Measurement of the jet radius and transverse momentum dependence of inclusive jet suppression in lead-lead collisions at $\sqrt{s_{NN}} = 2.78$ TeV with the ATLAS detector


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Measurement of the jet radius and transverse momentum dependence of inclusive jet suppression in lead–lead collisions at $\sqrt{s_{NN}} = 2.76$ TeV with the ATLAS detector

ATLAS Collaboration

A R T I C L E  C O N T E N T S

1. Introduction

Collisions of lead ions at the LHC are expected to create strongly interacting matter at the highest temperatures ever produced in the laboratory [1]. This matter may be deconfined with a high density of unscreened colour charges. High transverse momentum ($p_T$) quarks and gluons generated by hard-scattering processes have long been considered an important tool for probing the properties of the matter created in ultra-relativistic nuclear collisions. The energy loss of the partons propagating through the matter may provide direct sensitivity to the colour charge density and to the transport properties of the matter [2–4]. Indirect observations of substantial parton energy loss or “jet quenching” via suppressed single high-$p_T$ hadron yields [5–8] and disappearance of the dijet contribution to di-hadron correlations [9,10] have contributed to the conclusion that Au + Au collisions at RHIC produce a quark–gluon plasma [11,12]. Observations of highly asymmetric dijets in central Pb + Pb collisions at the LHC [13–15] can be understood in the context of “differential” jet quenching, where one parton produced from an initial hard-scattering loses significantly more energy than the other, possibly as a result of different path lengths of the partons in the matter [16]. However, the asymmetry is not sensitive to situations where the two jets in a dijet pair lose comparable amounts of energy, so other measurements are required to probe “inclusive” jet quenching.

The inclusive, per-event jet production rate provides such a measurement. Energy loss of the parent partons in the created matter may reduce or “suppress” the rate for producing jets at a given transverse momentum. Such energy loss is expected to increase with medium temperature and with increasing path length of the parton in the medium [17]. As a result, there should be more suppression in central Pb + Pb collisions, which have nearly complete overlap between the incident nuclei, and little or no suppression in peripheral collisions where the nuclei barely overlap. In the absence of energy loss, the jet production rate is expected to vary with Pb + Pb collision centrality approximately in proportion to $N_{coll}$, the number of nucleon–nucleon collisions that take place during a single Pb + Pb collision. The jet suppression may be quantified using the central-to-peripheral ratio, $R_{CP}$, the ratio of the per-event jet yields divided by the number of nucleon–nucleon collisions in a given centrality bin to the same quantity in a peripheral centrality bin. The quantity, $R_{CP}$, has the advantage that potentially large systematic uncertainties, especially those arising from systematic errors on the jet energy scale, largely cancel when evaluating the ratios of jet spectra within the same data set. The variation of the suppression with jet transverse momentum and with collision centrality will depend both on the energy loss mechanism and on the experimental definition of the jet. In the case of radiative energy loss, jet energies can be reduced by greater “out-of-cone” radiation, which should be more severe for
smaller jet radii [18–20]. Naively, collisional energy loss would result in a suppression that is independent of radius. However recent calculations suggest that collisional processes can also contribute to jet broadening [21]. A measurement of the radius dependence of jet suppression could further clarify the roles of radiative and collisional energy loss in jet quenching.

This Letter presents measurements of the inclusive jet $R_{CP}$ in Pb + Pb collisions at a nucleon–nucleon centre-of-mass energy of $\sqrt{s_{NN}} = 2.76$ TeV using data collected during 2010 corresponding to an integrated luminosity of approximately 7 $\mu$b$^{-1}$. Results are presented for jets reconstructed from energy deposits measured in the ATLAS calorimeters using the anti-$k_t$ jet-finding algorithm [22]. The anti-$k_t$ reconstruction was performed separately for four different values of the anti-$k_t$ distance parameter, $R$, that specifies the nominal radius of the reconstructed jets, $R = 0.2, 0.3, 0.4$ and 0.5. For the remainder of the Letter the term “radius” will refer to the distance parameter, $R$. The jet energy is functionally defined to be the total energy within the jet clustering algorithm above an uncorrelated underlying event. This jet definition may include medium response with is correlated with the jet. The underlying event contribution to each jet was subtracted on a per-jet basis, and the $R_{CP}$ values were calculated after unfolding the jet spectra for distortions due to intrinsic jet resolution and underlying event fluctuations.

2. Experimental setup and trigger

The measurements presented here were performed using the ATLAS calorimeter, inner detector, trigger, and data acquisition systems [23]. The ATLAS calorimeter system consists of a liquid argon (LAr) electromagnetic (EM) calorimeter covering $|\eta| < 3.2$, a steel-scintillator sampling hadronic calorimeter covering $|\eta| < 1.7$, a LAr hadronic calorimeter covering $1.5 < |\eta| < 3.2$, and two LAr electromagnetic and hadronic forward calorimeters (FCal) covering $3.2 < |\eta| < 4.9$. The hadronic calorimeter granularities or cell sizes in $\Delta R \times \Delta \phi$ are $0.1 \times 0.1$ for $|\eta| < 2.5$ and $0.2 \times 0.2$ for $2.5 < |\eta| < 4.9$. The EM calorimeters are longitudinally segmented into three compartments with an additional pre-sampler layer. The EM calorimeter has a granularity that varies with layer and pseudorapidity, but which is generally much finer than that of the hadronic calorimeter. The middle sampling layer, which typically has the largest energy deposit in EM showers, has a $\Delta R \times \Delta \phi$ granularity of $0.025 \times 0.025$ over $|\eta| < 2.5$.

Charged particles associated with the calorimeter jets were measured over the pseudorapidity interval $|\eta| < 2.5$ using the inner detector [24]. The inner detector is composed of silicon pixel detectors in the innermost layers, followed by silicon microstrip detectors and a straw-tube tracker, all immersed in a 2 T axial magnetic field provided by a solenoid. Minimum bias Pb + Pb collisions were identified using measurements from the zero-degree calorimeters (ZDCs) and the minimum-bias trigger scintillator (MBTS) counters. The ZDCs are located symmetrically at $z = \pm 140$ m and cover $|\eta| > 8.3$. In Pb + Pb collisions the ZDCs primarily measure “spectator” neutrons – neutrons from the incident nuclei that do not interact hadronically. The MBTS measures charged particles over $2.1 < |\eta| < 3.9$ using two sets of counters.

1 ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z-axis along the beam pipe. The x-axis points from the IP to the centre of the LHC ring, and the y-axis points upward. Cylindrical coordinates $(\rho, \phi)$ are used in the transverse plane, $\phi$ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle $\theta$ as $\eta = -\ln(\tan(\theta/2))$.

2 An exception is the third (outermost) sampling layer, which has a segmentation of $0.2 \times 0.1$ up to $|\eta| = 1.7$.

