The missing GeV γ-ray binary: searching for HESS J0632+057 with Fermi-LAT


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The missing GeV $\gamma$-ray binary: searching for HESS J0632+057 with Fermi-LAT

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ABSTRACT

The very high energy (VHE; $>100$ GeV) source HESS J0632+057 has been recently confirmed as a $\gamma$-ray binary, a subclass of the high-mass X-ray binary population, through the detection of an orbital period of 321 d. We performed a deep search for the emission of HESS J0632+057 in the GeV energy range using data from the Fermi Large Area Telescope (LAT). The analysis was challenging due to the source being located in close proximity to the bright $\gamma$-ray pulsar PSR J0633+0632 and lying in a crowded region of the Galactic plane where there is prominent diffuse emission. We formulated a Bayesian block algorithm adapted to work with weighted photon counts, in order to define the off-pulse phases of PSR J0633+0632. A detailed spectral-spatial model of a $5^\circ$ circular region centred on the known location of HESS J0632+057 was generated to accurately model the LAT data. No significant emission from the location of HESS J0632+057 was detected in the $0.1–100$ GeV energy range integrating over $\sim3.5$ yr of data, with a 95 per cent flux upper limit of $F_{0.1–100\text{GeV}} < 3 \times 10^{-8}$ ph cm$^{-2}$ s$^{-1}$. A search for emission over different phases of the orbit also yielded no significant detection. A search for source emission on shorter time-scales (days–months) did not yield any significant detections. We also report the results of a search for radio pulsations using the 100-m Green Bank Telescope. No periodic signals or individual dispersed bursts of a likely astronomical origin were detected. We estimated the flux density limit of $<90/40$ $\mu$Jy at 2/9 GHz. The LAT flux upper limits combined with the detection of HESS J0632+057 in the 136–400 TeV energy band by the MAGIC collaboration imply that the VHE spectrum must turn over at energies $<136$ GeV placing constraints on any theoretical models invoked to explain the $\gamma$-ray emission.

Key words: gamma-rays: general – X-rays: binaries.

1 INTRODUCTION

The population of $\gamma$-ray binaries is a subclass of the high-mass X-ray binary (HMXB) population; they are binary systems comprised of a high-mass ($>10\ M_\odot$) star and a compact object whose spectral energy distribution (SED) peaks above 1 MeV. Indeed, they...
Searching for γ-rays from HESS J0632+057

The LAT is the primary instrument on board Fermi; it is a pair conversion telescope sensitive to photons with energies from ~20 MeV to more than 300 GeV. The direction of an incident photon is derived by tracking the electron–positron pair in a high-resolution converter tracker, and the energy of the pair is measured with a hodoscopic CsI(Tl) crystal calorimeter. The Fermi-LAT has a peak on-axis effective area of 8000 cm², a 2.4 sr field of view and an angular resolution of ~0.6 at 1 GeV (for events converting in the front section of the tracker). Furthermore, an anticoincidence detector vetoes the background of charged particles (Atwood et al. 2009).

The observatory operates principally in survey mode; the telescope is rocked north and south on alternate orbits to provide more uniform coverage so that every part of the sky is observed for ~30 min every 3 h. The analysis presented here makes use of all data taken from 2008 August 4 until 2012 March 3.

2.1 Data reduction

The Fermi Science Tools v09r28p00 package was used to analyse the ‘source’ event class data. All events within a circular region of interest (ROI) of 15° radius centred on the known location
of the massive star MWC 148 associated with HESS J0632+057 [R.A.(J2000) = 98:247, Dec.(J2000) = 5:800] and in the energy range 0.1–100 GeV were extracted. The good time intervals are defined such that the ROI does not cross the γ-ray-bright Earth limb (defined at 100° from the zenith angle) and that the source is always inside the LAT field of view, namely within a cone angle of 66°. The 'pysource_v6' instrument response functions (IRFs) are used throughout the analysis.

When constructing spectral-spatial models of the ROI, the models used for the Galactic diffuse emission (gal_2yearP7V6_v0.fits) and isotropic backgrounds (iso_P7V6Source.txt) were those currently recommended by the LAT team1 and were used by the Fermi-LAT collaboration to build the second Fermi-LAT source catalog (hereafter 2FGL; Nolan et al. 2012).

3 DATA ANALYSIS: THE γ-RAY SKY AROUND HESS J0632+057

HESS J0632+057 lies at the Galactic latitude \( b = -1:44 \), in the active region of the Monoceros Loop. An image of this region in the γ-ray band is shown in Fig. 1. This is a background-subtracted count map of the \( 10^\circ \times 10^\circ \) region around HESS J0632+057. It is generated by selecting events with energies greater than 1 GeV, and subtracting the counts due to the Galactic diffuse emission and the isotropic emission as calculated using the respective models with the SCIENCE TOOL gtsrcprob. The spectral models of the pulsar and the nearby sources adopted to calculate the weights are those from the 2FGL catalogue. Finally, the pulse profile in Fig. 2 was obtained by summing the weights in each phase bin for the events within 2° of the pulsar position.

To minimize contamination from the strong emission of PSR J0633+0632 in the analysis of HESS J0632+057, the pulsar emission was 'gated'2 by only using events which occurred outside the pulse peaks, during phases 0.66–0.98 and 0.21–0.49 (hereafter out-of-peak phases). The off-pulse and intrapeak phases of PSR J0633+0632 were defined by applying a Bayesian block algorithm (Ray et al. 2011). A phase was associated with each event using the Fermi plug-in of the TEMPO2 software package (Hobbs, Edwards & Manchester 2006). A weight was also assigned to each event, corresponding to the probability that the γ-ray is emitted by the pulsar, rather than by a nearby source or by the diffuse emission (Kerr 2011b); weights are calculated using the SCIENCE TOOL gtsrcprob. The spectral models of the pulsar and the nearby sources adopted to calculate the weights are those from the 2FGL catalogue. Finally, the pulse profile in Fig. 2 was obtained by summing the weights in each phase bin for the events within 2° of the pulsar position.

3.1 PSR J0633+0632

PSR J0633+0632 is a bright, radio-quiet γ-ray pulsar discovered in the first six months of the Fermi mission (Abdo et al. 2009a). It has a flux (>100 MeV) of \((8.4\pm 1.2) \times 10^{-8} \text{ph cm}^{-2} \text{s}^{-1}\), and the pulse profile has the characteristic two narrow peaks separated by ∼0.5 in phase (Abdo et al. 2010).

The pulse profile obtained with the data set analysed in this paper is shown in Fig. 2. To coherently assign a phase to each photon over such a long observation period (∼3.5 yr) requires an up-to-date ephemeris for the pulsar. As PSR J0633+0632 is radio quiet, the new ephemeris has been calculated using the Fermi-LAT data with the method described by Ray et al. (2011). A phase was associated with each event using the Fermi plug-in of the TEMPO2 software package (Hobbs, Edwards & Manchester 2006). A weight was also assigned to each event, corresponding to the probability that the γ-ray is emitted by the pulsar, rather than by a nearby source or by the diffuse emission (Kerr 2011b); weights are calculated using the SCIENCE TOOL gtsrcprob. The spectral models of the pulsar and the nearby sources adopted to calculate the weights are those from the 2FGL catalogue. Finally, the pulse profile in Fig. 2 was obtained by summing the weights in each phase bin for the events within 2° of the pulsar position.

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3.2 Off-pulse and intrapeak definition

We found that an efficient way to define the off-pulse phase range is by subdividing the pulse profile into 'Bayesian blocks' and

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1 Descriptions of the models are available at the FSSC: http://fermi.gsfc.nasa.gov/.  
2 For a detailed example of pulsar gating see the FSSC tutorial: http://fermi.gsfc.nasa.gov/ssc/data/analysis/scitools/pulsar_gating_tutorial.html.
choosing the block with the lowest count rate. The intrapeak phase range is instead the lowest block between the two peaks. The Bayesian blocks algorithm was proposed by Scargle (1998): the algorithm defines the phase ranges of constant rate, called Bayesian blocks, on the basis of the likelihood ratio test applied to a Poissonian distribution. This is perfect for count light curves, but not for weighted pulse profiles, as in this case the weights are not integer values, and they do not follow a Poissonian distribution. Furthermore, the pulse profile is periodic, meaning that it does not have a well-defined start and end.

In order to manage weighted counts, in our method the Bayesian blocks are defined using the \( \chi^2 \) ratio test. Specifically, given a light curve of weighted counts, the null hypothesis is that the light curve is flat, and it is fitted by a single block. The associated \( \chi^2 \) value for a single block is defined as

\[
\tilde{\chi}^2_{\text{bl}} = \frac{1}{N-1} \sum_i \frac{(B_{w,i} - M_w)^2}{M_w},
\]

where \( B_{w,i} \) are the weighted source counts in the \( i \)th bin and \( M_w \) are the average weighted counts all along the light curve. \( N \) is the total number of bins. The test hypothesis is that two blocks fit better than a single block. The \( \tilde{\chi}^2 \) of the two blocks hypothesis is

\[
\tilde{\chi}^2_{\text{2bl}} = \frac{1}{N-2} \left[ \sum_i \frac{(B_{w1,i} - M_{w1})^2}{M_{w1}} + \sum_j \frac{(B_{w2,i} - M_{w2})^2}{M_{w2}} \right],
\]

where \( M_{w1}, B_{w1,i} \) and \( M_{w2}, B_{w2,i} \) are the average weighted counts and the weighted counts in the \( i \)th and \( j \)th bin of the first and second block, respectively. Finally, the significance of the improvement given by two blocks rather than a single block is calculated by means of the F-test.

If two blocks fit the data significantly better than a single block, then the division point between the blocks is recorded and the first of the new blocks (on the left) is tested for further subdivision. The process continues until further subdivision is not significant or the block length reaches a predefined minimum length. In this way, the shortest, furthest left significant block is defined. The subdivision restarts considering the remaining light curve on the right.

The issue relating to the periodicity of the pulse profile is solved by selecting the (starting and ending) point of the light curve to be the phase of the highest peak. This is the choice that least affects the definition of the off-pulse region.

The algorithm has three parameters: the number of bins of the pulse profile (100), the minimum possible length of a block (3 bins) and the threshold on the two blocks significance (1.5 \( \sigma \)). The values in parenthesis are the settings used for PSR J0633+0632. We tested this algorithm on all the \( \gamma \)-ray pulsars detected so far with the Fermi-LAT, finding that these settings give good results for all the bright pulsars.

### 3.3 Out-of-peak analysis

Having restricted the events to those selected from within the out-of-peak phases of PSR J0633+0632 (shaded phase ranges in Fig. 2), the analysis of the region around HESS J0632+057 was performed by means of the binned maximum-likelihood method (Mattox et al. 1996). Two different tools were used to perform the spatial and spectral analysis: SCIENCE TOOL $g\text{tlike}$ and pointlike $g\text{tlike}$ is the official tool supplied with the publicly available SCIENCE TOOLS to fit spectral models to LAT data. pointlike is an alternative binned likelihood implementation, optimized to localize the source and to characterize its extension. It operates quickly on long data sets making it better suited for iteratively constructing source models of a region. Results from pointlike have been extensively compared with $g\text{tlike}$ and found to be consistent (see, e.g., Kerr 2011a; Lande et al. 2012). The likelihood-ratio test statistic (TS) is employed to evaluate the significance of the \( \gamma \)-ray flux of a test source. The TS value is defined to be

\[
\text{TS} = -2\Delta \log \mathcal{L} = -2 \log \frac{\mathcal{L}(\text{Null})}{\mathcal{L}(\text{Test})},
\]

where \( \mathcal{L} \) is the likelihood function maximized by adjusting all parameters of the model with and without the test source (test and null hypothesis, respectively). A TS = 25 corresponds to a significance of just over 4\( \sigma \) when the testing source is modelled with a simple power law (Nolan et al. 2012).

Fig. 3 shows two TS maps of the 5° × 5° region centred on HESS J0632+057. They are generated by testing for an additional point-like source located at the centre of each pixel (of size 0.1° × 0.1°). In the left-hand panel of Fig. 3 is plotted a diffuse subtracted TS map, meaning that the spectral-spatial model for the null hypothesis includes only the Galactic and isotropic emission models. In contrast, the right-hand panel shows the residual TS map when the spectral-spatial model also includes the 2FGL sources within a radius of 20° from the ROI centre, in addition to the diffuse models. The 2FGL sources within 5° of HESS J0632+057 (11 sources) are modelled with their flux normalizations allowed to vary, while the other 45 sources had all their parameters fixed to the values given in the 2FGL catalogue. The spectral model of the isotropic diffuse emission had free normalization and the spectral model of the Galactic diffuse emission is scaled by a simple power law with normalization and index free to account for any local variations in the diffuse emission. In the final fit of the data the fitted value for the power-law index is consistent with 0 within the fit error, indicating that no adjustment to the local Galactic diffuse spectral shape was required.

HESS J0632+057 lies in a region of high Galactic diffuse emission. Since at low energies the diffuse emission is stronger, as well as its uncertainties, the following analysis was restricted to using data at \( E > 300 \text{ MeV} \).

The TS maps in Fig. 3 show that the region around HESS J0632+057 is not well modelled solely by the sources from the 2FGL catalogue. This is not unexpected considering that the 2FGL catalogue was constructed with 2 yr of Fermi-LAT data, while in this paper the data set is almost doubled (∼3.5 yr). Searching for emission from HESS J0632+057 requires an accurate modelling of the entire region, and hence, an improved source model of the region is required. The ROI source model was iteratively improved by adding a new point source at the location of the highest residual TS and then recalculating the residual TS map (see Section 3.3.3). Initially, the 3° region around HESS J0632+057 was fine-tuned before extending the radius to 5°. A detailed description of the process and results is given in the following sections.

#### 3.3.1 TS calculation of 2FGL sources

In the spectral-spatial model described in the previous section, an additional point source was added at the position of HESS J0632+057. Its spectrum is modelled with a simple power law, with flux and spectral index parameters set free. The TS values of all the 2FGL sources within 3° of HESS J0632+057 were calculated from a maximum likelihood fit with the $g\text{tlike}$ tool.
With the exception of PSR J0633+0632 and HESS J0633+057, only one source, 2FGL J0637.8+0737, had TS < 50. Four sources (2FGL J0636.0+0554, 2FGL J0631.7+0428, 2FGL J0634.3+0356c and 2FGL J0637.0+0416c) were found with TS values between 50 and 100 and only 2FGL J0631.6+0640 had TS > 100. As an initial step, the source with TS < 50 was deleted from the spectral-spatial model and those sources with 50 < TS < 100 were relocalized to their optimum positions using pointlike.

3.3.2 Source spectrum modelling

The residual off-pulse emission of PSR J0633+0632 is very weak, with TS = 23. As a consequence, the spectral model was changed from the superexponential power law, typical of γ-ray pulsars, to a simple power law.

After this change, we started the source spectrum modelling. In order to reach a stable solution for all the sources in the region, we first fit the brightest sources, keeping frozen the other sources to their nominal spectra. Then, we set free weaker and weaker sources, until in the last fit all sources had free spectral parameters. In practice, this consisted of three steps. First, a maximum likelihood fit was performed with gtlike setting free all the spectral parameters of the sources with TS > 90 within 3° of HESS J0632+057. In the second step, the fit was repeated setting free also the spectral parameters of the sources with TS between 50 and 100. Finally, in the third step, the spectral parameters of all the sources were set free. This established the best model of the region on the basis of the 2FGL catalogue source population from which additional, newly found sources could be assessed.

3.3.3 Adding new sources

An iterative procedure was implemented to model excesses in the residual TS map as additional point sources in the model. The process consisted of the following stages.

(i) Calculate the residual TS map using the current spectral-spatial model.
(ii) Add an additional point source to the model at the location of the highest excess in the TS map, modelled by a simple power law.
(iii) Fit the best position of the additional source with the pointlike tool.
(iv) Check the curvature of the spectrum of the additional source. If curvature is indicated, substitute the initial power-law model with a log parabola.
(v) Refit with gtlike the spectra of all the sources in the region, setting free also the spectral parameters of the additional source.

The procedure was repeated until the maximum value within the residual TS map was less than 15, below which the procedure would effectively be fitting only statistical noise. Five additional point sources were added within 3° of HESS J0632+057. Three more point sources were added when the procedure was repeated enlarging the radius to 5°. Table 1 summarizes the results obtained by modelling the region around HESS J0632+057; listed are all the sources included in the final spectral-spatial model within a radius of 5° from HESS J0632+057.

Of the eight additional sources, one source, ‘New1’, may be a result of resolving 2FGL J0634.3+0356c into two sources, as the nominal position of 2FGL J0634.3+0356c was shifted when it was re-localized in our fitting procedure. Note that this 2FGL catalogue source was denoted as a ‘c’ source and hence is found in a region of bright and/or possibly incorrectly modelled diffuse emission (Nolan et al. 2012). It is worth noting that ‘New2’ is 29 arcmin from 2FGL J0637.8+0737 (the source deleted from the model in the initial stage, see Section 3.3.1), although it lies outside the 95 per cent error radius for this catalogue source.

The majority of the additional sources are very weak. Due to this and the challenge of modelling the Galactic diffuse emission, it is very difficult to definitively characterize them as new γ-ray sources with any confidence rather than statistical fluctuations in the data or inaccurately modelled regions of the Galactic diffuse

Figure 3. TS map (left) and residual TS map (right; 2FGL sources are included in the model) of the 5° × 5° sky region centred on HESS J0632+057. The maps are calculated for $E > 300$ MeV using photons in the off-pulse and intrapeak phases of PSR J0633+0632. The 2FGL sources are labelled with green crosses, HESS J0632+057 with a diamond and the position of PSR J0633+0632 with the circle.
emission. Consequently, we only report one new γ-ray source, Fermi J0642.4+0648, that is detected at >7σ at a location of RA(J2000)=101.60, Dec.(J2000)=6.80 with a 95 per cent error radius of ∼7.8 arcmin. The source spectrum is modelled by a log parabola shape, 

$$\frac{dN}{dE} = N_0 \left( \frac{E}{E_0} \right)^{-\alpha + \beta \log(E/E_0)}$$

with parameters $N_0 = 0.54 \pm 0.08$, $\alpha = 1.5 \pm 0.4$, $\beta = 0.4 \pm 0.1$, and $E_0 = 2.3 \pm 0.6$.

### 4 HESS J0632+057

At no point in the analysis did the source that was added at the position of HESS J0632+057 have a significant TS value. Using the final source model, fitting for a source at the position of HESS J0632+057 yields a TS of 4.7, far below the significance threshold for claiming a source detection.

To cover the possibility that the source has a very soft spectrum in the LAT energy band, the entire analysis, including the iterative building of a source model, was repeated using all events at energies $E > 100$ MeV. The additional events incorporated into this analysis did not yield a statistically significant detection of HESS J0632+057. The region model for an $E > 100$ MeV analysis required the addition of four point sources in addition to the 2FGL catalogue sources. These four sources corresponded to four of the sources that were incorporated into the $E > 300$ MeV analysis and are indicated in Table 1.

### 4.1 Orbital variability and flaring analysis

In the analyses described so far, HESS J0632+057 was assumed to be a persistent source and we conclude that its average γ-ray emission is below the sensitivity of Fermi-LAT. All of the other LAT-detected γ-ray binaries show periodic orbital variability including the extreme case of the periodic transient emission seen from PSR B1259–63 (Abdo et al. 2011; Tam et al. 2011), which is only detectable at HE and VHE around periastron passage. To investigate the hypothesis that HESS J0632+057 is only detectable during a specific phase of the orbit, a further analysis was performed folding the Fermi-LAT data on the orbital period of 321 d and setting phase 0.0 to the epoch 548 57 MJD (Bongiorno et al. 2011). The orbit was subdivided into eight equally spaced phase bins, and for each of them a maximum likelihood fit was performed freezing the spectral parameters of all nearby sources in the model. No significant emission was found in any of the phase bins.

HESS J0632+057 has been routinely monitored by the LAT Collaboration. A standard analysis is performed on daily, weekly and 28 d time-scales to search for any flaring behaviour from the source; on these time-scales, the source has never been detected by the LAT.

### Table 1

The sources included in the final spectral-spatial model within a radius of 5° of HESS J0632+057. The table is split into four sections: 2FGL catalogue sources; 2FGL sources relocalized in our procedure (named Reloc instead of 2FGL); new sources introduced to model the TS residuals and HESS J0632+057. The flux is for $E > 300$ MeV in units of 10$^{-9}$ cm$^{-2}$ s$^{-1}$.

<table>
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</thead>
<tbody>
<tr>
<td>2FGL J0631.5+1035</td>
<td>97.88</td>
<td>10.60</td>
<td>PLEExpCutoff</td>
<td>17.4 ± 1.3</td>
<td>350</td>
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<td>2FGL J0631.6+0640</td>
<td>97.92</td>
<td>6.68</td>
<td>LogParabola</td>
<td>9.3 ± 1.2</td>
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<td>2FGL J0633.7+0633a</td>
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<td>6.556</td>
<td>PowerLaw</td>
<td>9.6 ± 2.2</td>
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<td>2FGL J0641.1+1006c</td>
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<td>10.10</td>
<td>LogParabola</td>
<td>12.5 ± 1.8</td>
<td>71</td>
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<td>2FGL J0642.9+0319</td>
<td>100.73</td>
<td>3.33</td>
<td>LogParabola</td>
<td>9.9 ± 2.0</td>
<td>30</td>
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<tr>
<td>2FGL J0643.2+0858</td>
<td>100.82</td>
<td>8.98</td>
<td>LogParabola</td>
<td>17.0 ± 1.6</td>
<td>170</td>
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<tr>
<td>Reloc J0631.7+0428b</td>
<td>97.70</td>
<td>4.80</td>
<td>LogParabola</td>
<td>9.1 ± 1.7</td>
<td>45</td>
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<tr>
<td>Reloc J0634.3+0356c</td>
<td>98.52</td>
<td>4.27</td>
<td>PowerLaw</td>
<td>13.2 ± 2.4</td>
<td>64</td>
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<tr>
<td>Reloc J0634.0+0554b</td>
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<td>6.10</td>
<td>LogParabola</td>
<td>11.3 ± 1.9</td>
<td>54</td>
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<tr>
<td>Reloc J0637.0+0416c</td>
<td>99.34</td>
<td>4.45</td>
<td>LogParabola</td>
<td>8.1 ± 2.0</td>
<td>27</td>
</tr>
<tr>
<td>Fermi J0642.4+0648d</td>
<td>101.60</td>
<td>6.80</td>
<td>LogParabola</td>
<td>8.9 ± 1.4</td>
<td>53</td>
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<tr>
<td>New1c</td>
<td>98.54</td>
<td>3.74</td>
<td>PowerLaw</td>
<td>10.1 ± 2.2</td>
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<td>New2e</td>
<td>99.65</td>
<td>7.17</td>
<td>PowerLaw</td>
<td>8.7 ± 1.8</td>
<td>34</td>
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<tr>
<td>New4f</td>
<td>96.55</td>
<td>8.80</td>
<td>PowerLaw</td>
<td>4.8 ± 1.2</td>
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<td>New5g</td>
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<td>PowerLaw</td>
<td>6.5 ± 1.6</td>
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<td>New6</td>
<td>98.66</td>
<td>5.27</td>
<td>PowerLaw</td>
<td>7.2 ± 1.8</td>
<td>28</td>
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<td>New7</td>
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<td>4.32</td>
<td>PowerLaw</td>
<td>6.1 ± 1.6</td>
<td>23</td>
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<tr>
<td>New8h</td>
<td>97.05</td>
<td>5.10</td>
<td>PowerLaw</td>
<td>6.2 ± 1.9</td>
<td>25</td>
</tr>
<tr>
<td>HESS J0632+057</td>
<td>98.247</td>
<td>5.800</td>
<td>PowerLaw</td>
<td>–</td>
<td>4.7</td>
</tr>
</tbody>
</table>

aPSR J0633+0632.
bSource position shifted from that found in the 2FGL catalogue beyond its 95 per cent confidence radius.
cThe original 2FGL J0634.3+0356c source has been resolved into two separate sources in this analysis.
dAlso detected in the >100 MeV iterative analysis.
e~0.5 from the catalogue position of 2FGL J0637.8+0737 which was removed from our initial source model (see the text for details).
4.2 Flux upper limits

As no significant source is identified at the location of HESS J0632+057, 95 per cent flux upper limits are derived in different energy bins using the Bayesian method developed by Helene (1983). With this method, the 95 per cent flux upper limits are found by integrating the likelihood profile (function of the source flux) starting from $F = 0$, without any assumption on its distribution. To calculate the values of the likelihood, the final spectral-spatial model is used with the flux normalization factors of all the sources with TS < 50 set free, while freezing all the parameters of the other sources.

We kept fixed the spectral index of HESS J0632+057 to the best-fitting value found in the previous analysis, $\Gamma = 2.9$, while its flux normalization factor was allowed to vary. Table 2 reports the 95 per cent upper limits calculated in different energy bins. Additional upper limits were calculated over a restricted range of orbital phases (0.250, 0.375) that correspond to when the X-ray and VHE fluxes are observed to peak over the orbit. In Fig. 4, the Fermi-LAT upper limits are plotted together with the detection points of MAGIC (Aleksić et al. 2012).

The uncertainties of the Fermi-LAT effective area and of the Galactic diffuse emission are the two primary sources of systematics that can affect the results derived in our analysis. The former systematic effect is estimated by repeating the upper limit calculations using modified IRFs that bracket the ‘P7SOURCE_V6’ effective areas. Specifically, they are a set of IRFs in which the effective area has been modified considering its uncertainty as a function of energy in order to maximally affect a specific spectral parameter (Ackermann et al. 2012a). The latter systematic effect is evaluated by substituting in the spectral-spatial model alternative models for the Galactic diffuse emission. Eight different models have been generated that vary the cosmic ray source distribution, the halo size and the HI spin temperature. In contrast to the standard model, the alternative models also have free independent normalization coefficients for the inverse Compton component, and for the components associated with H i and CO, which are subdivided in four independent Galactocentric annuli (de Palma et al. 2013). The effect of the systematic is included in the upper limits reported in Table 2.

Table 2. 95 per cent flux ($F_{95}$), energy flux ($G_{95}$) and $\nu F_{\nu}$ upper limits of HESS J0632+057 calculated over the LAT energy range. Greyed-out values indicate flux upper limits calculated over a restricted set of orbital phases that correspond to the peak emission in the X-ray & TeV energy bands.

<table>
<thead>
<tr>
<th>Energy range (GeV)</th>
<th>TS</th>
<th>$F_{95}$ ($10^{-9}$ ph cm$^{-2}$ s$^{-1}$)</th>
<th>$G_{95}$ ($10^{-12}$ erg cm$^{-2}$ s$^{-1}$)</th>
<th>$\nu F_{\nu}$ (95 per cent) ($10^{-12}$ erg cm$^{-2}$ s$^{-1}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.1–0.3</td>
<td>1.7</td>
<td>&lt;40</td>
<td>&lt;60</td>
<td>&lt;12</td>
</tr>
<tr>
<td>0.3–1.0</td>
<td>4.7</td>
<td>&lt;9.0</td>
<td>&lt;15</td>
<td>&lt;7.7</td>
</tr>
<tr>
<td>1.0–3.0</td>
<td>4.3</td>
<td>&lt;1.4</td>
<td>&lt;4.1</td>
<td>&lt;3.6</td>
</tr>
<tr>
<td>3.0–10</td>
<td>0.0</td>
<td>&lt;0.21</td>
<td>&lt;0.69</td>
<td>&lt;1.7</td>
</tr>
<tr>
<td>10–30</td>
<td>0.0</td>
<td>&lt;0.08</td>
<td>&lt;0.60</td>
<td>&lt;2.2</td>
</tr>
<tr>
<td>30–100</td>
<td>0.0</td>
<td>&lt;0.06</td>
<td>&lt;0.42</td>
<td>&lt;4.8</td>
</tr>
<tr>
<td>&gt;0.1</td>
<td>6.7</td>
<td>&lt;30</td>
<td>&lt;80</td>
<td>&lt;30</td>
</tr>
</tbody>
</table>

Figure 4. The HE SED of HESS J0632+057 showing the 95 per cent energy flux upper limits calculated from the LAT data using all phases (in black). Also shown are the 95 per cent energy flux upper limits calculated using orbital phases which correspond to the X-ray peak (in red), together with the MAGIC measurements (red data points) corresponding to the source detection in the same orbital phases (Aleksić et al. 2012).

5 RADIO OBSERVATIONS

We tested the hypothesis that the HESS J0632+057 system is a binary in which one of the two stars is an observable radio pulsar. With this aim, we searched for radio pulsations using the GBT. In order to mitigate the possible effects of scattering and eclipsing by the wind of a massive companion, especially near periastron, in analogy with PSR B1259–63 (Johnston et al. 1992), our strategy was to observe at the relatively high radio frequencies of 2 and 9 GHz. Since no orbital ephemeris was available when planning these observations, to try not to observe the system close to periastron, when obscuration and Doppler smearing of the radio pulses (due to orbital motion) are at their maximum, we made a series of six observations spread over the four-month period of 2008 December 24–2009 April 19 (see Table 3). All observations were performed with the GBT Pulsar Spigot (Kaplan et al. 2005).

Table 3. Summary of GBT Observations. The orbital phases are calculated with respect to the phase zero as defined by Casares et al. (2012) (at MJD = 548 570.0) and using an orbital period of 321 d. For reference, periastron and apastron occur at phases 0.967 and 0.467, respectively.

<table>
<thead>
<tr>
<th>Date</th>
<th>MJD</th>
<th>Obs. time (s)</th>
<th>Obs. freq. (GHz)</th>
<th>Orb. phase</th>
</tr>
</thead>
<tbody>
<tr>
<td>24-12-2008</td>
<td>548 24.140</td>
<td>3000</td>
<td>2</td>
<td>0.90</td>
</tr>
<tr>
<td>23-03-2009</td>
<td>549 13.853</td>
<td>5000</td>
<td>9</td>
<td>0.18</td>
</tr>
<tr>
<td>26-03-2009</td>
<td>549 16.863</td>
<td>5000</td>
<td>2</td>
<td>0.19</td>
</tr>
<tr>
<td>27-03-2009</td>
<td>549 17.947</td>
<td>3900</td>
<td>2</td>
<td>0.19</td>
</tr>
<tr>
<td>30-03-2009</td>
<td>549 20.814</td>
<td>5000</td>
<td>9</td>
<td>0.20</td>
</tr>
<tr>
<td>19-04-2009</td>
<td>549 40.832</td>
<td>8500</td>
<td>2</td>
<td>0.26</td>
</tr>
</tbody>
</table>
The centre frequencies were 1.95/8.90 GHz, with 768/1024 spectral channels spanning a bandwidth of 600/800 MHz. For both observing frequencies, the sampling time was 81.92 µs. The maximum dispersion measure (DM) along this line of sight, according to the NE2001 Galactic free electron density model of Cordes & Lazio (2002), is 182 pc cm\(^{-3}\). For reference, for a moderate distance of 3 kpc, the NE2001 model predicts a DM of \(\sim 100\) pc cm\(^{-3}\). Given the high spectral and time resolution of these data, the high observing frequencies and the relatively low expected DM along this line of sight, our searches were not significantly hampered by intrachannel dispersive smearing, even down to millisecond pulse periods.

We used both a Fourier-based acceleration search (Ransom, Eikenberry & Middleditch 2002) and a search for single dispersed pulses using the PRESTO suite of pulsar software. We used standard search procedures (e.g. Hessels et al. 2007) to clean the data of radio frequency interference and to dedisperse the data at a range of DMs from 0 to 364 pc cm\(^{-3}\) (up to twice the maximum expected total Galactic column density along this line of sight).

Inspection of the resulting pulsar candidates revealed no signals of likely astronomical origin. Under the assumption that the pulsar is beamed towards us and it is not significantly scattered by the intervening interstellar medium, we can place a simple flux density limit at our two observing frequencies using the modified radiometer equation (Dewey et al. 1985). At 2/9 GHz, we estimate a maximum flux density of \(< 90/40 \mu\text{Jy}\) (for \(P_{\text{up}} \geq 1\) ms, 1-h observation time, an assumed pulse width of 10 per cent and total system plus sky temperature of 40 K). This means that if the system does contain a pulsar, it must be significantly dimmer than the bulk of the known pulsar population, not beamed towards Earth, very distant and/or significantly obscured by the wind from its companion star.

5.1 Numerical simulations

Within the hypothesis that the compact object is a pulsar with its radio beam pointing towards us, we performed numerical simulations of observations at both 1.95 and 8.90 GHz in order to understand at which orbital phase the pulsar could be most likely visible, and we compared results with the GBT data. In particular, we calculated the orbital phase the pulsar could be most likely visible, and we performed numerical simulations of both 1.95 and 8.90 GHz in order to understand at which orbital phase the pulsar could be most likely visible, and we performed numerical simulations of both 1.95 and 8.90 GHz in order to understand at which orbital phase the pulsar could be most likely visible. We used both a Fourier-based acceleration search (Ransom, Eikenberry & Middleditch 2002) and a search for single dispersed pulses using the PRESTO suite of pulsar software. We used standard search procedures (e.g. Hessels et al. 2007) to clean the data of radio frequency interference and to dedisperse the data at a range of DMs from 0 to 364 pc cm\(^{-3}\) (up to twice the maximum expected total Galactic column density along this line of sight).

Of the orbit for almost all the explored geometrical configurations of the system; in particular, the GBT non-detections would always fall at orbital phases at which we would expect to observe 100 per cent of the total flux or so, thus indicating that the hypothesis of a truncated disc for this system is not compatible with the present data. In the case of an extended disc, none of the pulsar orbital phases can be selected as being safe for a possible pulsar detection. Largely because of our lack of information on the disc orientation and physical characteristics, we were always able to find a proper combination of geometrical parameters to put the eclipse interval at any possible orbital phase. If the pulsar signal were finally detected, we would then be able to put significant constraints on the system geometry.

6 DISCUSSION

A detailed analysis of \(\sim 3.5\) yr of Fermi-LAT data does not yield a significant detection of the \(\gamma\)-ray binary HESS J0632+057. Similarly, searches for persistent, orbitally variable or flaring emission did not yield any detection. This makes this object unique among known \(\gamma\)-ray binaries as the only member of the class eluding detection in the 0.1–100 GeV energy band, with an energy flux upper limit of \(< 3 \times 10^{-11} \text{ erg cm}^{-2} \text{ s}^{-1}\).

Radio observations were also conducted with GBT. No periodic signals or individual dispersed bursts of a likely astronomical origin were detected. We estimated the flux density limit of \(< 90\) and \(< 40 \mu\text{Jy}\) at frequencies of 2 and 9 GHz, respectively, assuming the pulsar is beamed towards Earth.

Some of the properties of the confirmed \(\gamma\)-ray binaries are summarized in Table 4. Upon first glance, the multiwavelength properties of HESS J0632+057 exhibit similarities to PSR B1259–63 and LS I + 61°303 as follows.

(i) All three systems have Be stars as the companion star.

(ii) The HESS J0632+057 orbital period lies in the range between LS I + 61°303 and PSR B1259–63.

(iii) The VHE and X-ray emissions are in phase in both HESS J0632+057 and LS I + 61°303.

LS I + 61°303 is one of the brightest persistent GeV sources detected by the LAT and is detectable throughout its orbit, while PSR B1259–63 is characterized by transient \(\gamma\)-ray emission, as evidenced by the LAT detecting the source only around periastron passage. Of all the known \(\gamma\)-ray binaries, HESS J0632+057 has the closest estimated distance to Earth and consequently it was anticipated as a strong candidate for detection at GeV energies by Fermi.

The non-detection is all the more intriguing in light of the recent report by Aleksić et al. (2012), who detected emission from HESS J0632+057 with the MAGIC telescope from 136 GeV to 4 TeV in 2011 February. They reported extending the spectral detection of the source into the 136–400 GeV band, and their spectral fit to the data indicated a pure power-law shape, consistent with the spectrum reported by the HESS collaboration (Aharonian et al. 2007), with

\[ n_e(r) = n_0 (r_0/r)^2, \]

\[ \beta = 2.55 \text{ and } n_0 = 10^{12} \text{ cm}^{-3}. \]

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Only a narrow region of the parameter space is consistent with a net decrease in the observed flux at the orbital phases corresponding to some of the non-detections. However, this only occurs at the lower frequency (1.95 GHz), assuming very high inclination angles (i \(\sim 90^\circ\)) and for an increased electron density of the isotropic wind with respect to Bogomazov (2005) best-fitting model.
no sign of a spectral turnover. The LAT upper limits calculated in the same phase range as the MAGIC observations are plotted in Fig. 4 in red, while the MAGIC measurements in black. The LAT 95 per cent flux upper limits are sufficiently constraining to indicate that the VHE spectrum observed by MAGIC and HESS cannot be extrapolated back by a simple power law to the LAT energy band and, consequently, there must be a turnover in the spectrum between 10–136 GeV in order for the source not to be detectable in the available LAT data set. The LAT upper limits calculated over all orbital phases are plotted in black in Fig. 4. With the current results, we cannot exclude that at HE the orbital flux modulation of HESS J0632+057 is a negligible fraction of its overall flux. Under this hypothesis, the comparison of the LAT upper limits calculated over all the orbits with the MAGIC measurements gives a more stringent constraint of the spectral turnover, which should happen close to 100 GeV.

In the context of the well-studied γ-ray binaries, LS 1+61°303 and LS 5039, a disconnection between the HE and VHE emission components is not entirely unexpected. The LAT spectra of LS 1+61°303 and LS 5039 are best described by a power law with an exponential cutoff that does not extrapolate up to the reported VHE spectra (Hadasch et al. 2012). Both sources have also shown indications of a HE tail that may connect to the VHE emission. However, it is clear that the GeV emission is dominated by the cutoff power-law spectral component which has led to the suggestion that two different emission processes may be in action to produce the HE and VHE emissions detected in these binary systems.

6.1 Striped wind emission from a pulsar

A potential explanation for the ‘missing’ HE emission from HESS J0632+057 has been suggested by Petri & Dubus (2011). They propose that the HE emission detected from PSR B1259–63 can be explained by inverse Compton up-scattering of stellar photons by particles within a ‘striped pulsar wind’. In this model, the HE emission is produced in the vicinity of the pulsar, explaining the ‘pulsar-like’ exponentially cutoff power-law spectrum observed by the LAT, while the radio/X-ray/VHE emission is produced beyond the termination shock of the pulsar wind. Such a model explains the disconnection between the HE and VHE emissions in the γ-ray binaries. For the HE emission to be detectable the Earth line of sight must pass through the ‘striped wind’. Assuming that the pulsar obliquities are isotropically distributed, this would only occur in 50 per cent of the γ-ray binaries.

7 CONCLUSIONS

A rigorous search for emission in the GeV energy band from the γ-ray binary HESS J0632+057 using ∼3.5 yr of Fermi-LAT data has not yielded a significant detection. There is no significant persistent, flaring or orbital phase-dependent emission. This makes HESS J0632+057 unique among the handful of known γ-ray binaries as the only source eluding detection by the LAT.

The 95 per cent LAT flux upper limit for the 0.1–100 GeV energy band of <3 × 10^{-6} cm^{-2} s^{-1} is consistent with the flux predicted by the one-zone synchrotron-self-Compton model shown in fig. 4 of Hinton et al. (2009) and Skilton et al. (2009), which assume a magnetic field of B = 70 mG and a radiation density of U_{rad} = 1 erg cm^{-3}. However, it is clear from the LAT data that there must be a turnover in the VHE spectrum of HESS J0632+057 below ∼136 GeV in order for the Fermi-LAT data to be compatible with the data obtained at higher energies by MAGIC and HESS. The ‘striped pulsar wind’ model invoked by Petri & Dubus (2011) to explain the emission from PSR B1259–63 could explain what is observed in HESS J0632+057; however, a pulsar nature for the compact object in this system has not been confirmed and hence other models may offer feasible explanations.

Numerical simulations were performed to calculate the radio flux fraction expected to be observed along the orbit by a pulsar with its radio beam pointing towards us, while absorbed by free-free interaction with the strong wind (polar and disc) from the companion. The results of the simulations lead us to exclude the truncation of the disc in favour of an indefinitely extended disc. Alternatively, we could exclude the hypothesis that the radio beam of the plausible pulsar is pointing towards the Earth.

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