Search for displaced muonic lepton jets from light Higgs boson decay in proton-proton collisions at $\sqrt{s} = 7$ TeV with the ATLAS detector


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Search for displaced muonic lepton jets from light Higgs boson decay in proton–proton collisions at $\sqrt{s} = 7$ TeV with the ATLAS detector

ATLAS Collaboration

1. Introduction

A search is presented for long-lived neutral particles decaying to final states containing collimated muon pairs in proton–proton collisions at $\sqrt{s} = 7$ TeV centre-of-mass energy. The event sample, collected during 2011 at the LHC with the ATLAS detector, corresponds to an integrated luminosity of 1.9 fb$^{-1}$. The model considered in this analysis consists of a Higgs boson decaying to a new hidden sector of particles which finally produce two sets of collimated muon pairs, but the search described is equally valid for other, distinct models such as heavier Higgs boson doublets or singlet scalars, produced through gluon fusion, that decay to a hidden sector and eventually produce collimated muon pairs.

Recently, evidence for the production of a boson with a mass of about 126 GeV has been published by ATLAS [1] and CMS [2]. The observation is compatible with the expected production and decay of the Standard Model (SM) Higgs boson at this mass. Testing the SM Higgs hypothesis is currently of utmost importance. To this end two effects may be considered: (i) additional resonances which arise in an extended Higgs sector found in many extensions of the SM, or (ii) rare Higgs boson decays which may deviate from those predicted by the SM. In this Letter we search for a scalar, produced through gluon fusion, that decays to a light hidden sector, according to the topology of Fig. 1, focusing on the 100 GeV to 140 GeV mass range.

The phenomenology of light hidden sectors has been studied extensively over the past few years [6–10]. Possible characteristic topological signatures of such extensions of the SM are “lepton jets”. A lepton jet is a cluster of highly collimated particles: electrons, muons and possibly pions [7,11–13]. These arise if light unstable particles with masses in the MeV to GeV range (for example dark photons, $\gamma_d$) reside in the hidden sector and decay predominantly to SM particles. At the LHC, hidden-sector particles may be produced with large boosts, causing the visible decay products to form jet-like structures. Hidden-sector particles such as $\gamma_d$ may be long-lived, resulting in decay lengths comparable to, or larger than, the detector dimensions. The production of lepton jets can occur through various channels. For instance, in supersymmetric models, the lightest visible superpartner may decay into the hidden sector. Alternatively, a scalar particle that couples to the visible sector may also couple to the hidden sector through Yukawa couplings or the scalar potential. This analysis is focused on the case where the Higgs boson decays to the hidden sector [14,15]. The SM Higgs boson has a narrow width into SM final states if $m_H < 2m_W$. Consequently, any new (non-SM) coupling to additional states, which reside in a hidden sector, may contribute significantly to the Higgs boson decay branching ratios. Even with new couplings, the total Higgs boson width is typically small, well below the order of one GeV. If a SM-like Higgs boson is confirmed, it will remain important to constrain possible rare decays, e.g. into lepton jets.

Neutral particles with large decay lengths and collimated final states represent, from an experimental point of view, a challenge both for the trigger and for the reconstruction capabilities of the detector. Collimated particles in the final state can be hard to disentangle due to the finite granularity of the detectors; moreover, in the absence of inner tracking detector information and a primary vertex constraint, it is difficult to reconstruct charged-particle tracks from decay vertices far from the interaction point (IP). The ATLAS detector [16] is equipped with a muon spectrometer (MS) with high-granularity tracking detectors that allow charged-particle tracks to be reconstructed in a standalone configuration using only the muon detector information (MS-only). This is a crucial feature for detecting muons not originating from the primary interaction vertex.
Fig. 1. Schematic picture of the Higgs boson decay chain, $H \rightarrow 2f_2 \rightarrow f_1 \gamma_d$. The Higgs boson decays to two hidden fermions ($f_2$). Each hidden fermion decays to a $\gamma_d$ and to a stable hidden fermion ($f_1$), resulting in two muon jets from the $\gamma_d$ decays in the final state.