3. Event selection and centrality definition

In the offline analysis, Pb + Pb collisions were required to have a primary vertex reconstructed from charged particle tracks with $p_T > 500$ MeV. The tracks were reconstructed from hits in the inner detector using the standard ATLAS track reconstruction algorithm [25] with settings optimized for the high hit density in heavy ion collisions [26]. Additional requirements of a ZDC coincidence, at least one hit in each MBTS counter, and a difference in times measured by the two sides of the MBTS detector of less than 3 ns were imposed. The combination of the ZDC and MBTS conditions and the primary vertex requirement efficiently eliminates both beam–gas interactions and photo-nuclear events [27]. These event selections yielded a total of 51 million minimum-bias Pb + Pb events. Previous studies [26] indicate that the combination of trigger and offline requirements select minimum-bias hadronic Pb + Pb collisions with an efficiency of 98 ± 2%.

The centrality of Pb + Pb collisions was characterized by $N_{FCal}$, the total transverse energy measured in the forward calorimeters. The distribution of $N_{FCal}$ was divided into intervals corresponding to successive 10% percentiles of the full centrality distribution after accounting for the missing 2% most peripheral events. A standard Glauber Monte Carlo analysis [28,29] was used to estimate the average number of participating nucleons, $\langle N_{part} \rangle$, and the average number of nucleon–nucleon collisions, $\langle N_{coll} \rangle$, for Pb + Pb collisions in each of the centrality bins. The results are shown in Table 1. The $R_{CP}$ measurements presented here use the 60–80% centrality bin as a common peripheral reference. The $R_{CP}$ calculation requires the ratio, $R_{CP} \equiv \langle N_{coll} \rangle / \langle N_{FCal} \rangle$, where $\langle N_{FCal} \rangle$ is the average number of collisions in the 60–80% centrality bin. The $R_{CP}$ uncertainties have been evaluated by changing the estimate of $R_{CP}$ due to variations of the minimum-bias trigger efficiency, parameters of the Glauber calculation, and parameters in the modelling of the $N_{FCal}$ distribution [26]. The $R_{CP}$ values and uncertainties are also reported in Table 1.

4. Monte Carlo samples

Three Monte Carlo (MC) samples [30] were used for the analysis in this Letter. A total of 1 million simulated minimum-bias Pb + Pb events were produced using version 1.38b of the HIJING event generator [31]. HIJING was run with default parameters except for the disabling of jet quenching. To simulate the effects of elliptic flow in Pb + Pb collisions, a parameterized centrality-,

<table>
<thead>
<tr>
<th>Centrality</th>
<th>$\langle N_{part} \rangle$</th>
<th>$\langle N_{coll} \rangle$</th>
<th>$R_{CP}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>0–10%</td>
<td>356 ± 2</td>
<td>1500 ± 115</td>
<td>57 ± 6</td>
</tr>
<tr>
<td>10–20%</td>
<td>261 ± 4</td>
<td>923 ± 68</td>
<td>35 ± 4</td>
</tr>
<tr>
<td>20–30%</td>
<td>186 ± 4</td>
<td>559 ± 41</td>
<td>21 ± 2</td>
</tr>
<tr>
<td>30–40%</td>
<td>129 ± 4</td>
<td>322 ± 24</td>
<td>12 ± 1</td>
</tr>
<tr>
<td>40–50%</td>
<td>86 ± 4</td>
<td>173 ± 14</td>
<td>6.5 ± 0.04</td>
</tr>
<tr>
<td>50–60%</td>
<td>53 ± 3</td>
<td>85 ± 8</td>
<td>3.2 ± 0.01</td>
</tr>
<tr>
<td>60–80%</td>
<td>23 ± 2</td>
<td>27 ± 4</td>
<td>1.0</td>
</tr>
</tbody>
</table>
pt-dependent $\cos 2\phi$ modulation based on previous ATLAS measurements [26] was imposed on the particles after generation [32]. The detector response to the resulting HIJING events was evaluated using GEANT4 [33] configured with geometry and digitization parameters matching those of the 2010 Pb + Pb run.

An “MC overlay” data set, intended specifically for evaluating jet performance, was obtained by overlaying GEANT4-simulated $\sqrt{s_{NN}} = 2.76$ TeV pp hard-scattering events on the HIJING events described above. The pp events were obtained from the ATLAS MC09 tune [34] of the PYTHIA event generator [35]. One million PYTHIA hard-scattering events were generated for each of five intervals of $p_T$, the transverse momentum of outgoing partons in the $2 \rightarrow 2$ hard-scattering, with boundaries 35, 70, 140, 280 and 560 GeV. The pp events for each $p_T$ interval were overlaid on the same sample of HIJING events.

A smaller sample of “data overlay” events was produced by overlaying 150k GEANT4-simulated PYTHIA pp events onto 150k minimum-bias Pb + Pb data events recorded during the 2011 LHC Pb + Pb run. Due to the different detector conditions in the 2010 and 2011 runs, the data overlay events cannot provide the corrections required for this analysis. However, they provide a valuable test of the accuracy of HIJING in describing the underlying event.

5. Jet reconstruction

Calorimeter jets were reconstructed from $\Delta \eta \times \Delta \phi = 0.1 \times 0.1$ towers using the anti-$k_T$ algorithm [22] in four-vector recombination mode with anti-$k_T$ distance parameters $R = 0.2, 0.3, 0.4$ and 0.5. The tower energies were obtained by summing energies, calibrated at the electromagnetic energy scale [36], of all cells in all layers within the $\eta$ and $\phi$ boundaries of the towers. Cells that span tower boundaries had their energy apportioned by the fraction of the cell contained within a given tower. The jet measurements presented here were obtained by performing the anti-$k_T$ reconstruction on the towers prior to underlying event (UE) subtraction and then evaluating and subtracting the UE from each jet at the calorimeter cell level. The subtraction procedure calculates a per-event average UE energy density excluding contributions from jets and accounting for effects of elliptic flow modulation on the UE [37]. The UE estimation and subtraction was performed using a two-step procedure that was identical for all calorimeter layers.

