Measurement of the high-mass Drell-Yan differential cross-section in pp collisions at $\sqrt{s} = 7$ TeV with the ATLAS detector


Published in:
Physics Letters B

DOI:
10.1016/j.physletb.2013.07.049

Link to publication

Citation for published version (APA):

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Measurement of the high-mass Drell–Yan differential cross-section in \( pp \) collisions at \( \sqrt{s} = 7 \) TeV with the ATLAS detector

ATLAS Collaboration

1. Introduction

At hadron colliders, the Drell–Yan (DY) process [1], proceeding at tree level via the s-channel exchange of a virtual photon or Z boson, can produce charged lepton pairs over a wide range of invariant mass. The differential cross-section as a function of the invariant mass is described by perturbative QCD (pQCD) calculations at next-to-next-to-leading order. Given the simple experimental signature and the low backgrounds, a small experimental uncertainty can be achieved on the measured invariant mass distribution allowing for a precision test of pQCD. The mass spectrum is also sensitive to the parton distribution functions (PDFs), in particular to the poorly known distribution of antiquarks at large \( x \) [2], where \( x \) can be interpreted, at leading order, as the fraction of the proton momentum carried by the interacting parton. Additionally, the production of DY dilepton pairs is a source of background for other Standard Model (SM) measurements, and the mass spectrum may be modified by new physics phenomena giving rise to, e.g., narrow resonances or an excess of high-mass pairs inconsistent with the known PDFs.

The differential cross-section for DY dilepton pair production in the high-mass range has been reported previously by the CMS [3], CDF [4] and D0 [5] Collaborations. With the ATLAS detector, total and differential cross-sections in a mass window of 66–116 GeV have been measured using the 2010 dataset [6]. In addition, searches for new physics in the high-mass range have been performed [7–9] and no deviations from the SM expectation were observed. This Letter reports an extension of these previous analyses by providing a measurement of the DY cross-section, fully corrected for detector effects, in the dielectron channel as a function of the \( e^+e^- \) invariant mass, \( m_{ee} \), up to 1500 GeV. To minimise model-dependent theoretical uncertainties, the cross-section is not extrapolated to the full phase space but is reported in a phase space only slightly extended with respect to the fiducial acceptance of the \( e^+ \) and \( e^- \). The results are compared to NNLO pQCD calculations with next-to-leading-order (NLO) electroweak corrections from the FEWZ 3.1 [10,11] framework and to the predictions from three event generators.

2. The ATLAS detector

The ATLAS detector is described in detail in Ref. [12]. The two systems most relevant to this analysis are the inner tracking detector, surrounded by a superconducting solenoid providing a 2 T axial magnetic field, and the calorimeter. Charged-particle tracks and vertices are reconstructed with silicon pixel and microstrip detectors covering the pseudorapidity range \( |\eta| < 2.5 \) and a strawtube transition-radiation tracker covering \( |\eta| < 2.0 \). Within the region \( |\eta| < 3.2 \), electromagnetic calorimetry is provided by barrel and endcap detectors consisting of lead absorbers and liquid argon (LAr) as the active material, with fine lateral and longitudinal...
segmentation within $|\eta| < 2.5$. The hadronic calorimeter is based on steel/scintillator tiles in the central region ($|\eta| < 1.7$) while the hadronic endcap calorimeters ($1.5 < |\eta| < 3.2$) use copper/LAr.

A three-level trigger system is used to select events. The first level is implemented in custom electronics and is followed by two software-based trigger levels. In 2011 the total output rate of events recorded for physics analysis was 200–300 Hz.

3. Simulated samples

Simulated data samples were generated in order to estimate backgrounds and correct the signal for the detector resolution, efficiency and acceptance. The PYTHIA 6.4.26 [13] and MC@NLO 4.02 [14] Monte Carlo (MC) generators were used to model the DY signal. In addition, SHERPA 1.3.1 [15] was used to produce signal samples with up to three additional partons, and the final result of the analysis is compared to the generator-level predictions from all three programs. MC@NLO was also used to simulate the $t\bar{t}$ background, while HERWIG 6.520 [16] was used for the diboson ($WW, WZ$ or $ZZ$) backgrounds. MC@NLO was interfaced to HERWIG to model parton showers and fragmentation processes, and to JIMMY 4.31 [17] for underlying event simulation. All event generators were interfaced to PHOTOS 3.0 [18] to simulate QED final-state radiation (FSR), except for SHERPA which uses the method of Ref. [19].

The PYTHIA and HERWIG samples were generated using the modified leading-order (LO**) PDF set MRSTMcAl1 [20] following the recommendations of Ref. [21], while the MC@NLO samples used the NLO CT10 [22] set. The SHERPA samples used the default CTEQ6L1 [23] PDF set of the generator.

