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Search for excited electrons and muons in $\sqrt{s} = 8$ TeV proton–proton collisions with the ATLAS detector

The ATLAS Collaboration

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E-mail: atlas.publications@cern.ch

Abstract. The ATLAS detector at the Large Hadron Collider is used to search for excited electrons and excited muons in the channel $pp \rightarrow \ell\ell^* \rightarrow \ell\ell\gamma$, assuming that excited leptons are produced via contact interactions. The analysis is based on 13 fb$^{-1}$ of $pp$ collisions at a centre-of-mass energy of 8 TeV. No evidence for excited leptons is found, and a limit is set at the 95% credibility level on the cross section times branching ratio as a function of the excited-lepton mass $m_{\ell^*}$. For $m_{\ell^*} \geq 0.8$ TeV, the respective upper limits on $\sigma B(\ell^* \rightarrow \ell\gamma)$ are 0.75 and 0.90 fb for the $e^*$ and $\mu^*$ searches. Limits on $\sigma B$ are converted into lower bounds on the compositeness scale $\Lambda$. In the special case where $\Lambda = m_{\ell^*}$, excited-electron and excited-muon masses below 2.2 TeV are excluded.
1. Introduction

Although the Standard Model (SM) of particle physics is very successful at describing a large range of phenomena, it does not provide an explanation for the generational structure and mass hierarchy of quarks and leptons. Fermion compositeness models [1–6] aim at reducing the number of fundamental matter constituents by describing SM fermions as bound states of more-elementary particles. The existence of excited states would then be a direct consequence of the fermion substructure.

This paper reports on searches for excited electrons (e*) and excited muons (µ*) using 13 fb$^{-1}$ of pp collision data at a centre-of-mass energy of $\sqrt{s} = 8$ TeV recorded in 2012 with the ATLAS detector at the Large Hadron Collider (LHC). Searches are based on a benchmark model [6] that describes excited-fermion interactions using an effective Lagrangian. Excited leptons (ℓ*) would be predominantly produced via four-fermion contact interactions, and are expected to decay into a lepton and a gauge boson, or a lepton and a pair of fermions. All unknown couplings of the model are set as in [6]. The contact interaction is then described by the Lagrangian

$$L_{\text{contact}} = \frac{2\pi}{\Lambda^2} j^\mu j^\nu, \quad j_\mu = \bar{f}_L \gamma_\mu f_L + \bar{f}_L^* \gamma_\mu f_L^* + \bar{f}_L \gamma_\mu f_L + \text{h.c.},$$

where $\Lambda$ is the compositeness scale, $j_\mu$ is the fermion current for ground states (f) and excited states (f*), and ‘h.c.’ stands for Hermitian conjugate. The gauge-mediated decays are given by the Lagrangian

$$L_{\text{GM}} = \frac{1}{2\Lambda} \bar{\ell}_R^* \sigma_\mu^a \left[ \frac{g}{2} W_\mu^a + \frac{g'}{2} B_\mu^a \right] \ell_L + \text{h.c.},$$

where $\ell$ and $\ell^*$ are the lepton and excited-lepton fields, respectively, $W_\mu^a$ and $B_\mu^a$ are the $SU(2)_L$ and $U(1)_Y$ field-strength tensors and $g$ and $g'$ are the corresponding gauge couplings. The searches described here focus on the single-production mechanism ($q\bar{q} \rightarrow \ell^+\ell^-\bar{\ell}^\tau\ell^\tau$) and the electromagnetic radiative decay mode $\ell^* \rightarrow \ell \gamma$. The signature thus consists of events containing two same-flavour, opposite-charge leptons and a photon ($\ell^+\ell^-\gamma$ final state). The kinematic properties of the signal are determined by the excited-lepton mass ($m_{\ell^*}$) and the compositeness
scale ($\Lambda$). Due to unitarity constraints on contact interactions [6, 7], the model does not apply in the regime $m_{\ell^*} > \Lambda$. For most of the parameter space, the presence of excited leptons would appear as a peak in the lepton–photon mass spectrum. However, for $m_{\ell^*} \simeq \Lambda$, the width of the resonance can become significantly larger than the experimental mass resolution of the detector. To avoid this complication as well as the lepton–photon pairing ambiguity, a search is performed for an excess in the $m_{\ell\ell\gamma}$ invariant mass ($m_{\ell\ell\gamma}$) spectrum.

Previous searches at LEP [8–11], HERA [12, 13] and the Tevatron [14–17] have found no evidence for excited leptons. For the case where $m_{\ell^*} = \Lambda$, $e^*$ and $\mu^*$ masses below 1.9 TeV have been excluded by both the ATLAS [18] and CMS [19] experiments using $\sqrt{s} = 7$ TeV data.

2. ATLAS detector

The ATLAS detector [20] consists of an inner tracking system surrounded by a superconducting solenoid, electromagnetic and hadronic calorimeters, and a muon spectrometer. It has a forward–backward symmetric cylindrical geometry$^1$ and nearly $4\pi$ coverage in solid angle. Charged-particle tracks and vertices are reconstructed in silicon-based pixel and microstrip tracking detectors that cover $|\eta| < 2.5$ and transition radiation detectors extending to $|\eta| < 2.0$. A hermetic calorimeter system, which covers $|\eta| < 4.9$, surrounds the solenoid. The liquid-argon electromagnetic calorimeter, which plays an important role in electron and photon identification and measurement, is finely segmented for $|\eta| < 2.5$ to provide excellent energy and position resolution. Hadron calorimetry is provided by an iron–scintillator tile calorimeter in the central pseudorapidity range $|\eta| < 1.7$ and a liquid-argon calorimeter with copper or tungsten as absorber material in the pseudorapidity range $1.5 < |\eta| < 4.9$. A spectrometer is installed outside the calorimeter to identify muons and measure their momenta with high precision. The toroidal magnetic field of the muon spectrometer is provided by three air-core superconducting magnet systems: one for the barrel and one per endcap, each composed of eight coils. Three layers of drift-tube chambers and/or cathode-strip chambers provide precise coordinate measurement in the bending plane ($r$–$z$) in the region $|\eta| < 2.7$. A system consisting of resistive-plate chambers for $|\eta| < 1.05$ and thin-gap chambers for $1.05 < |\eta| < 2.7$ provides measurement of the $\phi$ coordinate. It also provides triggering capability up to $|\eta| = 2.4$.