A first estimate of the UE average transverse energy density, $\rho_U(\eta)$, was evaluated in 0.1 intervals of $\eta$ from all cells in each calorimeter layer, $i$, within the given $\eta$ interval excluding those within “seed” jets. In the first subtraction step, the seeds are defined to be $R = 0.2 \text{ jets containing at least one tower with } E_T > 3 \text{ GeV and having a ratio of maximum tower transverse energy to average tower transverse energy, } E_T^{\text{max}}/E_T > 4$. Elliptic flow in Pb + Pb collisions can impose a $2v_2 \cos (2(\phi - \Psi_2))$ modulation on the UE. Here, $v_2$ is the second coefficient in a Fourier decomposition of the azimuthal variation of the UE particle or energy density, and the event plane angle, $\Psi_2$, determines the phase of the elliptic modulation. Standard techniques [26,37] were used to measure $v_2$.

$$\Psi_2 = \tan^{-1}\left(\frac{\sum_{k} w_k E_{T_k} \sin (2\phi_k)}{\sum_{k} w_k E_{T_k} \cos (2\phi_k)}\right),$$ (1)

where $k$ runs over cells in the FCal, $\phi_k$ represents the cell azimuthal angle, and $w_k$ represent per-cell weights empirically determined to ensure a uniform $\Psi_2$ distribution. An $\eta$-averaged $v_2$ was measured separately for each calorimeter layer according to

$$v_{2i} = \frac{\sum_{j \in \mathcal{J}_i} E_{T_j} \cos (2(\phi_j - \Psi_2))}{\sum_{j \in \mathcal{J}_i} E_{T_j}},$$ (2)

where $j$ runs over all cells in layer $i$. The UE-subtracted cell transverse energies were calculated according to

$$E_{T_j}^{\text{sub}} = E_{T_j} - A_j \rho_U(\eta_j) \left(1 + 2v_2 \cos (2(\phi_j - \Psi_2))\right),$$ (3)

where $E_{T_j}$, $\eta_j$, $\phi_j$, and $A_j$ represent the cell $E_T$, $\eta$ and $\phi$ positions, and area, respectively for cells, $j$, in layer $i$. The kinematics for $R = 0.2$ jets generated in this first subtraction step were calculated via a four-vector sum of all (assumed massless) cells contained within the jets using the $E_T$ values obtained from Eq. (3).

The second subtraction step starts with the definition of a new set of seeds using a combination of $R = 0.2 \text{ jets from the first subtraction step with } E_T > 25 \text{ GeV and track jets (defined below) with } p_T > 10 \text{ GeV. Using this new set of seeds, a new estimate of the UE, } \rho_U(\eta), \text{ was calculated excluding cells within } \Delta R = 0.4 \text{ of the new seed jets, where } \Delta R = \sqrt{(\eta_{\text{cell}} - \eta_j)^2 + (\phi_{\text{cell}} - \phi_j)^2}.\text{ New } v_2 \text{, values, } v_2', \text{ were also calculated excluding all cells within } \Delta R = 0.4 \text{ of any of the new seed jets. This exclusion largely eliminates distortions of the calorimeter } v_2 \text{ measurement in events containing high- } p_T \text{ jets. The background subtraction was then applied to the original cell energies using Eq. (3) but with } \rho_i \text{ and } v_2' \text{, replaced by the new values, } \rho_U(\eta) \text{ and } v_2'. \text{ New jet kinematics were obtained for all jet radii from a four-momentum sum of cells within the jets using the subtracted cell transverse energies. Jets generated in this second subtraction step having } E_T > 20 \text{ GeV were recorded for subsequent analysis.}\text{ A correction of typically a few per cent was applied to the reconstructed jets to account for incomplete exclusion of towers within jets from the UE estimate due, for example, to differences in direction between the seeds and the final jets. This correction was validated by applying the full heavy ion jet reconstruction procedure to } 2.76 \text{ TeV pp data collected by ATLAS in March 2011. The reconstructed jets were compared, jet-by-jet, to those obtained from the pp jet reconstruction procedure. After this last correction for incomplete exclusion of jets from the background, the energy scales of the heavy ion and pp reconstruction procedures agreed to better than 1% for } E_T > 25 \text{ GeV. A final correction depending on the jet } \eta, E_T, \text{ and } R \text{ was applied to obtain the correct hadronic energy scale for the reconstructed jets. The calibration constants were derived separately for the four jet radii using the same procedure applied to pp jet measurements [38].}\text{ In addition to the calorimeter jet reconstruction, track jets were reconstructed using the anti-}k_T\text{ algorithm with } R = 0.4 \text{ from charged tracks that have a good match to the primary vertex and that have } p_T > 4 \text{ GeV. This threshold suppresses contributions of the UE to the track jet measurement. Specifically, an } R = 0.4 \text{ track jet has an estimated likelihood of including an uncorrelated } p_T > 4 \text{ GeV charged track of less than 4% in the 0–10% centrality bin. The single track reconstruction efficiency is } \approx 80\%, \text{ approximately independent of centrality.}\text{ The fluctuating UE in Pb + Pb collisions can potentially produce reconstructed jets that do not originate from hard-scattering processes. In the remainder of this Letter such jets are referred to as “underlying event jets” or UE jets. A requirement that calorimeter jets match at least one track jet with } p_T > 7 \text{ GeV or an EM cluster reconstructed from cells in the electromagnetic calorimeter [38] with } p_T > 7 \text{ GeV was applied to exclude UE jets. The matching criterion for both track jets and EM clusters is that they lie within } \Delta R = 0.2 \text{ of the jet. Applying this matching requirement provides a factor of about 50 rejection against UE jets while inducing an additional } p_T\text{-dependent inefficiency in the jet measurement. To accommodate the use of track jets in the UE jet recor-
generated jets were obtained by running Separately, the presence and approximate kinematics of HIJING- and the 0–10% and 60–80% centrality bins after all event selection criteria, UE jet rejection, and the |η| < 2.1 cut have been applied. In that MC sample, the kinematics of the reference measuring unquenched jets was obtained using the MC over-

The primary evaluation of the combined performance of the ATLAS detector and the analysis procedures described above in measuring unquenched jets was obtained using the MC overlay sample. In that MC sample, the kinematics of the reference PYTHIA generator-level jets (hereafter called “truth jets”) were re-

The JER was found to be well described by a quadrature sum of three terms,

\[
\frac{\sigma[\Delta E_T]}{E_T^{\text{truth}}} = \frac{a}{E_T^{\text{truth}}} \oplus \frac{b}{E_T^{\text{truth}}} \oplus c, \tag{4}
\]

where \(a\) and \(c\) represent the usual sampling and constant contributions to calorimeter resolution. The term containing \(b\) describes the contribution of underlying event fluctuations, which do not depend on jet \(E_T\). Results of fitting the \(E_T\) dependence of the JER according to Eq. (4), using methods described below, are shown with curves in Fig. 1.

