Fine structure of the Bethe Ansatz for the spin-1/2 Heisenberg XXX model

Schoutens, C.J.M.; Essler, F.; Korepin, V.

Published in:
Journal of Physics. A, Mathematical and General

DOI:
10.1088/0305-4470/25/15/019

Citation for published version (APA):

General rights
It is not permitted to download or to forward/distribute the text or part of it without the consent of the author(s) and/or copyright holder(s), other than for strictly personal, individual use, unless the work is under an open content license (like Creative Commons).

Disclaimer/Complaints regulations
If you believe that digital publication of certain material infringes any of your rights or (privacy) interests, please let the Library know, stating your reasons. In case of a legitimate complaint, the Library will make the material inaccessible and/or remove it from the website. Please Ask the Library: http://uba.uva.nl/en/contact, or a letter to: Library of the University of Amsterdam, Secretariat, Singel 425, 1012 WP Amsterdam, The Netherlands. You will be contacted as soon as possible.
Fine structure of the Bethe ansatz for the spin-$\frac{1}{2}$ Heisenberg $XXX$ model

Fabian H L Essler, Vladimir E Korepin and Kareljan Schoutens
Institute for Theoretical Physics, State University of New York at Stony Brook, Stony Brook, NY 11794-3840, USA

Received 18 November 1991, in final form 20 March 1992

Abstract. We analyse the Bethe ansatz equations for the two-particle sector of the spin-$\frac{1}{2}$ Heisenberg $XXX$ model on a one-dimensional lattice of length $N$. We show that, beginning at a critical lattice length of $N = 21.86$, new pairs of real solutions develop, whereas complex solutions start to disappear. The integers (that appear in the logarithmic form of the Bethe equations) of the new solutions do not fit into the conventional classification scheme. The total number of solutions in the two-particle sector remains unchanged and is in agreement with the claim that the SU(2) extended Bethe ansatz gives a complete set of $2^N$ eigenstates.

1. Introduction

The one-dimensional isotropic spin-$\frac{1}{2}$ Heisenberg magnet ($XXX$ spin chain) was the first model solved by means of the Bethe ansatz method [1]. The model describes interacting spins, situated on the sites of a periodic lattice of length $N$. The Hamiltonian is given by

$$H = J \sum_{i=1}^{N} (S_i^x S_{i+1}^x + S_i^y S_{i+1}^y + S_i^z S_{i+1}^z - \frac{1}{4})$$

$$S_{N+1} = S_1$$

where $2S_i$ are the Pauli matrices. The Bethe ansatz provides the following set of eigenfunctions

$$|\Psi\rangle = \sum_{x_1 < x_2 < \ldots < x_M} \Psi(x_1, x_2, \ldots, x_M) S_{x_1}^x S_{x_2}^x \ldots S_{x_M}^x |0\rangle$$

$$\Psi(x_1, x_2, \ldots, x_M) = \sum_{P \in S_M} \exp \left( i \sum_{j=1}^{M} k_{P_j} x_j + i \sum_{\substack{j \neq i \quad \text{or} \quad P_j > P_i}} \phi_{P_j, P_i} \right)$$

Here $|0\rangle$ is the ferromagnetic state with all spins up and $M$ is the number of overturned spins. The wavefunction is symmetric under interchange of any two of the
coordinates \( x_j \) and \( k_1, \ldots, k_M \in [0, 2\pi] \) are the spectral parameters, \( P \) is a per-
mutation of \( M \) elements, and \( \phi_{j, l} \) is defined by

\[
2 \cot \left( \frac{\phi_{j, l}}{2} \right) = \cot \left( \frac{k_j}{2} \right) - \cot \left( \frac{k_l}{2} \right).
\]  

(3)

Imposing the periodic boundary conditions

\[
\Psi(x_1, x_2, \ldots, x_{M-1}, N+1) = \Psi(1, x_1, x_2, \ldots, x_{M-1})
\]  

(4)

leads to the following set of Bethe equations for the spectral parameters \( k_1, \ldots, k_M \)

\[
e^{ik_jN} = \prod_{j=1}^{M} e^{i\phi_{j,l}} \quad j = 1, 2, \ldots, M.
\]  