The search presented in this Letter focuses on neutral particles decaying to the simplest type of muon jets (MJs), containing only two muons; prompt MJ searches have been performed both at the Tevatron [17,18] and at the LHC [19]. Other searches for displaced decays of a light Higgs boson to heavy fermion pairs have also been performed at the LHC [20].

The benchmark model used for this analysis is a simplified scenario where the Higgs boson decays to a pair of neutral hidden fermions ($f_2$) each of which decays to one long-lived $\gamma_d$ and one stable neutral hidden fermion ($f_1$) that escapes the detector un-noticed, resulting in two lepton jets from the $\gamma_d$ decays in the final state (see Fig. 1). The mass of the $\gamma_d$ (0.4 GeV) is chosen to provide a sizeable branching ratio to muons [14].

2. The ATLAS detector

ATLAS is a multi-purpose detector [16] at the LHC, consisting of an inner tracking system (ID) embedded in a superconducting solenoid, which provides a 2 T magnetic field parallel to the beam direction, electromagnetic and hadronic calorimeters and a muon spectrometer using three air-core toroidal magnet systems.1 The trigger system has three levels [21] called Level-1 (L1), Level-2 (L2) and Event Filter (EF). L1 is a hardware-based system using information from the calorimeter and muon spectrometer, and defines one or more Regions of Interest (ROIs), geometrical regions of the detector, identified by ($\eta, \phi$) coordinates, containing interesting physics objects. L2 and the EF (globally called the High Level Trigger, HLT) are software-based systems and can access information from all sub-detectors. The ID, consisting of silicon pixel and micro-strip detectors and a straw-tube tracker, provides precision tracking of charged particles for $|\eta| \leq 2.5$. The electromagnetic and hadronic calorimeter system covers $|\eta| \leq 4.9$ and, at $\eta = 0$, has a total depth of 9.7 interaction lengths (22 radiations lengths in the electromagnetic part). The MS provides trigger information ($|\eta| \leq 2.4$) and momentum measurements ($|\eta| \leq 2.7$) for charged particles entering the spectrometer. It consists of one barrel and two endcap parts, each with 16 sectors in $\phi$, equipped with precision tracking chambers and fast detectors for triggering. Monitored drift tubes are used for precision tracking in the region $|\eta| \leq 2.0$ and cathode strip chambers are used for $2.0 \leq |\eta| \leq 2.7$. The MS detectors are arranged in three stations of increasing distance from the IP: inner, middle and outer. The air core toroidal magnetic field allows an accurate charged particle reconstruction independent of the ID information. The three planes of trigger chambers (resistive plate chambers in the barrel and thin gap chambers in the endcaps) are located in middle and outer (only in the barrel) stations.

The L1 muon trigger requires hits in the middle stations to create a low transverse momentum ($p_T$) muon ROI or hits in both the middle and outer stations for a high $p_T$ ROI. The muon ROIs have a spatial extent of $0.2 \times 0.2$ ($\Delta \eta \times \Delta \phi$) in the barrel and of $0.1 \times 0.1$ in the endcap. L1 ROI information seeds, at HLT level, the reconstruction of muon momenta using the precision chamber information. In this way sharp trigger thresholds up to 40 GeV can be obtained.

3. Signal and background simulation

The set of parameters used to generate the signal Monte Carlo samples is listed in Table 1. The Higgs boson is generated through the gluon–gluon fusion production mechanism which is the dominant process for a low mass Higgs boson. The gluon–gluon fusion Higgs boson production cross section in $pp$ collisions at $\sqrt{s} = 7$ TeV, estimated at the next-to-next-to-leading order (NNLO) [22], is $\sigma_{SM} = 24.0$ pb for $m_H = 100$ GeV and $\sigma_{SM} = 12.1$ pb for $m_H = 140$ GeV. The PYTHIA generator [23] is used, linked together with MadGraph4.4.2 [24] and BRIDGE [25], for gluon–gluon fusion production of the Higgs boson and the subsequent decay to hidden-sector particles.