The jet reconstruction efficiency decreases with decreasing jet \(E_T\) for \(E_T \lesssim 50\) GeV. The decrease starts at larger \(E_T\) and decreases more rapidly for larger jet radii and in more central col-

### Table 2

<table>
<thead>
<tr>
<th>R</th>
<th>(p_T &gt; 40) GeV 0–10%</th>
<th>(p_T &gt; 40) GeV 60–80%</th>
<th>(p_T &gt; 100) GeV 0–10%</th>
<th>(p_T &gt; 100) GeV 60–80%</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.2</td>
<td>112333 8068 2308 162</td>
<td>1287153 12629 3534 222</td>
<td></td>
<td></td>
</tr>
<tr>
<td>0.3</td>
<td>543444 15964 4974 277</td>
<td>710158 18573 7586 307</td>
<td></td>
<td></td>
</tr>
<tr>
<td>0.4</td>
<td>116231 8068 2308 162</td>
<td>1287153 12629 3534 222</td>
<td></td>
<td></td>
</tr>
<tr>
<td>0.5</td>
<td>710158 18573 7586 307</td>
<td>710158 18573 7586 307</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

Fig. 1. Results of MC evaluation of jet reconstruction performance in 0–10% and 60–80% collisions as a function of truth jet \(E_T\) for \(R = 0.2\) (left) and \(R = 0.4\) (right) jets. Top: jet energy resolution \(\sigma[\Delta E_T]/E_T^{\text{truth}}\) and jet energy scale closure, \(\langle \Delta E_T \rangle/E_T^{\text{truth}}\). Solid curves show parameterizations of the JER using Eq. (4). Bottom: Efficiencies, \(\epsilon\) and \(\epsilon'\), for reconstructing jets before and after application of UE jet removal (see text for explanation), respectively.
The inefficiency results primarily from the finite JER which causes jets with $E_{T}^{\text{truth}} > 20 \text{ GeV}$ to be measured with $E_{T}^{\text{rec}} < 20 \text{ GeV}$. The UE jet rejection causes an additional loss of jets but in a manner that reduces the centrality dependence of the inefficiency.

The accuracy of the MC overlay studies described above was evaluated using the data overlay sample analyzed using the same procedures that were applied to the MC overlay sample. The analysis yielded results for the JER, JES, and efficiency consistent with the MC overlay sample, although the JER in the data overlay sample was found to be slightly better than in the MC overlay sample. The JES in the data overlay sample was found to agree between peripheral and central collisions to better than 1% for $R = 0.4$ jets, and the reconstruction efficiency was found to differ by less than 5% on the rise of the efficiency curve.

A data-driven check of the HIJING description of UE fluctuations was performed by evaluating distributions of EM-scale summed $E_T$ in rectangular groups of towers within the interval $|y| < 2.8$. The groups were chosen to match the areas of jets used in this analysis: $3 \times 4$ and $7 \times 7$ for $R = 0.2$ and $R = 0.4$ jets, respectively. No attempt was made to exclude jets from the fluctuation analysis. The distributions of $E_{T}^{x=4}$ and $E_{T}^{x=7}$, the $\Sigma E_{T}$ for $3 \times 4$ and $7 \times 7$ groups of towers, are shown in Fig. 2 for a narrow range of $\Sigma E_{T}^{\text{Cal}}, 3.4 \leq \Sigma E_{T}^{\text{Cal}} < 3.5 \text{ TeV}$, that lies within the 0–1% centrality interval. These distributions have mean values, $(E_{T}^{x=4}) = 26 \text{ GeV}$ and $(E_{T}^{x=7}) = 105 \text{ GeV}$, subtracted and, thus, in principle represent the distribution of the residual contributions of the UE to jet energies after subtraction. However, the high tails of the distributions can be attributed to the presence of jets, which are not part of the UE. The corresponding distributions obtained from the HIJING MC sample, but with $(E_{T}^{x=4})$ and $(E_{T}^{x=7})$ obtained from data, are shown in Fig. 2 with filled histograms.

The shapes of the MC and data distributions in Fig. 2 (top) are very similar, but the MC result slightly over-predicts the positive fluctuations for all collision centralities. In central collisions the MC result also slightly over-predicts the size of negative fluctuations. In contrast, for non-central collisions (not shown here) the data has a broader distribution of negative fluctuations than the MC sample. These observations are demonstrated by Fig. 2 (bottom) which shows the standard deviations of the $E_{T}^{x=4}$ and $E_{T}^{x=7}$ distributions, $\sigma(E_{T}^{x=4})$ and $\sigma(E_{T}^{x=7})$, as a function of $\Sigma E_{T}^{\text{Cal}}$, obtained from both the data and the MC sample. The data and MC distributions have similar trends, but the MC $\sigma(E_{T}^{x=4})$ and $\sigma(E_{T}^{x=7})$ values are larger in central collisions by 2.5% and 5%, respectively. In non-central collisions, the broader spectrum of negative fluctuations in the data causes $\sigma(E_{T}^{x=4})$ and $\sigma(E_{T}^{x=7})$ to exceed the corresponding quantity in the HIJING MC sample by approximately the same percentages.

Consistency between the results of the fluctuation analysis and the evaluation of the JER described above has been established by fitting the $E_{T}$ dependence of the JER with the functional form given by Eq. (4) with fixed $b$ values obtained from the fluctuation analysis. The $b$ values for a given jet radius were determined by taking the standard deviation of the $\Sigma E_{T}$ distribution for the corresponding tower group averaged over centrality and corrected to the hadronic energy scale. The resulting $b$ values for $R = 0.2$ (0.4) jets are 5.62 (12.45) GeV and 1.15 (2.58) GeV for the 0–10% and 60–80% centrality bins respectively. The parameters $a$ and $c$ obtained from the fits are found to be independent of centrality within fit uncertainties, as expected, and to have values $a = 1.0$ (0.8), $c = 0.07$ (0.06) for $R = 0.2$ (0.4) jets with $E_{T}$ expressed in GeV. The accuracy of the fits in describing the $E_{T}$ dependence of the JER is demonstrated by the curves showing results for $R = 0.2$ and $R = 0.4$ jets in Fig. 1.

The contribution of UE jets to the measured jet spectrum after UE jet rejection is estimated to be $< 3\%$ approximately independent of jet $p_{T}$ for $40 < p_{T} < 60$ GeV and less than 1% for $p_{T} > 60$ GeV. This estimate was obtained by evaluating the rate of reconstructed jets in the HIJING MC sample which were not matched to HIJING truth jets and correcting for missing truth jets due to the $p_{T} > 4$ GeV requirement applied in the HIJING truth jet reconstruction.
7. Jet spectra and unfolding

