% for the $W \rightarrow \mu \nu$ and $W \rightarrow e\nu$ decay
channels, respectively.

### 8.2 Method based on the correlation between electron
and neutrino transverse momenta in $W \rightarrow e\nu$

In this section the correlation between the transverse mo-
menta of charged and neutral leptons from $W$ boson de-
cays is used to determine the $E_T^{\text{miss}}$ scale. The mean mea-
sured $E_T^{\text{miss}}$ is compared to the mean true $E_T^{\text{miss}}$ from sig-
nal MC events. The relative bias in the reconstructed $E_T^{\text{miss}}$,
$\langle (E_T^{\text{miss}}) - (E_T^{\text{miss, True}}) \rangle / \langle (E_T^{\text{miss, True}}) \rangle$, is studied as a function of $p_T^{\ell}$ because the MC simulation of the electron response is
more accurate than that for hadrons.

This method is shown for $W \rightarrow e\nu$ events by applying
selection criteria similar to the ones described in Sect. 3.3,
but with isolation requirements both on the electron track
and calorimeter signal. The $E_T^{\text{miss}}$ is required to be greater
than 20 GeV and no cut is applied on $m_T$.

MC samples are generated with MC@NLO [15]. A next-
to-leading-order (NLO) generator is used for this study be-
cause in this approach the $E_T^{\text{miss}}$ scale is validated on the

### Table 4 Results of $m_T$ fit in $W \rightarrow \ell \nu$ events. The second and third
columns show the scale and resolution parameters obtained.
The numbers of events for the signal, the electroweak and QCD backgrounds
obtained from the fit are shown in the fourth, fifth and sixth columns

<table>
<thead>
<tr>
<th>Channel</th>
<th>$\alpha - 1$ (%)</th>
<th>$k$</th>
<th>Signal</th>
<th>EW (fixed)</th>
<th>QCD</th>
<th>$\chi^2/\text{ndof}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$W \rightarrow \mu \nu$ data</td>
<td>5.1 ± 0.8</td>
<td>0.52 ± 0.01</td>
<td>164920 ± 840</td>
<td>14760</td>
<td>24870 ± 840</td>
<td>68/87</td>
</tr>
<tr>
<td>$W \rightarrow \mu \nu$ MC</td>
<td>5.5 ± 0.8</td>
<td>0.50 ± 0.01</td>
<td></td>
<td></td>
<td></td>
<td>70/78</td>
</tr>
<tr>
<td>$W \rightarrow e\nu$ data</td>
<td>-0.8 ± 1.6</td>
<td>0.49 ± 0.01</td>
<td>75660 ± 180</td>
<td>1210</td>
<td>980 ± 180</td>
<td>54/75</td>
</tr>
<tr>
<td>$W \rightarrow e\nu$ MC</td>
<td>1.8 ± 1.7</td>
<td>0.50 ± 0.01</td>
<td></td>
<td></td>
<td></td>
<td>38/54</td>
</tr>
</tbody>
</table>
basis of the known decay properties of the W boson. The correlation between $p_T^W$ and $p_T^\tau$ is important for this study, and is poorly described by leading-order generators such as PYTHIA, whereas it is much improved in MC@NLO. The MC events are weighted such that the true W boson transverse momentum, $p_T^W$, and pseudorapidity $\eta^W$ agree with that generated using the RESBOS [26] generator which is more accurate in describing $p_T^W$ at low values. Finally, an additional smearing is applied to the reconstructed electron momentum in the MC samples, to match the electron resolution measured in data, and the correction is propagated to $E_T^{miss}$.

A data-driven technique is used to estimate the impact of jet background, which is small (see Fig. 18 left) and concentrated at low $p_T^\tau$. $W \rightarrow \tau \nu$ events, where the $\tau$ decays to an electron, are the second largest background, but the impact on the mean value of $E_T^{miss}$ is found to be negligible.