(5)

The usual parametrization for the Bethe equations and eigenfunctions is obtained by

the substitution

\[
\Lambda_j = \cot(k_j/2).
\]  

(6)

In this parametrization we find that

\[
\Psi(x_1, x_2, \ldots, x_M) = \sum_{P \in \mathcal{S}_M} \prod_{j=1}^{M} \left( \kappa(\Lambda_{P_j}) \right)^{x_j} \prod_{i<j, P_i>P_j} \kappa \left( \frac{\Lambda_{P_i} - \Lambda_{P_j}}{2} \right)
\]  

(7)

and the Bethe equations take the form

\[
(\kappa(\Lambda_j))^N = \prod_{j=1}^{M} \kappa \left( \frac{\Lambda_j - \Lambda_l}{2} \right) \quad j = 1, 2, \ldots, M
\]  

(8)

\[
\kappa(\Lambda) = \frac{\Lambda + i}{\Lambda - i}.
\]

The energy of the state (7) is given by

\[
E(\Lambda_1, \ldots, \Lambda_M) = \sum_{j=1}^{M} \frac{-2J}{\Lambda_j^2 + 1}.
\]  

(9)

One fundamental assumption in investigating solutions of (8) is the so-called string
hypothesis [1-3], which states that in the large \( \mathcal{N} \) limit (for fixed \( M \)), any solution
\( \Lambda_1, \ldots, \Lambda_M \) of (8) consists of a set of strings of the form

\[
\Lambda_{\alpha,j} = \Lambda_{\alpha} + i(n + 1 - 2j) + O(\epsilon^{-\delta}) \quad j = 1, 2, \ldots, n
\]  

(10)

where \( n \geq 1 \) gives the length of the string, \( \alpha \) labels different strings of a given length,
\( j \) specifies the imaginary part of \( \Lambda \) (\( \Lambda_{\alpha} \) denotes the real part of \( \Lambda_{\alpha,j} \)) and \( \delta > 0 \).

The total number of strings of length \( n \) is denoted by \( M_n \). Assuming the validity of
this hypothesis, it is possible to reduce the Bethe equations (8) to a set of coupled
Fine structure of the Bethe ansatz for the spin-\(\frac{1}{2}\) Heisenberg 4117

equations involving only the real parts \(\Lambda^n_\alpha\). The logarithmic form of these equations is given by [2, 4]

\[
N \theta \left( \frac{\Lambda^n_\alpha}{n} \right) = 2\pi I^n_\alpha + \sum_{(m, \beta) \neq (n, \alpha)} \theta_{n,m}(\Lambda^n_\alpha - \Lambda^m_\beta)
\]  \hspace{1cm} (11)

where

\[
\theta(x) = 2 \arctan(x)
\]  \hspace{1cm} (12)

and

\[
\theta_{n,m}(x) = \theta \left( \frac{x}{|n - m|} \right) + 2\theta \left( \frac{x}{|n - m| + 2} \right) + \cdots
\]

\[
+ 2\theta \left( \frac{x}{n + m - 2} \right) + \theta \left( \frac{x}{n + m} \right) \hspace{1cm} \text{if } n \neq m
\]

\[
\theta_{n,m}(x) = 2\theta \left( \frac{x}{2} \right) + 2\theta \left( \frac{x}{4} \right) + \cdots
\]

\[
+ 2\theta \left( \frac{x}{2n - 2} \right) + \theta \left( \frac{x}{2n} \right) \hspace{1cm} \text{if } n = m.
\]  \hspace{1cm} (13)

Solutions of (11) are parametrized by integer (for \(N - M^n_n\) odd) or half odd integer (for \(N - M^n_n\) even) numbers \(I^n_\alpha\), which are distributed symmetrically around zero and should satisfy

\[
|I^n_\alpha| \leq \frac{1}{2} \left( N - \sum_{m=1}^{\infty} t_{n,m} M_m - 1 \right)
\]  \hspace{1cm} (14)

where \(t_{n,m} = 2 \text{Min}(m, n) - \delta_{m,n}\). It is generally believed that there is a one-to-one correspondence between solutions of the Bethe equations and sets of independent, non-repeating integers \(I^n_\alpha\). (By this we mean that \(I^n_\alpha \neq I^n_\beta\) for \(\alpha \neq \beta\), such that no two strings of the same length \(n\) have the same integer.) This is essential for all proofs of completeness of the (SU(2)-extended) Bethe ansatz solutions for the Heisenberg XXX model [1-4].

In this paper we investigate the validity of these assumptions in the two-particle sector. We find that counting the solutions of (8) using the integer method (based on the assumption that all strings are of the form (10)) [2-4] yields an incorrect result for lattice lengths \(N\) greater than \(N_{\text{crit}} = 21.86\). For \(N > N_{\text{crit}}\) there exist less complex string solutions than those predicted by counting the integers according to (14). These 'missing' strings are replaced by additional real solutions, which can be characterized by a pair of repeating integers. The total number of solutions thus remains unchanged, whereas their nature changes drastically. The change in type of solution to the Bethe equations had already been noticed by Bethe himself in [1] but for some reason has been overlooked in the more recent literature.

In section 2 we will give a graphical method for finding all complex two-particle solutions of the Bethe equations (8) and in section 3 we will apply a similar method to determine all real solutions in the two-particle sector. We will find a new type of real solution that is not present in the classification scheme based on (14) (and the assumption of non-repeating integers). In section 4 we classify the new type of real solution in the large \(N\) limit and show that the number of new solutions is of order \(\sqrt{N}\).
2. String solutions of the Bethe equations in the two-particle sector

Due to the symmetry of (8) under complex conjugation, complex solutions \( \Lambda_1, \Lambda_2 \) must be of the form \( \Lambda_1 = (\Lambda_2)^* \), where * denotes complex conjugation. This enables us to reduce the Bethe equation (8) to a single equation in the variables \( x \) and \( y \) with \( \Lambda_1 = x + iy \):

\[
\left( \frac{x + i(y + 1)}{x + i(y - 1)} \right)^N = \frac{y + 1}{y - 1}.
\]

Dividing (15) by its conjugate equation we obtain

\[
\left( \kappa \left( \frac{x}{y + 1} \right) \right)^N \left( \kappa \left( \frac{x}{1 - y} \right) \right)^N = 1.
\]

After taking the logarithm this yields the analogue to (11) for arbitrary string width \( y \):

\[
N \left[ \theta \left( \frac{x}{y + 1} \right) + \theta \left( \frac{y - 1}{x} \right) \right] = 2\pi I_a^2.
\]

Clearly this reduces to (11) for \( y = 1 \). The cut of \( \tan^{-1}(z) \) has been chosen from \( i \) to \( i\infty \) and from \( -i \) to \( -i\infty \) along the imaginary axis. Note that \( I_a^2 \) is integer (half odd integer) for even (odd) lattice length \( N \). Taking the magnitude of (15) we arrive at

\[
\left( \frac{x^2 + (y + 1)^2}{x^2 + (y - 1)^2} \right)^N = \left( \frac{y + 1}{y - 1} \right)^2.
\]

The curve \( B_N \), that solves (18) has two branches as shown in figure 1 and includes the whole \( x \)-axis \( (y = 0) \) as well. The 'inner' branch intersects the \( x \)-axis for \( N \gg 1 \) at

\[
x^{(1)} = \pm \sqrt{(N - 1)}.
\]

The 'outer' branch has the asymptotes [5]

\[
y = \pm \frac{1}{\sqrt{(N - 1)}} x.
\]

Figure 1. The curve \( B_N \) for \( N = 5 \).
To obtain solutions of (15) we also have to discuss the phases. Noting that the right-hand side of (15) is real we can write

$$\left( \frac{x+i(y+1)}{x+i(y-1)} \right)^N = \frac{y+1}{y-1} = \pm r^N \tag{21}$$

where \( r \in [0, \infty) \). Taking the \( N \)th root of this equation we arrive at

$$\left( \frac{x+i(y+1)}{x+i(y-1)} \right) = \left( \frac{y+1}{y-1} \right)^{1/N} r \quad m = 0, 1, \ldots, 2N-1. \tag{22}$$

We denote the solutions of (22) without restriction on \( r \) for given \( m \) by \( C_m \). Looking at (21) we see that there is a constraint on the allowed range of \( r \) in (22). Even (odd) \( m \) correspond to the + (−) sign on the right-hand side of (21) and thus to \(|y| > 1 \) (\(|y| < 1 \)). This constraint selects a certain part \( K_m \) of the curve \( C_m \) and only points \((x, y)\) that lie on \( K_m \) are solutions of (21).