All MC events were generated at $\sqrt{s} = 7$ TeV and include the full ATLAS detector simulation [24] based on GEANT4 [25]. Settings of MC parameters that describe properties of minimum bias events and the underlying event were chosen based on results from previous ATLAS measurements [26]. The effects of having on average nine interactions per bunch crossing ("pile-up") were accounted for by overlaying simulated minimum bias events. To match the measured instantaneous luminosity profile of the LHC, MC events were reweighted to yield the same distribution of the mean number of interactions per bunch crossing as measured in data.

Several independent corrections were applied to the simulated samples, for the detector response, missing higher order terms in the generation of the signal events and for the modelling of the transverse momentum spectrum of the lepton pair. The electron [2] energy resolution was corrected to match that observed in data, following Ref. [27]. In addition, the efficiencies for electrons to pass requirements on the trigger, the reconstruction, and the particle identification in the MC simulation were corrected by scale factors, defined as the ratio of the measured efficiency in data to that in the simulation. The PYTHIA and MC@NLO signal predictions were reweighted to a NNLO pQCD calculation with $m_{ee}$-dependent $K$-factors obtained from a modified version of PHOZPR [28]. Additionally, NLO electroweak corrections, calculated using HORACE 3.1 [29], were applied to the PYTHIA MC sample. The $t\bar{t}$ sample was rescaled to its inclusive near-NNLO cross-section prediction [30,31] and the diboson samples were normalised to NLO cross-sections calculated using MCFM [32]. The PYTHIA signal MC sample was reweighted at generator level to a version that used an ATLAS tune found to yield a good agreement with the transverse momentum distribution of the $Z$ boson observed in data [33]. This procedure gives an adequate description of the transverse momentum distribution for the high $m_{ee}$ region studied in this analysis.

4. Event selection

The analysis is based on the full 2011 data sample collected at $\sqrt{s} = 7$ TeV. The data were selected online by a trigger that required two electromagnetic (EM) energy deposits each with a transverse energy greater than 20 GeV. Applying trigger and data-quality requirements yields an integrated luminosity of $4.9 \pm 0.1 \text{ fb}^{-1}$. Events from these $pp$ collisions are selected by requiring a collision vertex with at least three associated tracks, each with transverse momentum greater than 400 MeV. Events are then required to have at least two electron candidates as defined below.

Electron candidates are reconstructed from the energy deposits in the calorimeter matched to inner-detector tracks. The electron energy is measured in the calorimeter and its direction from the associated track. The calorimeter energy resolution is between 1% and 2% for high-energy electrons [27]. An energy scale correction obtained from an in situ calibration, using $W/Z$ boson and $J$/ψ meson decays, following the recipe of Ref. [27], is applied to the data. The electron candidates are required to have a transverse energy $E_T > 25$ GeV and pseudorapidity $|\eta| < 2.47$, excluding the transition regions between the barrel and endcap calorimeters at $1.37 < |\eta| < 1.52$. They must satisfy the "medium" identification criteria based on shower shape, track-quality and track–cluster matching variables, which are inclusive of the shower shape criteria applied as part of the "loose" identification [27]. Additionally, the electron candidates must have an associated hit in the inner-most pixel layer to suppress background from photon conversions.

If an event contains more than two electron candidates passing the above selection, the two with highest $E_T$ are chosen. To further reduce the background from jet production, the leading (highest $E_T$) electron is required to be isolated by demanding that the sum of the transverse energy in the calorimeter cells in a cone of $\Delta R = 0.2$ around the electron direction is less than 7 GeV. This sum excludes the core of the electron energy deposition and is corrected for the $E_T$-dependent transverse shower leakage from the core, as well as for pile-up contributions.

After all selection requirements, a total of 26 844 candidate events are found in the $m_{ee}$ range considered. The dominant backgrounds are events containing one or two misidentified electron candidates, denoted $W +$ jets and dijet. Other backgrounds arise from events containing two real electrons, originating from the dileptonic decays of pair-produced top quarks (denoted $t\bar{t}$) and from diboson production processes.