3. Simulated samples

The simulation of the excited-lepton signal is based on calculations from [6]. Signal samples are generated at leading order (LO) with CompHEP 4.5.1 [21] using MSTW2008 LO [22] parton distribution functions (PDFs). CompHEP is interfaced with Pythia version 8 [23, 24] for the simulation of parton showers and hadronization. The emission of photons via initial-state radiation and final-state radiation (FSR) is handled by Pythia. Only the single production of excited leptons followed by a $\ell^* \rightarrow \ell\gamma$ decay is simulated.

For both the $e^*$ and $\mu^*$ searches, the dominant background arises from Drell–Yan processes accompanied by either a prompt photon from initial- or final-state radiation ($Z + \gamma$) or a jet misidentified as a photon ($Z + \text{jets}$). The $Z + \gamma$ background results in the same final state as the

$^1$ ATLAS uses a right-handed coordinate system with the $z$-axis along the beam pipe. The $x$-axis points to the centre of the LHC ring, and the $y$-axis points upward. Cylindrical coordinates ($r, \phi$) are used in the transverse plane, $\phi$ being the azimuthal angle. The pseudorapidity is defined in terms of the polar angle $\theta$ as $\eta = -\ln \tan(\theta/2)$. 

signal, whereas the $Z$ + jets background is suppressed by imposing stringent requirements on the quality and isolation of the reconstructed photon. Small contributions from $t\bar{t}$ and diboson ($WW$, $WZ$ and $ZZ$) production are also present in both channels. In the electron channel, the $W + \gamma +$ jets background contributes to the $ee\gamma$ selection when a jet is misidentified as an electron. Backgrounds from $W +$ jets and multi-jet events, including semileptonic decays of heavy-flavour hadrons, are suppressed by requiring the leptons and the photon to be isolated, and have negligible contribution to the total background after selection.

The $Z + \gamma$ sample is generated with SHERPA 1.4.1 [25] using CT10 [26] PDFs and includes the LO emission of up to three partons in the initial state. To avoid phase-space regions where matrix elements diverge, the angular separation between the photon and each lepton is required to be $\Delta R(\ell, \gamma) = \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2} > 0.1$ and the transverse momentum of the photon ($p_T^\gamma$) is required to be above 10 GeV. To ensure adequate statistics, 1.2 million events (equivalent to 37 fb$^{-1}$) were generated for each of the electron and muon channels.

The $Z +$ jets and $W + \gamma +$ jets backgrounds are generated with ALPGEN 2.13 [27] using CTEQ6L1 [28] PDFs. The $t\bar{t}$ background is produced with MC@NLO 3.41 [29] with CT10 PDFs. In both cases, JIMMY 4.31 [30] is used to describe multiple parton interactions and HERWIG 6.510 [31] is used to simulate the remaining underlying event, parton showers and hadronization. The diboson processes are generated with POWHEG [32] and PYTHIA using CT10 PDFs. For all these samples, FSR is handled by PHOTOS [33]. To remove overlaps between the $Z +$ jets and $Z + \gamma$ samples, $Z +$ jets events with prompt photons are rejected if $p_T^\gamma > 10$ GeV and $\Delta R(\ell, \gamma) > 0.1$. The predictions for $Z +$ jets and $W + \gamma +$ jets backgrounds are normalized using the data-driven techniques described in section 5. Cross sections for diboson processes are evaluated at next-to-LO [34] and the $t\bar{t}$ cross section is calculated at approximate-next-to-next-to-LO [35], with uncertainties of 5% and $+10%$ $-9%$, respectively.

The generated samples are processed using a detailed detector simulation [36] based on GEANT4 [37] to propagate the particles and account for the detector response. Monte Carlo (MC) minimum-bias events are overlaid on both the signal and background processes to simulate the effect of additional $pp$ collisions (pile-up). Simulated events are weighted so that the distribution of the expected number of interactions per event agrees with the data, with an average of 20 interactions per bunch crossing.

4. Data and selection

The data were collected between April and October 2012 during stable-beam periods of $\sqrt{s} = 8$ TeV $pp$ collisions, and correspond to an integrated luminosity of 13.0 fb$^{-1}$ for the electron channel and 12.8 fb$^{-1}$ for the muon channel [38]. For the $e^+$ search, a trigger relying only on calorimetric information is used to select events. It requires two electromagnetic clusters with transverse momentum ($p_T$) thresholds of 35 and 25 GeV for the leading and subleading clusters, respectively, with loose shower-shape requirements aiming to select electrons and photons. For the $\mu^+$ search, a single-muon trigger is used. It requires a track to be reconstructed in both the muon spectrometer and the inner detector with a combined track $p_T > 24$ GeV.

Offline, events are selected if they contain at least two lepton candidates and a photon candidate. A primary vertex with at least three associated charged-particle tracks with $p_T > 0.4$ GeV is also required. If several vertices fulfill this requirement, the vertex with the largest $\Sigma p_T^2$ is selected, where the sum is over all reconstructed tracks associated with the vertex.
Each electron candidate is formed from a cluster of cells in the electromagnetic calorimeter associated with a charged-particle track in the inner detector. For the $e^+$ search, two electron candidates are required. Their transverse momentum ($p_T^e$) must satisfy $p_T^e > 40\, \text{GeV}$ ($30\, \text{GeV}$) for the leading (subleading) electron. Both electrons must be reconstructed within the range $|\eta| < 2.47$ and not in the transition region $1.37 < |\eta| < 1.52$ between the barrel and endcap calorimeters. The ATLAS medium electron identification criteria [39] for the transverse shower shape, the longitudinal leakage into the hadronic calorimeter and the association with an inner-detector track are applied to the cluster. The electron energy is obtained from the calorimeter measurement, and its direction is given by the associated track. A hit in the innermost layer of the pixel detector is required (if an active pixel layer is traversed) to suppress background from photon conversions. To suppress background from jets, the highest-$p_T$ electron is required to be isolated by demanding that the sum of the transverse energies in the cells around the electron direction in a cone of radius $\Delta R = 0.2$ be less than 7 GeV. The core of the electron energy deposition is excluded and the sum is corrected for transverse shower leakage and pile-up from additional $pp$ collisions to make the isolation variable essentially independent of $p_T^e$. The electron trigger and reconstruction efficiencies are evaluated using tag-and-probe techniques with $Z \rightarrow ee$ events [39] for data and MC simulation. Correction factors are extracted in several $\eta \times p_T^e$ bins and applied to the simulation. In cases where more than two electrons are found to satisfy the above requirements, the pair with the largest invariant mass is chosen. No requirement is applied to the electric charge of the electrons, as it could induce an inefficiency in the signal selection for high-$p_T$ electrons due to charge misidentification.

