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KEPLER-68: THREE PLANETS, ONE WITH A DENSITY BETWEEN THAT OF EARTH AND ICE GIANTS


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ABSTRACT

NASA’s Kepler Mission has revealed two transiting planets orbiting Kepler-68. Follow-up Doppler measurements have established the mass of the innermost planet and revealed a third Jovian-mass planet orbiting beyond the two transiting planets. Kepler-68b, in a 5.4 day orbit, has \( M_P = 8.3_{-2.4}^{+2.2} M_{\oplus} \), \( R_P = 2.31_{-0.09}^{+0.06} R_{\oplus} \), and \( \rho_P = 3.32_{-0.98}^{+0.86} g \text{ cm}^{-3} \), giving Kepler-68b a density intermediate between that of the ice giants and Earth. Kepler-68c is Earth-sized, with a radius \( R_P = 0.953_{-0.042}^{+0.037} R_{\oplus} \) and transits on a 9.6 day orbit; validation of Kepler-68c posed unique challenges. Kepler-68d has an orbital period of 580 \pm 15 days and a minimum mass of \( M_P \sin i = 0.947 \pm 0.035 M_{J} \). Power spectra of the Kepler photometry at one minute cadence exhibit a rich and strong set of asteroseismic pulsation modes enabling detailed analysis of the stellar interior. Spectroscopy of the star coupled with asteroseismic modeling of the multiple pulsation modes yield precise measurements of stellar properties, notably \( T_{\text{eff}} = 5793 \pm 74 \text{ K} \), \( M_*= 1.079 \pm 0.051 M_{\odot} \), \( R_* = 1.243 \pm 0.019 R_{\odot} \), and \( \rho_* = 0.7903 \pm 0.0054 \text{ g cm}^{-3} \), all measured with fractional uncertainties of only a few percent. Models of Kepler-68b suggest that it is likely composed of rock and water, or has a H and He envelope to yield its density \( \sim 3 \text{ g cm}^{-3} \).

Key words: planetary systems – fundamental parameters – stars: individual (Kepler-68, KIC 11295426, 2MASS J19240775+4902249)

Online-only material: color figures

1. INTRODUCTION

The NASA Kepler Mission has presented a catalog of over 2300 stars with planet-like transit signatures (borucki et al. 2011; Batalha et al. 2013). Here, we report a detailed study of Kepler-68, a G-type main sequence star harboring a transiting planet, Kepler-68b, with a radius of \( \sim 2.5 R_{\oplus} \) and orbital period \( \sim 5.40 \) days. We describe the detection of a second transiting planet that is close in size to the Earth. We carry out multiple follow-up measurements of the star Kepler-68 (KIC 11295426), including additional Kepler photometry, ground-based spectroscopy and high resolution imaging, Spitzer Space Telescope photometry, and Doppler measurements. At Kepler magnitude, \( KP = 10.00 \), the star has high enough flux for asteroseismic analysis of its stellar properties using short cadence (SC; see Gilliland et al. 2010) Kepler photometry, offering correspondingly accurate measures of the stellar density, mass, and radius.
measurements to give planet parameters, as described in Section 7. We also discuss the properties of Kepler-68b from the standpoint of theoretical models in Section 7, especially regarding the planet’s composition. Its placement in a mass–radius diagram suggests a composition of large amounts of rock and significant amounts of volatiles to yield the observed density of 3 g cm$^{-3}$.

With the changing status of stars during the course of the Kepler Mission, as planet candidates are discovered and confirmed as planets, the nomenclature used and recognized by diverse analysis groups evolves. The star studied here is located at $\alpha = 19^h24^m07.75^s$, $\delta = +49^\circ02.25'$, and in the Kepler Input Catalog (KIC) was designated KIC 11295426. The previously existing Two Micron All Sky Survey (2MASS) ID was J19240775+4902249. The Kepler Object of Interest (KOI) name was given when it appeared in the Borucki et al. (2011) exoplanet candidate list as KOI-00246, or the more commonly appearing KOI-246, as used herein. KOI numbers were assigned per convention as KOI-246.01 with initial detection of the 5.4 day candidate transits, KOI-246.20 for the candidate detected based on non-Kepler input, and then KOI-246.02 for the second transiting planet candidate at 9.6 days. With validation and confirmation of planets, the star was given its final moniker, Kepler-68, and the planets Kepler-68b,c,d—from KOIs 246.01, 246.02 and 246.20, respectively.

2. KEPLER PHOTOMETRY

The Kepler instrument is described in Van Cleve & Caldwell (2009) and Argabright et al. (2008), while an overview of the performance is presented in Caldwell et al. (2010b) and Jenkins et al. (2010b). Here, we report the results using 12 quarters of Kepler data. The standard pipeline reduction of the photometry first yielded a transit signal with a period of 5.40 days, consistent with a planet with a size of approximately 2.5 $R_\oplus$ (Borucki et al. 2011; Batalha et al. 2013). Subsequent searches of the light curve alerted the Kepler team to a second transit signal with a period of ~9.61 days and approximately Earth size. Thus, the Kepler photometry and pipeline reduction from Kepler-68 reveals two periodic transit signals consistent with planets, hereafter called Kepler-68b and Kepler-68c.

Independent transit searches of the Kepler photometry have been carried out by Ofir & Dreizler (2012) and Huang et al. (2013), who also find evidence for Kepler-68b and Kepler-68c with the same period and transit depth within uncertainties. For Kepler-68c, Ofir & Dreizler (2012) find a period of 9.60538 ± 0.00026 days and a planet radius of 0.86 $R_\oplus$. They also find a single transit from a possible third planet with a transit duration of ~8 hr, implying a period of 970 ± 50 days and planet size about 2.4 $R_\oplus$. Thus, there is a possibility that a third planet transits Kepler-68. Kepler observations of Kepler-68 are ongoing, including acquisition of SC data. Inspection of data through Quarter 13 has not shown further evidence for the several hundred day planet candidate. At a period of nearly 1000 days, the next transit would not be expected until Quarter 16 in early 2013.

Raw flux light curves for each quarter (Jenkins et al. 2010a) are corrected for systematic errors, detrended, and stitched together to form contiguous time series, and are then searched for transit signals (Jenkins et al. 2010c). We remove systematic errors, outliers, and intra-quarter discontinuities by cotrending against the photometry of other stars in the vicinity of Kepler-68 using the Pre-Search Data Conditioning (PDC) pipeline module as described in Twicken et al. (2010), with updates as per Stumpe et al. (2012) and Smith et al. (2012).

Figure 1 shows the raw (SAP_FLUX, the result of simple aperture photometry) and corrected (PDCSAP_FLUX) flux time series for Kepler-68 during a representative quarter (Quarter 4). The largest remaining systematic errors in the PDC-MAP (Stumpe et al. 2012; Smith et al. 2012) processed data are minor offsets following thermal transients after monthly pointings changes to telemeter data to the ground (very small for Quarter 4). The slow variation with a period of ~50 days and amplitude of 0.0003 in this figure could either be intrinsic to the star or associated with imperfect removal of long-term drifts due to image motion (differential velocity aberration) that is present in the raw data. It is clear that variations in Kepler-68 are smaller than typical variations of the Sun, consistent with the slow rotation and advanced age argued for in Sections 6.1 and 6.2, respectively. After filtering out transit events, the measured relative standard deviation of the PDC-corrected, long-cadence light curve is 21 ppm per 6 hr interval (CDPP—the formal Combined Differential Photometric Precision—see Jenkins et al. 2010b; Christiansen et al. 2012a). An expected instrument + photon noise is computed for each flux in the timeseries. The mean of the per (29.4 minute) cadence noise estimates reported by the pipeline is 233 ppm. Both raw (simple aperture sums) and corrected (PDC-MAP) light curves are available at the Mikulski Archive for Space Telescopes17 at the Space Telescope Science Institute.

2.1. Transiting Planet Search

The Kepler-68b transits were identified by the Transiting Planet Search (TPS) pipeline module that identifies periodic reductions in flux with a duration of hours, each corresponding to a transit of a prospective planet. The algorithm is a wavelet-based, adaptive matched filter (Jenkins et al. 2010c). TPS then identifies a time series of single “events,” each having an associated “single event statistic (SES)” that represents the probability that a transit is present. The SES from each transit are combined into multiple event statistics (MES) by folding them at trial orbital periods ranging from 0.5 days to as long as half the data coverage interval.

Kepler-68b was identified by TPS in each quarter of data with an MES > 15$\sigma$.

Multi-quarter searching for transits was used. The transit depth, duration, period, and epoch are derived from physical modeling (see Section 7) using all of the available data. Kepler-68b is characterized as a 345.6 ± 1.5 ppm dimming lasting 3.459 ± 0.009 hr, with transit ephemeris of $T_{[BJD]} = 2455006.85729 \pm 0.00042 + N \times 5.938763 \pm 0.000004$ days. The longer-period transits of Kepler-68c were identified by non-pipeline inspections. Kepler-68c is characterized as a 53.1 ± 2.3 ppm dimming lasting 3.09 ± 0.09 hr and an ephemeris $T_{[BJD]} = 2454969.3805 \pm 0.0041 + N \times 9.60585 \pm 0.000072$ days.

3. DATA VALIDATION

Signals that mimic planet transits are also found by TPS. All “threshold crossing events” identified by TPS are subjected to assessment of standard vetting products that allow separate disposition of clear false positives before bestowing the KOI moniker. Most false positives can be identified by judicious assessment of the quality of the transit-planet model fit to the

17 http://archive.stsci.edu/kepler
photometry and by searching for astrometric displacements of the photocenter between times out of transit and in-transit. True transiting planets should exhibit photometry that is well fit by our transit model and they should show little, if any, variation and is the reason we did not use the Quarter 3 data. The upper panel shows the phase folded light curve for Kepler-68b, after de-trending and subtraction of the Kepler-68c transit signal. A common type of false positive is a background eclipsing binary blended with the target star; in such cases, a secondary eclipse is often seen. Evidence of a secondary eclipse is seen neither here, nor in automated searches. The lower left panel shows that the phase folded data for Kepler-68b are fit very well by our transit model. Another signature of a possible false positive associated with an eclipsing binary (either the target star itself as a grazing eclipse or as a blended background system) relies on searching for subtle differences of depth between alternate transits—a binary with twice the listed period and primary and secondary eclipses of slightly unequal depths are easily seen in data of this quality. For all aspects of lightcurve inspection, Kepler-68b is entirely clean.