Though jet reconstruction performance is naturally evaluated in terms of jet $E_T$, the physics measurements in this Letter were performed as a function of $p_T$ directly calculated from the jet four-momentum. The typical masses of the jets are sufficiently small that $E_T \approx p_T$ holds over the range of measured $p_T$ for all jet radii. The measured $p_T$ spectra of reconstructed jets passing UE jet rejection and having $|y| < 2.1$ were evaluated for each centrality bin using logarithmic $p_T$ bins spanning the range $38 < p_T < 210$ GeV. The correlations within and between $p_T$ bins arising from multi-jet events were quantified by the covariance, $C_{ij}$, between the number of jets measured in two bins, $i$ and $j$. The measured $R_{CP}$ was calculated as

$$R_{CP}(p_T)_{\text{cent}} = \left( \frac{N_{\text{cent}}(p_T)}{N_{\text{coll}}(p_T)} \right) / \left( \frac{N_{60-80}(p_T)}{N_{\text{coll}}(p_T)} \right),$$

where $N_{\text{cent}}(p_T)$ represents the measured jet yield in a given $p_T$ and centrality bin, and $N_{\text{cent}}$ and $N_{60-80}$ are the number of Pb + Pb collisions within the chosen and peripheral reference centrality intervals, respectively. Results for $R_{CP}(p_T)_{\text{0-10}}$ obtained from the measured spectra are shown in Fig. 3 for $R = 0.2$ and $R = 0.4$ jets. The $R_{CP}(p_T)_{\text{0-10}}$ for $R = 0.2$ jets is approximately equal to 0.5 over the measured $p_T$ range. The $R_{CP}(p_T)_{\text{0-10}}$ for $R = 0.4$ and $R = 0.2$ jets are consistent for $p_T > 120$ GeV, but at lower $p_T$, the $R = 0.4$ $R_{CP}(p_T)_{\text{0-10}}$ increases relative to the $R = 0.2$ values. The difference between $R = 0.2$ and $R = 0.4$ $R_{CP}(p_T)_{\text{0-10}}$ values can be mostly attributed to the difference in the size of the UE fluctuations for $R = 0.2$ and $R = 0.4$ jets shown in Fig. 1. The larger JER for $R = 0.4$ jets produces greater upward migration on the steeply falling jet $p_T$ spectrum in central collisions than in peripheral collisions, thus enhancing the measured $R_{CP}$. The drop in the $R = 0.4$ $R_{CP}(p_T)_{\text{0-10}}$ at low $p_T$ is due to the decrease in jet reconstruction efficiency between 60–80% and 0–10% centrality bins which, as noted above, largely results from the worse JER in central collisions.

To remove the effects of the bin migration, the jet spectra were unfolded using the singular value decomposition (SVD) technique as implemented in RooUnfoldOld [40]. The MC overlay samples were used to populate a response matrix, $A$, which describes the transformation of the true jet spectrum, $x$, to the observed spectrum, $b$, according to $b = Ax$. The truth and reconstructed jet $p_T$ were obtained from the MC overlap sample using the methods described in Sections 6 and 5, respectively, and the selection and matching of truth and reconstructed jet pairs was performed as described in Section 6. Using the weighting method suggested in Ref. [39], the unfolded spectrum is expressed as a set of weights $w$ multiplying the input spectrum $x_{\text{ini}}$ used to produce $A$. The SVD method expresses the solution for $w$ in terms of a least-square minimization problem that includes a prescription for regularizing the amplification of statistical fluctuations of the data that would result from the direct inversion of $A$. The regularization is controlled by a parameter $\tau$ such that contributions from singular values $s_k$ of the unfolding matrix with $s_k < \tau$ are suppressed. Inclusion of the $p_T$-dependent reconstruction efficiency in the response was found to strongly affect the spectrum of singular values of the matrix defining the SVD problem, so the efficiency correction was applied separately following the unfolding. The spectrum of MC truth jets was reweighted to provide a smooth, power-law initial spectrum, $x_{\text{ini}} \propto \epsilon^2(p_T)/p_T^\alpha$, where the power index was chosen to be $\alpha = 5$. An analysis of the optimal regularization in the SVD unfolding following the methods of Ref. [39] indicated that a regularization parameter fixed by the fifth singular value ($\tau = s_5^2$) of the SVD matrix was appropriate for all centralities and all $R$ values. The statistical uncertainties in the SVD unfolding due to statistical errors on the input spectrum were evaluated using the pseudo-experiment technique with 1000 separate stochastic variations of the input spectrum based on the full covariance matrix. The contributions of statistical fluctuations in the response matrix, $A$, were similarly evaluated using an equal number of stochastic variations of the response matrix. The two contributions to the statistical uncertainty were combined in quadrature.

Potential biases in the unfolding procedure were evaluated using two different methods. Each unfolded spectrum was re-folded with its corresponding response matrix and compared to the measured spectrum for self-consistency. In general, regularization can introduce differences between re-folded and measured spectra on the scale of statistical uncertainties on the measured spectra, while over-regularization can produce larger, systematic differences. For all of the unfolded spectra, the re-folding procedure yielded a typical difference between measured and re-folded spectra comparable to the statistical uncertainties on the measured spectra. A separate check was performed by unfolding the reconstructed MC spectrum for each centrality bin and each jet radius and comparing to the original MC truth jet spectrum. For this purpose, the MC data sets were divided in half and reconstructed spectra and response matrices were generated separately from each set. The unfolded and truth MC jet spectra typically agreed to better than 2%, though for the 0–10% centrality bin and for $R = 0.4$ and 0.5 $p_T$ jets, differences as large as 5% were observed in these lowest $p_T$ bins. These differences are covered by the unfolding systematic uncertainties described below.

The corrected $R_{CP}$ was evaluated according to

$$R_{CP}(p_T)_{\text{cent}} = \left( \frac{N_{\text{cent}}(p_T)}{N_{\text{coll}}(p_T)} \right) / \left( \frac{N_{60-80}(p_T)}{N_{\text{coll}}(p_T)} \right),$$

Fig. 3. Top: Measured and corrected $R_{CP}$ values for the 0–10% centrality bin as a function of jet $p_T$ for $R = 0.4$ and $R = 0.2$ jets. Bottom: Ratio of corrected to measured $R_{CP}$ values for both jet radii. The error bars on the points represent statistical uncertainties only.
where $\tilde{N}_{\text{jet}}$ represents the unfolded number of jets in the $p_T$ bin, and $\epsilon_{\text{cent}}^{pT}$ and $\epsilon_{60-80}^{pT}$ are the $p_T$-dependent jet reconstruction efficiencies after UE jet rejection for the indicated centrality bins. Fig. 3 shows the comparison of the corrected and measured $R_{\text{CP}}$ values as a function of jet $p_T$ for $R = 0.2$ and $R = 0.4$ jets in the 0–10% centrality bin. The unfolding has little effect on the $R = 0.2$ $R_{\text{CP}}$ due to the good energy resolution (relative to larger radii) for $R = 0.2$ jets even in central collisions. For the $R = 0.4$ jets, $R_{\text{CP}}$ is reduced by a factor of about two at the lowest $p_T$ values included in the analysis and is only slightly modified at the highest $p_T$. Because the unfolding provides a non-local mapping of the input jet $p_T$ spectrum onto the unfolded spectrum, the statistical uncertainties in the unfolded spectra have significant correlations between bins, and there is not a direct relationship between the statistical errors in the input spectrum and the unfolded spectrum. The regularization of the unfolding also suppresses statistical fluctuations in the unfolded spectrum, but the statistical uncertainties in the measured spectrum also contributes to the systematic uncertainties from the unfolding procedure.