Each muon candidate has to be reconstructed independently in both the inner detector and the muon spectrometer. Its momentum is determined from a combined fit to these two measurements. For the $\mu^+$ search, two muon candidates with a transverse momentum ($p_T^\mu$) above 25 GeV are required. Both muons must have a minimum number of hits in the inner detector and hits in each of the inner, middle and outer layers of the muon spectrometer. This requirement, which restricts the muon acceptance to $|\eta| < 2.5$, guarantees a precise momentum measurement. Muons with hits in the barrel–endcap overlap regions of the muon spectrometer ($1.05 \lesssim |\eta| \lesssim 1.4$) are discarded because of the limited coverage with drift-tube chambers in this angular range. To suppress background from cosmic rays, the muon tracks are required to have transverse and longitudinal impact parameters $|d_0| < 0.2$ and $|z_0| < 1\, \text{mm}$ with respect to the selected primary vertex. To reduce background from heavy-flavour hadrons, each muon is required to be isolated such that $\Sigma p_T / p_T^\mu < 5\%$, where the sum is over inner-detector tracks with $p_T > 1\, \text{GeV}$ that are contained in a cone of radius $\Delta R = 0.3$ surrounding the candidate muon track, the latter being excluded from the sum. The muon trigger and reconstruction efficiencies are evaluated using tag-and-probe techniques with $Z \rightarrow \mu\mu$ events [40], and $\eta$-dependent corrections to be applied to the simulation are determined. The two muons are additionally required to have opposite electric charge. In cases where more than one pair of muons are found to satisfy the above requirements, the pair with the largest invariant mass is considered.

Each photon candidate is formed from a cluster of cells in the electromagnetic calorimeter. A photon can be reconstructed either as an unconverted photon, with no associated track, or as a photon that converted to an electron–positron pair, associated with one or two tracks. The presence of at least one photon candidate with $p_T^\gamma > 30\, \text{GeV}$ and $|\eta| < 2.37$ is required in both channels. As for electrons, photons within the transition region between the barrel and endcap calorimeters are excluded. Photon candidates are required to satisfy the ATLAS...
tight photon definition [41]. This selection includes constraints on the energy leakage into the hadronic calorimeter as well as stringent requirements on the energy distribution in the first and second sampling layers of the electromagnetic calorimeter. These requirements increase the purity of the selected photon sample by rejecting most of the jet background, including jets with a leading neutral hadron (usually a $\pi^0$) that decays into a pair of collimated photons. The photon-identification efficiency and shower shapes in the electromagnetic calorimeter are studied using FSR photons from Z boson decays with loose lepton–photon separation requirements. Shower shapes are then adjusted in the simulation so that the resulting photon-identification efficiency matches the efficiency measured in data [42].

To further reduce background from misidentified jets, photon candidates are required to be isolated by demanding that either $E_T^{\text{iso}} < 10$ GeV or $E_T^{\text{iso}} / p_T^\gamma < 1\%$, where $E_T^{\text{iso}}$ is the sum of the transverse energies of the clusters within a cone of radius $\Delta R = 0.4$ surrounding the photon. As for the electron isolation, the clusters from the photon energy deposition are excluded and the sum is corrected for transverse shower leakage and pile-up. The relative-isolation criterion reduces the efficiency loss for high-$p_T$ photons ($p_T^\gamma > 1$ TeV). Since the photon and the leptons are expected to be well separated for the excited-lepton signal, only photons satisfying $\Delta R(\ell, \gamma) > 0.7$ are retained. This requirement is effective at suppressing Drell–Yan events with FSR photons that are typically highly collimated with the leptons. If more than one photon candidate in an event satisfy the above requirements, the one with the largest $p_T$ is used in the search.

Finally, two additional requirements are applied to drastically reduce the background level. The first one, referred to as the ‘Z veto’ in the following, requires the dilepton mass to satisfy $m_{\ell\ell} > 110$ GeV. The second is a variable lower bound on the dilepton–photon mass that defines the signal search region. As a result of optimization studies, the signal region for $m_{\ell\gamma} < 900$ GeV is $m_{\ell\gamma} > m_{\ell\gamma} + 150$ GeV. For $m_{\ell\gamma} \geq 900$ GeV, it is fixed to $m_{\ell\gamma} > 1050$ GeV. The signal efficiency for these two requirements is above 98% for $m_{\ell\gamma} \geq 200$ GeV.

The total signal acceptance times efficiency ($A \times \epsilon$) is shown in figure 1 as a function of the excited-lepton mass. For low values of $m_{\ell\gamma}$, the photon and the leptons tend to be produced more forward and have a softer $p_T$ spectrum than at high mass, which explains the decrease in $A \times \epsilon$. The lower geometrical acceptance in the muon channel is due to the requirement of hits in all three layers of precision chambers.