Figure 3 addresses Kepler-68c by showing vetting products for the time series analogous to Figure 2, but with the transit signal of Kepler-68b subtracted out. Again, there is no clear evidence for a secondary eclipse or depth differences for alternately averaged transits, either of which would be suggestive of a false positive interpretation. The question of a secondary eclipse feature will, however, be explored at depth in the BLENDER context in Section 5. This transit, which is only ~60% as deep as a true Earth analog, obviously appears at high signal-to-noise ratio (S/N) in the phase folded data and is well fit by our planet transit model.

Centroid analyses based on assessing the in- and out-of-transit signal positions can be particularly powerful (e.g., see Batalha et al. 2011 for application to Kepler-10b). The difference of average images taken out-of-transit, minus a similar average of images taken during transit for an unsaturated target, yields a point-spread function (PSF) at the inherent source position (coincident within errors with the target star for a clean candidate). Since Kepler-68 is saturated, the centroid analyses are much less discerning with the inherent loss of spatial information following saturation. For a strongly saturated target, changes of flux are concentrated at the end of bleed columns, while the central pixels in the bleed trail remain, well, saturated.

The KIC shows the location of a fainter (magnitude not available in KIC, but derived below in Section 4.2.1 from AO imaging) star offset by about 11" in such a way that it is almost precisely in the detector column direction from Kepler-68. In most quarters, Kepler-68 is roughly centered on a pixel in the column direction leading to bleeding that encompasses this secondary star, KIC 11295432, in which case using the Kepler data...
Figure 2. The upper panel shows the de-trended time series over Quarters 1–11 (without Quarter 3 as discussed in the text) after folding on the 5.39877 day period of Kepler-68b such that the transit falls at phase 0.75. The signal for Kepler-68c has also been subtracted. Green points show the data binned in 0.01 phase intervals, the red line is the best fitting transit model. The lower left panel provides detail on the phased and folded light curve at the position of the transit; solid dots indicate 30 minute averages and the solid line is the best fitting transit light curve fit. The lower right panel details Odd and Even numbered transits individually co-added. Red stars mark center-of-transit times.

(A color version of this figure is available in the online journal.)

Figure 3. Same as Figure 2, but for Kepler-68c after removal of the Kepler-68b transits.

(A color version of this figure is available in the online journal.)
Table 1

<table>
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<td>338</td>
<td>0.00</td>
<td>0.02</td>
<td>0.03</td>
</tr>
</tbody>
</table>

Notes. Column and row values refer to pixels on channel 59 of the *Kepler* detectors for Quarter 9. The Direct image has been normalized by the total electrons per cadence of $2.4 \times 10^7$. The difference images have been normalized by the same factor scaled by the known depths of 346 and 53 ppm for Kepler-68b and Kepler-68c respectively. For internal consistency, the sums over these normalized difference images should be $\sim 1$ (this is satisfied). Kepler-68 is nearly centered in these tabular domains, while KIC 11295432 (7 mag fainter companion—see text) is near the center of row 343, and columns 944–945 pixels as italicized in entries. Typical uncertainties for Kepler-68b entries are 0.004, and 0.04 for Kepler-68c.

To discern the true source of transits between the two stars is not possible. In Quarter 9 both Kepler-68 and the fainter neighbor KIC 11295432 are nearly centered between columns, and the bleeding terminates before reaching the latter. Table 1 details pixel values in both the direct out-of-transit image at the left, and for the difference images of Kepler-68b and Kepler-68c in the central and right blocks for Quarter 9 data. (Quarter 1 and Quarter 5 with the telescope at the same orientation also possess this feature, but the Kepler-68c difference images, which in the best of circumstance are low S/N, were unstable for these quarters.) The difference images for both Kepler-68b and Kepler-68c show maxima in the terminal saturation pixels, rather than in the pixel containing KIC 11295432. This proves that KIC 11295432 cannot be the source for either set of transits. Saturated pixels in these quarters span a rectangle of $2 \times 4$ pixels, or $8'' \times 16''$, for which centroid analysis does not rule out background contaminating sources. Deviations in the difference image along the row direction for saturated images likely result from imperfect correction for local detector electronics undershoot (Caldwell et al. 2010a), a signal-dependent offset to negative values along rows traced to an amplifier in the *Kepler* electronics. The high resolution imaging discussed in Section 4.2 will provide the primary constraints on potential background objects, which if these objects were eclipsing binaries could be the source of false positives.

In Figure 4, we present vetting evidence to further investigate the veracity of the single transit of the presumed long period reported in Ofir & Dreizler (2012) and shown in their Figure 3. The detrending of raw *Kepler* data adopted by Ofir & Dreizler (2012) results in an $\sim 8$ hr wide intensity drop that is adequately fit with a standard transit model. In the raw long cadence (LC, simple aperture photometry—no detrending) *Kepler* data, their

![Figure 4. The upper panel shows raw *Kepler* long cadence data spanning six days centered on the “monotransit” event shown in Figure 3 of Ofir & Dreizler (2012). The middle panel shows the same data after pipeline processing with PDC-MAP. The lower panel shows the same time period with short cadence raw data.](image)
event is easily seen at BJD \( - 2455000 = 403.2 \). However, the raw LC data shows a drop of intensity of about 0.0001 across this event which may easily be seen by drawing linear fits to the data before or after the event. Such behavior is commonly seen with sudden pixel sensitivity losses associated with radiation damage to a single pixel (see, e.g., discussion from discovery in Hubble Space Telescope (HST) ACS data by Gilliland & Bohlin (2007) and Christiansen et al. (2012b) for its presence in Kepler data)—and ultimate recovery of most, but not all of the depressed sensitivity. Such sensitivity drops can be particularly difficult to tell from transits in LC data that blur both sudden (spurious) drops and short ingress/egress events. The spurious signature may be more easily seen in SC data, as shown in the bottom panel of Figure 4—here, the event looks more like a sudden sensitivity loss and recovery than a transit. Note that the two occurrences of Kepler-68b transits about one day from both ends show symmetric responses about the transit centers, while the single-event behavior is quite asymmetric with shape, characteristic of a pixel sensitivity loss. The middle panel shows that the pipeline PDC-MAP treatment completely removes any suggestion of the single-event transit while nicely preserving the real events from Kepler-68b. These considerations do not prove the single-event transit as an instrument systematic false positive, but we favor this interpretation. Definitive proof would come if a difference image analysis across the transit could isolate the intensity drop to one pixel, while the stellar signal spans several pixels. We tested for this, but were only able to show that the source is within the set of saturated pixels for Kepler-68, leaving open an astrophysical source.

4. FOLLOW-UP OBSERVATIONS

Kepler-68 was identified as having a candidate planet in late 2009, prompting follow-up observations to confirm and characterize the planet and to secure more precise stellar parameters than are available in the KIC. We were particularly interested in checking for signs of a nearby eclipsing binary star system that might masquerade as a planet. We carried out spectroscopy of the host star Kepler-68 (Section 4.1) and high spatial resolution imaging to identify nearby stars in the photometric aperture (Section 4.2). Upon passing those gates, we carried out high-resolution, high S/N echelle spectroscopy with and without an iodine cell to measure atmospheric stellar parameters, magnetic activity, absorption line shape changes with time, and to make precise Doppler measurements. As described below, these follow-up observations revealed no evidence of a nearby eclipsing binary for either transiting planet, and they provide measurements and constraints on the masses of the transiting planets. In addition, a previously unknown non-transiting planet was revealed as discussed in Section 4.3. Bisector analyses of the high resolution spectroscopy are presented in Section 4.4. We obtained Spitzer Space Telescope photometry through transits of Kepler-68b and present the results in Section 4.5.

4.1. Reconnaissance Spectroscopy

Two reconnaissance spectra were obtained with the Tull Coude Spectrograph of the McDonald Observatory 2.7 m Harlan J. Smith Telescope on the nights of 2010 March 25 and 28, and a third at the TRES Echelle Spectrograph of the Tillinghast 1.5 m telescope on Mt. Hopkins, also on the night of 2010 March 25. These high spectral resolution, low signal/noise spectra showed no convincing evidence for RV variability at the 0.2 km s\(^{-1}\) level, and no hints of any contaminating spectra.
limits as shown in Figure 6. Within the full ARIES field of view, one additional source that was 6.1 and 5.7 mag fainter in $J$ and $K_s$, respectively, was identified. Using the 2MASS $J$ and $K_s$ magnitudes for Kepler-68 of 8.974 and 8.587 leads to estimates of 15.07 and 14.29, respectively, for the companion. Using the transformations in Appendix A of Howell et al. (2011) leads to a $K_p = 17.0 \pm 0.4$, or a $\delta$-magnitude of $7.0 \pm 0.4$ with respect to Kepler-68 in the *Kepler* bandpass. This AO identified source also shows up clearly in available UKIRT $J$ direct imaging (Lucas & Samuel 2009) and corresponds to the star KIC 11295432, which does not have a specified $K_p$ in the KIC. From a combination of *Kepler* data centroid consideration (Section 3) and the ARIES AO and UKIRT data, to roughly the exclusion limits reached in Figure 6, no other sources exist to relevant larger radii as well.

4.2.2. Speckle Imaging

Speckle imaging of Kepler-68 was obtained on the night of 2010 June 20 UT using the two-color speckle camera at the Wisconsin Indiana Yale NOAO (WIYN) 3.5 m telescope on Kitt Peak. The speckle camera simultaneously obtained 1000 40 ms images in two filters: $V$ (5620/400 Å) and $R$ (6920/400 Å). These data were reduced and processed to produce a final reconstructed speckle image for each filter. Figure 5 includes this speckle-reconstructed $R$ band image. The details of the two-color speckle camera are presented in Howell et al. (2011) for application in the *Kepler* follow-up program.

The speckle data allow the detection of a companion star within the approximately $2.76 \times 2.76$ arcsec box centered on the target. The speckle imaging can detect, or rule out, companions between 0.05 arcsec and 1.5 arcsec from Kepler-68. We found no companion star within the speckle image to the detection limits shown in Figure 6. When high quality near-IR AO imaging is available, as in this case, the WIYN speckle imaging is largely redundant, providing unique sensitivity for a small angular separation range due to its superior FWHM of $0.053$—a factor of two sharper, but generally shallower than the MMT AO imaging.

4.3. Precise Doppler Measurements of Kepler-68

We obtained 52 high resolution spectra of Kepler-68 between 2010 July 19 and 2012 August 12 using the HIRES spectrometer on the Keck I 10 m telescope (Vogt et al. 1994). We configured the spectrometer, took observations, and reduced the spectra with the same method used on thousands of nearby FGKM stars (Marcy et al. 2008). This technique yields a Doppler precision of $1.0–1.5$ m s$^{-1}$ for stars as faint as the 14th magnitude ($V$ band), depending on spectral type and rotational $v \sin i$. The HIRES fiber-feed was not used for these observations. An iodine cell was used to superimpose iodine lines on the stellar spectrum, providing empirical information for each exposure and each wavelength about the instantaneous wavelength scale and instrumental profile of the spectrometer.

The observations were made with the “C2 decker” entrance aperture, which projects to $0.87 \times 14.0'$, giving a resolving power of $70,000$ at $5500$ Å and enabling sky subtraction (typical seeing is $0.6–1.2'$). The average exposure was 11 minutes, giving a $S/N$ per pixel of 200.

The raw CCD images were reduced in the standard way, including the subtraction of background sky counts (mostly from moonlight) at each wavelength just above and below the stellar spectrum. We used Doppler analysis with the algorithm of Johnson et al. (2009). The internal Doppler errors (the weighted uncertainty in the mean of 400 spectral segments) are typically $1.0–1.5$ m s$^{-1}$. Our experience with hundreds of G-type main sequence stars shows that the actual errors are larger than the internal errors by $\sim 1.5$ m s$^{-1}$ for such stars. Thus, we added $1.5$ m s$^{-1}$ in quadrature to the internal uncertainties to yield our estimated uncertainty in the RV. The resulting velocities and uncertainties are given in Table 2 and shown in Figure 7 as a function of time. The error bars include the internal Doppler errors and an assumed jitter of $1.5$ m s$^{-1}$.

The center of mass velocity relative to the solar system barycenter (Gamma Velocity) for Kepler-68 is $-20.9 \pm 0.1$ km s$^{-1}$ (Table 4). This is a typical RV for a star in the Galactic disk, indicative of a middle-aged disk star. The near solar metallicity (Section 6.1), [Fe/H] = $+0.12 \pm 0.07$ mag, magnetic activity (Section 6.1), and asteroseismic age, $6.3 \pm 1.7$ Gyr (Section 6.2), also suggest Galactic disk membership.

The primary signal from the RVs is the $K = 19.9 \pm 0.75$ m s$^{-1}$ variation with a period of $580 \pm 15$ days, as shown first in Figure 7 and phased in Figure 8. This clear RV signature, coupled with the lack of bisector variations discussed in Section 4.4, provides discovery and confirmation of the Jovian-scale outer planet—Kepler-68b.

The velocities phased to the photometric period of Kepler-68b in Figure 8 show a clear, continuous, and nearly sinusoidal variation consistent with a nearly circular orbit of a planetary companion. Note that the 7 mag (see Sections 3 and 4.2.1) fainter companion KIC 11295432 blended into the *Kepler* photometry does not fall within the HIRES slit and cannot be the source of Kepler-68b RV variations. The lack of any discontinuities in the phased velocity plot argues against a background eclipsing binary star as the explanation. Such a binary with a period of 5.4 days would have orbital semi-amplitudes of tens of kilometers per second, so large that the spectral lines would completely separate from each other and separate from the lines of the main star. Such breaks in the spectral-line blends would cause discontinuities in the velocity variation, which are not
seen here. Thus, the chance that the 5.4 day periodicity exhibited independently in both the photometry and velocities might be caused by an eclipsing binary seems small.

Precision Doppler measurements are used to constrain the mass of Kepler-68b, as discussed in Section 7. The absence of a Doppler signal for Kepler-68c is used to compute an upper limit to the mass under the planet interpretation.

4.4. Bisector Analysis

From the Keck spectra, we computed a mean line profile and the corresponding mean line bisector. Time-varying line asymmetries are tracked by measuring the bisector spans—the velocity difference between the top and bottom of the mean line bisector—for each spectrum (Torres et al. 2005). When RV variations are the result of a blended spectrum between a star and an eclipsing binary, we expect the bisectors to reveal a phase-modulated line asymmetry (Queloz et al. 2001; Mandushev et al. 2005).

The bisector spans in the lower panel of Figure 7 show no correlated variation with the RVs and have a scatter of 7.8 m s\(^{-1}\), which is significantly less than the semi-amplitude of Kepler-68d. However, the uncertainties in the bisector measurements are larger than the semi-amplitude of the two smaller planets in the system and a bisector analysis is thus inconclusive with respect to Kepler-68b and Kepler-68c.

4.5. Photometry with the Spitzer Space Telescope

Kepler-68b was observed during two transits with Warm-Spitzer/IRAC (Werner et al. 2004; Fazio et al. 2004) at 4.5 \(\mu\)m (program ID 60028). The observations occurred on UT 2010 December 27 and on UT 2011 January 7. The visits lasted 7.6 hr. The data were gathered in subarray mode (32 \times 32 pixels) with an exposure time of 2 s per image that yielded 15,680 images per visit.

The images used are the Basic Calibrated Data (BCD) delivered by the Spitzer archive. These files are corrected for dark current, flat-fielding, detector non-linearity and converted into flux units. The method we used to produce photometric time series from the images is described in Désert et al. (2009). We first discard the first half-hour of observations, which is affected by a significant telescope jitter before stabilization. To facilitate the evaluation of the photometric errors, we then convert the pixel intensities to electrons using the information given in the detector gain and exposure time provided in the FITS headers. We convert to UTC-based BJD following the procedure developed by Eastman et al. (2010). We correct for transient pixels in each individual image using a 20-point sliding median filter of the pixel intensity versus time. We find the centroid position of the stellar PSF and perform aperture photometry using a circular aperture with a radius of 3.5 pixels on individual BCD images; we adopt the radius that provides the smallest errors. The final number of photometric measurements used is 13,146 data points for the first visit and 13,262 for the second one. The raw time series are presented in the top panels of Figure 9. We find a typical S/N 260 per image, which corresponds to about 90% of the theoretical S/N.

We used a transit light curve model multiplied by instrumental decorrelation functions to measure the transit parameters and their uncertainties from the Spitzer data, as described in Désert et al. (2011b). We compute the transit light curves with the IDL transit routine OCCULTSMALL from Mandel & Agol (2002). In the present case, this function depends on one parameter: the planet-to-star radius ratio \(R_p/R_\star\). The other transit parameters are fixed to the value derived from the Kepler lightcurves. The limb-darkening coefficients are set to zero, consistent with expected small values in the IR and insufficient photometric precision in the Spitzer light curves to matter.

The Spitzer/IRAC photometry is known to be systematically affected by the so-called pixel-phase effect (see, e.g., Charbonneau et al. 2005; Knutson et al. 2008). We decorrelated our signal in each channel using a linear function of time for the baseline (two parameters) and a quadratic function of the PSF position (four parameters) to correct the data for each channel. We performed a Markov Chain Monte Carlo (MCMC) analysis with six chains of length \(10^5\), each providing median depth values and errors. We allow asymmetric error bars spanning 34% of the points above and below the median of the distributions to

![Figure 7. Radial velocities versus time from the Keck-HIRES spectra over 2010–2012 are shown in the upper panel. Error bars include the internal uncertainties and the expected combined astrophysical and instrumental jitter of 1.5 m s\(^{-1}\), added in quadrature. The lower panel shows line bisectors derived from the same spectra.](Image)
derive the 1σ uncertainties for each parameter, as described in Désert et al. (2011a).

We measured the transit depth for Kepler-68b at 4.5 μm of 350 ± 70 ppm for the first visit and 560 ± 70 ppm for the second. The weighted mean of these two values provides a transit depth of 455 ± 50 ppm. The value for the first visit is in excellent agreement with the Kepler depth of 346 ppm, suggesting that the radius-ratio of the candidate Kepler-68b to its host star is a wavelength independent function, in agreement with a dark planetary object. However, visit two is 3σ from the Kepler value. Two possibilities for this behavior, other than the small statistical probability of two such discrepant values legitimately following from the same underlying distribution, are: (1) a physically different depth in the two epochs for the transit, and (2) an inconsistency in one of the Spitzer epochs that we have failed to uncover. To investigate the first possibility, we identified the position of the two Spitzer epochs relative to Kepler coverage. Alas, despite the excellent overall duty cycle with Kepler in excess of 90%, both epochs fell into the longest downtime experienced to date with Kepler—due to a safing event before Quarter 8. We have no data with which to challenge the hypothesis that the two epochs really do have
Briefly, we performed a systematic exploration of the different types of false positives that can mimic the signal by generating large numbers of synthetic blend light curves over a wide range of parameters and comparing each of them with the Kepler photometry in a $\chi^2$ sense. The photometry we used for the validation is the de-trended SC time series (one month from Quarter 2, plus the seven quarters of Quarters 5–11—886,638 measurements), which provides stronger constraints on the shape of the transit than the long cadence data, as shown later. We rejected blends that result in light curves inconsistent with the Kepler observations. We then estimated the frequency of the remaining blends by taking into account all of the available observational constraints from the follow-up observations mentioned above. Finally, we compared this frequency with the expected frequency of true planets (planet “prior”) to derive the odds ratio.

The types of false positives we considered include eclipsing systems falling within the Kepler aperture that are either in the background or foreground, or that are physically associated with the target in a hierarchical triple configuration. We allowed the object producing the eclipses to be either a star or a planet. To compute the blend frequencies, we used informed estimates of the number density of stars in the background from the Besancon Galactic structure model of Robin et al. (2003), rates of occurrence of eclipsing binaries in the Kepler field from the work of Slawson et al. (2011), and frequencies of planets involved in blends based on the catalog of planet candidates (KOIs) of Batalha et al. (2013), which was constructed using observations from Quarter 1 to Quarter 6. The same catalog was also used to estimate the planet prior. This list of KOIs is bound to contain some fraction of false positives (see, e.g., Morton & Johnson 2011; Morton 2012) and it is also likely incomplete, mainly due to the difficulty in detecting shallow signals especially with long periods. Consequently, we have applied corrections for these effects following a Monte Carlo procedure described by Fressin et al. (2012a), both in computing the blend frequencies and also for the planet prior.

The observational constraints used to further reduce the number of blends included the following: (1) the color of the star as reported in the KIC (Brown et al. 2011), which allows us to rule out any simulated blends resulting in a combined color that is significantly redder or bluer than the target; (2) limits from the centroid motion analysis on the angular separation of companions that could produce the signal (Section 3); (3) brightness and angular separation limits from high-resolution adaptive optics and speckle imaging (Section 4.2); (4) limits on the brightness of unresolved companions from high-resolution spectroscopy (Section 4.1). For eclipsing systems physically associated with the target, we also considered dynamical stability constraints in hierarchical triple configurations (Holman & Wiegert 1999).

The BLENDER simulations for Kepler-68c rule out all false positive scenarios involving eclipsing binaries physically associated with the target, as the predicted light curves invariably have the wrong shape to match a planetary transit. For the scenarios involving eclipsing binaries that are in the background or foreground, we find a blend frequency of $2.8 \times 10^{-6}$, and for those in which larger planets transit background stars, we estimate a much smaller frequency of $7.0 \times 10^{-8}$. Hierarchical triples (a larger planet transiting a stellar companion to the target) contribute a frequency of $6.7 \times 10^{-7}$. The total blend frequency is thus $3.5 \times 10^{-6}$. An illustration of the constraints on false positives resulting from BLENDER as well as those from the follow-up observations is given in Figure 10.
An estimate of the planet prior may be obtained by dividing the number of known planets of similar size as Kepler-68c from Batalha et al. (2013) by the total number of main-sequence Kepler targets observed during Quarters 1–6, which is 138,253. We find 71 KOIs that are in the same (3σ) radius range as the putative planet in Kepler-68c, of which we expect 9.4 to be false positives, following the procedures of Fressin et al. (2012a). We also expect such shallow (∼50 ppm) transit signals to be detectable around only about 9.7% of all Kepler targets, which we use to correct for incompleteness. We compute the planet prior as (71 − 9.4)/0.097/138,253 = 4.6 × 10^-3. This a priori planet frequency is 4.6 × 10^-3/3.5 × 10^-6 ≈ 1300 times larger than the estimated blend frequency, from which we conclude that Kepler-68c is validated as a true planet with a very high degree of confidence.

In the above calculations, we have not explicitly taken into consideration the period of the signal, which may be an important factor for small candidates such as Kepler-68c because such signals are relatively rare in the KOI list of Batalha et al. (2013). This may in principle influence both the planet prior and the blend frequencies we have just described, given that the latter also draw on the KOI list to estimate the rate of occurrence of larger planets involved in blends. Therefore, instead of allowing eclipsing binaries and transiting planets with any orbital period for the blend frequency calculations, we repeated the analysis with the more realistic approach of accepting only blends with periods near the measured period of Kepler-68c (within a factor of two). We did the same when computing the planet prior, for consistency. The total blend frequency is reduced in this way by about a factor of five, but the new planet prior is only 1.5 times smaller, resulting in a larger net odds ratio of ∼4300 that provides for an even stronger validation than before.

The above calculations neglect the fact that KOI-246.02 was found around a target that has a statistically validated transiting planet, KOI-246.01 = Kepler-68b. The planet priors used were averaged between single and multi-planet systems. Actual planet priors are about 30% smaller for single planets and more than an order of magnitude larger for multiple planets (Lissauer et al. 2012; J. J. Lissauer et al. 2013, in preparation). The presence of a known planet also increases the prior for

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Figure 10. BLENDER goodness-of-fit contours for Kepler-68c corresponding to the three different scenarios that contribute to the overall blend frequency: background eclipsing binaries (top), background or foreground stars transited by a planet (middle), and physical companions transited by a planet (bottom). Only blends inside the solid white contours produce light curves matching the Kepler photometry within acceptable limits (3σ, where σ is the significance level of the \( \chi^2 \) difference compared to a transit model fit; see Fressin et al. 2011). Lighter-colored areas (red, orange, yellow) mark regions of parameter space giving increasingly worse fits to the data (4σ, 5σ, etc.), and correspond to blends we consider to be ruled out. The axes in each panel represent two of the
physically associated blends, but by a smaller factor. The numbers quoted above for the likelihood of false positives are four times higher for background blends than for physically associated blends, so when the multiplicity of the system is accounted for, the odds ratios against the blends quoted above are increased by roughly an order of magnitude.

As an interesting test of the value of SC versus long cadence data for false positive discrimination, we repeated the calculations with the long-cadence time series (Quarters 1–2, Quarters 4–11, 34,556 measurements) including the period cut described above. We find an odds ratio of approximately 1500, about three times lower than when using SC, but still large enough to comfortably validate the signal. Therefore, at least for Kepler-68c, SC provides a clear advantage in ruling out blends. This is likely due to the better definition of the transit shape at ingress/egress, which is often where the main differences are between model light curves for blends and the model for a true planet. In this case, we find that the main improvement (decrease) in the blend frequencies when going from long to SC is in the scenarios involving hierarchical triples, with background eclipsing binary frequencies changing the least.

While in principle the above calculations provide a clear statistical validation of the Kepler-68c signal as a true planet, independent of the detection of the reflex motion of the star (RVs), we note that some of the blend scenarios involving background eclipsing binaries yield a fit to the *Kepler* photometry that is significantly better than that of a planet fit. This is a situation we have not encountered in previous validations of *Kepler* candidates. The light curves of these false positive scenarios all feature a very shallow (∼10 ppm) secondary eclipse that happens to match a similar dip present in the phase-folded photometry for Kepler-68c at phase 0.5. This shallow dip at phase 0.5 has a formal significance of 3σ–4σ and is the primary source for these favorable false positive fits.

The reality of the signal at phase 0.5 for Kepler-68c is in doubt; were this not so, then claiming validation in this case would be impossible. Figure 11 contrasts phase folded SC data for Kepler-68c at both the transit and offset by half phases. In assessing the reality of this putative secondary eclipse feature, we have examined the data in several ways. When averaged over widths comparable to the transit, the data show a few other excursions (both positive and negative) to deviations as large as the phase 0.5 feature; thus, the feature is obviously not highly significant. However, other aspects of data investigation do not support a straightforward dismissal of evidence consistent with the weak ∼10 ppm feature. Examination of medians over the phase bins shown in Figure 11 showed the same offsets, and thus the deviation does not result simply from a small subset of the data. Likewise, dividing the data into first and second halves before doing independent phase binnings shows evidence for a ∼10 ppm depression at phase 0.5 in both cases. In these two halves, there are again other instances of deviations of comparable width and depth; however, across the two halves, the only cases lining up are those at phase 0.5. While nothing here is convincing in terms of regarding the phase 0.5 offset as real, it is equally the case that the reality of the offset cannot be excluded. We therefore took the conservative approach of allowing BLENDER to be influenced by the apparent negative offset at phase 0.5, which still provides the favorable odds ratio required for formal validation of Kepler-68c as a planet.

Further complicating this interpretation, these data show a small degree of correlated noise, even after folding over 64 orbits of Kepler-68c. After binning to 0.001 phase bins, the first autocorrelation value is 0.25, falling to zero within about 0.005 in phase. We formed simulated time series consisting
only of Gaussian noise, stellar oscillations, and granulation (Gilliland et al. 2011a) appropriate to Kepler-68 and found that after folding this had a similar (although smaller) degree of autocorrelation, suggesting that inherent stellar variations may suffice to explain the modest correlated noise.

5.2. Validation of Kepler-68b

The robust detection of a Doppler signal at the period of Kepler-68b provides strong support for the planetary nature of that signal. There is, however, a small chance that a background eclipsing binary could mimic this spectroscopic signature of a planet (as well as the photometric signal), although this possibility was convincingly argued against in Section 4.3. The precision of our bisector span measurements in Section 4.4 is not quite high enough for a definite conclusion regarding this possibility, so we proceed here with a validation analysis along the lines of what was done for Kepler-68c.

The much deeper transit (~350 ppm) and higher S/N of Kepler-68b result in significantly stronger constraints on the shape of the signal, and consequently in a much reduced frequency of blends that give acceptable fits to the shape of the signal, and consequently in a much reduced precision of our bisector span measurements in Section 4.4. This process of ruling out by BLENDER was convincingly argued against in Section 4.3. The precision of our bisector span measurements in Section 4.4 is not quite high enough for a definite conclusion regarding this possibility, so we proceed here with a validation analysis along the lines of what was done for Kepler-68c.

To estimate the a priori likelihood of a true planet, we use the fact that there are 96 candidates in the list of Batalha et al. (2013) with a planetary radius similar to that implied by Kepler-68b (within 3σ), of which an estimated six should be false positives. Signals of this kind are expected to be detectable in 58.7% of all Kepler targets. The planet prior is then (96 − 6)/0.587/138,253 = 1.1 × 10^−3. With this, we obtain an odds ratio of nearly 79,000 in favor of the planet interpretation, i.e., a very clear validation of Kepler-68b. The above results used the long-cadence time series for simplicity, and included the period cut described above. Use of SC would likely result in more reasonable total errors of 74 K and 0.074 dex.

We measured the Ca II K emission (Isaacson & Fischer 2010), yielding a Mt. Wilson S value S = 0.139 and log R'HK = −5.15. These values suggest lower activity for Kepler-68 in comparison to mean solar reported as S = 0.178, log R'HK = −4.90 by Lockwood et al. (2007) and S = 0.171, log R'HK = −4.96 by Hall et al. (2009). Our own measure of solar activity using Ganshine as a proxy yielded S = 0.164 and log R'HK = −4.97 (Isaacson & Fischer 2010) for 2009 August 31 when the Sun was still at a low state of its cycle. Thus, Kepler-68 is a magnetically inactive star, consistent with its low rotational rate, vsin i = 0.5 km s^−1. Kepler-68 appears to be a middle-aged (age 2–10 Gyr), slowly rotating inactive star on the main sequence. This is consistent with the age derived from the asteroseismology analysis (Section 6.2).

The activity indices, S and log R'HK, have modestly lower than solar values, while Figure 1 suggests significantly lower than typical levels of photometric variability compared to the Sun. Hall et al. (2009) summarize 14 years of contemporaneous photometric and spectroscopic (for activity) measurements of 28 solar analog stars with precisions in the photometry capable of detecting changes at roughly half the level seen for the Sun. Although a strong correlation between photometric variability exists with activity, they find that stars with near-solar activity indices show a range of half (not well bounded) to twice solar in photometric variability. Kepler-68 seems consistent with this lower range of photometric change at a given activity level, although not knowing what phase of a Kepler-68 activity cycle was sampled limits fidelity. Multi-year Kepler data will eventually enable robust activity—photometric variability understanding.

6. STELLAR CHARACTERISTICS

6.1. Spectroscopic Parameters

We did a spectral synthesis analysis using SME (Valenti & Piskunov 1996; Valenti & Fischer 2005) of one of our high resolution template spectra from Keck-HIRES of Kepler-68 to derive an effective temperature, T_eff = 5754 ± 44 K, surface gravity, log g = 4.2 ± 0.1 (cgs), metallicity, [Fe/H] = +0.10 ± 0.04, and rotational vsin i = 0.5 ± 0.5 km s^−1.

The above effective temperature was used to constrain the fundamental stellar parameters derived via asteroseismic analysis (see Section 6.2). The asteroseismology analysis gave log g = 4.281 ± 0.008, which is 0.1 dex higher than the SME value. The asteroseismology analysis is likely superior because of the high sensitivity of the acoustic periods to stellar radius. Still, the asteroseismology result depended on adopting the value of T_eff from SME. We recomputed the SME analysis by freezing (adopting) the seismology value for log g. See Torres et al. (2012) for a recent discussion of such an iteration in the analogous context of high S/N transit light-curve analysis providing the log g “truth.” This iteration yielded values of T_eff = 5793 ± 44 K, [Fe/H] = +0.12 ± 0.04, and rotational vsin i = 0.5 ± 0.5 km s^−1. The revised effective temperature was then put back into the asteroseismic calculation to further constrain the stellar radius and gravity. This iterative process converged quickly, as the log g from seismology yielded an SME value for T_eff that was only slightly different from the original unconstrained determination, and the asteroseismic log g using the iterated T_eff remained at 4.281. Likely systematic errors on T_eff and [Fe/H] of 59 K and 0.062 dex have been derived by Torres et al. (2012) in comparing results across multiple spectroscopic packages for a large number of stars. Adding these additional errors in quadrature with the errors quoted above results in more reasonable total errors of 74 K and 0.074 dex.

The utility of asteroseismology for exoplanet interpretations fundamentally rests on recognizing that both asteroseismic and transit light curve modeling (at high S/N levels and with known or assumed eccentricity) provide constraints on the same stellar parameter—ρ⋆. Thus, high precision knowledge of ρ⋆, which commonly results when asteroseismology is feasible, provides for a natural means of tightening the exoplanet transit light curve solution by adopting this as a prior—see Gilliland et al. (2011b) and Nutzman et al. (2011). Kepler-68 at KP = 10.00, near solar temperature, and a KIC radius of 1.06 R_J was recognized early as an excellent candidate for asteroseismology, with a
predicted 99% chance of success with only two months of SC data (Chaplin et al. 2011). Kepler-68 was observed at SC (Gilliland et al. 2010) for one month in Quarter 2 as a KASC survey target, then added to the science team SC targets from Quarter 5 onward to support asteroseismology and fine analyses of the short transit ingress and egress.

The power of asteroseismology in setting estimates of the stellar radius, which determines the exoplanet radius via ratio with the transit depth depending on the square of this ratio, can be seen by recounting knowledge of $R_*$ for Kepler-68 at the time of initial analyses. Transit light curve solutions for Kepler-68b with stellar radius as a free parameter returned values of 1.63 $R_\odot$, compared to a KIC radius of 1.06 $R_\odot$ (with spectroscopic solutions favoring something like the KIC value). These two stellar radius values would lead to differences of a factor of 3.6 in the planetary density, emphasizing the need for better knowledge of the stellar radius.

Figure 12 shows the Kepler-68 power spectral density with input of 10 months of SC data, the near-evenly spaced peaks characteristic of solar-like oscillations are obvious. Even when first inspecting power spectra from single months of SC data, the spacing between modes was trivial to estimate at about 100 $\mu$Hz compared to a solar value of 135 $\mu$Hz. Since the mean stellar density is known to scale as the square of frequency spacing (Ulrich 1986), this allowed an early constraint that Kepler-68 was at 0.55 of the solar mean density, and assuming a solar mass (given the solar $T_{\text{eff}}$) provided an estimate of 1.22 $R_\odot$ at the back-of-the-envelope level of analysis.

The high S/N of the power spectrum shown in Figure 12 made peak-bagging (derivation of frequencies for individual modes) straightforward, and eight team members fit the modes using Maximum Likelihood Estimation, and MCMC approaches (e.g., see Fletcher et al. 2009; Handberg & Campante 2011; Appourchaux et al. 2012). The adopted frequencies listed in Table 3 came from the single set that most consistently was at the median over all eight.

The fitting technique has been reported in various versions for the analysis of HST observations of HD 17156 (Gilliland et al. 2011b) and Kepler observations of HAT-P-7 (Christensen-Dalsgaard et al. 2010) and Kepler-10 (Batalha et al. 2011). The underlying stellar evolution modeling is provided using the ASTEC code (Christensen-Dalsgaard 2008a), with eigenfrequency analyses coming from ADIPLS (Christensen-Dalsgaard 2008b).

### Table 3 Measured Frequencies $\nu_{\text{eff}}$ of Kepler-68 (in $\mu$Hz)

<table>
<thead>
<tr>
<th>$n$</th>
<th>$l = 0$</th>
<th>$l = 1$</th>
<th>$l = 2$</th>
</tr>
</thead>
<tbody>
<tr>
<td>14</td>
<td>$\ldots$</td>
<td>1661.02 ± 0.35</td>
<td>$\ldots$</td>
</tr>
<tr>
<td>15</td>
<td>1668.43 ± 0.29</td>
<td>1713.38 ± 0.15</td>
<td>1761.36 ± 0.32</td>
</tr>
<tr>
<td>16</td>
<td>1767.09 ± 0.29</td>
<td>1813.49 ± 0.19</td>
<td>1861.58 ± 0.37</td>
</tr>
<tr>
<td>17</td>
<td>1867.94 ± 0.17</td>
<td>1914.65 ± 0.20</td>
<td>1962.99 ± 0.20</td>
</tr>
<tr>
<td>18</td>
<td>1969.08 ± 0.14</td>
<td>2016.27 ± 0.11</td>
<td>2064.85 ± 0.11</td>
</tr>
<tr>
<td>19</td>
<td>2070.73 ± 0.09</td>
<td>2117.70 ± 0.07</td>
<td>2166.32 ± 0.23</td>
</tr>
<tr>
<td>20</td>
<td>2172.02 ± 0.13</td>
<td>2219.40 ± 0.14</td>
<td>2268.08 ± 0.20</td>
</tr>
<tr>
<td>21</td>
<td>2273.37 ± 0.15</td>
<td>2321.09 ± 0.14</td>
<td>2369.55 ± 0.60</td>
</tr>
<tr>
<td>22</td>
<td>2375.43 ± 0.25</td>
<td>2423.57 ± 0.29</td>
<td>2472.54 ± 0.64</td>
</tr>
<tr>
<td>23</td>
<td>2477.86 ± 0.36</td>
<td>2525.73 ± 0.30</td>
<td>$\ldots$</td>
</tr>
</tbody>
</table>

In the present case, the stellar parameters grid includes a few values of the mixing-length parameter $\alpha_{\text{ML}}$ in addition to the mass $M$ and the initial composition, characterized by the abundances $X$ and $Z$ by mass of hydrogen and heavy elements. Thus, the evolution sequences are characterized by a set of parameters $\{P_k\} = \{M, Z, X, \alpha_{\text{ML}}\}$. Details of the Kepler-68 grid are presented below.

The fit of a given model to the data is defined in terms of

$$
\chi^2 = \frac{1}{N-1} \sum_{i=1}^{N} \left( \frac{\nu_i^{(\text{obs})} - \nu_i^{(\text{mod})}}{\sigma_i} \right)^2,
$$

where $\nu_i^{(\text{obs})}$ and $\nu_i^{(\text{mod})}$ are the observed and model frequencies and $\sigma_i$ is the error in the observed frequencies. It is assumed that the degree and order of the observed frequencies have already been determined. In addition, an augmented fit

$$
\chi^2 = \chi^2_v + \left( \frac{T_{\text{eff}}^{(\text{obs})} - T_{\text{eff}}^{(\text{mod})}}{\sigma(T_{\text{eff}})} \right)^2,
$$

including the observed and model effective temperature $T_{\text{eff}}$, is formed.

For each evolution sequence, frequencies are calculated for selected models along the sequence (typically every fifth), but such that frequencies are also available at all models in the sequence in the vicinity of the model $M_{\text{min}}$ with the smallest $\chi^2_v$. Based on homology scaling, we then assume that the frequencies in the vicinity of $M_{\text{min}}$ can be obtained as $r \nu_i(M_{\text{min}})$, where
\( r = [r(R/M_{\text{min}})]^{-3/2}, R \) being the surface radius of a model intermediate between the actual timesteps in the evolution sequence. The best-fitting of such model is determined by minimizing
\[
\chi^2(r) = \frac{1}{N-1} \sum_{i=1}^{N} \left( \frac{v_i^{(\text{obs})} - r v_i(M_{\text{min}})}{\sigma_i} \right)^2
\]
as a function of \( r \). The resulting value \( r_{\text{min}} \) of \( r \) defines an estimate \( R_{\text{min}} \) of the radius of the best-fitting model along the given sequence. In this way, we ensure that the scaled model is intermediate between two successive timesteps in the evolution sequence for which frequencies have been calculated. Linear interpolation to \( R_{\text{min}} \) then defines the final best-fitting model \( M_{\text{min}}(P_k) \) (which obviously in general does not coincide with a timestep in the evolution sequence) for the given set of model parameters \( \{P_k\} \), and with corresponding \( \chi^2(r_{\text{min}}(P_k)) \) and \( \chi^2_{\text{min}}(P_k) \). To obtain the final best-fitting model, we find the parameter set corresponding to the smallest \( \chi^2(r_{\text{min}}(P_k)) \) (or \( \chi^2_{\text{min}}(P_k) \)) amongst all of the evolution sequences. The best-fitting frequencies, e.g., for comparison with the observations, are obtained by applying the appropriate scaling \( r_{\text{min}} \) to the frequencies of the model \( M_{\text{min}} \) in the minimizing sequence.

An important goal of the fit is obviously to obtain statistically well-characterized estimates of the stellar properties, in particular density, radius, mass, and age. In the present analysis, these were determined as averages and standard deviations of the properties of the models, \( M_{\text{min}}(P_k) \), over the parameters, \( \{P_k\} \), with the weights \( \chi^2_{\text{min}}(P_k) \).

The calculations used the latest OPAL equation of state tables (see Rogers et al. 1996) and OPAL opacities at temperatures above 10^4 K (Iglesias & Rogers 1996); at lower temperatures, the Ferguson et al. (2005) opacities were used. Nuclear reactions were calculated using the NACRE parameters (Angulo et al. 1999). Diffusion and settling of helium was treated using the Michaud & Proffitt (1993) approximations; diffusion and settling of heavy elements was not taken into account. Convection was treated using the Böhm-Vitense (1958) mixing-length formulation. Although some of the relevant models have a small convective core, core overshoot was not considered.

The spectroscopically determined \([\text{Fe/H}] = 0.12 \pm 0.04 \) (see Section 6.1) is related to the model quantities \( X \) and \( Z \) by
\[
[\text{Fe/H}] = \log \left( \frac{Z/X}{Z_\odot/X_\odot} \right),
\]
where “⊙” denotes solar values. To obtain the composition from this, one clearly needs to assume a value of \( Z_\odot/X_\odot \) and a value of \( X \) or a relation between \( Z \) and \( X \). For the former, the Grevesse & Noels (1993) value of \( Z_\odot/X_\odot = 0.0245 \) is used. We recognize that substantially lower values have been obtained in more recent solar spectroscopic analyses (see Asplund et al. 2009 and references therein). However, these determinations lead to solar models showing a substantial increase in the discrepancy with the helioseismically determined solar structure (e.g., Bahcall et al. 2005; Christensen-Dalsgaard et al. 2009), compared with models based on the Grevesse & Noels (1993) composition. Given that the reasons for this discrepancy are so far unknown, we prefer to use as reference a solar composition that provides reasonable agreement with the helioseismic inferences. As a model of Galactic chemical evolution, \( \Delta Y = 2\Delta Z \) has often been used, and hence (fixing the relation roughly to the Sun) \( X = 0.7679 - 3Z \). The grid in composition allows for a spread in \([\text{Fe/H}] \) using models with \([\text{Fe/H}] = 0.02, 0.1, \) and \( 0.18 \), consistent with the \([\text{Fe/H}] \) error of 0.074 after the inclusion of systematics. With the transformation discussed above, this corresponds to \((Z, X) = (0.0183, 0.7130), (0.0217, 0.7029), \) and \((0.0256, 0.6910)\). To avoid being restricted to a specific relation between \( X \) and \( Z \), we have computed models for all nine resulting combinations of \( X \) and \( Z \).

The values of the mixing length have been chosen as \( a_{\text{ML}} = 1.5, 1.8, \) and \( 2.1 \). Models were computed initially between 0.9 and 1.2 \( M_\odot \), with a step of 0.02 \( M_\odot \). The grid in mass was later refined, with a step of 0.01 \( M_\odot \), in the vicinity of the best-fitting models. A total of 282 evolution sequences, with typically 200–300 models in each, are considered in the fit.

The present use of adiabatic frequency calculations, and an inadequate modeling of the near-surface layers, cause errors in the resulting frequencies that must be taken into account in the fit. Here, we use a correction to the frequencies for these near-surface errors, of the form
\[
\delta v = a(v/v_0)^b
\]
(Kjeldsen et al. 2008), where \( b = 4.90 \) is obtained from a corresponding solar fit, \( v_0 = 2071.33 \mu\text{Hz} \) was fixed in the middle of the observed frequency range, and \( a = -1.527 \mu\text{Hz} \) was obtained from a fit of a suitable reference model to the observed frequencies.

With these procedures, the best-fitting model had \( \chi^2 = 5.0 \) (see Equation (2)). The quality of the fit is illustrated in the echelle diagram (Grec et al. 1983) shown in Figure 13. The stellar parameters and standard deviations obtained from the weighted averages over the evolution sequences are shown in Table 4.

As a consistency check, the global oscillation parameter values \( \Delta v = 101.51 \pm 0.09 \mu\text{Hz} \) and \( v_{\text{max}} = 2154 \pm 13 \mu\text{Hz} \) were derived from the best-fitting (peak bagging) frequencies and amplitudes of the most prominent peaks, providing values quite consistent with estimations from standard, automated detection codes (Hekker et al. 2010; Christensen-Dalsgaard et al. 2010; Verner et al. 2011). The resulting density, mass, and radius provided by grid-based solutions (Basu et al. 2010; Karoff et al. 2010) were consistent with our more detailed fit.
Table 4
Star and Planet Parameters for the Kepler-68 System

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value</th>
<th>Notes</th>
</tr>
</thead>
<tbody>
<tr>
<td><strong>Transit and orbital parameters: Kepler-68b</strong></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Orbital period $P$ (days)</td>
<td>5.398763 ± 0.000004</td>
<td>A</td>
</tr>
<tr>
<td>Midtransit time $E$ (BJD)</td>
<td>2455006.85729 ± 0.00042</td>
<td>A</td>
</tr>
<tr>
<td>Scaled semimajor axis $a/R_*$</td>
<td>10.68 ± 0.14</td>
<td>A</td>
</tr>
<tr>
<td>Scaled planet radius $R_p/R_*$</td>
<td>0.01700 ± 0.00046</td>
<td>A</td>
</tr>
<tr>
<td>Impact parameter $b$</td>
<td>0.45 ± 0.17</td>
<td>A</td>
</tr>
<tr>
<td>Orbital inclination $i$ (deg)</td>
<td>87.60 ± 0.90</td>
<td>A</td>
</tr>
<tr>
<td>Orbital semi-amplitude $K$ (m s$^{-1}$)</td>
<td>2.63 ± 0.71</td>
<td>B</td>
</tr>
<tr>
<td>Orbital eccentricity $e$</td>
<td>0</td>
<td>B</td>
</tr>
<tr>
<td>Center-of-mass velocity $y$ (km s$^{-1}$)</td>
<td>−20.9 ± 0.1</td>
<td>B</td>
</tr>
<tr>
<td><strong>Transit and orbital parameters: Kepler-68c</strong></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Orbital period $P$ (days)</td>
<td>9.605085 ± 0.000072</td>
<td>A</td>
</tr>
<tr>
<td>Midtransit time $E$ (BJD)</td>
<td>2454969.3805 ± 0.0041</td>
<td>A</td>
</tr>
<tr>
<td>Scaled semimajor axis $a/R_*$</td>
<td>15.68 ± 0.20</td>
<td>A</td>
</tr>
<tr>
<td>Scaled planet radius $R_p/R_*$</td>
<td>0.00703 ± 0.00025</td>
<td>A</td>
</tr>
<tr>
<td>Impact parameter $b$</td>
<td>0.84 ± 0.11</td>
<td>A</td>
</tr>
<tr>
<td>Orbital inclination $i$ (deg)</td>
<td>86.93 ± 0.41</td>
<td>A</td>
</tr>
<tr>
<td>Orbital semi-amplitude $K$ (m s$^{-1}$)</td>
<td>1.25 ± 0.05</td>
<td>B</td>
</tr>
<tr>
<td><strong>Observed stellar parameters</strong></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Effective temperature $T_{\text{eff}}$ (K)</td>
<td>5793 ± 74</td>
<td>C</td>
</tr>
<tr>
<td>Spectroscopic gravity log $g$ (cgs)</td>
<td>4.281 ± 0.06</td>
<td>C</td>
</tr>
<tr>
<td>Metallicity [Fe/H]</td>
<td>+0.12 ± 0.074</td>
<td>C</td>
</tr>
<tr>
<td>Projected rotation $v\sin i$ (km s$^{-1}$)</td>
<td>0.5 ± 0.5</td>
<td>C</td>
</tr>
<tr>
<td><strong>Fundamental Stellar Properties</strong></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Density $\rho_*$ (g cm$^{-3}$)</td>
<td>0.7903 ± 0.0054</td>
<td>D</td>
</tr>
<tr>
<td>Mass $M_* (M_\odot)$</td>
<td>1.079 ± 0.051</td>
<td>D</td>
</tr>
<tr>
<td>Radius $R_*(R_\odot)$</td>
<td>1.243 ± 0.019</td>
<td>D</td>
</tr>
<tr>
<td>Surface gravity log $g_*$ (cgs)</td>
<td>4.281 ± 0.008</td>
<td>D</td>
</tr>
<tr>
<td>Luminosity $L_*$ ($L_\odot$)</td>
<td>1.564 ± 0.141</td>
<td>D</td>
</tr>
<tr>
<td>Absolute $V$ magnitude $M_V$ (mag)</td>
<td>4.34 ± 0.09</td>
<td>D</td>
</tr>
<tr>
<td>Age (Gyr)</td>
<td>6.3 ± 1.7</td>
<td>D</td>
</tr>
<tr>
<td>Distance (pc)</td>
<td>135 ± 10</td>
<td>D</td>
</tr>
<tr>
<td><strong>Planetary parameters: Kepler-68b</strong></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Mass $M_p (M_\oplus)$</td>
<td>8.3$^{+2.2}_{-2.4}$</td>
<td>A,B,C,D,F</td>
</tr>
<tr>
<td>Radius $R_p (R_\oplus)$</td>
<td>2.31$^{+0.06}_{-0.09}$</td>
<td>A,B,C,D</td>
</tr>
<tr>
<td>Density $\rho_p$ (g cm$^{-3}$)</td>
<td>3.32$^{+0.08}_{-0.08}$</td>
<td>A,B,C,D,F</td>
</tr>
<tr>
<td>Surface gravity log $g_p$ (cgs)</td>
<td>3.14 ± 0.11</td>
<td>A,B,C,D,F</td>
</tr>
<tr>
<td>Orbital semimajor axis $a$ (AU)</td>
<td>0.06170 ± 0.00056</td>
<td>E</td>
</tr>
<tr>
<td>Equilibrium temperature $T_{\text{eq}}$ (K)</td>
<td>1280 ± 90</td>
<td>G</td>
</tr>
<tr>
<td><strong>Planetary parameters: Kepler-68c</strong></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Mass $M_p (M_\oplus)$</td>
<td>4.8$^{+2.5}_{-3.6}$</td>
<td>A,B,C,D,F</td>
</tr>
<tr>
<td>Radius $R_p (R_\oplus)$</td>
<td>0.953$^{+0.037}_{-0.042}$</td>
<td>A,B,C,D</td>
</tr>
<tr>
<td>Density $\rho_p$ (g cm$^{-3}$)</td>
<td>28$^{+13}_{-23}$</td>
<td>A,B,C,D,F</td>
</tr>
<tr>
<td>Orbital semimajor axis $a$ (AU)</td>
<td>0.09059 ± 0.00082</td>
<td>E</td>
</tr>
<tr>
<td><strong>Planetary parameters: Kepler-68d</strong></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Orbital period $P$ (days)</td>
<td>580 ± 15</td>
<td>F</td>
</tr>
<tr>
<td>Minimum mass $M_p \sin i (M_\oplus)$</td>
<td>0.947 ± 0.035</td>
<td>D,F</td>
</tr>
<tr>
<td>Orbital semi-amplitude $K$ (m s$^{-1}$)</td>
<td>19.9 ± 0.75</td>
<td>D,F</td>
</tr>
<tr>
<td>Orbital semimajor axis $a$ (AU)</td>
<td>1.40 ± 0.03</td>
<td>D,F</td>
</tr>
<tr>
<td>Orbital eccentricity $e$</td>
<td>0.18 ± 0.05</td>
<td>F</td>
</tr>
</tbody>
</table>

Notes. A: Based primarily on an analysis of the photometry. B: Based on a joint analysis of the photometry and radial velocities. C: Based on an analysis by D. Fischer of the Keck/HIRES template spectrum using SME (Valenti & Piskunov 1996). D: Based on asteroseismology analysis. E: Based on Newton’s revised version of Kepler’s Third Law and the results from D, F: Based on radial velocities, G: Calculated assuming a random distribution of Bond albedo over 0.0–0.5, and a random set ranging from zero to full redistribution of heat from day to night sides.

7. PLANET CHARACTERISTICS

7.1. Fits to Photometry and Radial Velocities

The physical and orbital properties of both transit signatures are derived by simultaneously fitting Kepler photometry and Keck RVs and by adopting the mean stellar density, $R_*$, and $M_*$ of the host star as determined by asteroseismology.

The Kepler photometry and Keck RVs are fit with non-interacting Keplerian orbits. The model parameters are the mean stellar density ($\rho_*$) and a flux and RV zero point, and the time of transit ($T_0$), orbital period ($P$), impact parameter ($b$), scaled planetary radius ($R_p/R_*$), RV amplitude ($K$) and eccentricity and argument of pericenter parameterized as $e\cos(w)$ and $e\sin(w)$ for each planet. The transit was modeled using the analytic formalization of Mandel & Agol (2002) to fit photometric observations of the transit. We use the quadratic parameterization of limb darkening also described by Mandel & Agol (2002) with coefficients (0.4096, 0.2602) calculated by Claret & Bloemen (2011) for the Kepler bandpass. Model presented here. Likewise, fits using the Asteroseismic Modeling Portal (see Metcalfe et al. 2009 for details) gives values very close to those presented here.

Our primary asteroseismic solutions used the formal spectroscopic errors of 44 K on $T_{\text{eff}}$ and allowed a spread of $±0.08$ dex on [Fe/H]. We have used grid-based solutions starting with these errors, and the more appropriate errors of 74 K on $T_{\text{eff}}$ and 0.074 dex on [Fe/H] to show that changes in the directly constrained stellar density are negligible, and that inferred values of stellar mass and radius change by 0.2$\sigma$ and 0.3$\sigma$, respectively, in comparison to the errors already quoted in Table 4. The associated errors on stellar properties also changed little in comparison to values already in use.

The asteroseismic solutions given here are based on only about half of the SC data now available. However, the quality of the frequency extractions for Kepler-68 with the data through Quarter 7 only is already sufficiently good that residual errors from the asteroseismic solution are negligible for the inference of exoplanet parameters. Indeed, the asteroseismology for Kepler-68 will be among the very best possible with Kepler, and further results concentrating on fine details of this star will appear in the future. As an example, some of the best fitting stellar evolution models indicated a small convective core, hence it would be good to explore the inclusion of convective core overshoot. Initial exploration of core overshoot indicated that this had negligible impact on the mean stellar density required for exoplanet inferences and was not further pursued. The quality of Kepler-68 data for asteroseismology will likely support inferences on the outer convection zone depth (Mazumdar et al. 2012) and obliquity of the rotation axis (Chaplin et al. 2013) in future analyses.

Kepler-68 is a near twin of a Cen A, with inferred values of $T_{\text{eff}}$, $R_*$, $M_*$, $L_*$, and age all within 1$\sigma$ of each other, despite the small error bars for both—see Bazot et al. (2012) and references therein. Of the fundamental parameters, only the [Fe/H] differs significantly, with a Cen A being more metal rich at 0.24 ± 0.02 (Neuorge–Verheeeke & Magain 1997) compared to Kepler-68 at 0.12 ± 0.074. Both stars have now been interpreted with the benefit of asteroseismology. The quality of asteroseismic constraints are superior for Kepler-68, with general astronomical knowledge being better for the nearby binary a Cen A. Differential analyses of these two interesting stars may prove fruitful.
fits to the Kepler-68b lightcurve yield an eccentricity that is consistent with zero \(e \cos(w) = 0.02 \pm 0.10; e \sin(w) = 20.13 \pm 0.20\), which is consistent with our expectations for tidal circularization (Mazeh 2008). Given the large orbital separation of the outer planet candidate, we can not assume its orbit to be circular based on tidal circularization. The duration of the transit for the outer planet candidate is consistent with a circular orbit, but the resulting upper limit is still significant \(e \approx 0.2\). For the remainder of our discussion, the models are constrained to zero eccentricity for both Kepler-68b and Kepler-68c. The RV variations are modeled by assuming non-interacting (Keplerian) orbits. With Kepler-68c validated, the relative inclination between the two orbits is likely less than 20 degrees, as larger relative inclinations would require a fortuitous alignment of the orbital nodes for both planets to transit (Ragozzine & Holman 2010).

We initially fit our observations by fixing \(\rho_s\) to its asteroseismic values (see Section 6.2). Model parameters are found by chi-squared minimization using a Levenberg-Marquardt prescription. We then use the best-fit values to seed an MCMC parameter search (Ford 2005) to fit all of the model parameters with \(\rho_s\) from the asteroseismic solution. We adopt the asteroseismic determined mean-stellar density as a prior of the overall solution.

### 7.2. Transit Timing Variation Prospects

For a nominal two-planet model (i.e., circular orbits with masses from Table 4), the predicted root mean square (rms) TTVs for Kepler-68b & Kepler-68c are both less than half a minute. The median timing uncertainties (based on long-cadence observations) are 3.3 and 20 minutes. Thus, it is not surprising that Kepler has not provided a TTV signal due to the interaction of Kepler-68b and Kepler-68c from initial searches. Even with a possible factor of two precision gain from SC data and an available long time-series providing the square root of the number of transits gain, predictions are that detection of a TTV signal on Kepler-68b remains marginal, and unlikely for Kepler-68c. Even increasing the masses of both planets by three times the upper “1\(\sigma\)” uncertainty above the estimates in Table 4, the predicted rms TTVs are less than a minute for both planets. Similarly, even models with unrealistically large eccentricities (\(e = 0.3\) and the nominal masses) can result in rms TTVs of less than a minute. Therefore, we have not performed a detailed TTV analysis of this system.

### 7.3. Composition of Kepler-68b

The synthesis of RV monitoring, transit photometry, and precise asteroseismic stellar characterization reveals that Kepler-68b’s mass, radius, and density are all intermediate between the properties of Earth and the solar system ice giants. Kepler-68b’s bulk density \(3.32^{+0.30}_{-0.39} \text{g cm}^{-3}\) is low enough to imply that volatiles (in the form of H/He or astrophysical ices) make a significant contribution to the total planet mass and volume. Kepler-68b cannot be composed of iron and silicates alone; an iron-poor silicate composition is too dense by more than 3\(\sigma\). Even a carbon-rich mineralogy—which may lead to solid planets with larger radii than the Earth-like mineralogy often assumed (Madhusudhan et al. 2012)—does not account for the planet density within 1\(\sigma\). Following Rogers & Seager (2010) and Rogers et al. (2011), we constrain the range of bulk compositions that are consistent with Kepler-68b’s measured mass and radius. Assuming an Earth-like rocky interior composition (consisting of 32% Fe and 68% silicate by mass), Kepler-68b would need between 0.07% and 0.6% of its mass in a H/He gas layer, or, alternatively, between 21% and 76% of its mass in a water vapor envelope. Intermediate compositions with a mixture of H/He and higher mean molecular weight material from ices are also possible.

Compared to Kepler-68b, the compositions of the other planets orbiting Kepler-68 are less well constrained because the bulk planet densities are unknown. Given its Jupiter-like minimum mass, the outermost planet (Kepler-68d) is likely to be dominated by H/He. Without a transit measurement of Kepler-68d’s radius, however, neither the dominant composition nor the more subtle proportion of heavy elements in the planet interior can be directly inferred. At the other extreme of the planetary mass scale, Earth-sized Kepler-68c has only a marginal RV detection. Planet interior structure models place more stringent constraints on the planet mass than the 2\(\sigma\) RV upper limit of 10.6 \(M_\oplus\). If Kepler-68c is a rocky body composed of iron and silicate, its mass would fall within 0.65 \(M_\oplus < M_p < 2.5 M_\oplus\)—even a pure iron configuration is allowed. A residual RV precision better than 70 cm s\(^{-1}\) is needed to constrain Kepler-68c’s make-up.

A striking feature of the Kepler-68 planetary system is that it harbors one of the most strongly irradiated volatile-rich mini-Neptunes detected to date. The stellar energy flux received by Kepler-68b is more than 412 ± 34 times larger than that received by the Earth. Among low-mass \((M_p < 20 M_\oplus\) planets with measured radii, only Kepler-10b, 55 Cnc e, CoRoT-7b, and Kepler-18b are more strongly irradiated by their host stars. These planets have higher densities than Kepler-68b, however, and are consistent with volatile-less compositions within the 1\(\sigma\) uncertainties on their measured masses and radii. Kepler-10b, CoRoT-7b, and Kepler-18b may be comprised solely of iron and silicates (with no H/He or astrophysical ices), and 55 Cnc e (although not dense enough to have a silicate configuration) could have a carbon-rich solid composition without a volatile envelope (Madhusudhan et al. 2012). Kepler-68b’s status as the most strongly irradiated mini-Neptune that unambiguously (when the 1\(\sigma\) uncertainties are taken into account) has a significant amount of volatiles makes it a valuable benchmark for planet mass-loss models.

Like many of the highly irradiated Kepler planets, mass loss has likely had an important influence sculpting Kepler-68b’s composition over its 6.3 Gyr lifetime. Indeed, Kepler-68b lies near the edge of the empirical mass loss destruction threshold noted by several authors (Lecavelier 2007; Ehrenreich & Désert 2011; Jackson et al. 2012; Lopez et al. 2012). At 8.3 \(M_\oplus\), Kepler-68b is close to the minimum mass of 6.5 \(M_\oplus\) predicted by mass loss in Lopez et al. (2012). In order to examine the vulnerability of Kepler-68b in greater detail, we employed the coupled thermal and mass loss evolution models of Lopez et al. (2012), assuming an Earth-like core and 10% mass loss efficiency.

Although, Kepler-68b is stable against mass loss today, it is possible that it underwent substantial mass loss early in its history when radii were larger and stellar XUV fluxes were over 100× higher (Ribas et al. 2005). In fact, if Kepler-68b had a H/He envelope then, it was vulnerable to a type of runaway mass loss that occurs when the mass loss timescale becomes short compared to the cooling timescale (Burrau et al. 2004; Lopez et al. 2012). Although less than 1% H/He today, Kepler-68b would need to have been ~80% H/He when it was 10 Myr old in order to have a residual primordial H/He envelope today. Moreover, models that undergo this type of extreme mass loss
almost always lead to a planet completely stripped of its H/He envelope (Lopez et al. 2012). The initial conditions must be carefully fine tuned in order to arrive at a planet that has such a small but non-zero H/He envelope today. This suggests that it is unlikely that Kepler-68b retains a primordial H/He envelope.

On the other hand, a steam envelope on Kepler-68b should be very stable against mass loss. If Kepler-68d is ∼50% water today, then it has only lost ∼1% of its initial water envelope since it was 10 Myr old. This suggests that it is possible that like Kepler-68d, 68b formed beyond the snow-line and migrated to its current orbit. However, another distinct possibility is that Kepler-68d does in fact have an H/He envelope, just not a primordial one. Elkins-Tanton & Seager (2008) showed that rocky planets could outgas up to 0.9% of their mass in H/He, more than sufficient to explain the envelope needed today.

On the whole, the Kepler-68 planetary system shares characteristics both with the compact Kepler multi-planet systems (e.g., Kepler-11 and Kepler-20) and with the solar system. Like Kepler-11 and Kepler-20, Kepler-68 has multiple transiting planets within 0.1 AU. In common with the solar system, Kepler-68 has a Jovian-mass planet residing outside (at greater orbital separations than) the smaller bodies in the inner system. Between 0.1 and 1.4 AU, there are no confirmed planets in the Kepler-68 system, although our census may be incomplete. The presence of a volatile-rich super Earth within 0.06 AU, combined with the presence of a luke-warm Jupiter inside the snow line, makes the Kepler-68 system an interesting case study for planet formation and migration theories.

8. SUMMARY

Two distinct sets of transit events were detected in the lightcurve of Kepler-68 constructed from ∼2 years of Kepler photometry. Physical models constrained by the asteroseismology-derived stellar parameters were simultaneously fit to the transit light curves and the precision Doppler measurements. Modeling produced tight constraints on the properties of Kepler-68b: \( \frac{M_p}{M_\oplus} = 8.3^{+2.2}_{-2.4} \), \( R_p / R_\oplus = 2.31^{+0.06}_{-0.09} \), \( p_p = 3.32^{+0.86}_{-0.99} \) g cm\(^{-3}\). Evaluation of these properties within a theoretical framework allowed us to draw conclusions about the planet’s composition, arguing that a simple iron and silicate structure is excluded. Kepler-68b must retain significant volatiles, even though it is highly irradiated.

The outer planet, Kepler-68d, is detected only in RVs for which an upper limit to the mass is approximately Jupiter in scale, and in a Mars-like orbit. Transits of Kepler-68d would not be expected if it were at the same inclination inferred from the impact parameters for Kepler-68b and Kepler-68c. At the two epochs of Kepler data in which transits of Kepler-68d would be seen if they existed, none are present. There is still a 1% chance that if Kepler-68d does transit, the transits could have been missed in minor data gaps. The presence of an outer giant planet further enriches the interpretive potential for the Kepler-68 system.

Kepler-68c, the intermediate planet, produces transits of only 55 ppm depth (less than 2/3 that of an Earth analog). But due to the brightness of Kepler-68, coupled with low stellar activity and the modest orbital period of 9.6 days, it is detected at high confidence from the Kepler photometry. RVs do not provide confirmation, although the formally inferred amplitude at a phase fixed from the transits appears at the 1σ level. The upper limit on mass from the RVs is not significant for inferring interesting aspects about the planet composition. BLENDER provided a sufficiently high odds ratio to assert that Kepler-68c is a planet. This validation came with an added complication in this case, however, in that some individual false positive scenarios (the best being for a 7 mag fainter background eclipsing binary) provided formally better fits to the phased light curve than did a simple planet transit model. With an odds ratio over 10,000 for the planet interpretation, it is proper to accept this, with perhaps some qualification reserved in this case. Were further data to more definitively provide evidence that a subtle feature at phase 0.5 in the light curve is properly interpreted as a secondary eclipse, then the statistical argument for “validation” would be over-ruled by “confirmation” as a false positive.

This qualifies as an interesting system even in the context of so many exciting and unique discoveries coming from the Kepler Mission—with a bright, quiet star providing exquisite asteroseismic constraints on stellar properties; RVs providing a precise mass for one transiting planet and supporting discovery of an outer planet; and the second transiting planet validated. That the innermost transiting planet has been shown to have a density intermediate between terrestrial and gas giant planets, with sufficient fidelity to inform theoretical models of its structure, further bolsters the assertion that the Kepler-68 exoplanet system is an important development in this rapidly expanding field.

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Facility: Kepler

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