8. Systematic uncertainties

Systematic uncertainties in the $R_{\text{CP}}$ measurement can arise due to errors on the jet energy scale (JES), the jet energy resolution (JER), jet finding efficiency, the unfolding procedure, and the $R_{\text{coll}}$ values. Uncertainties in jet $E_T$ and $p_T$ are assumed to be equal (i.e. $\delta p_T = \delta E_T$). Uncertainties in the JES and the JER influence the unfolding of the jet spectra. The resulting systematic uncertainties on the $R_{\text{CP}}$ values were estimated by producing new response matrices according to the procedures described below, generating unfolded spectra from these matrices, and calculating new $R_{\text{CP}}$ values. The resulting changes in the $R_{\text{CP}}$ values were taken to be estimates of $\delta R_{\text{CP}}^\text{sys}$. For uncertainties fully correlated in centrality, $\delta R_{\text{CP}}^\text{sys}$ was evaluated by simultaneously varying the chosen centrality bin and the 60–80% bin, while for other uncertainties, the chosen centrality bin and 60–80% centrality bins were varied separately and the variations in $R_{\text{CP}}$ combined in quadrature.

Overall JES uncertainties common to the different centrality bins cancel in the ratio of the spectra in $R_{\text{CP}}$, but centrality-dependent JES errors will produce a systematic shift in $R_{\text{CP}}$. Studies using the MC overlay sample discussed in Section 6 indicate a maximum difference in JES between the 0–10% and 60–80% centrality bins for the jet $p_T$ range included in this analysis of 0.5%, 1%, 1.5% and 2.5% for $R = 0.2$, 0.3, 0.4 and 0.5 jets. At lower $p_T$, the assigned systematic uncertainties increase linearly with decreasing $p_T$ such that they double in size between 70 GeV and 38 GeV. For other centrality bins, the systematic errors on the centrality dependence of the JES decrease smoothly from central to peripheral collisions. The resulting $\delta R_{\text{CP}}^\text{sys}$ values were evaluated using new response matrices generated by scaling the reconstructed $p_T$ to account for the above-quoted JES uncertainties. The JES systematic uncertainty is assumed to be fully correlated between different centrality bins and different $R$ values.

Systematic uncertainties in the JER due to inaccuracies in the MC description of the UE fluctuations were evaluated using results of the fluctuation analysis described above. The effects of those inaccuracies were evaluated by rescaling the jet $p_T$ to $p_T^{\text{truth}}$ values obtained from the MC study by factors that cover the differences between data and MC result. For each centrality and jet radius, a modified value of the $b$ parameter in Eq. (4) was evaluated and used to obtain new JER values, $\sigma'[\Delta E_T]$ from Eq. (4). Then a rescaled $\Delta p_T$ was obtained from

$$\Delta p_T = \Delta p_T \left( \frac{\sigma'}{\sigma} \right).$$

Since the discrepancies between the MC and the data were observed to be different for positive and negative fluctuations, the rescaling was applied separately for positive and negative $\Delta p_T$.

The $\Sigma E_T$ values in the MC study were found to have larger positive fluctuations than those in the data for all centralities by approximately 2.5%, 2.5%, 5%, and 7.5% for $R = 0.2$, 0.3, 0.4 and 0.5 jets, respectively, so for positive $\Delta p_T$, $b$ was reduced by these percentages. For the 0–10% centrality bin, the negative fluctuations were also larger in the MC study than in the data by the same approximate percentages, so for central collisions the same, modified $b$ value was used for negative $\Delta p_T$. For all other centrality bins, the negative fluctuations in the data were larger than in the MC by approximately twice the above-quoted percentages. Thus, for those centralities, the modified $b$ values were obtained for negative $\Delta p_T$ by increasing $b$ by 5%, 5%, 10%, and 15%, respectively, for $R = 0.2$, 0.3, 0.4 and 0.5 jets.
New response matrices were generated using the calculated \( \Delta p_L \) values according to \( p^\text{rec}_L = p^\text{true}_L + \Delta p_L \), and these modified response matrices were used to estimate the JER systematic uncertainties following the procedure described above. The systematic uncertainty on the spectra due to the JER for the 0–10% centrality bin was taken to be one-sided as all evaluations indicate that the MC simulations slightly overestimate UE fluctuations. Asymmetric errors were obtained for the other centrality bins by applying the positive and negative \( \Delta E_T \) scalings separately. The JER systematic uncertainties were assumed to be fully correlated between different jet \( R \) values but uncorrelated between different collision centralities, so the uncertainties on the spectra were combined in quadrature in evaluating \( \delta R_{CP} \). The conservative assumption that the JER uncertainties are fully uncorrelated between different centrality bins is based on the observation that the differences between data and the HIJING MC sample in the fluctuation analysis are not the same for all centralities.

The systematic uncertainties associated with the non-UE contributions to the JER (described by the \( a \) and \( c \) terms in Eq. (4)) were evaluated following procedures used by ATLAS in previous \( pp \) jet measurements [41]. New response matrices were generated by applying an additional stochastic smearing to the \( \Delta p_T \) values, and the systematic uncertainty was obtained by applying the procedure described above.

Systematic uncertainties on \( R_{CP} \) due to the unfolding were evaluated by changing the power index \( (n) \) in the functional form for \( x_{\text{lim}} \) by \( \pm 0.5 \) and by varying the regularization parameter. The \( \pm 0.5 \) change in the power law index was chosen because it produces a spectrum that changes relative to the default \( x_{\text{lim}} \) over the measured \( p_T \) range by a factor of about two – the typical suppression observed in central collisions. Thus, it covers the possibility that the true \( R_{CP} \) could increase to one or decrease to 0.25 over the measured \( p_T \) range. To evaluate the potential systematic uncertainty due to regularization, the unfolding was performed with regularization parameters obtained from the fourth and sixth singular values of the unfolding matrix, \( \tau = \frac{\sigma^4}{\lambda} \) and \( \tau = \frac{\sigma^6}{\lambda} \). Systematic uncertainties on the spectra were determined from the differences in the unfolded spectra. The resulting \( \delta R_{CP}^{\text{sys}} \) values were obtained assuming that the regularization uncertainties on the two spectra are uncorrelated.

The systematic uncertainty on the efficiency correction was evaluated by comparing MC overlay and data overlay samples where differences less than 5% were observed on the “turn on” part of the efficiency curve. A 5% uncertainty due to the efficiency correction was applied to \( R_{CP} \) for \( p_T < 100 \) GeV in the four most central bins. To check for biases introduced by the UE jet rejection, the analysis was repeated with a significantly weakened rejection criterion in which jets were required to match a single track with \( p_T > 4 \) GeV. No significant differences in the \( R_{CP} \) were observed except for \( p_T < 50 \) GeV where differences as high as 4% were found. These differences can be attributed to the contribution of additional UE jets.