5. Background determination

Most of the background predictions are estimated with MC samples normalized with calculated cross sections and the measured integrated luminosity of the data. Because the misidentification of jets as photons is not accurately modelled in the simulation, the $Z +$ jets background is instead normalized to the data using a control region defined as $70 < m_{\ell\ell} < 110$ GeV, where the contribution from signal events is at most 3% for $m_{\ell\gamma} \geq 200$ GeV. In this control region, the number of $Z +$ jets events is estimated by subtracting from the data all simulated backgrounds except $Z +$ jets. The normalization of the $Z +$ jets MC sample is corrected accordingly by a scale factor, separately determined to be $0.53 \pm 0.10$ for both the electron and muon channels. The quoted uncertainty combines the statistical uncertainties on the data and simulated backgrounds and the uncertainty on the $Z + \gamma$ cross section. Other sources of uncertainty including the integrated luminosity and the cross sections of the $t\bar{t}$ and diboson processes are negligible.
Figure 1. Acceptance times efficiency ($A \times \epsilon$) of the excited-lepton selection as a function of the excited-lepton mass ($m_{\ell^*}$), evaluated for a compositeness scale of 5 TeV. The uncertainties correspond to the sum in quadrature of the statistical uncertainty and systematic uncertainties associated with the lepton and photon efficiencies.

Scale factors were evaluated in different $p_T^\gamma$ bins, and results are consistent within statistical uncertainties.

In the electron channel, the $W + \gamma + \text{jets}$ background is also normalized to the data because of the imperfect modelling of the jet-to-electron fake rate. For this background only, the identification criteria were relaxed for one electron to increase the MC statistics. The $W + \gamma + \text{jets}$ normalization is derived using a likelihood template fit to the data, in the same control region as for $Z + \text{jets}$. The fit is simultaneously performed on transverse mass distributions $m_T(e_1, E_T^{\text{miss}})$ and $m_T(e_2, E_T^{\text{miss}})$, where $E_T^{\text{miss}}$ denotes the magnitude of the missing transverse momentum, which is calculated [43] from calorimeter cells with $|\eta| < 4.9$ using the local energy calibration of electrons, photons, hadronically decaying $\tau$-leptons and jets. Cells belonging to clusters not associated with such reconstructed objects as well as cells associated with a muon candidate are also included. The transverse mass is $m_T(e_i, E_T^{\text{miss}}) = \sqrt{2p_T^{e_i}E_T^{\text{miss}}(1 - \cos \Delta\phi)}$, where $\Delta\phi$ is the angle between the transverse momentum of electron $i$ ($p_T^{e_i}$) and the missing transverse momentum. The only floating parameter in the fit is the scale factor of the $W + \gamma + \text{jets}$ background, which is found to be $0.22^{+0.25}_{-0.22}$ (stat $\oplus$ syst). Systematic uncertainties account for correlations between the two $m_T$ variables, the choice of control region in which the fit is performed, the loosening of the electron identification criteria, and the dependence on the $Z + \text{jets}$ scale factor.

The numbers of events in the control region ($70 < m_{\ell\ell} < 110 \text{ GeV}$) and the numbers after the $Z$ veto ($m_{\ell\ell} > 110 \text{ GeV}$) are shown in table 1 after scaling the $Z + \text{jets}$ background, as well as the $W + \gamma + \text{jets}$ background in the electron channel. In the control region, by construction, the total background is equal to the number of events in data. After the $Z$ veto, the observed data are found to be consistent with the background prediction. Good agreement is also observed between data and background in the control region for the lepton and photon kinematic...
Table 1. Data yields and background expectations in the control region and after the Z veto. The Z+jets and W+γ+jets backgrounds are scaled as described in the text. The uncertainties shown are purely from MC statistics, except for Z+jets and W+γ+jets where the statistical uncertainty on the associated scale factor is reported.

<table>
<thead>
<tr>
<th>Regions (GeV)</th>
<th>70 &lt; m_ee &lt; 110</th>
<th>m_ee &gt; 110</th>
<th>70 &lt; m_μμ &lt; 110</th>
<th>m_μμ &gt; 110</th>
</tr>
</thead>
<tbody>
<tr>
<td>Z + γ</td>
<td>1235 ± 25</td>
<td>208 ± 10</td>
<td>1067 ± 22</td>
<td>131 ± 8</td>
</tr>
<tr>
<td>Z + jets</td>
<td>371 ± 48</td>
<td>25 ± 7</td>
<td>334 ± 43</td>
<td>12 ± 3</td>
</tr>
<tr>
<td>t¯t, diboson</td>
<td>18 ± 1</td>
<td>19 ± 2</td>
<td>16 ± 1</td>
<td>6 ± 1</td>
</tr>
<tr>
<td>W + γ + jets</td>
<td>9 ± 9</td>
<td>21 ± 21</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>Total MC</td>
<td>1633 ± 55</td>
<td>273 ± 24</td>
<td>1417 ± 48</td>
<td>149 ± 8</td>
</tr>
<tr>
<td>Data</td>
<td>1633</td>
<td>263</td>
<td>1417</td>
<td>147</td>
</tr>
</tbody>
</table>

Figure 2. Distributions of the transverse momentum of the photon (p_T) for the electron (a) and muon (b) channels, in the control region defined by the dilepton mass range 70 < m_ℓℓ < 110 GeV. The background uncertainty corresponds to the sum in quadrature of the statistical uncertainties and the uncertainty in the data-driven Z+jets normalization.

distributions. In particular, figure 2 shows that the background prediction for the photon p_T distribution matches the data for both the e^+ and μ^+ searches.

Because only a small fraction of the simulated background events survive the m_ℓℓ > 110 GeV requirement, the m_ℓℓγ distributions of dominant backgrounds are separately fitted with an exponential function and extrapolated to the high-mass region. The binned results of these fits are used as final background estimates in the statistical analysis. The same operation is performed for the m_ℓγ distribution of each background, although in this case, the fit results are not used in any numerical analysis. The resulting background estimates are shown in figures 3 and 4 as functions of the invariant mass of the ℓγ and ℓℓγ systems, respectively. For table 1 and figures 2–4, the lower bound on m_ℓℓγ described in section 4 is not applied.
Figure 3. Distributions of the $\ell\ell\gamma$ invariant mass ($m_{\ell\ell\gamma}$) for the electron (a) and muon (b) channels after requiring the dilepton mass to satisfy $m_{\ell\ell} > 110\text{ GeV}$. Combinations with both the leading and subleading leptons are shown. The binned results of exponential fits are used for all backgrounds. The background uncertainty corresponds to the sum in quadrature of the statistical and systematic uncertainties. The last bin contains the sum of all entries with $m_{\ell\ell\gamma} > 875\text{ GeV}$.