As discussed in the introduction, the signal is chosen to enable a study of rare, non-SM, Higgs boson decays in the (possibly extended) Higgs sector. To do so we choose two points which envelope a mass range covering the 126 GeV resonance. The lower mass point, $m_H = 100$ GeV, is chosen to be compatible with the decay-mode-independent search by OPAL at LEP [26]. The higher mass point, $m_H = 140$ GeV, is chosen well below the $WW$ threshold, where a sizeable branching ratio into a hidden sector may be naturally achieved. The masses of $f_2$ and $f_3$ are chosen to be light relative to the Higgs boson mass, and far from the kinematic threshold at $m_{f_2} + m_{\gamma_d} = m_{f_3}$. For the chosen dark photon mass (0.4 GeV), the $\gamma_d$ decay branching ratios are expected to be [14]: 45% $e^+e^-$, 45% $\mu^+\mu^-$, 10% $\pi^+\pi^-$. Thus 20% of the Higgs $H \rightarrow \gamma_d\gamma_d + X$ decays are expected to have the required four-muon final state.

The mean lifetime $\tau$ of the $\gamma_d$ (expressed throughout this Letter as $\tau$ times the speed of light $c$) is a free parameter of the model. In the generated samples $\tau$ is chosen so that a large fraction of the decays occur inside the sensitive ATLAS detector volume, i.e. up to 7 m in radius and 13 m along the $z$-axis, where the trigger chambers of the middle stations are located. The detection efficiency can then be estimated for a range of $\gamma_d$ mean lifetimes through re-weighting of the generated samples.

Potential backgrounds include all the processes which lead to real prompt muons in the final state such as the SM processes $W +$ jets, $Z +$ jets, $t\bar{t}$, WW, WZ, and ZZ. However, the main contribution to the background is expected from processes giving a high production rate of secondary muons which do not point to the primary vertex, such as decays in flight of $K/\pi$ and heavy flavour decays in multi-jet processes, or muons due to cosmic rays. The prompt lepton background samples are generated using PYTHIA ($W +$ jets, and $Z +$ jets) and MC@NLO [27] ($t\bar{t}$, WW, WZ, and ZZ).

<table>
<thead>
<tr>
<th>Higgs mass [GeV]</th>
<th>$m_{f_3}$ [GeV]</th>
<th>$m_{f_1}$ [GeV]</th>
<th>$\gamma_d$ mass [GeV]</th>
<th>$\tau$ [mm]</th>
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</thead>
<tbody>
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<td>100</td>
<td>5.0</td>
<td>2.0</td>
<td>0.4</td>
<td>47</td>
</tr>
<tr>
<td>140</td>
<td>5.0</td>
<td>2.0</td>
<td>0.4</td>
<td>36</td>
</tr>
</tbody>
</table>

1 ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the $z$-axis coinciding with the beam pipe axis. The $x$-axis points from the IP to the centre of the LHC ring, and the $y$-axis points upward. Cylindrical coordinates ($r, \phi$) are used in the transverse plane, $\phi$ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle $\theta$ as $\eta = -\ln \tan(\theta/2)$. 
The generated Monte Carlo events are processed through the full ATLAS simulation chain based on GEANT4 [28, 29]. Additional pp interactions in the same and nearby bunch crossings (pile-up) are included in the simulation. All Monte Carlo samples are re-weighted to reproduce the observed distribution of the number of interactions per bunch crossing in the data. For the multi-jet background evaluation a data-driven method is used. The cosmic-ray background is also evaluated from data.

