Elucidating the backbone conformation of photoswitchable foldamers using vibrational circular dichroism

Published in: Physical Chemistry Chemical Physics

DOI: 10.1039/c3cp53243g

Citation for published version (APA):
Elucidating the backbone conformation of photoswitchable foldamers using vibrational circular dichroism

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The backbone conformation of amphiphilic oligo(azobenzene) foldamers is investigated using vibrational circular dichroism (VCD) spectroscopy on a mode involving the stretching of the \(N=\)N bonds in the backbone. From denaturation experiments, we find that the VCD response in the helical conformation arises mainly from through-space interaction between the \(N=\)N-stretch transition-dipole moments, so that the coupled-oscillator model can be used to predict the VCD spectrum associated with a particular conformation. Using this approach, we elucidate the origin of the VCD signals in the folded conformation, and can assign the observed partial loss of VCD signals upon photo-induced unfolding to specific conformational changes. Our results show that the \(N=\)N-stretch VCD response provides an excellent probe of the helical conformation of the \(N=\)N bonds in this type of switchable molecular system.

1 Introduction

Molecular systems with the ability to undergo conformational changes triggered by light have inspired the design of light-activated molecular machinery.\(^1\)\(^-\)\(^6\) An interesting class of such artificial molecular architectures is formed by photoswitchable amphiphilic oligo(azobenzene) foldamers, which can be manipulated to undergo a helix–coil conformational transition upon photoisomerization around the \(N=\)N bonds, with extraordinary folding–unfolding efficiencies.\(^7\)\(^-\)\(^9\) The backbone of these foldamers consists of azobenzene monomeric units, which can be assembled to obtain molecular systems of various lengths (Fig. 1). The photo-induced \(\text{trans-} \rightarrow \text{cis}\) isomerization triggers a conformational change in the backbone of the \(X_n\) azo helix (where \(X\) denotes the number of phenylene units, and \(n\) indicates the number of azobenzene moieties), causing it to unfold.\(^7\)\(^-\)\(^9\) Refolding to the helical conformation can be achieved by irradiation with visible light or alternatively by thermal relaxation.\(^7\)\(^-\)\(^9\)

As yet, the folding and unfolding of synthetic foldamers are generally investigated using electronic circular dichroism (CD),\(^1\)\(^1\) which provides a rather indirect probe of the helical backbone conformation. Here, we investigate the conformation and photo-induced conformational changes in azo-foldamers using vibrational circular dichroism (VCD) spectroscopy on a vibrational mode involving the \(N=\)N-stretch and in-plane vibrations of the aromatic rings of the azobenzene units in the foldamer backbone. In the last decade VCD has emerged as a powerful structure-resolving tool to study chiral molecules. In most cases, the interpretation of VCD spectra relies on complementary density-functional theory (DFT) calculations, and combined with such calculations VCD provides detailed information on the configuration and conformation of chiral molecules in the condensed phase.\(^1\)\(^4\)\(^-\)\(^2\)\(^1\) In particular, studies that combine VCD spectroscopy and \(\textit{ab initio}\) calculations show convincing results regarding the preferred secondary structures of foldamers.\(^2\)\(^4\)\(^-\)\(^2\)\(^6\) As an alternative to spectral interpretation using DFT calculations,
the simpler, semiquantitative coupled-oscillator method, in which the molecular vibrations are treated as localized, interacting transition dipoles, is attractive because of its transparency and computational ease,\textsuperscript{27–30} in particular for larger systems such as foldamers. Since we find that the N–N-stretch VCD signal of azo-foldamers arises predominantly from through-space interaction between the N–N-stretch transition-dipole moments, we can use this coupled-oscillator approach to interpret the observed foldamer VCD spectra. We find that the N=N groups give rise to unique VCD signatures, which not only provide direct access to the twist sense of the helical structure, but also contain more structural information than the conventional (electronic) CD spectrum.

2 Materials and methods

The synthesis and purification of the foldamers have been reported elsewhere.\textsuperscript{7} Fourier-transform infrared (FTIR) and VCD spectra were obtained using a Bruker Vertex 70 spectrometer in combination with a PMA 50 module for polarization modulation measurements. Samples of the 14\textsubscript{7} foldamer were prepared in CD\textsubscript{3}CN and CD\textsubscript{2}Cl\textsubscript{2} and kept in sealed infrared cells with CaF\textsubscript{2} windows. Photo-switching of the 14\textsubscript{7} foldamer was performed using a 450 W Xe lamp (Müller Elektronik) and a notch filter with $T = 12\%$ at $\lambda = 358$ nm (FWHM = 45 nm). The cartesian coordinates of the nitrogen atoms in the backbone of the 14\textsubscript{7} foldamer were retrieved from an optimized structure obtained using the MM2\textsuperscript{10} force field (Fig. 1).

3 Mode assignment

The folded conformation of a 14\textsubscript{7} azo-foldamer is structurally organized as follows: an alternating sequence of seven azobenzene units is connected through ethynylene bridges. The side chains consist of oligo(ethylene glycol) moieties, which are attached to the main chain azobenzene phenyl rings and thus provide the necessary solvophobic driving force for folding in a polar medium.\textsuperscript{7} In addition, the presence of the side-chains causes a symmetry breaking around the N–N bond. This lower symmetry as compared to that of non-substituted azobenzene gives rise to infrared activity of the vibrational mode involving mainly the stretching of the N–N bond. For simplicity, we will from now on refer to this mode (which also involves the phenyl ring vibration) as the N–N-stretch mode. The infrared activity of the N–N-stretching mode in such non-symmetrically substituted trans-azoazobenzene systems has been convincingly demonstrated for a series of compounds in a recent combined experimental and theoretical study.\textsuperscript{31} We confirmed our assignment of the IR band at $\approx 1450$ cm$^{-1}$ to this mode by quantum-chemical calculations on an isolated substituted azobenzene unit (Fig. 2). The geometry optimization and calculation of harmonic vibrational frequencies were performed at the DFT/BLYP/6-31G(d)\textsuperscript{32–37} level of theory for a simplified model of the system, where we substitute the oligo(ethylene glycol) branches by ethyl-ester groups. Fig. 2 displays the calculated displacement vectors and the direction of the transition-dipole moment for the N–N-stretching mode. We find the transition dipole to be approximately collinear ($\approx 15^\circ$) with the N–N bond.

4 N=N-stretch VCD spectra

Initially we investigate to what extent the observed VCD signal arises from through-space interactions between the N–N-stretch transition-dipole moments. Previous studies with electronic CD have shown that the helical content of 14\textsubscript{7} in CD\textsubscript{3}CN solution progressively vanishes upon addition of CDCl\textsubscript{3} to the solution due to the reduction of the solvophobic effect. We therefore measured the IR and VCD spectra of a 10$^{-3}$ M solution of 14\textsubscript{7} in CD\textsubscript{3}CN and CD\textsubscript{2}Cl\textsubscript{2} (Fig. 3; CD\textsubscript{2}Cl\textsubscript{2} rather than CDCl\textsubscript{3} was used as a solvent because of strong infrared absorption of CDCl\textsubscript{3} in the N–N-stretch region). We find no N=N-stretch VCD signal arising from the denatured, optically-inactive conformation of the foldamer in apolar media. This proves that the VCD signals observed for the folded conformation (the black solid line in Fig. 3) arise mainly from the helical arrangement and through-space coupling of the N–N-stretch vibrational chromophores.

4.1 Coupled-oscillator model for the N=N-stretch modes

To calculate the vibrational frequencies and the vibrational circular dichroism of a helical azo-foldamer containing $n$ N=N groups using the coupled-oscillator model, we use the formalism derived by Diem et al.\textsuperscript{32} Due to the interaction between the N–N-stretch modes, the N–N-stretch normal modes will be
delocalized over the helix, i.e., they will involve the vibrational motion of more than one chromophore. These delocalized normal modes are the eigenstates of the excitonic Hamiltonian

\[
H = \begin{bmatrix}
  \nu_0 & J_{12} & J_{13} & \ldots & J_{1n} \\
  J_{21} & \nu_0 & J_{23} & \ldots & J_{2n} \\
  J_{31} & J_{32} & \nu_0 & \ldots & J_{3n} \\
  \vdots & \vdots & \vdots & \ddots & \vdots \\
  J_{n1} & J_{n2} & J_{n3} & \ldots & \nu_0
\end{bmatrix}
\]

(1)

where \(\nu_0\) is the frequency of an unperturbed \(N=\text{N}\)-stretch mode, \(J_{ij}\) the coupling between \(N=\text{N}\)-stretch modes \(i\) and \(j\). The Hamiltonian \(H\) is a symmetric (Hermitian and real) matrix so that \(J_{ij} = J_{ji}\). The interaction between the \(N=\text{N}\) groups is described by transition-dipole coupling:

\[
J_{ij} = \frac{1}{4\pi\varepsilon_0} \left[ \frac{\vec{\mu}_i \cdot \vec{r}_{ij}}{|\vec{r}_{ij}|^3} - \frac{3(\vec{\mu}_i \cdot \vec{\mu}_j)(\vec{\mu}_j \cdot \vec{r}_{ij})}{|\vec{r}_{ij}|^5} \right],
\]

(2)

where \(\vec{r}_{ij}\) is the distance vector between dipoles \(\vec{\mu}_i\) and \(\vec{\mu}_j\) (Fig. 4). Note that the meta-linkages between the individual azobenzene repeated units prevent any through-bond coupling due to cross-conjugation.

The dipole–dipole interaction lifts the degeneracy of the \(N=\text{N}\)-stretching frequencies, and gives rise to as many delocalized normal modes as there are interacting \(N=\text{N}\) oscillators. The vibrational frequencies of these delocalized normal modes (excitons) are the eigenvalues of the excitonic Hamiltonian matrix \(H\). The dipole \((D)\) and rotational \((R)\) strengths of the \(k\)th exciton are given by

\[
D_k = \sum_{i=1}^{N} \sum_{j=1}^{N} b_{ik} b_{jk} (\vec{\mu}_i \cdot \vec{\mu}_j)
\]

(3)

and

\[
R_k = -(\pi\nu_0/c) \sum_{i=1}^{N} \sum_{j=1}^{N} b_{ik} b_{jk} |\vec{r}_{ij}| \vec{\mu}_i \times \vec{\mu}_j.
\]

(4)

where \(c\) is the velocity of light in vacuum, and \(b_{ij}\) are the components of the eigenvectors of the Hamiltonian. The infrared absorption and vibrational circular dichroism spectra are calculated using eqns (3) and (4), respectively, and convoluted with a Lorentzian band profile.

4.2 147 azo-foldamer

In Fig. 5a and b (solid lines) we show experimental FTIR and VCD spectra of the 147 azo-foldamer, respectively, in the \(N=\text{N}\)-stretch mode. In Fig. 5c and d (dashed lines) we show the calculated infrared absorption and vibrational circular dichroism spectra of the 147 azo-foldamer, respectively, in the \(N=\text{N}\)-stretch mode. The observed spectra are fitted with the calculated spectra to determine the vibrational frequencies and intensities. The observed and calculated spectra show good agreement, indicating that the model is accurate.
spectral region. Using the experimental infrared absorption spectrum (Fig. 5a, solid line), we determine the magnitude of $\vec{\mu}$ from the integrated absorption coefficient $A = \int \epsilon(\nu) d\nu$. This value is used to calculate the coupling terms in eqn (2) and to construct the Hamiltonian with solely $v_0$ as a free parameter. Fig. 5c and d show simulated infrared absorption and VCD spectra (with $|\vec{\mu}| = 0.12$ Debye), respectively, for the N=N-stretch modes calculated using the coupled-oscillator formalism described above. The experimental VCD pattern, which exhibits a characteristic bisignate signal (positive couplet), is nicely reproduced in the simulated spectrum. The three most IR-intense calculated eigenmodes have vibrational frequencies of 1452, 1451 and 1450 cm$^{-1}$. The remaining four predicted eigenmodes have extremely weak intensities (less than 5% of the total intensity) and are therefore neglected in the following analysis.

Fig. 6 displays the three most intense normal modes as color coding on schematic representations of the 14-foldamers backbone structure. The color code is used to illustrate the relative amplitudes (eigenvector coefficients $b_{ij}$) of the N=N oscillators in the eigenmodes that constitute the VCD spectrum (Fig. 5d). Eigenmode I (Fig. 6-I), which gives rise to the negative VCD peak ($\approx 1452$ cm$^{-1}$), has its origin in a strong coupling (with negative amplitudes) of the terminal and central N=N units when compared with that of the 14-foldamer (with 7 N=N units). Moreover, in Fig. 5d (dashed blue curve) we show the simulated VCD spectrum of a 10-foldamer obtained with $|\vec{\mu}| = 0.12$ Debye as an input parameter in the model. The signal magnitudes are extremely weak compared with those of the 14-foldamer, which confirms our conclusion that the VCD arises mainly from the triple-stacked N=N bonds.

Clearly, the coupled-oscillator formalism applied to the backbone N=N units successfully describes the observed features in the VCD spectrum of the 14-azo-foldamer. The good agreement of simulated and observed spectral features confirms the conclusions from the previously reported ECD spectra regarding the native-folded conformation of the azo-foldamer backbone.7–9

4.3 Photo-induced 14, unfolding

To determine the effect of azobenzene $E \rightarrow Z$ photoisomerization on the VCD response we irradiated a 10$^{-4}$ M solution of 14, in CD$_3$CN at $\lambda = 358$ nm using a 500 W Xe lamp equipped with a notch filter for a period of 24 h to assure reaching the photosationary state (PSS). Upon irradiation, the 14-azo-foldamer undergoes a helix-coil transition, leading to a PSS in which a fraction of the N=N units is in the Z-conformation, leading to partial denaturation of the folded helical state. Previous studies have shown that in the PSS approximately 40% of the N=N units that are in the Z-form belong to the termini.7 The IR and VCD spectra of the PSS are shown in Fig. 5a and b (dashed lines). While the IR spectrum shows only a slight red-shift, the VCD coupllet is reduced in intensity by approximately 50%. This decrease in signal magnitude corroborates the previous conclusions based on electronic circular dichroism experiments;7 where the decrease in signal intensity was assigned to the loss of excess helicity of the azo-foldamer. The observed decrease in intensity of the VCD bisignate signal (Fig. 5b) is consistent with the predicted role of eigenmodes I and II as local probes in the coupling of the N=N termini and core units, which is partially disrupted upon photoisomerization. Moreover, the partial denaturation of the foldamer, i.e., the reorientation of the termini and subsequent loss of helical content, effectively leads to the
formation of a 10<sub>-</sub>-type foldamer, in which the backbone still adopts a helical structure, but now with fewer transoid backbone N–N units. This observation is in agreement with the unfolding mechanism proposed in which unfolding occurs predominantly from the termini of the foldamer.  

5 Conclusion

Our results demonstrate that VCD can be used to probe conformational changes during the folding and unfolding of azobenzene foldamers. The fact that the measurements can be interpreted using a simple coupled-oscillator approach ensures a straightforward relationship between observed VCD spectra and the foldamer conformation. Much more detailed information can be obtained from VCD measurements on foldamers in which specific N=N bonds have been isotope-labeled. These labeled N=N bonds will have a different local-mode stretching frequency, and will be spectrally isolated from the other N=N bonds. They can therefore be described by a smaller (ideally 2 × 2, if two N=N bonds are labeled) Hamiltonian. Such experiments should make it possible to probe conformational changes at specific sites in the foldamer backbone, and the interpretation should be just as straightforward as for the unlabeled foldamers studied here.

Acknowledgements

The authors gratefully acknowledge Michiel F. Hilbers, Hans J. Sanders, Jan-Hein Hooijshuur and Freek Ariese for technical support. We would also like to thank Wybren Jan Buma for critically reading the manuscript. S.R.D. acknowledges financial support from the Portuguese Foundation for Science and Technology (FCT) under the fellowship SFRH/BD/48295/2008. S.H. would like to thank the Deutsche Akademischer Bund der Naturforscher Leopoldina – German National Academy of Sciences for a Leopoldina research fellowship (grant number LPDS 2011-18). S.H. gratefully acknowledges generous support from the German Research Foundation (DFG via SFB 658), the European Research Council (ERC) via ERC-2012-STG 308117 (Light4Function), and the European Science Foundation (ESF via P2M). S.W. acknowledges the European Research Council (ERC) for funding through Grant No. 210999.

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