Signal predictions for three different values of the excited-lepton mass ($m_{\ell^*}$) with a compositeness scale ($\Lambda$) of 10 TeV are also shown.

Figure 4. Distributions of the $\ell\ell\gamma$ invariant mass ($m_{\ell\ell\gamma}$) for the electron (a) and muon (b) channels after requiring the dilepton mass to satisfy $m_{\ell\ell} > 110\text{ GeV}$. The binned results of exponential fits are used for the $Z + \gamma$, $Z + \text{jets}$, $t\bar{t}$, and $W + \gamma + \text{jets}$ backgrounds. The background uncertainty combining the statistical and systematic uncertainties is displayed as the hatched area. The last bin contains the sum of all events with $m_{\ell\ell\gamma} > 1350\text{ GeV}$. Signal predictions for three different values of the excited-lepton mass ($m_{\ell^*}$) with a compositeness scale ($\Lambda$) of 10 TeV are also shown.
Table 2. Dominant uncertainties on the expected numbers of events for the lowest-mass search region, $m_{\ell\ell\gamma} > 350 \text{ GeV}$. The theory uncertainty reported for the background corresponds to the uncertainty on the $Z + \gamma$ cross section only.

<table>
<thead>
<tr>
<th>Source</th>
<th>$e^\star$ (%)</th>
<th>$\mu^\star$ (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Theory</td>
<td>1 $^{+25}_{-16}$</td>
<td>1 $^{+25}_{-16}$</td>
</tr>
<tr>
<td>Statistics</td>
<td>- 18</td>
<td>- 21</td>
</tr>
<tr>
<td>Luminosity</td>
<td>3 3</td>
<td>3 3</td>
</tr>
<tr>
<td>Efficiencies</td>
<td>5 5</td>
<td>5 5</td>
</tr>
</tbody>
</table>

### 6. Systematic uncertainties

The most important sources of uncertainty are discussed below and summarized in table 2. A large part of the background uncertainty comes from the $Z + \gamma$ cross-section calculation. It includes the renormalization and factorization scale uncertainties, obtained by varying independently each scale by a factor of two, as well as uncertainties in the PDFs and the strong coupling constant $\alpha_s$. These uncertainties are evaluated by generating $Z + \gamma$ SHERPA samples for the 52 CT10 eigenvector PDF sets, the four CT10.AS PDF sets corresponding to $\alpha_s = 0.116, 0.117, 0.119$ and 0.120, and the four combinations of scales. For $m_{\ell\ell\gamma} > 350 \text{ GeV}$ ($m_{\ell\ell\gamma} > 1050 \text{ GeV}$), the resulting uncertainty is $\pm^{25\%}_{16\%}$ ($^{+32\%}_{-18\%}$) for both channels. Cross-section uncertainties for the $t\bar{t}$ and diboson processes have a negligible impact on the total background uncertainty.

The statistical uncertainties associated with the $m_{\ell\ell\gamma}$ fits contribute to the background uncertainty at a comparable level at low mass, and become increasingly important at high mass. The sum in quadrature of fit uncertainties, including uncertainties on data-driven scale factors for the relevant backgrounds, increases from about $\pm 20\%$ for $m_{\ell\ell\gamma} > 350 \text{ GeV}$ in both channels to $\pm 65\%$ ($^{+200\%}_{-60\%}$) for $m_{\ell\ell\gamma} > 1050 \text{ GeV}$ in the $e^\star (\mu^\star)$ search. The main contributions come from the $Z + \gamma$ and $Z + \text{jets}$ backgrounds, as well as the $W + \gamma + \text{jets}$ background in the electron channel.

Experimental systematic uncertainties that affect both the signal and background yields include the uncertainty on the luminosity measurement and uncertainties in particle reconstruction and identification as described below.

The uncertainty on the integrated luminosity is 2.8%. It is based on a preliminary calibration of the luminosity scale derived from beam-separation scans [38] performed in November 2012.

The total uncertainty on the photon reconstruction and identification efficiencies is 4% [42]. The combination of uncertainties on the electron trigger, reconstruction, identification and isolation efficiencies results in a 2% uncertainty on both the signal efficiency and background level. The combined uncertainty on the trigger, reconstruction and identification efficiencies for muons is estimated to increase linearly as a function of $m_{\ell\star}$ to about 2% for $m_{\ell\star} = 2 \text{ TeV}$. This uncertainty is dominated by the impact of large energy loss from muon bremsstrahlung in the calorimeter. The sum in quadrature of the lepton and photon uncertainties for the lowest $m_{\ell\ell\gamma}$ threshold is shown in table 2. Uncertainties on the energy scale and resolution for final-state objects have a negligible effect on signal and background selection efficiencies.
Table 3. Data yields and background expectation as a function of a lower bound on $m_{ee\gamma}$ for the $e^*$ search. The uncertainties represent the sum in quadrature of the statistical and systematic uncertainties.

<table>
<thead>
<tr>
<th>$m_{ee\gamma}$ region (GeV)</th>
<th>$Z + \gamma$</th>
<th>Total background</th>
<th>Data</th>
</tr>
</thead>
<tbody>
<tr>
<td>$&gt;350$</td>
<td>53$^{+16}_{-14}$</td>
<td>69$^{+18}_{-16}$</td>
<td>60</td>
</tr>
<tr>
<td>$&gt;450$</td>
<td>27$^{+9}_{-7}$</td>
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<td>$&gt;550$</td>
<td>14$^{+5}_{-4}$</td>
<td>17$^{+7}_{-6}$</td>
<td>12</td>
</tr>
<tr>
<td>$&gt;650$</td>
<td>7.0$^{+3.6}_{-3.3}$</td>
<td>8.7$^{+4.8}_{-3.4}$</td>
<td>7</td>
</tr>
<tr>
<td>$&gt;750$</td>
<td>3.5$^{+2.5}_{-1.9}$</td>
<td>4.4$^{+3.4}_{-2.0}$</td>
<td>3</td>
</tr>
<tr>
<td>$&gt;850$</td>
<td>1.8$^{+1.8}_{-1.1}$</td>
<td>2.2$^{+2.5}_{-1.1}$</td>
<td>2</td>
</tr>
<tr>
<td>$&gt;950$</td>
<td>0.9$^{+1.2}_{-0.6}$</td>
<td>1.1$^{+1.7}_{-0.6}$</td>
<td>1</td>
</tr>
<tr>
<td>$&gt;1050$</td>
<td>0.4$^{+0.8}_{-0.3}$</td>
<td>0.5$^{+1.2}_{-0.4}$</td>
<td>1</td>
</tr>
</tbody>
</table>

Table 4. Data yields and background expectation as a function of a lower bound on $m_{\mu\mu\gamma}$ for the $\mu^*$ search. The uncertainties represent the sum in quadrature of the statistical and systematic uncertainties.