4. The kinematics of the signal

The main kinematic characteristics of the signal sample are:

- The $\gamma\gamma$ pair are emitted approximately back-to-back in $\phi$, with an angular spread of the distribution due to the emission of the $f_{\gamma\gamma}$.
- The average $p_T$ of the $\gamma\gamma$ in the laboratory frame is about 20 GeV for $m_H = 100$ GeV and 30 GeV for $m_H = 140$ GeV; due to the small mass of the $\gamma\gamma$ large boost factors in the decay length should be expected.
- Fig. 2 shows the distribution of $\Delta R = \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2}$ between the two muons from the $\gamma\gamma$ decay. The $\Delta R$ is computed at the decay vertex of the $\gamma\gamma$ from the vector momenta of the two muons. Due to the small mass of the $\gamma\gamma$ the $\Delta R$ is almost always below 0.1.

Since the two $f_{\gamma\gamma}$ are, like the two $\gamma\gamma$, emitted back-to-back in $\phi$, the observed missing transverse momentum $E_T^{miss}$, computed at the event-generator level, is small and cannot be used as a discriminating variable against the background.

5. Data samples and trigger selection

The dataset used for this analysis was collected at a centre-of-mass energy of 7 TeV during the first part of 2011, where a low level of pile-up events in the same bunch-crossing was present (an average of $\approx 6$ interactions per crossing). Only periods in which all ATLAS subdetectors were operational are used. The total integrated luminosity used is $1.94 \pm 0.07$ fb$^{-1}$ [30, 31]. All events are required to have at least one reconstructed vertex along the beam line with at least three associated tracks, each with $p_T \geq 0.4$ GeV. The primary interaction vertex is defined to be the vertex whose constituent tracks have the largest $\sum p_T^2$. This analysis deals with displaced $\gamma\gamma$ decays with final states containing only muons. Signal events are therefore characterized by a four-muon final state with the four muons coming from two displaced decay vertices. Due to the relatively low $p_T$ of the muons and due to the displaced decay vertex, a low-$p_T$ multi-muon trigger with muons reconstructed only in the MS is needed. In order to have an acceptably low trigger rate at a low $p_T$ threshold, a multiplicity of at least three muons is required. Candidate events are collected using an unprescaled HLT trigger with three reconstructed muons of $p_T \geq 6$ GeV, seeded by a L1-accept with three different muon ROIs. These muons are reconstructed only in the MS, since muons originating from a neutral particle decaying outside the pixel detector will not have a matching track in the ID tracking system. The trigger efficiency for the Monte Carlo signal samples, defined as the fraction of events passing the trigger requirement with respect to the events satisfying the analysis selection criteria (described in Section 6) is 0.32 $\pm$ 0.01 (for $m_H = 100$ GeV) and 0.31 $\pm$ 0.01 (for $m_H = 140$ GeV).

6. Muon jet reconstruction and event selection

MJs from displaced $\gamma\gamma$ decays are characterized by a pair of muons in a narrow cone, produced away from the primary vertex of the event. Consequently tracks reconstructed in the MS with a good quality track fit [32] are used. MJs are identified using a simple clustering algorithm that associates all the muons in cones of at least $0.2$, starting with the muon with highest $p_T$. The size of the cone takes into account the multiple scattering of the muons in the calorimeters. All the muons found in the cone are associated with a MJ. After this procedure, if any muons are unassociated with a MJ the search is repeated for this remainder, starting again with the highest $p_T$ muon. This continues until all possible MJs
are formed. The MJ direction and momentum are obtained from the vector sum over all muons in the MJ. Only MJs with two reconstructed muons are accepted and only events with two MJs are kept for the subsequent analysis. In order to keep the search as model independent as possible no requirement on the muon momenta has been introduced.

The possible contribution to the background of SM processes which lead to real prompt muon pairs in the final state is evaluated using simulated samples. After the trigger and the requirement of having two MJs in the event, their contributions have been found to be negligible. The only significant background sources are expected to be from processes giving a high production rate of secondary muons which do not point to the primary vertex, such as decays in flight of $K/\pi$ and heavy flavour decays in multi-jet production, or cosmic-ray muons not pointing to the primary vertex.

