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Search for a Standard Model Higgs Boson in the $H \rightarrow ZZ \rightarrow \ell^+\ell^-\nu\bar{\nu}$ Decay Channel with the ATLAS Detector

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A search for a heavy standard model Higgs boson decaying via $H \rightarrow ZZ \rightarrow \ell^+\ell^-\nu\bar{\nu}$, where $\ell = e, \mu$, is presented. It is based on proton-proton collision data at $\sqrt{s} = 7$ TeV, collected by the ATLAS experiment at the LHC in the first half of 2011 and corresponding to an integrated luminosity of 1.04 fb$^{-1}$. The data are compared to the expected standard model backgrounds. The data and the background expectations are found to be in agreement and upper limits are placed on the Higgs boson production cross section over the entire mass window considered; in particular, the production of a standard model Higgs boson is excluded in the region $340 < m_H < 450$ GeV at the 95% confidence level.


The search for the standard model (SM) Higgs boson [1–3] is one of the most important aspects of the Large Hadron Collider (LHC) physics program. Direct searches at the CERN LEP $e^+e^-$ collider have set a lower limit of 114.4 GeV on the Higgs boson mass, $m_H$, at 95% confidence level [4]. Searches by the CDF and D0 experiments at the Fermilab Tevatron $p\bar{p}$ collider have explored the mass range up to 200 GeV and exclude the additional region $156 < m_H < 177$ GeV [5]. For $m_H$ greater than twice the Z boson mass, $m_Z$, a significant fraction of Higgs bosons decay to two Z bosons. The $ZZ \rightarrow \ell^+\ell^-\nu\bar{\nu}$ decay channel offers a substantial branching fraction in combination with a good separation from potential background processes owing to the high transverse momentum, $p_T$, of the electron or muon pair from the leptonic Z decay and the high missing transverse momentum, $E_T^{\text{miss}}$, from the Z decaying to neutrinos.

The first cross section limits for a SM Higgs boson in the mass region $200 < m_H < 600$ GeV were set by the ATLAS and CMS collaborations in Refs. [6,7]. This letter extends the $H \rightarrow ZZ \rightarrow \ell^+\ell^-\nu\bar{\nu}$ results therein, with a 30-fold increase in the integrated luminosity, as well as a significant improvement in the event reconstruction and background rejection.

The data sample considered in this search was recorded by the ATLAS experiment during the first half of the 2011 LHC run at a center-of-mass energy $\sqrt{s} = 7$ TeV. The integrated luminosity of the data sample, considering only data-taking periods where all relevant detector subsystems were operational, is 1.04 fb$^{-1}$.

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The cross sections for Higgs boson production, the associated branching fractions [24], as well as their uncertainties, are compiled in Ref. [25]. They correspond to next-to-next-to-leading order in QCD for the gluon fusion [26–31] and the vector-boson fusion [32] processes. In addition, QCD soft-gluon resummations up to next-to-next-to-leading logarithms are available for the gluon fusion process [33], while next-to-leading order electro-weak corrections are applied to both the gluon fusion [34,35] and the vector-boson fusion [36,37] processes. These cross section calculations do not account for the width of the Higgs boson, which is implemented through an ad hoc Breit-Wigner line shape applied at the event generator level. Recent studies [25,38] have indicated that effects due to off-shell Higgs boson production and interference with other SM processes may become sizeable at the highest available precision calculations for the relevant process. To suppress backgrounds from top, W, and QCD multijet production, the dilepton invariant mass, the dilepton mass distribution is shown in Fig. 1. Inclusive Z boson production is the dominant background at this stage of the analysis. To reduce the background from events with fake E_T^{miss} due to mismeasured jets, events are rejected if the azimuthal angle between the missing transverse momentum vector, \( \vec{p}_T^{miss} \), and the leading jet in the event satisfies \( \Delta \phi (\vec{p}_T^{miss}, \vec{p}_T^{jet}) < 0.3 \). To reduce the background from top quark production, events with one or more b-tagged jets are rejected, where the b tagging is based on a single

Jets are used in this analysis to reject backgrounds from events with heavy quark decays or from events with fake E_T^{miss} due to mismeasured jets. For this purpose jets are reconstructed from clusters of energy deposits in the calorimeters using the anti-k_t algorithm [42] with a radius parameter R = 0.4. Only jets with \( p_T > 25 \text{ GeV} \) and \(|\eta| < 2.5\) are considered.

To remove leptons associated with jets, such as those originating from semileptonic decays of b hadrons, leptons are not considered in the analysis if the sum of inner detector track momenta in a cone \( \Delta R < 0.2 \) around the lepton direction is greater than 10% of the \( p_T \) of the lepton itself or if the lepton is within a distance \( \Delta R < 0.4 \) of the nearest jet.

The missing transverse momentum is measured as the (negative) vectorial sum of the transverse momenta of all clusters in the calorimeters within \(|\eta| < 4.5\) and all selected muons in the event. Calorimeter deposits associated with muons are subtracted to avoid double counting.

Events are required to contain a reconstructed primary vertex formed from at least 3 tracks and exactly two oppositely charged electrons or muons, consistent with originating from the primary vertex. The dilepton mass distribution is shown in Fig. 1. Inclusive Z boson production is the dominant background at this stage of the analysis. To suppress backgrounds from top, W, and QCD multijet production, the dilepton invariant mass, \( m_{\ell \ell} \), is required to satisfy \(|m_Z - m_{\ell \ell}| < 15 \text{ GeV} \).

To reduce the background from events with fake E_T^{miss} due to mismeasured jets, events are rejected if the azimuthal angle between the missing transverse momentum vector, \( \vec{p}_T^{miss} \), and the leading jet in the event satisfies \( \Delta \phi (\vec{p}_T^{miss}, \vec{p}_T^{jet}) < 0.3 \). To reduce the background from top quark production, events with one or more b-tagged jets are rejected, where the b tagging is based on a single
discriminating combining information from both the impact parameter with respect to the primary vertex of tracks associated to the jet and the presence of displaced secondary vertices associated to the jet’s tracks. The chosen cut achieves an efficiency of about 70% for identifying real b jets, with a light-quark jet rejection of about 80% [43].

To exploit the mass dependent kinematic features of $H \rightarrow ZZ \rightarrow \ell^+ \ell^- \nu \bar{\nu}$ production, the search is subdivided into a low Higgs boson mass ($m_H < 280$ GeV) and a high Higgs boson mass ($m_H \approx 280$ GeV) search region, where dedicated cuts are applied to two important discriminating variables used to reduce the background contributions: $E_T^{\text{miss}}$ and the azimuthal angle between the two leptons, $\Delta \phi(\ell, \ell)$. Figure 2 shows the distributions of these variables after the application of the $m_{\ell\ell}$ window cut. Since inclusive Z production gives rise to a steeply falling $E_T^{\text{miss}}$ distribution, systematic uncertainties on the $E_T^{\text{miss}}$ reconstruction are particularly important to estimate this background correctly. The dominant contributions to the $E_T^{\text{miss}}$ uncertainty come from the knowledge of the jet energy scale and the modelling of inclusive Z production. Figure 2 shows that a good agreement within systematic uncertainties is observed between data and the combined background expectation. In the low $m_H$ region, events are required to satisfy $E_T^{\text{miss}} > 66$ GeV, while in the high $m_H$ region the requirement is $E_T^{\text{miss}} > 82$ GeV. These cuts reduce significantly the backgrounds from processes with no or modest genuine missing transverse momentum originating from unobserved neutrinos.

The boost of the Z bosons originating from a Higgs boson decay increases with $m_H$, thus reducing the expected opening angle between the leptons. In the low $m_H$ region this boost is expected to be modest and a cut $1 < \Delta \phi(\ell, \ell) < 2.64$ is applied. In the high $m_H$ region an upper limit $\Delta \phi(\ell, \ell) < 2.25$ is required.

Finally, in the high $m_H$ region, events are also rejected if the azimuthal angle between the missing transverse momentum vector and the direction of the $Z \rightarrow \ell \ell$ boson candidate is $\Delta \phi(p_T^{\text{miss}}, \vec{p}_T^{\ell\ell}) < 1$. The efficiency of the event selection is very similar in the electron and muon channels, ranging from 3% for $m_H = 200$ GeV to about 48% for $m_H = 600$ GeV.

SM pair production of Z bosons has a final state identical to the signal, and is therefore expected to survive most of the applied selection criteria and form a continuum in the transverse mass distribution (defined below). The normalization for this background is obtained from a calculation including next-to-leading order terms [44] with an additional 6% term to account for missing quark-box diagrams (gg → ZZ) [45]. A 11% normalization uncertainty is assigned to this background, estimated from scale, PDF and model uncertainties. WW and WZ backgrounds are normalized in a similar way.

The background from inclusive Z production is derived from MC, after checking that the simulation describes well the data in samples selected by requiring the presence of a lepton pair. The background from top events is also taken from the MC prediction. This prediction is verified to agree with data, within systematic uncertainties, in two independent control samples: the first one requires at least one identified b-jet, while the second selects events containing electron-muon pairs.

Additional backgrounds can arise from QCD multijet events or inclusive W production due to heavy flavour decays or jets faking leptons. The normalization of the W background is obtained from the ratio between data and MC in control samples of like-sign electron-electron and electron-muon events with high $E_T^{\text{miss}}$. The QCD multijet background in the electron channel is determined using a data sample based on a loosened electron selection, thus dominated by jets; this sample is scaled to describe the tails of the $m_{\ell\ell}$ distribution. In the muon channel, the background from heavy flavour decays is studied using simulation, whereas other muon sources from multijet events are constrained using a sample of like-sign muon pairs in data. In both cases the background is found to be negligible.

The signal efficiencies and overall background expectations are similar in the electron and the muon channels.

FIG. 2 (color online). The $E_T^{\text{miss}}$ (left) and $\Delta \phi(\ell, \ell)$ (right) distributions for events with exactly two oppositely charged electrons or muons inside the Z mass window. The insets at the bottom show the ratio between data and the combined background expectations as well as a band corresponding to the combined systematic uncertainties of the analysis.
therefore only combined results are presented. The numbers of candidate $H \rightarrow ZZ \rightarrow \ell^+ \ell^- \nu \bar{\nu}$ events selected in data and the expected yields from signal and background processes are shown in Table I.

The systematic uncertainties include experimental uncertainties related to the selection and calibration of electrons, muons, jets and $b$ jets, which are also explicitly propagated to the $E_T^{miss}$ calculation. Shape uncertainties for the signal and for the single $Z$ and $ZZ$ backgrounds are estimated using PYTHIA as an alternative MC generator.

Normalization uncertainties for signal (gluon fusion $+14\%$ and VBF $-10\%$) and diboson backgrounds ($11\%$) are obtained from theory [25]; uncertainties for the inclusive $Z$ boson production ($2.5\%$), top quark production ($9\%$), inclusive $W$ boson production ($100\%$) and QCD multijet production in the electron channel ($50\%$) are estimated from data. A $3.7\%$ luminosity uncertainty [46] is included for those processes for which the normalization is not obtained from the data. The dominant systematic uncertainties in the analysis are the $E_T^{miss}$ uncertainties for the $Z$ background, the $b$-tagging uncertainty for the top background and the normalization uncertainties for the signal and the $W$ and diboson backgrounds.

After the event selection, the Higgs boson search is performed by looking for an excess of data over the SM background expectation in the transverse mass distribution of the selected $ee\nu\nu$ and $\mu\mu\nu\nu$ events. The transverse mass is calculated from the lepton pair and the $p_T^{miss}$ vector as

$$m_T^2 = \left[ \sqrt{m_Z^2 + \left| p_T^{miss} \right|^2} + \sqrt{m_Z^2 + \left| p_T^{miss} \right|^2} \right]^2 - \left| p_T^{miss} + \vec{p}_T^{miss} \right|^2.$$ 

**Table I.** The expected number of background and signal events for the Higgs boson search in the $H \rightarrow ZZ \rightarrow \ell^+ \ell^- \nu \bar{\nu}$ channel, along with the observed numbers of candidates in data, for an integrated luminosity of $1.04 \text{ fb}^{-1}$. The quoted uncertainties are statistical and systematic, respectively. Signal to background ratios are also given for various masses (see text).

<table>
<thead>
<tr>
<th>Source</th>
<th>low $m_H$ search</th>
<th>high $m_H$ search</th>
</tr>
</thead>
<tbody>
<tr>
<td>$Z$</td>
<td>$19.1 \pm 2.6 \pm 0.9$</td>
<td>$6.0 \pm 1.4 \pm 1.8$</td>
</tr>
<tr>
<td>$W$</td>
<td>$8.5 \pm 2.3 \pm 8.5$</td>
<td>$3.1 \pm 1.0 \pm 3.1$</td>
</tr>
<tr>
<td>top</td>
<td>$29.9 \pm 1.3 \pm 6.0$</td>
<td>$14.9 \pm 0.8 \pm 3.1$</td>
</tr>
<tr>
<td>multijet</td>
<td>$0.4 \pm 0.4 \pm 0.2$</td>
<td>$0.0 \pm 0.0 \pm 0.0$</td>
</tr>
<tr>
<td>$ZZ$</td>
<td>$17.6 \pm 0.4 \pm 2.1$</td>
<td>$14.7 \pm 0.4 \pm 1.7$</td>
</tr>
<tr>
<td>$WZ$</td>
<td>$16.7 \pm 0.6 \pm 2.0$</td>
<td>$12.1 \pm 0.5 \pm 1.4$</td>
</tr>
<tr>
<td>$WW$</td>
<td>$12.4 \pm 0.4 \pm 1.5$</td>
<td>$4.6 \pm 0.3 \pm 0.5$</td>
</tr>
<tr>
<td>Total</td>
<td>$104.6 \pm 3.8 \pm 16.0$</td>
<td>$55.3 \pm 2.0 \pm 7.8$</td>
</tr>
</tbody>
</table>

$\mu$-decay channels are combined. Figure 3 shows the $m_T$ distribution in the high $m_H$ search region. Signal to background ratios for different $m_H$ values, determined in a $m_T$ window defined to enclose 95% of the corresponding signal events, are listed in Table I.

The number and distribution of candidate $H \rightarrow ZZ \rightarrow \ell^+ \ell^- \nu \bar{\nu}$ events observed in the data agree with the expected backgrounds within the uncertainties, with no indication of an excess. Upper limits are set on the Higgs boson production cross section relative to its predicted SM value as a function of $m_H$. The limits are extracted from a maximum likelihood fit to the $m_T$ distribution following the $CL_s$ modified frequentist formalism with the profile likelihood test statistic [47,48]. All systematic uncertainties are taken into account.

Figure 4 shows the expected and observed limits at the 95% confidence level. The expected limit is lowest around $m_H = 380 \text{ GeV}$ where it is 1.1 times the SM Higgs boson cross section. Fluctuations in the background can lead to better or worse expected limits. Over the entire mass range the observed limits agree with the expectations within the $\pm 2\sigma$ band. A SM Higgs boson in the range $340 \text{ GeV} < m_H < 450 \text{ GeV}$ is excluded at the 95% confidence level.

In summary, results of a search for a heavy SM Higgs boson with a mass in the range $200 < m_H < 600 \text{ GeV}$ decaying to $ZZ \rightarrow \ell^+ \ell^- \nu \bar{\nu}$ have been presented. These results are based on a data sample corresponding to an integrated luminosity of $1.04 \text{ fb}^{-1}$, recorded with the ATLAS detector at the LHC. No evidence for a signal is observed and cross section limits are placed over the entire mass range, excluding the production of a SM Higgs boson in the region $340 < m_H < 450 \text{ GeV}$ at the 95% confidence level.

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FIG. 4 (color online). Observed and expected 95% confidence level upper limits on the Higgs boson production cross section divided by the SM prediction. The green and yellow bands indicate the divided by the SM prediction. The green and yellow bands indicate the ±1σ and ±2σ fluctuations, respectively, around the median sensitivity. The limits are based on 1.04 fb⁻¹ of data at √s = 7 TeV.

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