Of the dijet and $W +$ jets background, the dijet component additionally contains multi-jet, heavy-flavour quark and $γ +$ jet production. The $W +$ jets includes pair-produced top quarks and single-top-quark production, where at least one electron comes from the misidentification of a jet or a heavy quark. A data-driven method is used to evaluate the sum of these components. The probability for a jet to be misidentified as an electron (the fake rate) is determined in an $E_T$- and $η$-dependent way from nine background-enriched samples recorded by different inclusive jet triggers. These triggers had $E_T$ thresholds in the range 20–240 GeV, each with a different predefined rate achieved via the automatic rejection of a certain fraction of events, such that the nine samples were needed to collect sufficient background events over the full $E_T$ range. In each of these jet-triggered samples, the fake rate is calculated as the fraction of electron candidates passing the "loose" identification requirement that also pass the "medium" requirement. Events containing electron candidates from $W$ or $Z$ boson decays are first removed by dedicated cuts in order to avoid bias from real electron contamination: $W$ candidates are rejected by
requiring low missing transverse energy and low transverse mass; and \(Z\) candidates are rejected if they contain two “medium” electrons. A weighted average of the fake rates obtained from the nine jet samples is then calculated. To estimate the total dijet plus jets background, a factor derived from the averaged fake rate is applied to events that pass the signal selection but with one or both electron candidates passing only the “loose” identification requirement and failing the “medium” requirement.

The \(t\bar{t}\) and diboson backgrounds are estimated from MC simulation and account for up to 5% and 9% of the selected events, respectively. The overall level of agreement between data and the sum of the signal and background predictions is shown in Fig. 1.

5. Cross-section measurement

The differential cross-section, \(d\sigma/dm_{ee}\), is measured in 13 bins of \(m_{ee}\) from 116 GeV to 1500 GeV in a fiducial region in which both electrons have transverse momentum \(p_{T} > 25\) GeV and lie within \(|\eta| < 2.5\). The cross-section and fiducial region are determined for two conventions regarding QED FSR corrections. For the Born-level result, the true (meaning without detector simulation) \(m_{ee}\) and electron kinematics are defined by the electrons originating from the \(Z/\gamma^{*}\) decay before FSR. At the dressed level, true final-state electrons after FSR are recombined with radiated photons within a cone of \(\Delta R = 0.1\).

The cross-section is calculated from

\[
\frac{d\sigma}{dm_{ee}} = \frac{N_{\text{data}} - N_{\text{bkg}}}{C_{\text{DY}} L_{\text{int}} T_{\text{bin}}} 1
\]

where \(N_{\text{data}}\) is the number of candidate events observed in a given bin of \(m_{ee}\) (of width \(T_{\text{bin}}\)), \(N_{\text{bkg}}\) is the total background in that bin and \(L_{\text{int}}\) is the integrated luminosity. The correction factor, \(C_{\text{DY}}\), takes into account the efficiency of the signal selection and bin migration effects. It also includes the small extrapolation (about 10% to 13%) over the small region in \(|\eta|\) that is excluded for reconstructed electron candidates (1.37 < \(|\eta|\) < 1.52 and 2.47 < \(|\eta|\) < 2.5). The correction factor is defined as the number of MC-generated events that pass the signal selection in a bin of reconstructed \(m_{ee}\), divided by the total number of generated events within the fiducial region, at the Born or dressed level, in the corresponding bin of true \(m_{ee}\). It is obtained from the \textsc{pythia} MC signal sample and corrected for differences in the reconstruction, identification and trigger efficiencies between data and MC simulation. The value of \(C_{\text{DY}}\) varies from 0.55 (0.57) in the lowest bin to 0.70 (0.73) in the highest bin at the Born (dressed) level.

The \(m_{ee}\) resolution varies from approximately 3% at low \(m_{ee}\) to 1% at high \(m_{ee}\). The purity, defined as the fraction of simulated events reconstructed in a given \(m_{ee}\) bin that have true \(m_{ee}\) in the same bin, ranges from 79% (82%) to 98% (98%) at the Born (dressed) level.

6. Systematic uncertainties

The main contributions to the systematic uncertainties are given in Table 1 and described below.

6.1. Background estimation

In the estimation of the dominant dijet and \(W + \text{jets}\) background, a systematic uncertainty of 11% is assigned to the \(E_{T}\) - and \(\eta\)-dependent fake rate, corresponding to the spread of this quantity as measured in the nine independent jet samples, in order to cover any possible bias introduced in the triggering of these background events. A further uncertainty on the fake rate of up to 11% arises due to the presence of remaining signal contamination in the background-enriched sample. The total systematic uncertainty on the fake rate combines with a smaller effect (around 5%) from signal contamination in the sample where the fake rate is applied, to give a total uncertainty on the resulting background estimate of up to 16%. An additional systematic uncertainty can arise if the fake rate differs for different sources of fake electrons and the relative contribution of the different sources is not the same in the data sample where the fake rate is measured and the sample of events to which it is applied. It is found that \(b\)-jets have a higher fake rate than jets initiated by gluons or light quarks, but that the fraction of \(b\)-jets is small and similar in both samples. Conservatively taking this additional source of uncertainty into account, the overall uncertainty on the background is enlarged to 20%.