The distribution of $p_T^\tau$ is shown in Fig. 18. The distribution from data after event selection is fitted by varying the normalization of signal MC and QCD background distributions. A satisfactory description of data is achieved except for the first bin, which is excluded from the fit. For each $p_T^\tau$ bin, the corrected distribution of $E_T^{miss}$ is obtained by subtracting that of the background sample (after normalizing it according to the fit) from the data distribution. The largest impact of background corresponds to $p_T^\tau = 20$ GeV, with an effect of about 2 GeV on the mean value of $E_T^{miss}$; the effect decreases quickly to 0.2 GeV at $p_T^\tau = 30$ GeV.

Since a cut on $E_T^{miss}$ is used for the event selection and the $E_T^{miss}$ resolution is finite, the results are biased. To correct for the bias in signal MC events the requirement of correlation between $p_T^\nu$ and $p_T^\tau$ is important for this study, and the bias in signal MC events alone, the resulting average bias in $E_T^{miss}$ is 2.9 $\pm$ 0.1%. The method relies on simulation to derive the correlation between $E_T^{miss, True}$ and $p_T^\tau$, so it can be sensitive to details of the simulation. In particular, the jet factorization and renormalization scales, as well as the choice of PDF, can affect the results, but all these also change the $p_T^W$ distribution. Therefore the shape of the $p_T^W$ distribution was distorted by $\pm$10%, justified by the comparison of a recent measurement of the $p_T^W$ distribution [27] with RESBOS predictions, and the relative bias was calculated again.

A systematic uncertainty on the relative $E_T^{miss}$ scale bias of $\pm 2\%$ is evaluated. The results for the average $E_T^{miss}$ scale are summarized in Table 5. These results agree within errors with the values of $\alpha - 1$ shown in Table 4.

9 Conclusion

The missing transverse momentum ($E_T^{miss}$) has been measured in minimum bias, di-jet, $Z \rightarrow \ell\ell$ and $W \rightarrow \ell\nu$ events in 7 TeV $pp$ collisions recorded with the ATLAS detector in 2010.

The value of $E_T^{miss}$ is reconstructed from calorimeter cells in topological clusters, with the exception of electrons and photons for which a different clustering algorithm is used, and from reconstructed muons. The cells are calibrated according to their parent particle type. The scheme yielding the best performance is evaluated to be that in which electrons are calibrated with the default electron calibration and photons are used at the EM scale, the $\tau$-jets and jets are calibrated with the local hadronic calibration (LCW), the jets...
Table 5  Average relative $E^\text{miss}_T$ scale bias obtained from data and MC simulation from the electron-neutrino correlation method. The statistical and the systematic uncertainties are given for data

<table>
<thead>
<tr>
<th>source</th>
<th>scale bias (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>data</td>
<td>2.0 ± 0.1 ± 2.0</td>
</tr>
<tr>
<td>MC</td>
<td>2.9 ± 0.1</td>
</tr>
</tbody>
</table>

with $p_T$ greater than 20 GeV are scaled to the jet energy scale, and the contribution from topoclusters not associated to high-$p_T$ objects is calculated with LCW calibration combined with tracking information.

Monte Carlo simulation is found to describe the data in general rather well. No large tails are observed in combined with tracking information.