The different contributions to the total \( \delta R_{CP}^{\text{sys}} \) are shown in Fig. 4 for \( R = 0.4 \) jets in the 0–10% centrality bin. The JES and \( x_{\text{lim}} \) uncertainties are approximately independent of \( p_T \), while the JER uncertainty decreases with increasing \( p_T \). The regularization uncertainty grows with increasing \( p_T \) due to the poorer statistical precision of the high-\( p_T \) points. The systematic uncertainties for the other radii show similar \( p_T \) and centrality dependence, with the JES and JER uncertainties increasing with jet radius as expected.

### 9. Results

Fig. 5 shows the \( R_{CP} \) values obtained for \( R = 0.2 \) and \( R = 0.4 \) jets as a function of \( p_T \) in four bins of collision centrality with three different error contributions: statistical uncertainties, partially correlated systematic uncertainties, and fully correlated uncertainties. The \( R_{CP} \) values for all centralities and for both jet radii are observed to have at most a weak variation with \( p_T \). For the 0–10% centrality bin the \( R_{CP} \) values for both jet radii show a factor of about two suppression in the \( 1/N_{\text{cell}} \)-scaled jet yield. For more peripheral collisions, \( R_{CP} \) increases at all \( p_T \) relative to central collisions, with the \( R_{CP} \) values reaching 0.9 for the 50–60% centrality bin. A more detailed evaluation of the centrality dependence...
Fig. 6. $R_{CP}$ values as a function of $N_{\text{part}}$ for $R = 0.4$ anti-$k_t$ jets in six $p_T$ bins. The error bars indicate statistical errors from the unfolding; the shaded boxes indicate point-to-point systematic errors that are only partially correlated. The solid lines indicate systematic errors that are fully correlated between all points. The horizontal errors indicate systematic uncertainties on $N_{\text{part}}$.

Fig. 7. Left: $R_{CP}$ in the 0–10% centrality bin as a function of jet radius for four bins of jet $p_T$. Right: $R_{CP}$ as a function of jet radius for four centrality bins for the $p_T$ interval $89 < p_T < 103$ GeV. The error bars indicate statistical errors from the unfolding; the shaded boxes indicate point-to-point systematic errors that are only partially correlated. The solid lines indicate systematic errors that are fully correlated between all points. The horizontal width of the systematic error band is chosen for presentation purposes only. Dotted lines indicate $R_{CP} = 0.5$, and the dashed lines on the top panels indicate $R_{CP} = 1$.

The dependence of $R_{CP}$ on jet radius is shown in Fig. 7 for the 0–10% centrality bin in four jet $p_T$ intervals (left) and for different centrality bins in the $89 < p_T < 103$ GeV bin (right). For this figure, the shaded boxes indicate the combined contribution of systematic uncertainties due to regularization, $x_{\text{ini}}$, and efficiency, which are only partially correlated between points. All other systematic errors are fully correlated and are indicated by solid lines. The results in Fig. 7 show a weak variation of $R_{CP}$ with $R$, that is nonetheless significant when taking into account the correlations in the errors between the different $R$ values.

To demonstrate this conclusion more clearly, Fig. 8 shows the ratio of $R_{CP}$ values between $R = 0.3, 0.4$ and 0.5 jets and $R = 0.2$ jets, $R_{CP}^{0.3}/R_{CP}^{0.2}$, as a function of $p_T$ for the 0–10% centrality bin. When evaluating the ratio, there is significant cancellation between the correlated systematic uncertainties. Statistical correlations between the jet yields for the different radii were evaluated in the measured spectra and tracked through the unfolding procedure separately for the 0–10% and 60–80% centrality bins. Those correlations were then included when evaluating the statistical errors on $R_{CP}^{0.3}/R_{CP}^{0.2}$ shown in Fig. 8. The results in that figure indicate a significant dependence of $R_{CP}$ on jet radius. For $p_T < 100$ GeV the $R_{CP}^{0.3}/R_{CP}^{0.2}$ values for both $R = 0.4$ and $R = 0.5$ differ from one beyond the statistical and systematic uncertainties. The deviation persists for $R = 0.5$ above 100 GeV. A similar, but weaker dependence is observed in the 10–20% centrality bin. In more peripheral bins, no significant radial dependence is observed. The differences of $R_{CP}$ for $R = 0.4$ jets is presented in Fig. 6, which shows $R_{CP}$ vs $N_{\text{part}}$ for six jet $p_T$ bins. $R_{CP}$ decreases monotonically with increasing $N_{\text{part}}$ for all $p_T$ bins. The lower $p_T$ bins, for which the data are more statistically precise, show a variation of $R_{CP}$ with $N_{\text{part}}$ that is most rapid at low $N_{\text{part}}$. Trends similar to those shown in Figs. 5 and 6 are observed for all jet radii.
per nucleon at the LHC. The measurements were performed over
$R$ dependence of jet suppression, characterized by the inclusive jet
suppression. The suppression varies monotonically with collision
$T$, except for the lowest two $p_T$ bins. However, direct comparisons of $R_{CP}$ between different jet radii at low $p_T$ should be treated with care as the same jets measured using smaller radii will tend to appear in lower $p_T$ bins than when measured with a larger radius.

10. Conclusions

This Letter presents results of measurements of the centrality
dependence of jet suppression, characterized by the inclusive jet
central-to-peripheral ratio, $R_{CP}$, in Pb + Pb collisions at 2.76 TeV
per nucleon at the LHC. The measurements were performed over
the $p_T$ range $38 < p_T < 210$ GeV for anti-$k_T$ jets of radii $R = 0.2, 0.3, 0.4$ and 0.5. The inclusive jet yield is observed to be sup-
pressed by a factor of about two in central collisions relative to
peripheral collisions with at most a weak $p_T$ dependence to the
suppression. The suppression varies monotonically with collision
centrality over the measured $p_T$ range and for all jet radii. The
$R_{CP}$ at fixed $p_T$ is observed to vary with $p_T$ increasing gradu-
ally from $R = 0.2$ to $R = 0.5$. That variation is most significant
for $p_T < 100$ GeV where more than a 30% variation is observed.
These results provide the first direct measurement of inclusive jet
suppression in heavy ion collisions. The substantial suppression of
the jet yield observed at all $p_T$ values complements the previous
measurements of dijet transverse energy imbalance in Pb + Pb col-
lisions at the LHC [13–15].