This 20% is added in quadrature to the statistical uncertainty of the sample to which the fake rate is applied; the latter uncertainty dominates in the highest two \(m_{ee}\) bins. The resulting overall uncertainty on the cross-section from the dijet and \(W + \text{jets}\) background varies between 1.3% and 7.9%, depending on \(m_{ee}\).

Two alternative methods to estimate the dijet and \(W + \text{jets}\) background are considered as cross-checks. The first of these is similar to the baseline method but uses fake rates derived from loosely selected electrons collected by the EM signal trigger. Here the background-enriched sample is derived by employing a tag-and-probe technique selecting, among other requirements to suppress real electron contamination, a jet-like tag and a probe with the same charge. This method, being correlated to the baseline method due to the overlap of electron candidates passing the EM and jet triggers, yields very similar predictions with comparable systematic uncertainties. In the third method, the combined dijet plus \(W + \text{jets}\) background is estimated by performing a template fit to the isolation of the leading versus sub-leading electron. The background templates are obtained from data by reversing some of the identification requirements on one or both of the electrons, and the signal templates are made from the \textsc{pythia} DY sample. No additional systematic uncertainty is assigned from the two cross-checks, as their results are in agreement with the baseline method.

The uncertainties on the diboson and \(t\bar{t}\) background expectations include the theoretical uncertainties on their cross-sections, 5% for the dibosons and 10% for \(t\bar{t}\). At high \(m_{ee}\), the statistical uncertainties on the simulated samples dominate.
exceeding 50% in the highest bin for both processes. The resulting uncertainty on the cross-section is small compared to the data-driven 6.3. Energy scale and resolution

event yields. Following the prescription of Ref. [27], the scale factors to account for the difference in the EM signal-trigger efficiency between data and simulation are obtained by measuring the efficiency in data and MC events using a tag-and-probe method. The Z → e⁺e⁻ events are tagged by selecting events passing a single-electron trigger, thus providing one electron probe free of trigger bias to test against the signal-trigger requirements. The scale factors are very close to unity, and the electron isolation are obtained from data by reversing some of the requirements applied in the electron identification criteria. The isolation quantity is defined in a similar way to that used in the selection of the leading electron in the signal sample. The background isolation templates are then normalised to data in the tail of the distributions where no contribution from signal is expected, both before and after applying the identification requirements, in order to estimate the background fraction in the probe sample. The identification efficiencies are found to be consistent with those obtained by the method of Ref. [27] in the common measurement range, and are stable for electrons with Et up to 500 GeV.

The differences between the measured reconstruction and identification efficiencies and their values in MC simulation are taken as nτ- and Et-dependent scale factors with which the MC-derived CDY is corrected. An additional scale factor for the isolation requirement on the leading electron is also applied. Varying the scale factors for the electron reconstruction (identification) within their systematic uncertainties results in a change in the cross-section of up to 1.7% (2.6%).

6.4. Bin-by-bin correction

The results obtained from the bin-by-bin correction are cross-checked using an iterative Bayesian approach [34] and found to be consistent. In addition, a consistency test is performed by correcting the MC@NLO signal sample using the PYTHIA-derived CDY factor. The discrepancy between the sample corrected in this way and the true MC@NLO sample is about 1.5%. This is due to the slightly different shapes of the mττ distribution from the two generators, considered to represent the possible shape difference between data and the PYTHIA simulation. This is conservatively added as a systematic uncertainty on the cross-section in all mττ bins.

6.5. Trigger efficiency

Scale factors to account for the difference in the EM signal-trigger efficiency between data and simulation are obtained by measuring the efficiency in data and MC events using a tag-and-probe method. The Z → e⁺e⁻ events are tagged by selecting events passing a single-electron trigger, thus providing one electron probe free of trigger bias to test against the signal-trigger requirements. The scale factors are very close to unity, and the effect on the cross-section of varying them within their systematic uncertainties is approximately 1%.

6.6. MC statistics and MC modelling

The finite number of events in the MC samples from which the CDY factor is derived contribute an uncertainty of up to 2.4% on CDY and the computed cross-section. Systematic uncertainties are associated with the use of the K-factors and with the reweighting of the PYTHIA signal MC events in order to better match the transverse momentum distribution of the Z bosons and the mean number of interactions per bunch crossing in the data. The effect of a further reweighting of the vertex position distribution in the z direction, not applied by default when calculating CDY, is also taken as an uncertainty. These uncertainties enter into the calculation of CDY and result in an overall uncertainty on the cross-section of less than 1%. Excellent agreement in the FSR predictions between photos and sant [35,36] has been shown [37] and uncertainties related to the modelling of the detector response to low-energy photons from FSR are negligible.