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University at Albany, Albany NY, United States of America
2Department of Physics, University of Alberta, Edmonton AB, Canada
3(a)Department of Physics, Ankara University, Ankara; (b)Department of Physics, Dumlupinar University, Kutahya;
(c)Department of Physics, Gazi University, Ankara; (d)Division of Physics, TOBB University of Economics and
Technology, Ankara; (e)Turkish Atomic Energy Authority, Ankara, Turkey
4LAPP, CNRS/IN2P3 and Université de Savoie, Annecy-le-Vieux, France
5High Energy Physics Division, Argonne National Laboratory, Argonne IL, United States of America
6Department of Physics, University of Arizona, Tucson AZ, United States of America
7Department of Physics, The University of Texas at Arlington, Arlington TX, United States of America
8Physics Department, University of Athens, Athens, Greece
9Physics Department, National Technical University of Athens, Zografou, Greece
10Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
11Institut de Física d’Altes Energies and Universitat Autònoma de Barcelona and ICREA, Barcelona, Spain
12(a)Institute of Physics, University of Belgrade, Belgrade; (b)Vinca Institute of Nuclear Sciences, Belgrade, Serbia
13Department for Physics and Technology, University of Bergen, Bergen, Norway
14Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley CA, United States of
America
15Department of Physics, Humboldt University, Berlin, Germany
16Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern,
Switzerland
17School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
18(a)Department of Physics, Bogazici University, Istanbul; (b)Division of Physics, Dogus University, Istanbul;
(c)Department of Physics Engineering, Gaziantep University, Gaziantep; (d)Department of Physics, Istanbul Technical
University, Istanbul, Turkey
19(a)INFN Sezione di Bologna; (b)Dipartimento di Fisica, Università di Bologna, Bologna, Italy
20Physikalisches Institut, University of Bonn, Bonn, Germany
21Department of Physics, Boston University, Boston MA, United States of America
22Department of Physics, Brandeis University, Waltham MA, United States of America
23(a)Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro; (b)Federal University of Juiz de Fora (UFJF),
Juiz de Fora; (c)Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei; (d)Instituto de Fisica, Universidade de
Sao Paulo, Sao Paulo, Brazil
24Physics Department, Brookhaven National Laboratory, Upton NY, United States of America
25(a)National Institute of Physics and Nuclear Engineering, Bucharest; (b)University Politehnica Bucharest, Bucharest;
(c)West University in Timisoara, Timisoara, Romania
26Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
27Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
28Department of Physics, Carleton University, Ottawa ON, Canada
29CERN, Geneva, Switzerland
30Enrico Fermi Institute, University of Chicago, Chicago IL, United States of America
31(a)Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; (b)Departamento de Física, Universidad
Técnica Federico Santa María, Valparaíso, Chile
32(a)Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; (b)Department of Modern Physics,
University of Science and Technology of China, Anhui; (c)Department of Physics, Nanjing University, Jiangsu; (d)High
Energy Physics Group, Shandong University, Shandong, China
33Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Aubiere
Cedex, France
34Nevis Laboratory, Columbia University, Irvington NY, United States of America
35Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark
36(a)INFN Gruppo Collegato di Cosenza; (b)Dipartimento di Fisica, Università della Calabria, Arcavata di Rende, Italy
37Faculty of Physics and Applied Computer Science, AGH-University of Science and Technology, Krakow, Poland
38The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland
39Physics Department, Southern Methodist University, Dallas TX, United States of America
40Physics Department, University of Texas at Dallas, Richardson TX, United States of America
Also at Università di Napoli Parthenope, Napoli, Italy
Also at Institute of Particle Physics (IPP), Canada
Also at Department of Physics, Middle East Technical University, Ankara, Turkey
Also at Louisiana Tech University, Ruston LA, United States of America
Also at Group of Particle Physics, University of Montreal, Montreal QC, Canada
Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
Also at Department of Physics, Middle East Technical University, Ankara, Turkey
Also at Louisiana Tech University, Ruston LA, United States of America
Also at Group of Particle Physics, University of Montreal, Montreal QC, Canada
Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany
Also at Manhattan College, New York NY, United States of America
Also at School of Physics and Engineering, Sun Yat-sen University, Guanzhou, China
Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan
Also at High Energy Physics Group, Shandong University, Shandong, China
Also at Section de Physique, Université de Genève, Geneva, Switzerland
Also at Departamento de Fisica, Universidade de Minho, Braga, Portugal
Also at Department of Physics and Astronomy, University of South Carolina, Columbia SC, United States of America
Also at FKF Research Institute for Particle and Nuclear Physics, Budapest, Hungary
Also at California Institute of Technology, Pasadena CA, United States of America
Also at Institute of Physics, Jagiellonian University, Krakow, Poland
Also at Department of Physics, Oxford University, Oxford, United Kingdom
Also at Institute of Physics, Academia Sinica, Taipei, Taiwan
Also at Department of Physics, The University of Michigan, Ann Arbor MI, United States of America
Also at DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat a l’Energie Atomique), Gif-sur-Yvette, France
Also at Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France
Also at Department of Physics, Nanjing University, Jiangsu, China
*Deceased