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Royal Society and Leverhulme Trust, United Kingdom; DOE and
NSF, United States of America.
Department of Physics and Astronomy, University of New Mexico, Albuquerque, NM, United States

Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen, Nijmegen, Netherlands

Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands

Department of Physics, Northern Illinois University, DeKalb, IL, United States

Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia

Department of Physics, New York University, New York, NY, United States

Ohio State University, Columbus, OH, United States

Faculty of Science, Okayama University, Okayama, Japan

Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, OK, United States

Department of Physics, Oklahoma State University, Stillwater, OK, United States

Prague University, KTPM, Olomouc, Czech Republic

Center for High Energy Physics, University of Oregon, Eugene, OR, United States

LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France

Graduate School of Science, Osaka University, Osaka, Japan

Department of Physics, University of Oslo, Oslo, Norway

Department of Physics, Oxford University, Oxford, United Kingdom

INFN Sezione di Pavia, (b) Dipartimento di Fisica, Università di Pavia, Pavia, Italy

Department of Physics, University of Pennsylvania, Philadelphia, PA, United States

Petersburg Nuclear Physics Institute, Gatchina, Russia

INFN Sezione di Pisa, (b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy

Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA, United States

Laboratorio de Instrumentacao e Fisica Experimental de Particulas – LIP, Lisboa, Portugal; (b) Departamento de Fisica Teorica y del Cosmos and CAFPE, Universidad de Granada, Granada, Spain

Institute of Physics, Academy of Sciences of the Czech Republic, Prague, Czech Republic

Faculty of Mathematics and Physics, Charles University in Prague, Prague, Czech Republic

Czech Technical University in Prague, Prague, Czech Republic

State Research Center Institute for High Energy Physics, Protvino, Russia

Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom

Physics Department, University of Regina, Regina, SK, Canada

Gifu University, Gifu, Gifu, Japan

INFN Sezione di Roma I; (a) Dipartimento di Fisica, Università La Sapienza, Roma, Italy

INFN Sezione di Roma Tor Vergata; (b) Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy

INFN Sezione di Roma Tre; (b) Dipartimento di Fisica, Università Roma Tre, Roma, Italy

Faculté des Sciences Ain Chock, Réséau Universitaire de Physique des Hautes Énergies – Université Hassan II, Casablanca; (b) Centre National de l’Energie des Sciences Techniques Nucléaires, Rabat; (c) Faculté des Sciences Semlalia, Université Cadi Ayyad, LPTEA, Marrakech; (d) Faculté des Sciences, Université Mohamed Premier et LPTPM, Oujda; (e) Faculté des Sciences, Université Mohammed V-Agdal, Rabat, Morocco

PSM/IFRIF (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat a l’Energie Atomique), Gif-sur-Yvette, France

Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, CA, United States

Department of Physics, University of Washington, Seattle, WA, United States

Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom

Department of Physics, Shinyu University, Nagano, Japan

Fachbereich Physik, Universität Siegen, Siegen, Germany

Department of Physics, Simon Fraser University, Burnaby, BC, Canada

SLAC National Accelerator Laboratory, Stanford, CA, United States

Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom

Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava; (b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic

Department of Physics, University of Johannesburg, Johannesburg; (b) School of Physics, University of the Witwatersrand, Johannesburg, South Africa

Department of Physics, Stockholm University, The Oskar Klein Centre, Stockholm, Sweden

Physics Department, Royal Institute of Technology, Stockholm, Sweden

Departments of Physics and Astronomy and Chemistry, Stony Brook University, Stony Brook, NY, United States

Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom

School of Physics, University of Sydney, Sydney, Australia

Institute of Physics, Academia Sinica, Taipei, Taiwan

Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel

Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel

Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece

International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan

Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan

Department of Physics, Tokyo Institute of Technology, Tokyo, Japan

Department of Physics, University of Toronto, Toronto, ON, Canada

TRIUMF, Vancouver, BC, (b) Department of Physics and Astronomy, York University, Toronto, ON, Canada

Institute of Pure and Applied Sciences, University of Tsukuba, 1-1-1 Tennoji, Tsukuba, Ibaraki 305-8571, Japan

Science and Technology Center, Tsu University, Medford, MA, United States

Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia

Department of Physics and Astronomy, University of California Irvine, Irvine, CA, United States

INFN Gruppo Collegato di Udine; (b) ITP Trieste; (b) Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy

Department of Physics, University of Illinois, Urbana, IL, United States

Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden

Istituto di Fisica Corpuscular (IFIC) and Departamento de Fisica Atomica, Molecular y Nuclear and Departamento de Ingenieria Electronica and Instituto de Microelectonica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain

Department of Physics, University of British Columbia, Vancouver, BC, Canada

Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada

Department of Physics, University of Warwick, Coventry, United Kingdom

Waseda University, Tokyo, Japan

Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel

Department of Physics, University of Wisconsin, Madison, WI, United States

Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany

Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany

Department of Physics, Yale University, New Haven, CT, United States

177 Yerevan Physics Institute, Yerevan, Armenia
178 Domaine Scientifique de la Doua, Centre de Calcul CNRS/IN2P3, Villeurbanne Cedex, France

Also at Laboratorio de Instrumentacao e Fisica Experimental de Particulas – LIP, Lisboa, Portugal.

Also at Faculdade de Ciencias and CFPUL, Universidade de Lisboa, Lisboa, Portugal.

Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom.

Also at TRIUMF, Vancouver, BC, Canada.

Also at Department of Physics, California State University, Fresno, CA, United States.

Also at Novosibirsk State University, Novosibirsk, Russia.

Also at Fermilab, Batavia, IL, United States.

Also at Department of Physics, University of Coimbra, Coimbra, Portugal.

Also at Department of Physics, UASLP, San Luis Potosi, Mexico.

Also at Università di Napoli Parthenope, Napoli, Italy.

Also at Institute of Particle Physics (IPP), Canada.

Also at Department of Physics, Middle East Technical University, Ankara, Turkey.

Also at Louisiana Tech University, Ruston, LA, United States.

Also at Dep. Fisica and CEFITEC of Faculdade de Ciencias e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal.

Also at Department of Physics and Astronomy, University College London, London, United Kingdom.

Also at Group of Particle Physics, University of Montreal, Montreal, QC, Canada.

Also at Department of Physics, University of Cape Town, Cape Town, South Africa.

Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.

Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany.

Also at Manhattan College, New York, NY, United States.

Also at School of Physics, Shandong University, Shandong, China.

Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France.

Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China.

Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan.

Also at INFN Sezione di Roma I, Italy.

Also at DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l'Univers), CEA Saclay (Commissariat a l'Energie Atomique), Gif-sur-Yvette, France.

Also at Departamento de Fisica, Universidade de Minho, Braga, Portugal.

Also at Department of Physics and Astronomy, University of South Carolina, Columbia, SC, United States.

Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.

Also at California Institute of Technology, Pasadena, CA, United States.

Also at LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France.

Also at Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom.

Also at Department of Physics, Oxford University, Oxford, United Kingdom.

Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.

Also at Department of Physics, The University of Michigan, Ann Arbor, MI, United States.

Deceased.