In order to separate the signal from the background, a number of discriminating variables have been studied. The multi-jet background can be significantly reduced by using calorimeter isolation requirements around the MJ direction. The calorimeter isolation variable $E_{\text{isol}}^{\text{ID}}$ is defined as the difference between the transverse calorimetric energy $E_T$ in a cone of $\Delta R = 0.4$ around the highest $p_T$ muon of the MJ and the $E_T$ in a cone of $\Delta R = 0.2$; a cut $E_{\text{isol}}^{\text{ID}} \leq 5$ GeV keeps almost all the signal. The isolation modelling is validated for real isolated muons with a sample of muons coming from $Z \to \mu\mu$ decays. To further improve the signal-to-background ratio, two additional discriminating variables are used: $\Delta \phi$ between the two MJs and $\sum p_T^{\text{ID}}$ for the MJ, defined as the scalar sum of the transverse momentum of the tracks, measured in the ID, inside a cone $\Delta R = 0.4$ around the direction of the MJ. The muon tracks of the MJ in the ID, if any, are not removed from the isolation sum, so that prompt muons, which give a reconstructed track in both the ID and MS, will contribute to the $\sum p_T^{\text{ID}}$. As a consequence a cut on $\sum p_T^{\text{ID}}$ of a few GeV will remove prompt MJs or MJs with very short decay length.

For the background coming from cosmic-ray muons (mainly pairs of almost parallel cosmic-ray muons crossing the detector) a cut on the impact parameters of the muon tracks with respect to the primary interaction vertex is used.

The final set of selection criteria used is the following:

- Topology cut: events are required to have exactly two MJs, $N_{\text{MJs}} = 2$.
- MJ isolation: require MJ isolation with $E_{\text{isol}}^{\text{ID}} \leq 5$ GeV for both MJs in the event.
- Require $|\Delta \phi| \geq 2$ between the two MJs.
- Require opposite charges for the two muons in a MJ ($Q_{\text{MJs}} = 0$).
- Require a cut on the transverse and longitudinal impact parameters of the muons with respect to the primary vertex: $|d_0| < 200$ mm and $|z_0| < 270$ mm.
- Require $\sum p_T^{\text{ID}} < 3$ GeV for both MJs.

The distributions of the relevant variables used in the selection before each step of the cut flow are shown in Fig. 4. The results are summarized in Table 2. No events survive the selection in the data sample whereas the expected signals from Monte Carlo simulation, assuming the Higgs boson SM production cross section, 100% branching ratio for $H \to \gamma\gamma$ and $H \to Z\gamma$ and the parameters given in Table 1, are 75 or 48 events for Higgs boson masses of 100 GeV and 140 GeV respectively. The method used to estimate the cosmic-ray and multi-jet background yields, quoted in Table 2, is discussed in Section 7.

The resulting single $\gamma\gamma$ reconstruction efficiency for the mean lifetimes given in Table 1 is shown in Fig. 5 as a function of $\eta$, the $\Delta R$ separation of the two muons from the $\gamma\gamma$ decay and the decay length in the transverse plane, $L_{xy}$, of the $\gamma\gamma$. The efficiency is defined as the number of $\gamma\gamma$ passing the offline selection divided by the number of $\gamma\gamma$ passing the spectrometer acceptance ($|\eta| \leq 2.4$) with both muons having $p_T \geq 6$ GeV. The low reconstruction efficiency at very short $L_{xy}$ is a consequence of the $\sum p_T^{\text{ID}}$ cut.

The systematic uncertainty on the reconstruction efficiency is evaluated using a tag-and-probe method by comparing the reconstruction efficiency $E_{\text{rec}}^{\text{ID}}$ for $j/\psi \to \mu^+\mu^-$ samples from collision data and $j/\psi \to \mu^+\mu^-$ Monte Carlo simulation. The tag-and-probe definitions and the cut on $\Delta R < 0.1$ between the two muons are the same as in Section 5. To measure the reconstruction efficiency the ID probe track is associated with a MS-only muon track, different from the one associated with the tag. The result is shown in Fig. 6.