6.7. Theoretical uncertainties

Several theoretical uncertainties apply to the extrapolation of the cross-section in |y| from the measured region to the fiducial

Table 1

<table>
<thead>
<tr>
<th>Source of uncertainty</th>
<th>Uncertainty [%] in mττ bin</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>116–130 GeV</td>
</tr>
<tr>
<td>Total background estimate (stat.)</td>
<td>0.1</td>
</tr>
<tr>
<td>Total background estimate (syst.)</td>
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</tr>
<tr>
<td>Electron energy scale &amp; resolution</td>
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<tr>
<td>Electron identification</td>
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<tr>
<td>Electron reconstruction</td>
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<tr>
<td>Bin-by-bin correction</td>
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<tr>
<td>Trigger efficiency</td>
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<tr>
<td>MC statistics (CDY stat.)</td>
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<tr>
<td>MC modelling</td>
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<tr>
<td>Theoretical uncertainty</td>
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<tr>
<td>Total systematic uncertainty</td>
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<tr>
<td>Luminosity uncertainty</td>
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<tr>
<td>Data statistical uncertainty</td>
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Table 2

<table>
<thead>
<tr>
<th>mττ [GeV]</th>
<th>dσ/dmττ (Born)</th>
<th>dσ/dmττ (dressed)</th>
<th>Stat. err. [%]</th>
<th>Syst. err. [%]</th>
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<tr>
<td>116–130</td>
<td>2.24 × 10⁻¹</td>
<td>2.15 × 10⁻¹</td>
<td>1.1</td>
<td>4.2</td>
</tr>
<tr>
<td>130–150</td>
<td>1.02 × 10⁻¹</td>
<td>0.94 × 10⁻¹</td>
<td>1.4</td>
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<td>150–170</td>
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<td>4.93 × 10⁻²</td>
<td>2.0</td>
<td>4.6</td>
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<tr>
<td>170–190</td>
<td>2.84 × 10⁻²</td>
<td>2.76 × 10⁻²</td>
<td>2.7</td>
<td>4.7</td>
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<tr>
<td>190–210</td>
<td>1.87 × 10⁻²</td>
<td>1.82 × 10⁻²</td>
<td>3.0</td>
<td>5.3</td>
</tr>
<tr>
<td>210–230</td>
<td>1.07 × 10⁻²</td>
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<tr>
<td>230–250</td>
<td>8.23 × 10⁻³</td>
<td>7.98 × 10⁻³</td>
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<tr>
<td>250–300</td>
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<td>4.52 × 10⁻³</td>
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<td>300–400</td>
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<td>400–500</td>
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<td>5.7</td>
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<tr>
<td>700–1000</td>
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<td>7.5</td>
</tr>
<tr>
<td>1000–1500</td>
<td>2.88 × 10⁻⁶</td>
<td>2.76 × 10⁻⁶</td>
<td>50</td>
<td>9.8</td>
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region and thus contribute to an additional uncertainty on $C_{D\gamma}$. To evaluate the effect of the choice of PDF, the calculation of $C_{D\gamma}$ using PYTHIA with its default PDF (MRSTMCa1) is compared to that obtained after reweighting to CT10 (NLO) and HERAPDF1.5 [38] (NLO). The largest difference between the reweighted results and the default is taken as the systematic uncertainty, and amounts to 0.2%. A further systematic uncertainty is calculated using the MC@NLO sample reweighted to the 52 CT10 eigenvector error sets, the result being 0.5% at most. Finally, comparisons are made between PYTHIA reweighted to the CT10 PDF and MC@NLO (which uses as default CT10), and cross-checked using FEWZ 2.1 at NLO using the CT10 PDF. The effect is at most 0.3%. These systematic uncertainties, which each have a different dependence on $m_{ee}$, are added in quadrature and together give a 0.2–0.5% uncertainty on the cross-section.

The contributions from the above sources of systematic uncertainty to the uncertainty on the measured cross-section are summarised in Table 1 for the lowest and highest bin in the $m_{ee}$ range considered. The overall systematic uncertainty, excluding the luminosity uncertainty of 1.8% [39], rises from 4.2% in the lowest $m_{ee}$ bin to 9.8% in the highest $m_{ee}$ bin. The data statistical uncertainties increase from 1.1% to 50%.

7. Results and comparison to theory

The cross-sections obtained in the fiducial region (electron $p_T > 25$ GeV and $|\eta| < 2.5$) at the Born and dressed levels are given in Table 2. The difference between the two results is at most 4%. The precision of the measurement is limited by the statistical uncertainty on the data for $m_{ee} > 400$ GeV.

Fig. 2 shows the results at the dressed level, where they are compared to the predictions from PYTHIA, MC@NLO and SHERPA. No corrections have been applied to the generator-level predictions; instead, the prediction of each generator has been scaled globally to match the total number of events observed in data. The resulting scale factors are 1.23 for PYTHIA, 1.08 for MC@NLO and 1.39 for SHERPA. As expected, the only prediction at NLO in pQCD, from the MC@NLO generator, yields the scale factor closest to unity. The overall shape of the $m_{ee}$ distribution from all three generators is consistent with the data.

Fig. 3 shows the differential cross-section at the Born level compared to calculations in the FEWZ 3.1 framework using various recent NNLO PDFs. The FEWZ 3.1 framework allows the (N)NLO QCD corrections to lepton pair production to be combined with the NLO electroweak corrections. It has been verified at NLO in QCD that the choice of the electroweak scheme, $G_\mu$ or $\alpha(m_Z)$ as introduced in Ref. [40], has an effect of at most 0.4% on the calculated cross-section after applying NLO electroweak corrections. The electroweak-corrected NNLO QCD predictions shown are calculated using the $G_\mu$ scheme. The electroweak corrections include a positive contribution from the irreducible, non-resonant photon-induced background, i.e., $\gamma\gamma \rightarrow e^+e^-$. This contribution is estimated at leading order (LO) using the MRST2004qed [41] PDF, currently the only set available that includes QED corrections to the proton PDF, by taking the average of the predictions obtained under the current and constituent quark mass schemes. The symmetric difference between the average and either scheme is assigned as the corresponding uncertainty on this additive correction, being approximately 50% and representing a 3% uncertainty on the cross-section prediction in the highest $m_{ee}$ bin. The electroweak and photon-induced corrections were verified by SANC [35,36]. An additional small correction arises from single-boson production in which the final-state charged lepton radiates a real $W$ or $Z$ boson [42]. This is estimated using MADGRAPH 5 [43], following the prescription outlined in Ref. [42], to be at most 2%, in the highest $m_{ee}$ bin.

It can be seen in Fig. 3 that the deviations between the MSTW2008 [2] and the CT10 [22], HERAPDF1.5 [38] and NNPDF2.3 [44] predictions are covered by the total uncertainty band assigned to the MSTW2008 prediction, which is dominated by the combined 68% confidence level (CL) PDF and $\alpha_s$ variation. At low $m_{ee}$ the ABM11 [45] prediction lies above this theoretical uncertainty band, in part due to the ABM11 PDF set using a value of $\alpha_s$ outside of the 68% CL variation. The renormalisation and factorisation scale uncertainties contribute at most 1% to the theoretical uncertainty band in the highest $m_{ee}$ bin, having been evaluated by varying both scales up or down together by a factor of two, using VRAP [46]. The size of the photon-induced contribution is similar to the sum of the PDF, $\alpha_s$ and scale uncertainties as can be seen in the lower panel of Fig. 3(left), where the nominal calculation using the MSTW2008 PDF set is compared to the case where this contribution is not taken into account.

In the region where the precision of the measurement is limited by systematic uncertainties, $m_{ee} < 400$ GeV, the data generally lie above the FEWZ calculations. However, assuming that all systematic uncertainties, except those of statistical origin on the background and on $C_{D\gamma}$ (Table 1), are fully correlated bin-to-bin, the comparison between data and the different predictions over the full mass range yields chi-squared values of 13.9 for MSTW2008, 18.9 for CT10, 13.5 for HERAPDF1.5, 14.7 for ABM11 and 14.8 for NNPDF2.3, for the 13 data points, indicating compatibility between the theory and data.

8. Summary

Using 4.9 fb$^{-1}$ of data from $pp$ collisions at a centre-of-mass energy of $\sqrt{s} = 7$ TeV, the invariant mass distribution of $e^+e^-$ pairs from $WZ$ production has been measured at ATLAS in the range $116 < m_{ee} < 1500$ GeV, for electrons with $p_T > 25$ GeV and $|\eta| < 2.5$ and combined with the total number of events observed in data.
$|\eta| < 2.5$. Comparisons have been made to the predictions of the PYTHIA, MC@NLO and SHERPA MC generators, after scaling them globally to match the total number of events observed in data. The MC predictions are consistent with the shape of the measured $m_{\ell\ell}$ distribution. The predictions of the FEWZ 3.1 framework using five PDF sets at NNLO have also been studied. The framework combines calculations at NNLO QCD with NLO electroweak corrections, to which LO photon-induced corrections and real $W$ and $Z$ boson emission in single-boson production have been added. The resulting predictions for all PDFs are consistent with the measured differential cross-section, although the data are systematically above the theory. The data have the potential to constrain PDFs, in particular for antiquarks at large $x$, in the context of a PDF fit involving the world data sensitive to the proton structure.

Acknowledgements

We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently.

We acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWF and FWF, Austria; ANAS, Azerbaijan; STFC, Belgium; CNPq and FAPESP, Brazil; NSERC, NRC and CFI, Canada; CERN; CONICYT, Chile; CAS, MOST and NSFC, China; COFECYT/CONICET, Colombia; MSMT CR, MPO CR and VSC CR, Czech Republic; DDIR, DLR and RAL, Germany; INFN, Italy; MEXT and JSPS, Japan; CNRS, CEA and IRFU, France; CEA, Morocco; FOM and NWO, Netherlands; BRF and RCN, Norway; MNiSW, Poland; GRICES and FCT, Portugal; MNCH and STFC, Switzerland; STFC, the Royal Society and Leverhulme Trust, United Kingdom; DOE and NSF, United States of America.

The crucial computing support from all WLCG partners is acknowledged gratefully, in particular from CERN and the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CC-IN2P3 (France), KIT/GridKA (Germany), INFN–CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (UK) and BNL (USA) and in the Tier-2 facilities worldwide.

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References

21 Physikalisches Institut, University of Bonn, Bonn, Germany
22 Department of Physics, Boston University, Boston, MA, United States
23 Department of Physics, Brandeis University, Waltham, MA, United States
24 (a) Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro; (b) Federal University of Juiz de Fora (UFJF), Juiz de Fora; (c) Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei; (d) Instituto de Fisica, Universidade de Sao Paulo, Sao Paulo, Brazil
25 Physics Department, Brookhaven National Laboratory, Upton, NY, United States
26 (a) National Institute of Physics and Nuclear Engineering, Bucharest; (b) National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj Napoca; (c) Polytechnica Bucharest, Bucharest; (d) West University in Timisoara, Timisoara, Romania
27 Departamento de Fisica, Universidad de Buenos Aires, Buenos Aires, Argentina
28 Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
29 Department of Physics, Carleton University, Ottawa ON, Canada
30 CERN, Geneva, Switzerland
31 Enrico Fermi Institute, University of Chicago, Chicago, IL, United States
32 (a) Departamento de Fisica, Pontificia Universidad Catolica de Chile, Santiago; (b) Departamento de Fisica, Universidad Tecnica Federico Santa Maria, Valparaiso, Chile
33 (a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; (b) Department of Modern Physics, University of Science and Technology of China, Anhui; (c) Department of Physics, Nanjing University, Jiangsu; (d) School of Physics, Shandong University, Shandong; (e) Physics Department, Shanghai Jiao Tong University, Shanghai, China
34 Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France
35 Nevis Laboratory, Columbia University, Irvington, NY, United States
36 Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark
37 (a) INFN Gruppo Collegato di Cosenza; (b) Dipartimento di Fisica, Università della Calabria, Rende, Italy
38 (a) AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow; (b) Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland
39 The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland
40 Physics Department, Southern Methodist University, Dallas, TX, United States
41 Physics Department, University of Texas at Dallas, Richardson, TX, United States
42 DESY, Hamburg and Zeuthen, Germany
43 Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
44 Institut für Kern- und Teilchenphysik, Technical University Dresden, Dresden, Germany
45 Department of Physics, Duke University, Durham, NC, United States
46 SUPA – School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
47 INFN Laboratori Nazionali di Frascati, Frascati, Italy
48 Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany
49 Section de Physique, Université de Genève, Geneva, Swiss-Ph"{a}s
50 (a) INFN Sezione di Genova; (b) Dipartimento di Fisica, Università di Genova, Genova, Italy
51 (a) E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi; (b) High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
52 Il Physikalisches Institut, Justus-Liebig-Universität Gießen, Gießen, Germany
53 SUPA – School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
54 Il Physikalisches Institut, Georg-August-Universität, Göttingen, Germany
55 Laboratoire de Physique Subatomique et de Cosmologie, Université Joseph Fourier and CNRS/IN2P3 and Institut National Polytechnique de Grenoble, Grenoble, France
56 Department of Physics, Hampton University, Hampton, VA, United States
57 Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, MA, United States
58 (a) Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; (b) ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany
59 Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
60 Department of Physics, Indiana University, Bloomington, IN, United States
61 Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
62 University of Iowa, Iowa City, IA, United States
63 Department of Physics and Astronomy, Iowa State University, Ames, IA, United States
64 Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia
65 KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
66 Graduate School of Science, Kobe University, Kobe, Japan
67 Faculty of Science, Kyoto University, Kyoto, Japan
68 Kyushu University of Education, Kyotanabe, Japan
69 Department of Physics, Kyushu University, Fukuoka, Japan
70 Instituto de Fisica La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
71 Physics Department, Lancaster University, Lancaster, United Kingdom
72 (a) INFN Sezione di Lecce; (b) Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy
73 Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom
74 Department of Physics, Jožef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia
75 School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom
76 Department of Physics, Royal Holloway University of London, Egham, United Kingdom
77 Department of Physics and Astronomy, University College London, London, United Kingdom
78 Louisiana Tech University, Ruston, LA, United States
79 Laboratoire de Physique Nucl"{e}aire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France
80 Fysiska institutionen, Lund universitet, Lund, Sweden
81 Departamento de Fisica Teorica C-15, Universidad Autonoma de Madrid, Madrid, Spain
82 Institut für Physik, Universität Mainz, Mainz, Germany
83 School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
84 CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
85 Department of Physics, University of Massachusetts, Amherst, MA, United States
86 Department of Physics, McGill University, Montreal, QC, Canada
87 School of Physics, University of Melbourne, Victoria, Australia
88 Department of Physics, The University of Michigan, Ann Arbor, MI, United States
89 Department of Physics and Astronomy, Michigan State University, East Lansing, MI, United States
90 (a) INFN Sezione di Milano; (b) Dipartimento di Fisica, Università di Milano, Milano, Italy
Gif-sur-Yvette, France
National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Belarus
Department of Physics, Massachusetts Institute of Technology, Cambridge, MA, United States
Group of Particle Physics, University of Montreal, Montreal, QC, Canada
P.N. Lebedev Institute of Physics, Academy of Sciences, Moscow, Russia
Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia
Moscow Engineering and Physics Institute (MSPU), Moscow, Russia
D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
Nagasaki Institute of Applied Science, Nagasaki, Japan
Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan
INFN Sezione di Napoli, (b) Dipartimento di Scienze Fisiche, Università di Napoli, Napoli, Italy
Department of Physics and Astronomy, University of New Mexico, Albuquerque, NM, United States
Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
Department of Physics, Northern Illinois University, DeKalb, IL, United States
Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia
Department of Physics, New York University, New York, NY, United States
Ohio State University, Columbus, OH, United States
Faculty of Science, Okayama University, Okayama, Japan
Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, OK, United States
Department of Physics, Oklahoma State University, Stillwater, OK, United States
Palacký University, Olomouc, Czech Republic
Center for High Energy Physics, University of Oregon, Eugene, OR, United States
LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France
Graduate School of Science, Osaka University, Osaka, Japan
Department of Physics, University of Oslo, Oslo, Norway
Department of Physics, Oxford University, Oxford, United Kingdom
INFN Sezione di Pavia, (b) Dipartimento di Fisica, Università di Pavia, Pavia, Italy
Department of Physics, University of Pennsylvania, Philadelphia, PA, United States
Petersburg Nuclear Physics Institute, Gatchina, Russia
INFN Sezione di Pisa, (b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA, United States
Laboratorio de Instrumentación e Física Experimental de Partículas – LIP, Lisboa, Portugal,
Departamento de Física Teórica y del Cosmos y CAIFPE, Universidad de Granada, Granada, Spain
Institute of Physics, Academy of Sciences of the Czech Republic, Prague, Czech Republic
Czech Technical University in Prague, Prague, Czech Republic
Faculty of Mathematics and Physics, Charles University in Prague, Prague, Czech Republic
State Research Center Institute for High Energy Physics, Protvino, Russia
Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
Physics Department, University of Regina, Regina, SK, Canada
Ritsumeikan University, Kusatsu, Shiga, Japan
INFN Sezione di Roma I, (b) Dipartimento di Fisica, Università La Sapienza, Roma, Italy
INFN Sezione di Roma Tor Vergata, (b) Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
INFN Sezione di Roma Tre, (b) Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy
Faculté des Sciences Ain Chock, Université de Sciences et de Techniques d’Alger, Algiers, Algeria,
Centre National de l’Énergie des Sciences Nucléaires, Abidjan, Côte d’Ivoire,
Laboratoire de Physique des États Condensés, Université de Liège, Belgium
Department of Physics and Astronomy, University of Cape Town, Cape Town, South Africa
Institute of Physics, Academy of Sciences of the Slovak Republic, Bratislava, Slovakia
Department of Physics, University of the Witwatersrand, Johannesburg, South Africa
School of Physics, University of the Witwatersrand, Johannesburg, South Africa
Physics Department, Stockholm, Sweden
The Oskar Klein Centre, Stockholm, Sweden
Physics Department, Royal Institute of Technology, Stockholm, Sweden
Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook, NY, United States
Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom
School of Physics, University of Sydney, Sydney, Australia
Institute of Physics, Academia Sinica, Taipei, Taiwan
Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel
Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan
Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
Department of Physics, University of Toronto, Toronto, ON, Canada
TRIUMF, Vancouver, BC, (b) Department of Physics and Astronomy, York University, Toronto, ON, Canada
Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan
Department of Physics and Astronomy, Tufts University, Medford, MA, United States
Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia