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THE X-RAY FLARING PROPERTIES OF Sgr A* DURING SIX YEARS OF MONITORING WITH SWIFT

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Abstract

Starting in 2006, Swift has been targeting a region of ∼21′ × 21′ around Sagittarius A* (Sgr A*) with the onboard X-Ray Telescope. The short, quasi-daily observations offer a unique view of the long-term X-ray behavior of the supermassive black hole. We report on the data obtained between 2006 February and 2011 October, which encompasses 715 observations with a total accumulated exposure time of ∼0.8 Ms. A total of six X-ray flares were detected with Swift, which all had an average 2–10 keV luminosity of \( L_X \sim (1 - 3) \times 10^{35} \) erg s\(^{-1}\) (assuming a distance of 8 kpc). This more than doubles the number of such bright X-ray flares observed from Sgr A*. One of the Swift-detected flares may have been softer than the other five, which would indicate that flares of similar intensity can have different spectral properties. The Swift campaign allows us to constrain the occurrence rate of bright \( (L_X \gtrsim 10^{35} \) erg s\(^{-1}\)\) X-ray flares to be \( \sim 0.1\)–0.2 day\(^{-1}\), which is in line with previous estimates. This analysis of the occurrence rate and properties of the X-ray flares seen with Swift offers an important calibration point to assess whether the flaring behavior of Sgr A* changes as a result of its interaction with the gas cloud that is projected to make a close passage in 2013.

Key words: accretion, accretion disks – black hole physics – Galaxy: center – X-rays: individual (Sgr A*)

Online-only material: color figure

1. INTRODUCTION

Understanding accretion onto supermassive black holes and the associated feedback to their environment lies at the basis of understanding their formation, growth and evolution, the chemical enrichment of the interstellar medium, galaxy evolution, and the formation of large-scale structures in the universe. Sagittarius A* (Sgr A*) is a supermassive black hole that forms the dynamical center of the Milky Way (e.g., Reid & Brunthaler 2004; Ghez et al. 2008; Gillessen et al. 2009). It is the most nearby Galactic nucleus, located at ∼8 kpc, and therefore allows for an unparalleled study of the fueling process of supermassive black holes. 6

The bolometric luminosity of Sgr A* is ∼8–9 orders of magnitude lower than expected for Eddington-limited accretion onto a supermassive black hole with a mass of ∼4 \( \times 10^6 \) M\(_{\odot}\). Its puzzling sub-luminous character has been explained in terms of a very low accretion rate, radiatively inefficient accretion flows, and outflows ejecting the in-falling matter (for reviews, see Melia & Falcke 2001; Genzel et al. 2010; Morris et al. 2012). Investigation of the fueling process of Sgr A* may help to understand a large class of dim galactic nuclei and to place these objects into context with brighter supermassive black holes (e.g., Ho 1999; Nagar et al. 2005).

Although Sgr A* is currently quiescent, the detection of the Fermi bubbles and time-variable fluorescent emission from nearby gas clouds suggests that it may have been more active and considerably brighter in the past (e.g., Sunyaev et al. 1993; Koyama et al. 1996; Muno et al. 2007; Su et al. 2010; Terrier et al. 2010; Ponti et al. 2010, 2012). Furthermore, it was recently discovered that an ionized gas cloud of ∼3 M\(_{\odot}\) is approaching the supermassive black hole and is projected to pass as close as ∼2200 R\(_s\) (where R\(_s\) = 2GM/c\(^2\) is the Schwarzschild radius) around 2013 September (Gillessen et al. 2012, 2013). The interaction with the gas cloud may enhance the activity of Sgr A* in the next few years–decades (e.g., Anninos et al. 2012; Burkert et al. 2012; Gillessen et al. 2013; Mościbrodzka et al. 2012; Schartmann et al. 2012).

In the soft X-ray band, Sgr A* emits a steady X-ray luminosity of \( L_X \sim 2 \times 10^{33} \) erg s\(^{-1}\) (2–10 keV; Baganoff et al. 2001, 2003). However, occasionally the X-ray emission flares up by a factor of ∼5−150 for tens of minutes to a few hours. To date, a few dozen such X-ray flares have been detected by Chandra and XMM-Newton (e.g., Baganoff et al. 2001, 2003; Goldwurm et al. 2003; Bélangier et al. 2005; Porquet et al. 2003, 2008; Trap et al. 2011; Nowak et al. 2012). Most observed flares have an intensity of \( L_X \lesssim 10^{35} \) erg s\(^{-1}\) (2–10 keV), but on four occasions bright flares with \( L_X \approx (1 - 5) \times 10^{35} \) erg s\(^{-1}\) have been detected (Baganoff et al. 2001; Porquet et al. 2003, 2008; Nowak et al. 2012). The bright events can be modeled with a power law of photon index \( \Gamma \approx 2\), but the spectral slope of the weaker flares is less well constrained.

The total duration and short-timescale variability observed during the X-ray flares suggests that the emission originates close to the black hole, within ∼10R\(_s\). Therefore, the flares allow investigation of the inner accretion flow and offer a new view of the accretion processes at work in the Galactic nucleus. Different mechanisms have been proposed as the origin of the flares, ranging from magnetic reconnection and electron acceleration models (e.g., Markoff et al. 2001; Yuan et al. 2003; Liu et al. 2004, 2006), to the infall of gas clumps or the disruption of small bodies, such as comets or asteroids (e.g., Čadež et al. 2006; Tagger & Melia 2006).

The proposed flare models usually invoke synchrotron or inverse Compton radiation processes. Apart from addressing
the nature of the engine driving the flares, much effort has therefore gone into trying to understand the underlying emission mechanism, particularly by studying the flares simultaneously at different wavelengths (Eckart et al. 2003, 2006; Yusef-Zadeh et al. 2006, 2008; Hornstein et al. 2007; Aharonian et al. 2008; Marrone et al. 2008; Dodds-Eden et al. 2009; Trap et al. 2010, 2011). Constraining the repetition rate and spectral properties of the flares in the X-ray band is an important aspect of understanding the cause and emission mechanism of the flares (e.g., Porquet et al. 2008; Trap et al. 2010).

Our current knowledge about the X-ray flares from Sgr A* is based on observations with Chandra and XMM-Newton. However, the Swift satellite is also equipped with a sensitive X-Ray Telescope (XRT) that provides sufficient spatial resolution (comparable to that of XMM-Newton) to study the X-ray flares. In this work, we report on the results obtained during a six-year long monitoring campaign of the Galactic center carried out with Swift.

2. DESCRIPTION OF THE SWIFT CAMPAIGN

The Swift observatory (Gehrels et al. 2004) carries an XRT (Burrows et al. 2005) that is sensitive in the 0.5–10 keV range, has a field of view (FOV) of \(\approx 21' \times 21'\), and a spatial resolution of \(\approx 3'' - 5''\) (90% confidence level). The XRT can be operated either in the photon counting (PC) mode, in which a two-dimensional image is acquired, or the windowed timing (WT) mode, in which the CCD columns are collapsed into one dimension to reduce the frame time. During standard operations, the XRT automatically selects the best mode based on the brightness of sources in the field (Hill et al. 2004).

In 2006 February, the Swift satellite embarked on a program to monitor the central \(\approx 21'\times 21'\) of our Galaxy with the onboard XRT (Kennea & The Swift/XRT Team 2006; Degenaar & Wijnands 2009, 2010). This campaign consists of short (\(\approx 1\) ks) pointed X-ray observations that are carried out in the months February–October.7 In 2006–2008, the exposures were taken approximately once every day, whereas the cadence was lowered to one observation every three days from 2009 onward (Table 1).

We use the Swift monitoring campaign of the Galactic center to study the long-term X-ray behavior of Sgr A*. We restrict ourselves to the PC mode data, so that spatial information is available. Between 2006 and 2011, a total of 715 observations were carried out, amounting to a total exposure time of \(\approx 0.8\) Ms (Table 1). Depending on the exact spectral model, the typical sensitivity in a 1 ks exposure is \(L_X \approx 5 \times 10^{33}\) erg s\(^{-1}\) (2–10 keV) for a point source located in the Galactic center region. Single XRT observations usually consist of a number of short data segments with an exposure time of \(\approx 100–500\) s that are separated by one satellite orbit (\(\approx 1.5\) hr).

3. DATA REDUCTION AND ANALYSIS

3.1. XRT Data Reduction

We processed all data of the Galactic center using the Swift tools ver. 3.8 within heasoft ver. 6.11, and using the calibration data released on 2011 August (ver. 3.8). As an initial step, all observations were reprocessed using the tool xrtpipeline. For events of interest, X-ray spectra and light curves were extracted with XSELECT. Ancillary response files (arfs) were created using the task xrtmkarf, taking into account exposure map corrections. The appropriate version of the response matrix files were obtained from the calibration database (CALDB).

The X-ray spectral data were fitted within the XSPEC package (Arnaud 1996). In all spectral fits, we took into account interstellar absorption by using the TBABS model (Wilms et al. 2000).

Throughout this work, we give the absorbed 2–10 keV fluxes (\(F_X\), because these are the least subject to systematic uncertainties in the spectral fitting procedure and allow for a direct comparison with previously reported X-ray flares. Absorption-corrected luminosities (\(L_X\)) were calculated by assuming a distance of \(D = 8\) kpc. Quoted errors for spectral parameters and fluxes refer to 90% confidence intervals (in one interesting parameter). We note that all events detected from Sgr A* had count rates of \(\lesssim 0.2\) counts s\(^{-1}\), and therefore pile-up is not an issue.

3.2. Continuum X-Ray Emission around Sgr A*

To obtain X-ray light curves and spectra, we use a circular region with a radius of 10" centered on the radio position of Sgr A* (Yusef-Zadeh et al. 1999). Chandra observations providing subarcsecond resolution suggest that a 10" extraction region not only contains Sgr A*, but also contributions from (persistent) faint X-ray point sources and a strong diffuse emission component (Baganoff et al. 2003; Porquet et al. 2008; Munuza et al. 2009).

We investigated the continuum X-ray emission detected at the position of Sgr A* in order to have a background reference for studying the X-ray flares (Section 3.4). To aid this, we extracted accumulated spectra for each year of the campaign after removing possible flare events (Section 4). We fitted the spectral data with a simple absorbed power-law model

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7 The Galactic center cannot be observed with Swift between November and January due to Sun-angle constraints.
Figure 1. Long-term XRT light curve of Sgr A∗, binned per gti interval (0.3–10 keV). The solid horizontal line indicates the mean count rate observed in 2006–2011, whereas the dashed line indicates the 3σ level. The six confirmed X-ray flares are numbered and indicated by light gray triangles (Table 2). Flares 5 and 6 were both detected in two subsequent gti intervals.

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(TBABS * PEGPWRLW). We found that the spectral parameters and fluxes inferred for the different years were consistent within the errors. This indicates that the year-averaged sky emission detected by Swift is stable.

A simultaneous fit to the year-averaged spectra with all parameters tied yields $N_H = (9.1 \pm 0.5) \times 10^{22}$ cm$^{-2}$, $\Gamma = 2.5 \pm 0.1$, and a (absorbed) 2–10 keV flux of $F_X = (2.1 \pm 0.1) \times 10^{-12}$ erg cm$^{-2}$ s$^{-1}$. This is consistent with previous results obtained with Chandra when using a similar extraction region (Baganoff et al. 2003). The quiescent emission of Sgr A∗ contributes only $\approx 10\%$ to the total flux observed within our XRT extraction region (Baganoff et al. 2001, 2003).

3.3. Searching for X-Ray Flares from Sgr A∗

To search the Swift data for X-ray flares, we created a long-term XRT light curve of Sgr A∗. We extract the count rate for each single XRT good time interval (gti; generally, these represent different satellite orbits, see Section 2). Occasionally, data segments of limited exposure were obtained, e.g., when an observation was interrupted by a trigger of Swift’s Burst Alert Telescope, or due to an automated switch to the WT mode triggered by activity of one of the bright X-ray transients that are located in the FOV (e.g., Degenaar & Wijnands 2010; Degenaar et al. 2012). Since very short exposures do not generate reliable results, we chose to reject intervals with an exposure time of <100 s. After applying this selection, we were left with a total of 1349 gts that were searched for occurrences of X-ray flares. The resulting long-term Swift/XRT light curve is shown in Figure 1.

We searched the XRT light curve for occurrences of events that had an intensity of >3σ above the mean and marked these as potential X-ray flares. We subsequently manually inspected the images and light curves of these observations to investigate whether these could indeed contain X-ray flares from Sgr A∗ (see Figures 2 and 3 for examples). We analyzed the X-ray spectra of all events that were positively identified as X-ray flares and yielded sufficient source counts.

3.4. Spectral Analysis of the X-Ray Flares

The short XRT exposures typically provided a few tens of photons per event. Therefore, we chose to fit the unbinned X-ray spectra without background subtraction using Cash statistics (Cash 1979). We modeled the flares as a simple power law (PEGPWRLW), which was included as an additional component to the underlying continuum emission. All the parameters for the continuum were fixed to the values obtained from simultaneously fitting the year-averaged spectra (Section 3.2), so that the index and normalization (i.e., the flux) of the flare spectrum were the only free fit parameters. The additional power law describing the flares is thus absorbed with $N_H = 9.1 \times 10^{22}$ cm$^{-2}$ fixed.

To allow for a model-independent test of the overall spectral shape of the X-ray flares, we also determined their hardness ratio (HR). For the present purpose, we defined HR as the ratio of counts in the 4–10 and 2–4 keV energy bands.

4. RESULTS

The mean XRT count rate of Sgr A∗ inferred from the 2006–2011 data set is $1.1 \times 10^{-2}$ counts s$^{-1}$, with a standard
deviation of $6.7 \times 10^{-3}$ counts s$^{-1}$. The long-term light curve contained six events that had an intensity of $>4\sigma$ above the mean, and collected enough photons to allow for further study. Investigation of the images, light curves, and spectra suggests that all six can be identified as X-ray flares from Sgr A$^*$. These are numbered and indicated by light gray triangles in Figure 1. The basic properties of the six flares are summarized in Table 2, while the results of the spectral analysis are presented in Table 3. One of the flares (flare 6) is briefly discussed in Section 4.1. The observational setup (i.e., short exposures of $\lesssim 1$ ks) that had an intensity of a factor of $3–4$ above the mean intensity ($\lesssim 10^{-3}$ counts s$^{-1}$), i.e., a factor of $\lesssim 10$ ($\lesssim 13.5\sigma$) above the continuum count rate. In the second orbit ($\gtrsim 0.6$ ks) that occurred $\lesssim 1.4$ hr after the first, the count rate from Sgr A$^*$ was still $\gtrsim 2.5\sigma$ above the mean intensity ($\gtrsim 2.8 \times 10^{-2}$ counts s$^{-1}$). Comparing the XRT image of the second orbit with observations of similar exposure (i.e., $\lesssim 0.6$ ks) reveals an excess of photons that suggests that the flare was still detected in this gti. The intensity averaged over the full observation (i.e., 2 orbits) is $\gtrsim 6.1 \times 10^{-2}$ counts s$^{-1}$ and a total of 65 counts were obtained. The average X-ray spectrum is plotted in Figure 4, along with the spectral fit. The first and the second orbit collected 48 and 17 photons, respectively, which allow to fit the two spectra separately. The intensity inferred from the second orbit was a factor $\gtrsim 3$ lower than during the first. This suggests that the second orbit likely caught the tail of the X-ray flare, whereas the first may have detected it closer to the peak. Spectral analysis suggests that both intervals had similar spectra, which is supported by the inferred HRs (Table 3). There is thus no indication of spectral evolution during the flare, although the errors bars are large.

### 4.1. Flare 6: 2010 June 12

The sixth X-ray flare was detected in an observation that had a total exposure time of $\gtrsim 1.1$ ks and consisted of two orbits. The light curve of this observation is shown in Figure 2. Figure 3 compares the XRT image of the 1.1 ks flare observation with that of the proceeding 1.0 ks pointing, which was taken three days earlier.

Enhanced activity is clearly detected in the first orbit (exposure time of $\gtrsim 0.5$ ks), during which the intensity of Sgr A$^*$ is $\gtrsim 10 \times 10^{-1}$ counts s$^{-1}$, i.e., a factor of $\gtrsim 10$ ($\gtrsim 13.5\sigma$) above the continuum count rate. In the second orbit ($\gtrsim 0.6$ ks) that occurred $\lesssim 1.4$ hr after the first, the count rate from Sgr A$^*$ was still $\gtrsim 2.5\sigma$ above the mean intensity ($\gtrsim 2.8 \times 10^{-2}$ counts s$^{-1}$). Comparing the XRT image of the second orbit with observations of similar exposure (i.e., $\gtrsim 0.6$ ks) reveals an excess of photons that suggests that the flare was still detected in this gti. The intensity averaged over the full observation (i.e., 2 orbits) is $\gtrsim 6.1 \times 10^{-2}$ counts s$^{-1}$ and a total of 65 counts were obtained. The average X-ray spectrum is plotted in Figure 4, along with the spectral fit. The first and the second orbit collected 48 and 17 photons, respectively, which allow to fit the two spectra separately. The intensity inferred from the second orbit was a factor $\gtrsim 3$ lower than during the first. This suggests that the second orbit likely caught the tail of the X-ray flare, whereas the first may have detected it closer to the peak. Spectral analysis suggests that both intervals had similar spectra, which is supported by the inferred HRs (Table 3). There is thus no indication of spectral evolution during the flare, although the errors bars are large.

### 4.2. The Source of the Swift X-Ray Flares

We determined the XRT position of the six confirmed X-ray flares utilizing the point-spread fitting method of the algorithm.
Notes. The spectral data of the X-ray flares were fitted using an absorbed power-law model with \( N_H = 9.1 \times 10^{22} \text{ cm}^{-2} \) fixed, and applying Cash statistics. The hardness ratio (HR) is defined as the ratio of counts in the 4–10 and 2–4 keV bands. \( F_X \) gives the observed (i.e., not corrected for absorption) flux in the 2–10 keV band. \( L_X \) represents the average 2–10 keV luminosity corrected for absorption and assuming \( D = 8 \) kpc. Flares 5 and 6 were both detected in two subsequent orbits. We give the spectral and color information both for the separate orbits (denoted by “a” and “b”), as well as the values averaged over the two intervals.

(A color version of this figure is available in the online journal.)

Figure 4. XRT spectrum of the flare detected from Sgr A* on 2010 June 12 (6; black bullets), compared to the summed spectra of flares 1–5 (red squares). For representation purposes, the spectral data were rebinned to contain 10 photons bin\(^{-1}\). The solid lines indicate fits to a combination of two absorbed power laws, with all parameters for the one representing the continuum emission fixed (see Section 3.4).

(A color version of this figure is available in the online journal.)

devolved by Evans et al. (2009). The coordinates of the individual flares are all consistent within their 90% confidence errors of \( \pm 3\!.6–3\!.9 \). The Swift positions are offset by \( \pm 1\!.9–5\!.6 \) from the radio position of Sgr A* (Yusef-Zadeh et al. 1999).

Stacking the data of the six flares to improve the statistics yields R.A. = 17:45:40.13, decl. = -29:00:27.0 (J2000), with a 90% confidence uncertainty of 3\!.5. This localization is within \( \pm 1\!.6 \) from the coordinates of the radio counterpart of Sgr A* and fully consistent within the errors.

Given the spatial resolution of the XRT, we cannot formally exclude that the six X-ray flares were caused by another X-ray source located within our 10\( \prime \)\! extraction region. However, long-term Chandra monitoring has shown that all the faint persistent X-ray sources located within this region display a stable X-ray luminosity; none of them is known to exhibit X-ray variation of the observed magnitude (Muno et al. 2009).

Objects such as accreting white dwarfs or X-ray binaries may display transient X-ray events, but these typically last much longer (days–months) than the short (hours) events that we detect. There is one known transient X-ray binary located within 10\( \prime \) of Sgr A* (Porquet et al. 2005b; Muno et al. 2005a, 2005b). However, its position is 4\!.2 away from the XRT position, outside the 90% confidence error, inferred for the six flares (Muno et al. 2005a).

The above considerations suggest that it is unlikely that another X-ray source was responsible for causing the observed flares. The spectral parameters and intensities observed with the XRT are very similar to the X-ray flares seen with XMM-Newton and Chandra (e.g., Baganoff et al. 2001; Goldwurm et al. 2003; Bélander et al. 2005; Porquet et al. 2003, 2008; Nowak et al. 2012). This further strengthens the association with Sgr A*. We therefore conclude that we have detected six new bright X-ray flares from Sgr A* with Swift.

4.3. Different Types of X-Ray Flares?

The HRs of the flares suggest that flare 6 may have been spectrally softer than the other five (Figure 5). To investigate whether the spectral differences could be biased by statistics, we co-added the spectra and weighted response files of flares 1–5. Combining these yields 122 counts and a net count rate of \( \pm 4.5 \times 10^{-2} \) counts s\(^{-1}\). We obtain a photon index of \( \Gamma = 2.0 \pm 0.6 \) and HR = 1.7 \pm 0.2. This is indeed harder than flare 6 (\( \Gamma = 3.0 \pm 0.8 \) and HR = 0.7 \pm 0.2), while their 2–10 keV absorbed fluxes are similar (Table 2).

Figure 4 compares the averaged spectrum of flares 1–5 to that of flare 6. This shows that the main difference between the spectra occurs at energies \( \lesssim 3 \) keV. This suggests that possibly the absorption or dust scattering along our line of sight was different during flare 6. Comparing the continuum spectra at different epochs does not indicate that the background emission was different at that time. We also fitted the spectrum of flare 6 while subtracting the spectrum of the preceding observation (ObsID 90416020) as a background. This yielded an index of \( \Gamma = 2.6 \pm 1.0 \) (with \( N_H = (9.1 \pm 0.5) \times 10^{22} \text{ cm}^{-2} \) fixed; Section 3.2). While this remains softer than the average value

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Table 3: Results from Spectral Analysis of the X-Ray Flares Detected with Swift/XRT

<table>
<thead>
<tr>
<th>No.</th>
<th>HR</th>
<th>( \Gamma )</th>
<th>( \text{Cstat} ) (dof)</th>
<th>( F_X ) (erg cm(^{-2}) s(^{-1}))</th>
<th>( L_X ) (erg s(^{-1}))</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>2.1 \pm 1.0</td>
<td>1.7 \pm 1.5</td>
<td>18.89 (19)</td>
<td>(1.0 \pm 0.6) \times 10^{-11}</td>
<td>(1.4 \pm 0.5) \times 10^{35}</td>
</tr>
<tr>
<td>2</td>
<td>1.3 \pm 0.5</td>
<td>1.9 \pm 1.5</td>
<td>35.03 (26)</td>
<td>(0.7 \pm 0.4) \times 10^{-11}</td>
<td>(0.9 \pm 0.3) \times 10^{35}</td>
</tr>
<tr>
<td>3</td>
<td>2.5 \pm 1.2</td>
<td>3.1 \pm 2.4</td>
<td>21.62 (18)</td>
<td>(1.5 \pm 0.9) \times 10^{-11}</td>
<td>(2.9 \pm 1.2) \times 10^{35}</td>
</tr>
<tr>
<td>4</td>
<td>1.9 \pm 0.9</td>
<td>1.8 \pm 1.8</td>
<td>12.11 (16)</td>
<td>(1.2 \pm 0.6) \times 10^{-11}</td>
<td>(1.5 \pm 0.9) \times 10^{35}</td>
</tr>
<tr>
<td>5</td>
<td>1.9 \pm 0.7</td>
<td>1.9 \pm 1.2</td>
<td>29.23 (31)</td>
<td>(1.3 \pm 0.5) \times 10^{-11}</td>
<td>(1.8 \pm 0.7) \times 10^{35}</td>
</tr>
<tr>
<td>5a</td>
<td>2.8 \pm 1.5</td>
<td>1.6 \pm 1.6</td>
<td>26.28 (16)</td>
<td>(1.5 \pm 1.1) \times 10^{-11}</td>
<td>(2.0 \pm 0.9) \times 10^{35}</td>
</tr>
<tr>
<td>5b</td>
<td>1.3 \pm 0.6</td>
<td>2.7 \pm 2.5</td>
<td>14.06 (15)</td>
<td>(1.0 \pm 0.5) \times 10^{-11}</td>
<td>(1.8 \pm 0.8) \times 10^{35}</td>
</tr>
<tr>
<td>1–5</td>
<td>1.7 \pm 0.2</td>
<td>2.0 \pm 0.6</td>
<td>97.73 (105)</td>
<td>(1.1 \pm 0.3) \times 10^{-11}</td>
<td>(1.5 \pm 0.3) \times 10^{35}</td>
</tr>
<tr>
<td>6</td>
<td>0.7 \pm 0.2</td>
<td>3.0 \pm 0.8</td>
<td>50.52 (59)</td>
<td>(0.9 \pm 0.2) \times 10^{-11}</td>
<td>(1.8 \pm 0.3) \times 10^{35}</td>
</tr>
<tr>
<td>6a</td>
<td>0.8 \pm 0.2</td>
<td>3.0 \pm 0.9</td>
<td>48.03 (43)</td>
<td>(1.6 \pm 0.4) \times 10^{-11}</td>
<td>(3.1 \pm 0.5) \times 10^{35}</td>
</tr>
<tr>
<td>6b</td>
<td>0.5 \pm 0.4</td>
<td>3.2 \pm 2.6</td>
<td>11.23 (13)</td>
<td>(3.4 \pm 2.0) \times 10^{-12}</td>
<td>(0.7 \pm 0.3) \times 10^{35}</td>
</tr>
</tbody>
</table>
inferred for flares 1–5, we cannot exclude that the apparent differences are caused by changes in the background emission.

Flare 6 stands out in the long-term XRT light curve as the brightest event (Figure 1), although it was detected in two orbits and its averaged intensity was similar to that of the other five flares (Table 2). Spectral differences may possibly arise if we caught flare 6 near its peak, whereas the other five events trace the rising or decaying parts of the flares. However, we did not find signs of spectral evolution during flare 6 (detailed analysis of other bright X-ray flares detected with Chandra and XMM-Newton did not reveal evidence for spectral evolution either; Porquet et al. 2008; Nowak et al. 2012). Therefore, we do not consider it likely that the different spectral properties of the Swift flares are the result of tracing different parts of the X-ray flares.

The position of flare 6 (R.A. = 17:45:40.27, decl. = −29:00:24.1 with a 90% confidence error of 3′) is ≃5″ offset from the radio position of Sgr A*, but consistent with the location of the other five Swift flares. We do not consider it likely that flare 6 was caused by a different source, although we cannot formally exclude this possibility (Section 4.2).

5. DISCUSSION

We have used six years of Swift/XRT monitoring data of the Galactic center to search for X-ray flares from the supermassive black hole. Within a 10″ extraction region centered on Sgr A*, we detect a continuum emission level of ≃1.1 × 10^{-2} counts s^{-1}. The inferred corresponding 2–10 keV luminosity is $L_X \simeq 2 \times 10^{35}$ erg s^{-1}. This is a factor of ≃10 above the quiescent X-ray emission of the supermassive black hole and is dominated by contributions from diffuse emission and faint X-ray point sources.

We searched the long-term Swift light curve of Sgr A* for occurrences during which the XRT intensity increased by a factor >3σ above the mean (i.e., above a count rate of ≃3 × 10^{-2} counts s^{-1}). This chosen threshold makes us sensitive to X-ray flares that reach a 2–10 keV intensity of $L_X \gtrsim 7 \times 10^{34}$ erg s^{-1} (i.e., a factor of ≳35 above the quiescent level). This is at the high end of the luminosity distribution of those observed with Chandra and XMM-Newton, implying that we are sensitive only to the brightest X-ray flares from Sgr A* (e.g., Baganoff et al. 2001, 2003; Bélanger et al. 2005; Porquet et al. 2003, 2008; Nowak et al. 2012).

Investigation of the images, light curves, and spectra allowed us to identify six X-ray flares from Sgr A*. They have inferred 2–10 keV mean luminosities of $L_X \simeq (1 − 3) \times 10^{35}$ erg s^{-1} (Table 2). Only four such similarly bright events have been reported before: two were seen with Chandra (Baganoff et al. 2001; Nowak et al. 2012) and two with XMM-Newton (Porquet et al. 2003, 2008). Using ≲0.8 Ms of data, the Swift monitoring campaign has thus more than doubled the number of bright (i.e., $L_X > 10^{35}$ erg s^{-1}) X-ray flares detected from Sgr A*, setting the current counter to 10.

The multitude of observations provided by the Swift campaign clearly has the advantage of increasing the detection likelihood of X-ray flares. This is also illustrated by the fact that we detected five X-ray flares between 2006 and 2008 when the Galactic center was observed every day, while only one flare was detected in 2009–2011 when the cadence was lowered to one pointing every three days (see Section 5.2).

5.1. Properties of the X-Ray Flares Detected with Swift

Prior studies with Chandra and XMM-Newton suggested that the brightest ($L_X \gtrsim 10^{35}$ erg s^{-1}) X-ray flares had different spectral slopes (cf. Baganoff et al. 2001; Porquet et al. 2003, 2008). However, it was recently established that these differences were likely due to instrumental effects (primarily pile-up) and spectral modeling (Nowak et al. 2012).

The Swift campaign has detected multiple bright X-ray flares with the same instrument. This allows for a comparative study of their properties unbiased by instrumental effects, which are induced when comparing flares detected with different instruments (Porquet et al. 2008; Nowak et al. 2012). The photon index inferred from Swift flares 1–5 ($Γ = 2.0 \pm 0.6$) is comparable to that of the bright flares detected by Chandra and XMM-Newton ($Γ = 2.0^{+0.7}_{−0.6}, 2.3 ± 0.3$, and 2.4^{+0.3}_{−0.3}; Nowak et al. 2012). However, Swift flare 6 was softer ($Γ = 3.0 ± 0.8$).

Although we cannot exclude that this event was caused by a different source or that changes in the background emission play a role (Section 4.3), the Swift results may indicate that flares of similar intensity can have different spectral shapes. This could possibly be explained in terms of different emission mechanisms (e.g., synchrotron versus inverse Compton radiation), different emitting regions (e.g., local versus global events), or different Lorentz factors. One model that can account for different spectral properties is the stochastic acceleration formalism proposed by Liu et al. (2004). Within this framework, harder flares can be produced by local magnetohydrodynamic processes, such as magnetic reconnection, whereas softer flares are expected to result from global fluctuations.

5.2. Constraints on the Flaring Repetition

We can use the results of the Swift campaign to estimate the occurrence rate of bright X-ray flares from Sgr A*, i.e., those that have a luminosity of $L_X \gtrsim 10^{35}$ erg s^{-1} (2–10 keV). Between 2006 February 24 and 2011 October 25, the Galactic center was observed with the XRT on 715 days. A total of six X-ray flares were detected during those observations. We cannot put firm constraints on the length of the flares from the Swift campaign.
data, but we will assume that the duration of all events was the same as similarly bright flares observed with \textit{Chandra} and XMM-\textit{Newton}, i.e., $\approx 1$–3 hr (Baganoff et al. 2001; Porquet et al. 2003, 2008; Bélangé et al. 2005; Nowak et al. 2012).

Making the simplest assumption that the occurrence is uniform, we can estimate the flaring rate by dividing the number of events ($N$) by the number of days that were observed ($N_{\text{obs}} = 715$), and multiplying this with the time in a day ($t_d = 24$ hr) divided by the duration of a flare ($t_f = 1$–3 hr). Thus, we calculate the number of flares per day as \(N/N_{\text{obs}} \times t_d/t_f\). For a total of six X-ray flares the flaring rate is then $\approx 0.1$–0.2 day$^{-1}$. This implies that a bright flare with an intensity of $L_X \gtrsim 10^{34}$ erg s$^{-1}$ occurs approximately every 5–10 days. We note that this flaring rate is consistent with the fact that several more flares were detected with Swift in 2006–2008 than in 2009–2011 (Figure 1). This difference is simply explained by the larger number of observations (i.e., the higher sampling rate) during the first three years of the campaign (Table 1).

The flaring rate that we obtain from the \textit{Swift} data is in line with previous estimates. Based on 0.6 Ms of \textit{Chandra} data, Baganoff (2003) infers an occurrence rate of $\approx 0.3$–0.9 flares day$^{-1}$ for events that increase a factor $>10$ above the quiescent luminosity of Sgr A*. \textit{Chandra} has also exposed several faint X-ray flares that do not become brighter than $L_X \lesssim 10^{34}$ erg s$^{-1}$, and are thought to occur more frequently (e.g., Baganoff et al. 2003; Eckart et al. 2004, 2006, 2008). However, such weak flares are beyond the reach of \textit{Swift}.

### 5.3. Outlook

It was recently discovered that an ionized gas cloud with an estimated mass of $\lesssim 3 \, M_\odot$ is rapidly moving toward Sgr A*, and is expected to make a close passage around 2013 September (Gillessen et al. 2012, 2013). The cloud is expected to become disrupted due to its gravitational interaction with the supermassive black hole. Parts of the shredded gas may become accreted, enhancing the persistent emission of Sgr A* for the next few years–decades (e.g., Anninos et al. 2012; Burkert et al. 2012; Gillessen et al. 2013; Mościbrodzka et al. 2012; Schartmann et al. 2012).

Possibly, the interaction of the gas cloud with the shock at the Bondi radius can generate instabilities that produce more powerful and more frequent X-ray flares of Sgr A*. Mapping out the occurrence rate and X-ray properties of the flares detected with \textit{Swift} in 2006–2011 provides a calibration point that allows us to study whether the flaring rate of the supermassive black hole will be affected as a result of the approaching gas cloud.

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