The relative difference between the result obtained from the $j/\psi \to \mu^+\mu^-$ data and the $j/\psi \to \mu^+\mu^-$ Monte Carlo sample in the same range of $\Delta R \leq 0.1$, as for the signal, is taken as the systematic uncertainty on the reconstruction efficiency and amounts to 13%.

7. Multi-jet and cosmic-ray background evaluation

To estimate the multi-jet background contamination in the signal region we use a data-driven ABCD method slightly modified to cope with the problem of the very low number of events in the control regions. The ABCD method assumes that two variables can be identified, which are relatively uncorrelated, and which can each be used to separate signal and background. It is assumed that the multi-jet background distribution can be factorized in the MJ $E_{\text{isol}}^{\text{ID}} - |\Delta \phi|$ plane. The region A is defined by $E_{\text{isol}}^{\text{ID}} \leq 5$ GeV and $|\Delta \phi| < 2$; the region B, defined by $E_{\text{isol}}^{\text{ID}} \leq 5$ GeV and $|\Delta \phi| > 2$, is the signal region. The regions C and D are the anti-isolated regions ($E_{\text{isol}}^{\text{ID}} > 5$ GeV) and they are defined by $|\Delta \phi| < 2$ and $|\Delta \phi| > 2$, respectively. Neglecting the signal contamination in regions A, C and
The following effects are considered as possible sources of systematic uncertainty:

- **Luminosity**
  The overall normalization uncertainty of the integrated luminosity is 3.7% [30,31].

- **Muon momentum resolution**
  The systematic uncertainty on the muon momentum resolution for MS-only muons has been evaluated by smearing and shifting the momenta of the muons by scale factors derived from $Z \rightarrow \mu\mu$ data-Monte Carlo comparison, and by observing the effect of this shift on the signal efficiency. The overall effect of the muon momentum resolution uncertainty is negligible.

- **Trigger**
  The systematic uncertainty on the single $\gamma d$ trigger efficiency, evaluated using a tag-and-probe method is 17% (see Section 5).

- **Reconstruction efficiency**
  The systematic uncertainty on the reconstruction efficiency, evaluated using a tag-and-probe method for the single $\gamma d$ reconstruction efficiency, is 13% (see Section 6).

- **Effect of pile-up**
  The systematic uncertainty on the signal efficiency related to the effect of pile-up is evaluated by comparing the number of signal events after imposing all the selection criteria on the signal Monte Carlo sample increasing the average number of interactions per crossing from $\approx 6$ to $\approx 16$. This systematic uncertainty is negligible.
uncertainties are statistical only. Evaluated with a similar tag-and-probe method, is also shown. The uncertainties are statistical only.

Fig. 6. Tag-and-probe reconstruction efficiency $\varepsilon_{\text{rec}}$ as a function (a) of $\eta$, (b) of $\Delta R$ and (c) of the transverse decay length of the $\gamma_3$ for $m_H = 100$ GeV and $m_H = 140$ GeV and for the mean lifetimes given in Table 1. The reconstruction efficiency is defined as the number of $\gamma_3$ passing the offline selection divided by the number of $\gamma_3$ in the spectrometer acceptance ($|\eta| \leq 2.4$) with both muons having $p_T \geq 6$ GeV. The uncertainties are statistical only.