<table>
<thead>
<tr>
<th>$m_{\mu\mu\gamma}$ region (GeV)</th>
<th>$Z + \gamma$</th>
<th>Total background</th>
<th>Data</th>
</tr>
</thead>
<tbody>
<tr>
<td>$&gt;350$</td>
<td>33$^{+1.1}_{-0.9}$</td>
<td>40$^{+11}_{-10}$</td>
<td>32</td>
</tr>
<tr>
<td>$&gt;450$</td>
<td>17$^{+6}_{-5}$</td>
<td>21$^{+7}_{-6}$</td>
<td>12</td>
</tr>
<tr>
<td>$&gt;550$</td>
<td>8.7$^{+3.8}_{-3.6}$</td>
<td>11$^{+5}_{-4}$</td>
<td>3</td>
</tr>
<tr>
<td>$&gt;650$</td>
<td>4.4$^{+2.4}_{-2.2}$</td>
<td>5.9$^{+3.5}_{-2.6}$</td>
<td>3</td>
</tr>
<tr>
<td>$&gt;750$</td>
<td>2.2$^{+1.5}_{-1.3}$</td>
<td>3.2$^{+2.6}_{-1.6}$</td>
<td>2</td>
</tr>
<tr>
<td>$&gt;850$</td>
<td>1.1$^{+1.0}_{-0.8}$</td>
<td>1.7$^{+2.0}_{-0.9}$</td>
<td>1</td>
</tr>
<tr>
<td>$&gt;950$</td>
<td>0.6$^{+0.7}_{-0.6}$</td>
<td>0.9$^{+1.5}_{-0.6}$</td>
<td>0</td>
</tr>
<tr>
<td>$&gt;1050$</td>
<td>0.3$^{+0.5}_{-0.2}$</td>
<td>0.5$^{+1.0}_{-0.3}$</td>
<td>0</td>
</tr>
</tbody>
</table>

The impact of the $\ell^*$ decay width on the signal selection efficiency was also investigated. The decay width is computed with formulas given in [6]. It increases with $m_{\ell^*}$ and decreases with $\Lambda$, and over the $\Lambda- m_{\ell^*}$ region accessible in these searches, it ranges from $\simeq 1$ MeV to $\simeq 200$ GeV. Signal efficiencies were computed at the generator level for different values of $\Lambda$, and efficiency variations were observed to be at most 1%, which is negligible compared to the other uncertainties in the selection efficiency.

7. Results

The $m_{\ell\ell\gamma}$ distributions are shown in figure 4 for the data, the expected backgrounds, and three signal predictions. The expected and observed numbers of events in each of the search regions, used for the statistical analysis, are shown in tables 3 and 4 for the electron and muon
Figure 5. Upper limits at 95% CL on the cross section times branching ratio ($\sigma B$) as a function of the excited-lepton mass ($m_\ell^*$), for the electron (a) and muon (b) channels. LO signal predictions with uncertainties from renormalization and factorization scales and PDFs are shown for three different compositeness scales ($\Lambda$).

Figure 6. Exclusion limits in the compositeness scale ($\Lambda$) versus excited-lepton mass ($m_\ell^*$) parameter space for the electron (a) and muon (b) channels. The filled area is excluded at 95% CL. No limits are set in the dark shaded region $m_\ell^* > \Lambda$ where the model is not applicable.

channels, respectively. The uncertainties include both the statistical and systematic contributions as described earlier. The data are consistent with the background expectation, and no significant excess is observed in the signal region.

An upper limit on the cross section times branching ratio $\sigma (pp \to \ell\ell^*) \times B(\ell^* \to \ell\gamma)$ is determined for each channel and each $m_\ell^*$ hypothesis at the 95% credibility level (CL) using a Bayesian approach [44] with a flat positive prior for $\sigma B$. Systematic uncertainties are incorporated into the limit calculation as nuisance parameters with Gaussian priors.
Uncertainties in particle reconstruction and identification efficiencies as well as the uncertainty on the luminosity are fully correlated between signal and backgrounds. All other uncertainties are uncorrelated. The expected limit is evaluated as the median of the upper-limit distribution obtained with a set of background-only pseudo-experiments. Figure 5 shows the 95% CL expected and observed limits on $\sigma B$ for the $e^*$ and $\mu^*$ searches. For $m_{\ell^*} \geq 800$ GeV, the observed upper limits are 0.75 and 0.90 fb for the electron and muon channels, respectively. The sensitivity to the prior for $\sigma B$ was studied using a reference prior [45], resulting in 20–25% better limits for both channels. Theoretical predictions of $\sigma B$ for three different values of $\Lambda$ are also displayed in figure 5, along with the uncertainties from renormalization and factorization scales and PDFs. These uncertainties are shown for illustrative purpose only and are not used when setting limits.

For each $m_{\ell^*}$ hypothesis, the limit on $\sigma B$ is then translated into a lower bound on the compositeness scale. This bound corresponds to the value of $\Lambda$ for which the theoretical prediction $\sigma B(m_{\ell^*}, \Lambda)$ is equal to the upper limit on $\sigma B$. The excluded region in the $\Lambda$–$m_{\ell^*}$ plane is shown in figure 6 for both the $e^*$ and $\mu^*$ searches. For $m_{\ell^*} = \Lambda$, excited-electron and excited-muon masses are both excluded at 95% CL up to 2.2 TeV. The limits obtained with $\sqrt{s} = 7$ TeV data by ATLAS [18] and CMS [19] are also shown.

8. Conclusions

The results of a search for excited electrons and excited muons with the ATLAS detector at the LHC are reported, using a sample of $\sqrt{s} = 8$ TeV $pp$ collisions corresponding to an integrated luminosity of 13 fb$^{-1}$. The observed data are consistent with SM background expectations. An upper limit is set at 95% CL on the cross section times branching ratio $\sigma B(\ell^* \rightarrow \ell \gamma)$ as a function of the excited-lepton mass. For $m_{\ell^*} \geq 0.8$ TeV, the respective limits on $\sigma B$ are 0.75 and 0.90 fb for the $e^*$ and $\mu^*$ searches. These upper limits are converted into lower bounds on the compositeness scale $\Lambda$. In the special case where $\Lambda = m_{\ell^*}$, excited-electron and excited-muon masses below 2.2 TeV are excluded.

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The ATLAS Collaboration

8 Department of Physics, The University of Texas at Arlington, Arlington, TX, USA
9 Physics Department, University of Athens, Athens, Greece
10 Physics Department, National Technical University of Athens, Zografou, Greece
11 Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
12 Institut de Física d’Altes Energies and Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain
13a Institute of Physics, University of Belgrade, Belgrade, Serbia
13b Vinca Institute of Nuclear Sciences, University of Belgrade, Belgrade, Serbia
14 Department for Physics and Technology, University of Bergen, Bergen, Norway
15 Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, CA, USA
16 Department of Physics, Humboldt University, Berlin, Germany
17 Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
18 School of Physics and Astronomy, University of Birmingham, Birmingham, UK
19a Department of Physics, Bogazici University, Istanbul
19b Department of Physics, Dogus University, Istanbul
19c Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey
20a INFN Sezione di Bologna, Bologna, Italy
20b Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy
21 Physikalisches Institut, University of Bonn, Bonn, Germany
22 Department of Physics, Boston University, Boston, MA, USA
23 Department of Physics, Brandeis University, Waltham, MA, USA
24a Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil
24b Federal University of Juiz de Fora (UFJJ), Juiz de Fora, Brazil
24c Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei, Brazil
24d Instituto de Física, Universidade de Sao Paulo, Sao Paulo, Brazil
25 Physics Department, Brookhaven National Laboratory, Upton, NY, USA
26a National Institute of Physics and Nuclear Engineering, Bucharest, Romania
26b Physics Department, National Institute for Research and Development of Isotopic and Molecular Technologies, Cluj Napoca, Romania
26c University Politehnica Bucharest, Bucharest, Romania
26d West University in Timisoara, Timisoara, Romania
27 Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
28 Cavendish Laboratory, University of Cambridge, Cambridge, UK
29 Department of Physics, Carleton University, Ottawa, ON, Canada
30 CERN, Geneva, Switzerland
31 Enrico Fermi Institute, University of Chicago, Chicago, IL, USA
32a Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile
32b Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
33a Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China
33b Department of Modern Physics, University of Science and Technology of China, Anhui, China
33c Department of Physics, Nanjing University, Jiangsu, China
33d School of Physics, Shandong University, Shandong, China
33e Physics Department, Shanghai Jiao Tong University, Shanghai, China
34 Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France
35 Nevis Laboratory, Columbia University, Irvington, NY, USA
36 Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark
37a INFN Gruppo Collegato di Cosenza, Italy
37b Dipartimento di Fisica, Università della Calabria, Rende, Italy
38a AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland
38b Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland
39 The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland
40 Physics Department, Southern Methodist University, Dallas, TX, USA
41 Physics Department, University of Texas at Dallas, Richardson, TX, USA
42 DESY, Hamburg and Zeuthen, Germany
43 Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
44 Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
45 Department of Physics, Duke University, Durham, NC, USA
46 SUPA—School of Physics and Astronomy, University of Edinburgh, Edinburgh, UK
47 INFN Laboratori Nazionali di Frascati, Frascati, Italy
48 Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany
49 Section de Physique, Université de Genève, Geneva, Switzerland
50a INFN Sezione di Genova, Genova, Italy
50b Dipartimento di Fisica, Università di Genova, Genova, Italy
51a E. Andronikashvili Institute of Physics, Ivane Javakhishvili Tbilisi State University, Tbilisi, Georgia
51b High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
52 II Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
53 SUPA—School of Physics and Astronomy, University of Glasgow, Glasgow, UK
54 II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany
55 Laboratoire de Physique Subatomique et de Cosmologie, Université Joseph Fourier and CNRS/IN2P3 and Institut National Polytechnique de Grenoble, Grenoble, France
56 Department of Physics, Hampton University, Hampton, VA, USA
57 Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, MA, USA
58a Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
58b Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
58c ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany
59 Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
60 Department of Physics, Indiana University, Bloomington, IN, USA
61 Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
62 University of Iowa, Iowa City, IA, USA
63 Department of Physics and Astronomy, Iowa State University, Ames, IA, USA
64 Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia
65 KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
66 Graduate School of Science, Kobe University, Kobe, Japan
67 Faculty of Science, Kyoto University, Kyoto, Japan

104 Department of Physics and Astronomy, University of New Mexico, Albuquerque, NM, USA
105 Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
106 Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
107 Department of Physics, Northern Illinois University, DeKalb, IL, USA
108 Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia
109 Department of Physics, New York University, New York, NY, USA
110 Ohio State University, Columbus, OH, USA
111 Faculty of Science, Okayama University, Okayama, Japan
112 Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, OK, USA
113 Department of Physics, Oklahoma State University, Stillwater, OK, USA
114 Palacký University, RCPTM, Olomouc, Czech Republic
115 Center for High Energy Physics, University of Oregon, Eugene, OR, USA
116 LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France
117 Graduate School of Science, Osaka University, Osaka, Japan
118 Department of Physics, University of Oslo, Oslo, Norway
119 Department of Physics, Oxford University, Oxford, UK
120a INFN Sezione di Pavia, Italy
120b Dipartimento di Fisica, Università di Pavia, Pavia, Italy
121 Department of Physics, University of Pennsylvania, Philadelphia, PA, USA
122 Petersburg Nuclear Physics Institute, Gatchina, Russia
123a INFN Sezione di Pisa, Pisa, Italy
123b Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
124 Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA, USA
125a Laboratorio de Instrumentacao e Fisica Experimental de Particulas—LIP, Lisboa, Portugal
125b Departamento de Fisica Teorica y del Cosmos and CAFPE, Universidad de Granada, Granada, Spain
126 Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic
127 Czech Technical University in Prague, Praha, Czech Republic
128 Faculty of Mathematics and Physics, Charles University in Prague, Praha, Czech Republic
129 State Research Center Institute for High Energy Physics, Protvino, Russia
130 Particle Physics Department, Rutherford Appleton Laboratory, Didcot, UK
131 Physics Department, University of Regina, Regina, SK, Canada
132 Ritsumeikan University, Kusatsu, Shiga, Japan
133a INFN Sezione di Roma I, Roma, Italy
133b Dipartimento di Fisica, Università La Sapienza, Roma, Italy
134a INFN Sezione di Roma Tor Vergata, Roma, Italy
134b Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
135a INFN Sezione di Roma Tre, Roma, Italy
135b Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy
136a Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies—Université Hassan II, Casablanca, Morocco
136b Centre National de l’Energie des Sciences Techniques Nucleaires, Rabat, Morocco
136c Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech, Morocco
Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda, Morocco

DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat à l’Energie Atomique et aux Energies Alternatives), Gif-sur-Yvette, France

Department of Physics, University of California Santa Cruz, Santa Cruz, CA, USA

Department of Physics, University of Washington, Seattle, WA, USA

Department of Physics and Astronomy, University of Sheffield, Sheffield, UK

Department of Physics, Shinshu University, Nagano, Japan

Fachbereich Physik, Universität Siegen, Siegen, Germany

Department of Physics, Simon Fraser University, Burnaby, BC, Canada

SLAC National Accelerator Laboratory, Stanford, CA, USA

Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava, Slovakia

Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic

Department of Physics, University of Cape Town, Cape Town, South Africa

Department of Physics, University of Johannesburg, Johannesburg, South Africa

School of Physics, University of the Witwatersrand, Johannesburg, South Africa

Department of Physics, Stockholm University, Sweden

The Oskar Klein Centre, Stockholm, Sweden

Physics Department, Royal Institute of Technology, Stockholm, Sweden

Departments of Physics and Astronomy and Chemistry, Stony Brook University, Stony Brook, NY, USA

Department of Physics and Astronomy, University of Sussex, Brighton, UK

School of Physics, University of Sydney, Sydney, Australia

Institute of Physics, Academia Sinica, Taipei, Taiwan

Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel

Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel

Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece

International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan

Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan

Department of Physics, Tokyo Institute of Technology, Tokyo, Japan

Department of Physics, University of Toronto, Toronto, ON, Canada

TRIUMF, Vancouver, BC, Canada

Department of Physics and Astronomy, York University, Toronto, ON, Canada

Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan

Department of Physics and Astronomy, Tufts University, Medford, MA, USA

Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia

Department of Physics and Astronomy, University of California Irvine, Irvine, CA, USA

INFN Gruppo Collegato di Udine, Udine, Italy

ICTP, Trieste, Italy

Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy

Department of Physics, University of Illinois, Urbana, IL, USA

Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden

Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain

Department of Physics, University of British Columbia, Vancouver, BC, Canada

Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada

Department of Physics, University of Warwick, Coventry, UK

Waseda University, Tokyo, Japan

Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel

Department of Physics, University of Wisconsin, Madison, WI, USA

Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany

Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany

Department of Physics, Yale University, New Haven, CT, USA

Yerevan Physics Institute, Yerevan, Armenia

Centre de Calcul de l’Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France

Also at Department of Physics, King’s College London, London, UK

Also at Laboratorio de Instrumentacao e Fisica Experimental de Particulas—LIP, Lisboa, Portugal

Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan

Also at Faculdade de Ciencias and CFNUL, Universidade de Lisboa, Lisboa, Portugal

Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, UK

Also at TRIUMF, Vancouver, BC, Canada

Also at Department of Physics, California State University, Fresno, CA, USA

Also at Novosibirsk State University, Novosibirsk, Russia

Also at Department of Physics, University of Coimbra, Coimbra, Portugal

Also at Università di Napoli Parthenope, Napoli, Italy

Also at Institute of Particle Physics (IPP), Canada

Also at Department of Physics, Middle East Technical University, Ankara, Turkey

Also at Louisiana Tech University, Ruston, LA, USA

Also at Dep Fisica and CEFITEC of Faculdade de Ciencias e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal

Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France

Also at Department of Physics and Astronomy, Michigan State University, East Lansing, MI, USA

Also at Department of Financial and Management Engineering, University of the Aegean, Chios, Greece

Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain

Also at Department of Physics, University of Cape Town, Cape Town, South Africa

Also at CERN, Geneva, Switzerland

Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany

Also at Manhattan College, New York, NY, USA

Also at Institute of Physics, Academia Sinica, Taipei, Taiwan

Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China

Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan

Also at Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France

Also at School of Physical Sciences, National Institute of Science Education and Research, Bhubaneswar, India

Also at Dipartimento di Fisica, Università La Sapienza, Roma, Italy

Also at DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat à l’Energie Atomique et aux Energies Alternatives), Gif-sur-Yvette, France

Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia

Also at section de Physique, Université de Genève, Geneva, Switzerland

Also at Departamento de Fisica, Universidade de Minho, Braga, Portugal

Also at Department of Physics, The University of Texas at Austin, Austin, TX, USA

Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary

Also at DESY, Hamburg and Zeuthen, Germany

Also at International School for Advanced Studies (SISSA), Trieste, Italy

Also at Department of Physics and Astronomy, University of South Carolina, Columbia, SC, USA

Also at Faculty of Physics, M V Lomonosov Moscow State University, Moscow, Russia

Also at Nevis Laboratory, Columbia University, Irvington, NY, USA

Also at Physics Department, Brookhaven National Laboratory, Upton, NY, USA

Also at Department of Physics, Oxford University, Oxford, UK

Also at Department of Physics, The University of Michigan, Ann Arbor, MI, USA

Also at Discipline of Physics, University of KwaZulu-Natal, Durban, South Africa

Deceased

References


[45] Casadei D 2012 J. Instrum. 7 1012