A parametrization of \( C_m \) is given by

$$x(r) = \frac{-2r \sin(\pi m/N)}{1 + r^2 - 2r \cos(\pi m/N)} \quad y(r) = \frac{1 - r^2}{1 + r^2 - 2r \cos(\pi m/N)} \tag{23}$$

where \( r \in [0, \infty) \). The constraint mentioned earlier has to be imposed on these solutions. The curves \( C_m \) intersect the \( x \)-axis at \((r = 1)\)

$$x_m^{(2)} = \pm \cot(\pi m/2N). \tag{24}$$

Solutions of the Bethe equations (15) are given by intersection points of the curves \( K_m \) with the curves \( B_N \). In figures 3 and 4 we show the resulting curves for lattice lengths \( N = 5 \) and \( N = 21 \). Inspection of figure 3 shows that there exist two string solutions on the lattice with five sites. This is in agreement with the counting according to (14). Numerically one finds two string solutions with spectral parameters

$$\Lambda_1^{2,1} = -0.63 + 1.04i = -\Lambda_2^{2,2}$$

$$\Lambda_1^{2,2} = -0.63 - 1.04i = -\Lambda_2^{2,1} \tag{25}$$
which is in perfect agreement with figure 3 and the standard string picture. We see that for even and sufficiently small \( N \) every curve \( K_m \) with \( 1 < m < N/2 \) and \( 3N/2 < m < 2N - 1 \) has exactly two intersection points with \( B_N \), which leads to one string solution for any such \( m \) and, in addition, we have one string solution at zero \((x, y) = (0, \pm 1)\), which can be shown to correspond to the wavefunction

\[
|\psi\rangle = \sum_{y=1}^{N} (-1)^y S_y^- S_{y+1}^- |0\rangle.
\]

(Clearly this state only makes sense on a lattice with an even number of sites.) This gives a total of \( N - 3 \) solutions.

For odd and small enough \( N \) every curve \( K_m \) with \( 1 < m < (N + 1)/2 \) and with \( (3N - 1)/2 < m < 2N - 1 \) gives one solution which also adds up to a total of \( N - 3 \).

Using the standard string counting (14) we find that \( |I_2^3| \leq (N - 4)/2 \) which leads to a total of \( N - 3 \) solutions for \( N \) either even or odd. Thus one would think that (14) indeed predicts the correct number of solutions to the Bethe equations. This is not true, however, as can be seen in figure 5 which shows all string solutions for \( N = 25 \). Clearly the curves \( K_3 \) and \( K_{47} \) do not lead to any intersection points and we are thus left with a total of \( 25 - 5 = 20 \) solutions, two less than predicted by (14). The situation for 21 sites (see figure 4) is different, because the curves \( K_3 \) and \( K_{39} \) still intersect with the curve \( B_{21} \) (although the width of the corresponding strings is visibly much smaller than the 'ideal' value of 2).

Thus the following picture emerges: The strings corresponding to intersection points of \( K_3 \) and \( K_{2N-3} \) with \( B_N \) get narrower and narrower until they disappear for \( N \geq 22 \). We find that for \( 22 \leq N \leq 61 \) there are only \( N - 5 \) string solutions to the Bethe equations. Starting at \( N = 62 \) we 'lose' an additonal two solutions. Using (19) and (24) we see that whenever \( |x_m^{(2)}| \geq |x^{(1)}| \) for any odd \( m \) with \( 3 \leq m \leq \lfloor N/2 \rfloor \) (where \( \lfloor z \rfloor \) denotes the integer part of \( z \)), we lose two string solutions of the Bethe equations, i.e. we find (an additional) two fewer solutions than predicted by (14).

Figure 3. String solutions for lattice length \( N = 5 \).
3. Real solutions of the Bethe equations in the two-particle sector

The Bethe equations (5) in the two-particle sector are of the form

\[ e^{i k_2 N} = e^{i \phi_{2,1}} \]

and

\[ e^{i (k_1 + k_2) N} = 1 \]  \hspace{1cm} (27)

where \( \phi_{1,2} \) is given by (3) and the second equality is obtained by multiplying the equations for \( k_1 \) and \( k_2 \). Taking the logarithm and solving for \( k_1 \) one arrives at

\[ k_1 = 2 \cot^{-1} (\cot(k_2/2) - 2 \cot(k_2 N/2)) \]

\[ k_1 = -k_2 + (2 \pi l / N) \mod 2 \pi \hspace{1cm} l = 1, \ldots, N. \]  \hspace{1cm} (28)

Plotting \( k_1 \) as a function of \( k_2 \) and taking into account that the \( k_s \) range between 0 and \( 2 \pi \), one arrives at the graphs shown in figures 6 and 7 for \( N = 4 \).
The real solutions of the periodic boundary conditions for two particles are given by the intersection points of the graphs in figures 6 and 7. It is a well known fact that the wavefunction (2) vanishes if any two spectral parameters coincide (‘Pauli principle’). Therefore we can ignore intersection points on the line $k_1 = k_2$, and, due to the symmetry of the problem under interchange of $k_1$ and $k_2$, restrict our attention to the region $k_2 > k_1$. The requirement of all $\Lambda s$ to be finite [2, 3] leads to the exclusion of the axis $k_1 = 0$. Figure 8 shows the resulting graph for $N = 4$, and we can clearly identify one solution, as predicted by the integer method ((14) gives $|J_\alpha| < \frac{1}{2}(4 - 2 - 1) = \frac{1}{2}$ and thus there are the two permitted values $\pm 1$ for the two integers $I_{1,2}$, thus one solution).

The qualitative behaviour of the graphs remains unchanged for $N \leq 21$ and the integer method predicts the right number, which is $\binom{N-2}{2}$, of solutions. A close investigation of the situation for a larger lattice length like $N = 25$ reveals an interesting development, however. The dotted line in figure 9 is the line $k_1 = k_2$. As mentioned earlier there is a clear symmetry in the graph with respect to the interchange of $k_1$ and $k_2$. Due to this symmetry and the ‘Pauli principle’ (the vanishing of the wavefunction for $k_1 = k_2$) we only have to consider the region below the dotted line in figure 9. As can be seen in figure 9, the line with $l = 3$ of (28) approaches the $\cot^{-1}$ curve rather closely in the regions $k_1 \approx k_2 \approx 0.4$ and $k_1 \approx k_2 \approx 5.9$. 

Figure 6. Plot of $k_1 = -k_2 + 2\pi l/N \mod 2\pi$ against $k_2$ for $N = 4$.

Figure 7. Plot of $k_1 = 2\cot^{-1}(\cot(k_2/2) - 2\cot(k_2N/2))$ against $k_2$ for $N = 4$.

Figure 8. Real solutions for $N = 4$.

Figure 9. Real solutions for $N = 25$. 
Magnification of these regions reveals that instead of one intersection point with \( k_1 = k_2 \) (as was the case for \( N \leq 21 \)), there now exist two additional intersection points (related by the interchange \( k_1 \leftrightarrow k_2 \), so that we must count only one of them) in each region (see figure 10), so that we obtain a total of 255 real solutions. The integer method predicts two real solutions less, as (14) gives \( |I^2_1| \leq \frac{1}{2} (25-2-1) = 11 \) and thus there are 23 permitted values for \( I^1_{1,2} \), which leads to \((23)^2 = 253 \) solutions.

A numerical solution of the Bethe equations (8) for \( N = 25 \) (see appendix I) is in complete agreement with the situation described here and the new solutions are indeed found in the regions predicted by the graphical method.

The critical lattice length \( N_{\text{crit}} \), for which new solutions start to appear is found by equating the first derivatives of the two intersecting curves. The numerical value is found to be \( N_{\text{crit}} = 21.86 \). From figure 9 it becomes clear that for large lattice lengths \( N \) further additional solutions will develop, as the situation that occurs for the second 'branch' of the \( \cot^{-1} \) curve and the line with \( l = 3 \) at \( N_{\text{crit}} = 21.86 \), will eventually have an analogue in the third, fourth, etc., 'branch' of the \( \cot^{-1} \) curve. Numerical studies reveal that the next two additional real solutions appear for \( N = 61.34 \) (the line with \( l = 5 \) intersecting the third branch of the \( \cot^{-1} \) curve).

4. Large-\( N \) analysis

The new real solutions always develop in the vicinity of the curve \( k_1 = k_2 \). Therefore one can solve the equations (28) perturbatively in powers of \( 1/N \) around this line. One finds that for a given \( N \)

\[
\begin{align*}
k_1 &= \frac{2\pi(n-1)}{N} \left[ 1 + \frac{2n}{N} + \frac{4n \left( n^2 + (n-1)^2 \right)}{N^2} + O \left( N^{-3} \right) \right] \\
k_2 &= \frac{2\pi n}{N} \left[ 1 - \frac{2(n-1)}{N} + \frac{4(n-1) \left( n^2 + (n-1)^2 \right)}{N^2} + O \left( N^{-3} \right) \right]
\end{align*}
\]

\( n = 2, 3, \ldots, n_{\text{max}} = \left[ \frac{\sqrt{N}}{\pi} + \frac{1}{2} \right] \)

where \([x]\) denotes the integer part of \( x \). This means that there exist \( n_{\text{max}} - 1 \) pairs of new real solutions for a given \( N \). The 'critical' lattice length, at which new real
solutions of the form (29), with parameter \( n \), first develop is given by the solution of the equation

\[
N_{\text{crit}}^{(n)} \sin^2 \left( \frac{\pi (n - \frac{1}{2})}{N_{\text{crit}}} \right) = 1.
\]  

(30)

Perturbatively one finds that

\[
N_{\text{crit}}^{(n)} = 2(n - \frac{1}{2})^2.
\]  

(31)

The results for \( n = 2, 3 \) are in very good agreement with the numerically found values given earlier. Using (6), (11) and (29) we can compute the integers for the new real solutions. Elementary calculations yield for the solution with parameter \( n \)

\[
I_1^1 = I_2^1 = \frac{N + 1}{2} - n.
\]  

(32)

Thus the new real solutions are characterized by a pair of repeating integer numbers. This is a new phenomenon, which contradicts the usual assumption that all integers are independent and non-repeating.

5. Conclusions

In this paper we have shown that the total number of solutions of the Bethe equations in the two-particle sector is in agreement with the result obtained by counting integers according to (14) as carried out in [2-4]. We saw, however, that the integer method fails to predict the correct distribution of solutions among real and complex solutions. Furthermore we demonstrated that the 'Pauli principle for the integers', i.e. the assumption that all solutions of the Bethe equations can be characterized by a set of independent non-repeating (half-odd) integer numbers, is invalid. The new states we found are described by two identical integers and the physical effects of such excitations above the ferromagnetic vacuum can be studied along the lines described in [6].

Acknowledgments

We thank S Dasmahapatra for discussions and B Sauer for plotting figure 9. This work was supported in part by the National Science Foundation under research grants PHY91-07261 and NSF91-08054.

Appendix. Some numerical results for \( N = 25 \)

In this appendix we give parts of the numerical solution of the Bethe equations in the two-particle sector. Table 1 lists the two additional real solutions as well as their integers, which are found to be repeating. Table 3 lists all (complex) string solutions together with their integers. We find that the solutions with 'integer' \( \pm 9.5 \) are missing. Tables 2 and 4 give samples of the numerical results for the 'regular' real solutions and the corresponding integers.
Fine structure of the Bethe ansatz for the spin-$\frac{1}{2}$ Heisenberg

Table 1. Extra real solutions.

<table>
<thead>
<tr>
<th>$\Lambda_1$</th>
<th>$\Lambda_2$</th>
<th>$(I_1^1, I_1^2)$</th>
</tr>
</thead>
<tbody>
<tr>
<td>6.076310351764672</td>
<td>4.604119967471151</td>
<td>(11,11)</td>
</tr>
<tr>
<td>-6.076310351764672</td>
<td>-4.604119967471151</td>
<td>(-11,-11)</td>
</tr>
</tbody>
</table>

Table 2. Real solutions with integers of the form $(-I_0^1, I_0^2)$.

<table>
<thead>
<tr>
<th>$\Lambda_1$</th>
<th>$\Lambda_2$</th>
<th>$(I_1^1, I_1^2)$</th>
</tr>
</thead>
<tbody>
<tr>
<td>-0.131652483159000</td>
<td>0.131652483159000</td>
<td>(-0.9999999,0.9999999)</td>
</tr>
<tr>
<td>-0.267949207552647</td>
<td>0.267949207552647</td>
<td>(-2.0000000,2.0000000)</td>
</tr>
<tr>
<td>-0.414213554923667</td>
<td>0.414213554923667</td>
<td>(-3.3,3)</td>
</tr>
<tr>
<td>-0.577350271101900</td>
<td>0.577350271101900</td>
<td>(-4.4,4)</td>
</tr>
<tr>
<td>-0.767326987742179</td>
<td>0.767326987742179</td>
<td>(-5.5,5)</td>
</tr>
<tr>
<td>-1.000000000101941</td>
<td>1.000000000101941</td>
<td>(-6.6,6)</td>
</tr>
<tr>
<td>-1.303225372381047</td>
<td>1.303225372381047</td>
<td>(-7.7,7)</td>
</tr>
<tr>
<td>-1.732050807568888</td>
<td>1.732050807568888</td>
<td>(-8.8,8)</td>
</tr>
<tr>
<td>-2.414213562373096</td>
<td>2.414213562373096</td>
<td>(-9.9,9)</td>
</tr>
<tr>
<td>-3.73205080756863</td>
<td>3.73205080756863</td>
<td>(-9.9,9)</td>
</tr>
<tr>
<td>-7.595754112725638</td>
<td>7.595754112725638</td>
<td>(-11,11)</td>
</tr>
</tbody>
</table>

Table 3. Spectral parameters and integers of the string solutions.

<table>
<thead>
<tr>
<th>$\Lambda^{a,b}$</th>
<th>$J_0^a$</th>
</tr>
</thead>
<tbody>
<tr>
<td>7.550697954902596 ± 1.674220479211304</td>
<td>10.5</td>
</tr>
<tr>
<td>3.603532496872190 ± 1.060262279024803</td>
<td>8.5</td>
</tr>
<tr>
<td>2.76045192520186 ± 0.989331795228972</td>
<td>7.5</td>
</tr>
<tr>
<td>2.129094407420760 ± 1.000733494880967</td>
<td>6.5</td>
</tr>
<tr>
<td>1.654575100754238 ± 0.999974181007365</td>
<td>5.500000000000001</td>
</tr>
<tr>
<td>1.26923865274572 ± 1.000000336229817</td>
<td>4.500000000000002</td>
</tr>
<tr>
<td>0.941128564704973 ± 0.9999998926892</td>
<td>3.49999999999998</td>
</tr>
<tr>
<td>0.649839392364719 ± 1.00000000000355</td>
<td>2.49999999999996</td>
</tr>
<tr>
<td>0.381520404437133 ± 1.000000000000000</td>
<td>1.49999999999998</td>
</tr>
<tr>
<td>0.1258293800000</td>
<td>1.000000000000000</td>
</tr>
<tr>
<td>-0.1258293800000</td>
<td>1.000000000000000</td>
</tr>
<tr>
<td>-0.381520404437133 ± 1.000000000000000</td>
<td>1.500000000000017</td>
</tr>
<tr>
<td>-0.649839392364719 ± 1.000000000000000</td>
<td>-2.49999999999998</td>
</tr>
<tr>
<td>-0.941128564704973 ± 0.9999998926892</td>
<td>-3.49999999999998</td>
</tr>
<tr>
<td>-1.26923865274572 ± 1.000000336229817</td>
<td>-4.500000000000002</td>
</tr>
<tr>
<td>-1.654575100754238 ± 0.999974181007365</td>
<td>-5.500000000000001</td>
</tr>
<tr>
<td>-2.129094407420760 ± 1.000733494880967</td>
<td>-6.5</td>
</tr>
<tr>
<td>-2.76045192520186 ± 0.989331795228972</td>
<td>-7.5</td>
</tr>
<tr>
<td>-3.603532496872190 ± 1.060262279024803</td>
<td>-8.5</td>
</tr>
<tr>
<td>-7.550697954902596 ± 1.674220479211304</td>
<td>-10.5</td>
</tr>
</tbody>
</table>
Table 4. Real solutions with one integer (but not both) equal to 11.

<table>
<thead>
<tr>
<th>$\Lambda_1^1$</th>
<th>$\Lambda_2^1$</th>
<th>$(I_1^1, I_2^1)$</th>
</tr>
</thead>
<tbody>
<tr>
<td>2.660818949240510</td>
<td>6.937316937832509</td>
<td>(10, 11)</td>
</tr>
<tr>
<td>1.884097488224177</td>
<td>7.077265200079445</td>
<td>(9, 11)</td>
</tr>
<tr>
<td>1.416202427893359</td>
<td>7.14242244443372</td>
<td>(8, 11)</td>
</tr>
<tr>
<td>1.092353837732522</td>
<td>7.181455030337195</td>
<td>(7, 11)</td>
</tr>
<tr>
<td>0.84800820447725</td>
<td>7.20823937845300</td>
<td>(6, 11)</td>
</tr>
<tr>
<td>0.651308925902454</td>
<td>7.228334310008103</td>
<td>(5, 11)</td>
</tr>
<tr>
<td>0.484695845135505</td>
<td>7.24430628292553</td>
<td>(4, 11)</td>
</tr>
<tr>
<td>0.337407004916716</td>
<td>7.258009100128192</td>
<td>(3, 11)</td>
</tr>
<tr>
<td>0.202216795753967</td>
<td>7.269970770157027</td>
<td>(2, 11)</td>
</tr>
<tr>
<td>0.073793040214152</td>
<td>7.280916005878445</td>
<td>(1, 11)</td>
</tr>
<tr>
<td>-0.052243320518647</td>
<td>7.291284166897977</td>
<td>(0, 11)</td>
</tr>
<tr>
<td>-7.595754112725151</td>
<td>7.595754112725151</td>
<td>(-11, 11)</td>
</tr>
<tr>
<td>-3.784235077366884</td>
<td>7.492364916365747</td>
<td>(-10, 11)</td>
</tr>
<tr>
<td>-2.46807197228016</td>
<td>7.438644066117615</td>
<td>(-9, 11)</td>
</tr>
<tr>
<td>-1.782616957891460</td>
<td>7.404677027965097</td>
<td>(-8, 11)</td>
</tr>
<tr>
<td>-1.350629918256678</td>
<td>7.380566510373989</td>
<td>(-7, 11)</td>
</tr>
<tr>
<td>-1.045151822554743</td>
<td>7.362036505894714</td>
<td>(-6, 11)</td>
</tr>
<tr>
<td>-0.811201425781150</td>
<td>7.346914476556023</td>
<td>(-5, 11)</td>
</tr>
<tr>
<td>-0.620884146211136</td>
<td>7.333960813623058</td>
<td>(-4, 11)</td>
</tr>
<tr>
<td>-0.458329409816131</td>
<td>7.322397412814130</td>
<td>(-3, 11)</td>
</tr>
<tr>
<td>-0.313613511426339</td>
<td>7.311689634578599</td>
<td>(-2, 11)</td>
</tr>
<tr>
<td>-0.179952252749314</td>
<td>7.301432693193145</td>
<td>(-1, 11)</td>
</tr>
</tbody>
</table>

References