Fig. 7. Table 3 shows the ranges in which the model of Fig. 1, as a function of the long-lived particle mean lifetime, with $c_\tau$ ranging from 0 to 700 mm. In this way the number of expected signal events is predicted as a function of the $\gamma_3$ mean lifetime. These numbers, together with the expected number of background events (multi-jet and cosmic rays) and taking into account the zero data events surviving the selection criteria described above, is used as input to obtain limits at the 95% confidence level (CL). The CLs method [33] is used to set 95% CL upper limits on the cross section times branching ratio ($\sigma \times BR$) for the process $H \rightarrow gamma_{\gamma} + X$, according to the model of Fig. 1. Here the branching ratio of $\gamma_3 \rightarrow \mu \mu$ is set to 45% with the $\gamma_3$ mass set to 0.4 GeV, as previously discussed. The $\sigma \times BR$ is given as a function of the $\gamma_3$ mean lifetime, expressed as $c_\tau$ for $m_H = 100$ GeV and $m_H = 140$ GeV. These limits are shown on Fig. 7. Table 3 shows the ranges in which the $\gamma_3$ ct is excluded at the 95% CL for $H \rightarrow \gamma_3 \gamma + X$ branching ratios of 100% and 10%.

9. Results and interpretation

The efficiency of the selection criteria described above is evaluated for the simulated signal samples (see Table 1) as a function of the mean lifetime of the $\gamma_3$. The signal Monte Carlo events are weighted by the detection probability of the two $\gamma_3$ in the various parts of the detector, generating their decay points according to a chosen value of the $\gamma_3$ lifetime, with $c_\tau$ ranging from 0 to 700 mm. In this way the number of expected signal events is predicted as function of the $\gamma_3$ mean lifetime. These numbers, together with the expected number of background events (multi-jet and cosmic rays) and taking into account the zero data events surviving the selection criteria in 1.9 fb$^{-1}$, are used as input to obtain limits at the 95% confidence level (CL). The CLs method [33] is used to set 95% CL upper limits on the cross section times branching ratio ($\sigma \times BR$) for the process $H \rightarrow gamma_{\gamma} + X$, according to the model of Fig. 1. Here the branching ratio of $\gamma_3 \rightarrow \mu \mu$ is set to 45% with the $\gamma_3$ mass set to 0.4 GeV, as previously discussed. The $\sigma \times BR$ is given as a function of the $\gamma_3$ mean lifetime, expressed as $c_\tau$ for $m_H = 100$ GeV and $m_H = 140$ GeV. These limits are shown on Fig. 7. Table 3 shows the ranges in which the $\gamma_3$ ct is excluded at the 95% CL for $H \rightarrow \gamma_3 \gamma + X$ branching ratios of 100% and 10%.

10. Conclusions

The ATLAS detector at the LHC was used to search for a light Higgs boson decaying into a pair of hidden fermions ($f_{\Delta f}$), each of which decays to a $\gamma_3$ and to a stable hidden fermion ($f_{\Delta f}$), resulting in two muon jets from the $\gamma_3$ decay in the final state. In a 1.9 fb$^{-1}$ sample of $\sqrt{s} = 7$ TeV proton–proton collisions no events consistent with this Higgs boson decay mode are observed. The observed data are consistent with the Standard Model background expectations.

Limits are set on the $\sigma \times BR$ to $H \rightarrow gamma_{\gamma} + X$, according to the model of Fig. 1, as a function of the long-lived particle mean lifetime for $m_H = 100$ GeV and 140 GeV with the chosen $\gamma_3$ mass that gives a decay branching ratio of 45% for $\gamma_3 \rightarrow \mu \mu$. Assuming the SM production rate for a 140 GeV Higgs boson, its branching ratio to two hidden-sector photons is found to be below 10%, at 95% CL for hidden photon $c_\tau$ in the range 7 mm $\leq c_\tau \leq 82$ mm. Bounds...
on the $\sigma \times BR$ of a 126 GeV Higgs boson may be conservatively extracted using the corresponding 140 GeV exclusion curve.

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References


Granada, Spain

INAF Sezione di Napoli; (b) Dipartimento di Scienze Fisiche, Università di Napoli, Napoli, Italy

Department of Physics and Astronomy, University of New Mexico, Albuquerque, NM, United States

Department of Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands

Nikhef Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands

Department of Physics, Northern Illinois University, DeKalb, IL, United States

Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia

Department of Physics, New York University, New York, NY, United States

Ohio State University, Columbus, OH, United States

Faculty of Science, Okayama University, Okayama, Japan

Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, OK, United States

Department of Physics, Oklahoma State University, Stillwater, OK, United States

Palacký University, RCPTM, Olomouc, Czech Republic

Center for High Energy Physics, University of Oregon, Eugene, OR, United States

LAL, Université Paris-Sud et CNRS/IN2P3, Orsay, France

Graduate School of Science, Osaka University, Osaka, Japan

Department of Physics, University of Oslo, Oslo, Norway

Department of Physics, Oxford University, Oxford, United Kingdom

INFN Sezione di Pavia; (b) Dipartimento di Fisica, Università di Pavia, Pavia, Italy

Department of Physics, University of Pennsylvania, Philadelphia, PA, United States

Petersburg Nuclear Physics Institute, Gatchina, Russia

INFN Sezione di Pisa; (b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy

Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA, United States

Laboratorio di Instrumentazione e Fisica Eperimentale de Particelle – LIP, Lisboa, Portugal; (b) Departamento de Física Teórica y del Cosmos and CAFF, Universidad de Granada, Granada, Spain

Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic

Faculty of Mathematics and Physics, Charles University in Prague, Praha, Czech Republic

Czech Technical University in Prague, Praha, Czech Republic

State Research Center Institute for High Energy Physics, Protvino, Russia

Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom

Physics Department, University of Regina, Regina, SK, Canada

Nagasaki Institute of Applied Science, Nagasaki, Japan

Physics Department, University of Regensburg, Regensburg, Germany

INFN Sezione di Roma 1; (b) Dipartimento di Fisica, Università La Sapienza, Roma, Italy

INFN Sezione di Roma Tor Vergata; (b) Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy

INFN Sezione di Roma Tre; (b) Dipartimento di Fisica, Università Roma Tre, Roma, Italy

Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies – Université Hassan II, Casablanca; (b) Centre National de l’Energie des Sciences Techniques Nucléaires, Rabat; (c) Faculté des Sciences Semlalia, Université Cadi Ayyad, LPEA, Marrakech; (d) Faculté des sciences, Université Mohamed Premier et LPTPM, Oujda; (e) Faculté des sciences, Université Mohammed V – Agdal, Rabat, Morocco

DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat a l’Energie Atomique), Gif-sur-Yvette, France

Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, CA, United States

Department of Physics, University of Washington, Seattle, WA, United States

Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom

Department of Physics, Shinshu University, Nagano, Japan

Fachbereich Physik, Universität Siegen, Siegen, Germany

Department of Physics, Simon Fraser University, Burnaby, BC, Canada

CERN Accelerator School Laboratory, Stanford, CA, United States

Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava; (b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic

Department of Physics, University of Johannesburg, Johannesburg; (b) School of Physics, University of the Witwatersrand, Johannesburg, South Africa

Department of Physics, Stockholm University; (b) The Oskar Klein Centre, Stockholm, Sweden

Physics Department, Royal Institute of Technology, Stockholm, Sweden

Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook, NY, United States

Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom

School of Physics, University of Sydney, Sydney, Australia

Institute of Physics, Academia Sinica, Taipei, Taiwan

Department of Physics, Technion - Israel Institute of Technology, Haifa, Israel

Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel

Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece

International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan

Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan

Department of Physics, Tokyo Institute of Technology, Tokyo, Japan

Department of Physics, University of Toronto, Toronto, ON, Canada

TRIUMF, Vancouver, BC; (b) Department of Physics and Astronomy, York University, Toronto, ON, Canada

Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan

Department of Physics and Astronomy, Tufts University, Medford, MA, United States

Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia

Department of Physics and Astronomy, University of California Irvine, Irvine, CA, United States

INFN Gruppo Collegato di Udine; (b) ITP, Trieste; (c) Dipartimento di Chimica, Fisica e Ambientale, Università di Udine, Udine, Italy

Department of Physics, University of Illinois, Urbana, IL, United States

Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden

Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica e Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain