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An X-ray–UV correlation in Cen X-4 during quiescence

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ABSTRACT
Quiescent emission from the neutron star low-mass X-ray binary Cen X-4 is seen to be variable on time-scales from hundreds of seconds to years, suggesting that at least in this object, low-level accretion is important during quiescence. Here, we present results from recent XMM–Newton and Swift observations of Cen X-4, where the X-ray flux (0.5–10 keV) varies by a factor of 6.5 between the brightest and faintest states. We find a positive correlation between the X-ray flux and the simultaneous near-ultraviolet (UV) flux, where as there is no significant correlation between the X-ray and simultaneous optical (V, B) fluxes. This suggests that while the X-ray and UV emitting regions are somehow linked, the optical region originates elsewhere. Comparing the luminosities, it is plausible that the UV emission originates due to reprocessing of the X-ray flux by the accretion disc, with the hot inner region of the disc being a possible location for the UV emitting region. The optical emission, however, could be dominated by the donor star. The X-ray/UV correlation does not favour the accretion stream impact point as the source of the UV emission.

Key words: stars: neutron – X-rays: binaries – X-rays: individual: Cen X-4.

1 INTRODUCTION
Low-mass X-ray binaries (LMXBs) are often transient, cycling through periods of outburst and quiescence where the X-ray luminosity varies by over four orders of magnitude between the two. In quiescence, the accretion rate is at less than $10^{-4}$ of the Eddington limit. The transient behaviour of X-ray binaries (and dwarf novae) is broadly described by the disc instability model (see Lasota 2001, 2008 for a review) where a thermal–viscous instability causes an outburst. During an outburst accretion likely occurs via an optically thick, geometrically thin disc (e.g. Shakura & Sunyaev 1973). During quiescence, the picture is not quite as clear, and accretion is thought to occur in a very different manner. In order to describe both the long recurrence times for transients and the observed quiescent X-ray binaries (and dwarf novae) is broadly described by the disc instability model (see Lasota 2001, 2008 for a review) where a thermal–viscous instability causes an outburst. During an outburst accretion likely occurs via an optically thick, geometrically thin disc (e.g. Shakura & Sunyaev 1973). During quiescence, the picture is not quite as clear, and accretion is thought to occur in a very different manner. In order to describe both the long recurrence times for transients and the observed quiescent X-ray emission a standard quiescent X-ray emission a standard thin disc must be truncated at large radii (e.g. Lasota 1996; Dubus, Hameury & Lasota 2001). Moreover, as the accretion rate drops, it is expected that the inner disc evaporates, leaving a hole (Meyer & Meyer-Hofmeister 1994).

Two component models for the quiescent accretion flow were therefore developed, where an outer truncated disc has an advection-dominated accretion flow (ADAF; Narayan & Yi 1994, 1995) inside (Narayan, McClintock & Yi 1996; Esin, McClintock & Narayan 1997; Narayan, Barret & McClintock 1997). Such models successfully explain the quiescent X-ray properties of black hole LMXBs, which are seen to display simple power-law emission. For neutron stars, there is the added complication of a stellar surface. The ADAF models predict that neutron stars should have a higher quiescent luminosity compared to black holes (with the same orbital period) as the hot flow can heat the neutron star surface, where as in black holes the energy will be advected through the event horizon (e.g., Garcia et al. 2001, and references therein). While such a difference is generally seen (though see Jonker et al. 2007 for a remarkable exception) it is also expected that quiescent neutron stars should be hot due to heating of the crust during the outburst (‘deep crustal heating’; Brown, Bildsten & Rutledge 1998). Moreover, it is even possible that energy may go into a jet rather than being advected (Fender, Gallo & Jonker 2003).

The exact geometry and details of how the hot, radiatively inefficient accretion flows work during quiescence is still not clear. There is uncertainty as to whether the flow is an ADAF (Narayan & Yi 1994, 1995) or an advection-dominated flow with strong outflows (ADIOS; Blandford & Begelman 1999), and what the value of the standard viscosity parameter, $\alpha$, is in quiescence (Menou et al. 2000;
The truncation radius is usually assumed to be between $10^2$ and $10^4$ Schwarzschild radii (Esin et al. 1997), but it too is uncertain (Menou et al. 1999). How matter is accreted on to the neutron star is also not clear. The quiescent luminosities of neutron star LMXBs are low enough to require that much of the accreted matter is prevented from reaching the neutron star surface, which could happen via the propeller effect (Menou et al. 1999; Menou & McClintock 2001). Alternatively, the hot accretion flow may not be an ADAF (or ADIOS), but a hot settling flow instead (Medvedev & Narayan 2001; Narayan & Medvedev 2003). In this hot settling flow scenario, hot quasi-spherical accretion on to a spinning neutron star happens subsonically, and the flow essentially settles on to the rotating neutron star without a shock. The luminosity of this flow is then mostly generated as rotational energy is extracted from the neutron star, rather than from the mass accretion rate. Thus, there exists a range of models to explain how accretion happens during quiescence, and even the basic parameters of such models are not well constrained, at least in the case of neutron star LMXBs.

The mechanism behind accretion in quiescence therefore remains elusive. Yet, one signature that it must be occurring is sporadic X-ray variability. Such variability in quiescent emission from neutron stars has been seen in several objects on all time-scales studied – from hundreds of seconds to years (Campana et al. 1997, 2004; Rutledge et al. 2001, 2002; Cackett et al. 2005, 2010, 2011; Fridriksson et al. 2011). X-ray spectra of quiescent neutron stars typically show two components – both a thermal, blackbody-like component and a non-thermal power-law component. A recent study of long-term quiescent variability in Cen X-4 (Cackett et al. 2010) confirmed that the thermal component, as well as the power-law component, is variable. Interestingly, the power law and thermal fluxes appear to vary in tandem, contributing roughly the same fraction of the total flux at all epochs. This suggests a clear link between the two components during quiescence. Compton upscattering of soft photons from the neutron star surface in the hot accretion flow would seem to be an obvious link between these two components. However, it would produce a power law that is only a few per cent of the 0.5–10 keV flux (Menou & McClintock 2001), while we observe the components to be roughly equal (Cackett et al. 2010). The power law arising from shock emission between a radio pulsar wind and inflowing matter from the companion star (Campana & Stella 2003) also seems to be ruled out as the variability in Cen X-4 cannot be explained by correlated changes in column density and the power-law component (Cackett et al. 2010). The accretion flow getting all the way down to the neutron star surface may therefore be the most realistic mechanism for the synchronous change in both the neutron star surface temperature and the power-law component as the accretion rate varies (e.g. Stella et al. 1994; Menou & McClintock 2001).

In order to further investigate the nature of quiescent variability in neutron star LMXBs on time-scales of weeks to months we observed Cen X-4 with XMM–Newton four times between 2010 August and 2011 January. Importantly, we also obtain optical and ultraviolet (UV) photometry with XMM–Newton to investigate the nature of broad-band variability during quiescence. Combining these four new XMM–Newton observations with one archival XMM–Newton observation and two Swift observations we show a clear correlation between X-ray and UV flux during quiescence. Cen X-4 is one of very few objects where such an X-ray/UV study can be performed given its proximity and relatively low extinction. In Section 2, we describe the data reduction and analysis before describing our results in Section 3 and discussing their implications in Section 4.

## 2 DATA REDUCTION AND ANALYSIS

The four XMM–Newton observations were obtained between the end of 2010 August and 2011 January with each observation lasting between 10 and 20ks. The ObsIDs are 0654470201, 0654470301, 0654470401, 0654470501 and their respective start dates (dd/mm/yy) are 25/08/10, 04/09/10, 24/01/11 and 31/01/11. The X-ray detectors were operated in full window mode with a medium filter. The Optical Monitor (OM; Mason et al. 2001) was operated such that data were collected in multiple optical/UV filters during the X-ray observations. In Table 1, we give details of all four observations along with exposure times for each instrument and filter used. The data were reduced with the XMM–Newton Science Analysis Software, version 11.0.0, using the latest calibration files.

In addition to the XMM–Newton data, we also analyse two Swift observations of Cen X-4 that also have both X-ray and optical/UV data. The ObsIDs 00035324001 and 00035324002 were performed on 03/09/06 and 01/05/12. The X-ray detector was operated in photon counting mode. In the first observation, only a UVW1 exposure was taken, while as in the second observation exposures in B, U, UVW1, UVM2 and UVM2 were taken. Table 2 gives details of the Swift observations. The data were reduced using HEASOFT, version 6.11.1. In the following sections, we detail the data reduction for the X-ray and optical/UV detectors.

### 2.1 X-ray data reduction

#### 2.1.1 XMM–Newton

We created calibrated event files for each observation from the Observation Data Files using the emproc and epproc commands for the MOS and pn detectors. We check for periods of high background by

<table>
<thead>
<tr>
<th>Detector or filter</th>
<th>0654470201</th>
<th>0654470301</th>
<th>0654470401</th>
<th>0654470501</th>
</tr>
</thead>
<tbody>
<tr>
<td>MOS</td>
<td>Exp. (ks)</td>
<td>11.2</td>
<td>10.4</td>
<td>4.0</td>
</tr>
<tr>
<td></td>
<td>Net rate (10^{-12} c s^{-1})</td>
<td>0.75 ± 0.03</td>
<td>3.07 ± 0.07</td>
<td>6.2 ± 0.6</td>
</tr>
<tr>
<td>pn</td>
<td>Exp. (ks)</td>
<td>16.1</td>
<td>9.2 ± 0.7</td>
<td>8.7 ± 1.0</td>
</tr>
<tr>
<td></td>
<td>AB flux (10^{-12} c s^{-1})</td>
<td>3.73 ± 0.05</td>
<td>2.86</td>
<td>5.5 ± 1.4</td>
</tr>
<tr>
<td>V</td>
<td>Exp. (ks)</td>
<td>11.3</td>
<td>23.1 ± 1.8</td>
<td>1.86</td>
</tr>
<tr>
<td></td>
<td>AB flux (10^{-12} c s^{-1})</td>
<td>1.20 ± 0.03</td>
<td>2.51 ± 0.18</td>
<td>1.20 ± 0.18</td>
</tr>
<tr>
<td>B</td>
<td>Exp. (ks)</td>
<td>11.3</td>
<td>2.1 ± 0.06</td>
<td>1.76</td>
</tr>
<tr>
<td></td>
<td>AB flux (10^{-12} c s^{-1})</td>
<td>3.07 ± 0.20</td>
<td>12.8 ± 2.1</td>
<td>10.2 ± 1.2</td>
</tr>
<tr>
<td>U</td>
<td>Exp. (ks)</td>
<td>17.6</td>
<td>8.8 ± 0.9</td>
<td>1.76</td>
</tr>
<tr>
<td></td>
<td>AB flux (10^{-12} c s^{-1})</td>
<td>1.20 ± 0.18</td>
<td>8.8 ± 0.9</td>
<td>9.7 ± 1.2</td>
</tr>
</tbody>
</table>

ND indicates that the source was not detected in that filter. MOS and pn net count rates are given in the 0.5–10 keV energy range. The count rate is given for MOS 1 only, the MOS 2 rate is always very similar.

Table 1. Details of the new XMM–Newton observations of Cen X-4, including count rates and optical/UV fluxes. AB flux densities have units of erg s^{-1} cm^{-2} Å.
creating light curves from the entire detectors with 100 s time binning, filtering for events with energies > 10 keV and PATTERN = 0 for the MOS, and 10 < E < 12 keV and PATTERN = 0 for the pn.

The background was high throughout the entire first observation (0654470201), with the MOS 1 count rate ranging from 0.2–0.30 counts s\(^{-1}\). We use the XRTMKARF tool to create the ancillary response file, and use the appropriate response matrix based on the epoch of the observation and the observing mode. Given the variability response file, and use the appropriate response matrix based on the epoch of the observation and the observing mode. Given the variability of the source, we were able to obtain images in the following filters during this observation: the pn filter (ks) (10\(^{-1}\)–10\(^{-1}\) cm\(^{-2}\) cm\(^{-2}\) Å).

Table 2. Details of the Swift observations of Cen X-4, including count rates and optical/UV fluxes. AB flux densities have units of erg s\(^{-1}\) cm\(^{-2}\) Å. X-ray Telescope (XRT) net count rates are given in the 0.5–10 keV energy range.

<table>
<thead>
<tr>
<th>Detector</th>
<th>00035324001</th>
<th>00035324002</th>
</tr>
</thead>
<tbody>
<tr>
<td>Filter</td>
<td>Exp. Net rate (10(^{-1}) cs(^{-1}))</td>
<td>Exp. Net rate (10(^{-1}) cs(^{-1}))</td>
</tr>
<tr>
<td>XRT</td>
<td>4.5 0.38 ± 0.03</td>
<td>3.8 0.76 ± 0.05</td>
</tr>
<tr>
<td>V</td>
<td>– –</td>
<td>– 0.30 0.75 ± 0.07</td>
</tr>
<tr>
<td>B</td>
<td>– –</td>
<td>0.30 1.11 ± 0.09</td>
</tr>
<tr>
<td>U</td>
<td>– –</td>
<td>0.30 1.35 ± 0.09</td>
</tr>
<tr>
<td>U1V1</td>
<td>6.25 0.32 ± 0.01</td>
<td>15.0 ± 0.5</td>
</tr>
<tr>
<td>U1V2</td>
<td>– –</td>
<td>0.95 0.29 ± 0.02</td>
</tr>
</tbody>
</table>

Table 3. XMM–Newton obsID 0144900101, optical/UV details. AB flux densities have units of erg s\(^{-1}\) cm\(^{-2}\) Å.

<table>
<thead>
<tr>
<th>Filter</th>
<th>Exp. (ks)</th>
<th>Net rate (10(^{-17}) erg s(^{-1}) cm(^{-2}))</th>
<th>AB flux (10(^{-17}) erg s(^{-1}) cm(^{-2}) Å)</th>
</tr>
</thead>
<tbody>
<tr>
<td>V</td>
<td>4.0</td>
<td>8.3 ± 0.5</td>
<td>20.8 ± 1.2</td>
</tr>
<tr>
<td>B</td>
<td>4.0</td>
<td>13.9 ± 0.5</td>
<td>17.3 ± 0.7</td>
</tr>
<tr>
<td>U</td>
<td>4.0</td>
<td>6.6 ± 0.3</td>
<td>12.7 ± 0.7</td>
</tr>
<tr>
<td>U1V1</td>
<td>12.2</td>
<td>3.7 ± 0.1</td>
<td>18.0 ± 0.5</td>
</tr>
<tr>
<td>U1V2</td>
<td>19.8</td>
<td>0.81 ± 0.05</td>
<td>17.9 ± 1.0</td>
</tr>
<tr>
<td>U1V2</td>
<td>25.0</td>
<td>0.25 ± 0.07</td>
<td>14.5 ± 3.8</td>
</tr>
</tbody>
</table>

Given: V (5407 Å), B (4334 Å), U (3472 Å), U1V1 (2905 Å) and U1V2 (2298 Å).

We reduce the OM data using the omichain tool. This tool processes the data with the latest calibration files, and then performs source detection and aperture photometry. The output files include both the images for each filter along with a combined source list containing detailed parameters for each source detected in all exposures, including the count rate (corrected for instrumental effects such as detector dead time) and AB flux densities (in erg s\(^{-1}\) cm\(^{-2}\) Å\(^{-1}\)). We visually identified Cen X-4 in all images before finding the corresponding source in the source list. In observation 0654470301 Cen X-4 was not automatically detected in the U1V2 filter, though a faint excess in counts did seem apparent on visual inspection. We therefore ran the omdetect tool (part of the omichain pipeline) on its own with the significance reduced to only requiring 1\(\sigma\) above the background; however, the source was still not detected (note the default is a 3\(\sigma\) threshold).

For both the MOS and pn detectors, we extract the source spectrum from a circular region of radius 40 pixels, and the background spectrum from a nearby, source-free 2 arcmin region. The response files are generated with the arfgen and rmfgen tools, and the resulting spectrum is binned to a minimum of 25 counts per bin. We give net count rates in Table 1.

2.1.2 Swift

Calibrated event files were created by reprocessing the data using the xrtpipeline tool and applying the standard (default) screening criteria. We analyse data taken in photon counting mode. We extracted the source spectrum using xselect, and a circular extraction region with a 20 pixel radius. The background spectrum was extracted from an annulus with inner radius of 40 pixels and outer radius of 120 pixels. We use the xrtmkarf tool to create the ancillary response file, and use the appropriate response matrix based on the epoch of the observation and the observing mode. Given the shorter exposures and lower count rates compared to xmm–newton we only bin spectra to 10 counts per bin.

2.2 Optical/UV data

2.2.1 XMM–Newton

We observed Cen X-4 with the OM in imaging mode. However, rather than using the standard setup, we chose to more efficiently obtain exposures of Cen X-4 in five filters by using a Science User Defined mode whereby only a single window is used to observe just the centre of the field of view (normally a mosaic is created from multiple windows to cover a much larger field of view). In this way, we were able to obtain images in the following filters during the short X-ray exposures of Cen X-4 (effective wavelengths are different, to the xmm–newton filters. The first observation used the u1v1 filter exclusively, where as the second observation used V (5468 Å), B (4392 Å), U (3465 Å), U1V1 (2600 Å), U1V2 (2246 Å) and U1V2 (1928 Å) filters (wavelengths from poole et al. 2008).

Note that both the U1V1 and U1V2 filters are broader than the U1V2 filter, and hence are more sensitive. However, U1V1 and U1V2 also have red leaks, meaning that their sensitivity stretches to


longer (redder) wavelengths than \textit{UVM2}. We used the \textit{UVOTMAGHIST} tool to perform aperture photometry, using a 3 arcsec radius circular extraction region, and a circular, source-free nearby background region with a 12 arcsec radius. This tool performs photometry on all separate exposures for a given filter, applies the aperture correction and determines the \textit{AB} flux. Having determined the count rates and fluxes for each individual exposure, we calculate the exposure-weighted average count rate and flux, as given in Table 2.

2.3 X-ray spectral analysis

We fit the X-ray spectra using \texttt{xspec} version 12.7.0 (Arnaud 1996). Each observation is fitted separately, but when fitting the individual \textit{XMM–Newton} observations we jointly fit the MOS 1, MOS 2 and pn spectra, with the parameters being the same for all detectors. We fit the spectra using a neutron star atmosphere plus a power law, all modified by Galactic photoelectric absorption. We use the nsatmos model for the neutron star atmosphere (Heinke et al. 2006). This is the same model we used to fit six previous observations of Cen X-4 (Cackett et al. 2010), thus, allowing for a direct comparison with these previous results.\footnote{This previous analysis looked at data from a range of different missions (1 ASCA, 2 Chandra, 2 XMM and 1 Suzaku).} For the same reason, we choose to fix \(N_\text{H} = 4.9 \times 10^{20} \text{ cm}^{-2}\), the value we found from jointly fitting six previous observations in Cackett et al. (2010). Also, to allow a direct comparison with our previous results, we fix the distance to 1.2 kpc (Chevalier et al. 1989) in the spectral fits. In Cackett et al. (2010), we investigated which parameters were variable, finding that the thermal component can be fitted equally well with either a variable temperature or a variable radius, and it is only the flux that we are mostly concerned with here. We assume a neutron star with \(R = 10 \text{ km}\) and \(M = 1.4 \text{ M}_\odot\), and that the entire surface is emitting. We fit the spectra in the 0.5–10 keV range and best-fitting parameters are given in Table 4. Uncertainties are quoted at the 1\(\sigma\) level throughout.

3 RESULTS

The long-term X-ray quiescent light curve of Cen X-4 is shown in Fig. 1, where we include the observations from Cackett et al. (2010) as well as the four new \textit{XMM–Newton} and two \textit{Swift} observations presented here. Notably, significant variability is seen on the weeks to months time-scales probed by these new observations. Most significant is the factor of 3 drop in flux between the third and fourth observations taken at the end of 2011 January and separated by only 7 d.

It is clear from looking at the OM count rates and fluxes (Table 1) that Cen X-4 is also highly variable in the optical/UV part of the spectrum during quiescence. In Fig. 2, we plot the optical/UV fluxes against the 0.5–10 keV unabsorbed X-ray flux. In order to compare the energetics, we convert the flux densities (erg s\(^{-1}\) cm\(^{-2}\)) that are the output from the \texttt{OMCHAIN} to fluxes by multiplying by the full width at half-maximum of the filter bandpass. We use: 684 Å (\(V\)), 976 Å (\(B\)), 810 Å (\(U\)), 620 Å (\(UVW1\)), 439 Å (\(UVW2\) for the \textit{XMM–Newton} filters, and 769 Å (\(V\)), 975 Å (\(B\)), 785 Å (\(U\)), 693 Å (\(UVW1\)), 498 Å (\(UVW2\)) for the \textit{Swift} filters. We also de-redden the UV/optical fluxes. We use the gas-to-dust ratio from Güver & Özel (2009) in order to convert from equivalent hydrogen column to \(E(B-V)\). We then use the interstellar extinction curve of Cardelli, Clayton & Mathis (1989) to estimate the extinction correction at the wavelengths of each filter.

We test for a correlation between the optical/UV and X-ray fluxes using a simple Pearson’s linear correlation test. The correlation

\begin{table}
\centering
\caption{Spectral fit parameters. The flux is given in units of erg s\(^{-1}\) cm\(^{-2}\). The column density was fixed at \(N_\text{H} = 4.9 \times 10^{20} \text{ cm}^{-2}\) in all spectral fits. We also assumed a neutron star radius of 10 km and mass of 1.4 M\(_\odot\), and a distance to Cen X-4 of 1.2 kpc. The power-law normalization is defined as photons keV\(^{-1}\) cm\(^{-2}\) s\(^{-1}\) at 1 keV. We define that the thermal fraction is the ratio of the unabsorbed 0.5–10 keV thermal flux to the total unabsorbed 0.5–10 keV flux.
\label{tab:fits}
\begin{tabular}{lcccc}
\hline
Parameter & \textit{XMM–Newton} & & \textit{Swift} & \\
 & 0654470201 & 0654470301 & 0654470401 & 00035324001 \\
\hline
\(kT_{\text{eff}}\) (eV) & 55.9 \pm 0.7 & 56.5 \pm 0.3 & 60.8 \pm 0.3 & 49.0 \pm 0.6 \\
Power-law index, \(\Gamma\) & 1.77 \pm 0.21 & 1.62 \pm 0.10 & 1.38 \pm 0.10 & 1.94 \pm 0.19 \\
Power-law norm (10\(^{-5}\)) & 6.6 \pm 1.4 & 6.3 \pm 0.7 & 5.5 \pm 0.7 & 3.6 \pm 0.7 \\
Unabs. 0.5–10 keV flux (10\(^{-12}\)) & 0.92 \pm 0.08 & 1.00 \pm 0.02 & 1.31 \pm 0.03 & 0.44 \pm 0.01 \\
Thermal fraction (per cent) & 58 & 56 & 62 & 59 \\
\chi^2 (dof) & 0.91 (106) & 0.93 (168) & 0.99 (129) & 1.27 (71) \\
\hline
\end{tabular}
\end{table}
coefficients we find are $r = 0.554, 0.779, 0.982, 0.973, 0.976$ for the $V, B, U, UVW1$ and $UVM2$ filters, respectively. Given the number of data points, and a two-tail test (no a priori knowledge of positive or negative correlation), this corresponds to a positive correlation at the 0.746, 0.880, 0.999 52, 0.999 77 and 0.976 confidence levels. Thus, the $U$ and $UVW1$ fluxes are correlated with the X-ray flux at greater than 3$\sigma$, while the $UVM2$ flux is correlated at greater than 2$\sigma$. There is no significant correlation between the $V$ or $B$ fluxes and the X-ray flux, indicating that this is a UV only correlation. We also look to see if there is a correlation between $UVW1$ and $V$ band fluxes, finding $r = 0.63$, which corresponds to a positive correlation at the 0.816 confidence level. Thus, there is no significant correlation between the $V$ band and $UVW1$ fluxes.

We now concentrate on the $UVW1$ versus X-ray correlation further, given that it is the most significant correlation. Obviously, given the high linear correlation coefficient, a simple straight line fits the $UVW1$ – X-ray correlation well. We get the following best-fitting parameters: $f_{UV} = (0.14 \pm 0.01)f_X - (0.50 \pm 0.18)$, giving a reduced $\chi^2 = 1.35$. We also test fitting a power law of the form $f_{UV} = af_Xb$, where $f_{0X}$ is a constant to allow a non-zero X-ray flux when $f_{UV} = 0$, or vice versa. The best fit is very close to a linear relationship, with $b = 1.00^{+0.34}_{-0.14}$, and giving a reduced $\chi^2 = 1.69$. If we fit a power law with the index, $b$ fixed at 0.5 (as expected for reprocessing by van Paradijs & McClintock 1994), we get a worse fit (reduced $\chi^2 = 4.5$). However, with only seven data points we can clearly not make any strong conclusions about the form of the correlation, especially as there is only one observation at a high luminosity, where a deviation from a linear relationship would become apparent. We also caution that we are comparing $UVW1$ fluxes with both $Swift$ and $XMM–Newton$ – the filter responses and wavelength range are slightly different between the missions which could lead to a small offset between the two. Further data are needed to strengthen and define the shape of the correlation.

We also look for whether there is an X-ray/UV correlation on shorter time-scales by looking at the UV and X-ray light curve during the two $Swift$ observations. This is possible with $Swift$ because the observations consist of multiple ~1 ks exposures, each giving separate UV images. With $XMM–Newton$, however, there is only one exposure in each filter during the entire observation and so we cannot look for X-ray–UV correlations during the $XMM–Newton$ observations. The first $Swift$ observation consists of six $Swift$ orbits of data, and the $UVW1$ filter was used for all of them. We can therefore compare the X-ray and UV light curves on time-scales of a few thousand seconds. The second $Swift$ observation comprised three orbits, with a $UVW1$ exposure during each orbit. For both observations, we create a background-subtracted $Swift$/X-ray Telescope (XRT) light curves, using the $XRTLCCORR$ tool to perform count rates corrections for vignetting, point spread function, etc. The $UVW1$ count rates are determined using the $UDOTMAGPHI$ tool (as described above). We show the light curves in Fig. 3. The XRT and $UVW1$ count rates are generally well correlated during both observations.

The $XMM–Newton$ observations only perform one exposure in each filter during a given observation. Therefore, while the X-ray flux we obtain is from the entire observation, the optical/UV fluxes are only obtained during some fraction of the observation. The X-ray exposures range from approximately 10–20 ks, where as the optical/UV exposures are between 1.76 and 5 ks. Variability during the observation could therefore add scatter to the UV and X-ray correlations. Examining the X-ray light curves from the $XMM–Newton$ observations there is variability present (as has also been seen during previous $XMM–Newton$ and $Chandra$ observations; see, e.g., Campana et al. 2004; Cackett et al. 2010). However, any subset of 2 ks or longer will not deviate too far from the mean X-ray flux, and thus not add a large amount of scatter.

**4 DISCUSSION AND CONCLUSIONS**

Our observations of Cen X-4 during quiescence have shown that the X-ray and near-UV emission are correlated, where as there is no significant correlation between the X-ray and optical emission. The X-ray/UV correlation is seen both on time-scales of weeks–years, as well as during observations spanning less than one day. In order to investigate the physical origin of this correlation, it is important to consider the energetics – if the correlation is due to reprocessing of the X-ray emission, then the near-UV flux must be only a fraction of the X-ray flux.

We find that the $UVW1$ flux is approximately 10 per cent of the 0.5–10 keV flux. However, much of the X-ray flux from the hot
emission is in more than just the UVW flux when extending the energy range. Of course, the near-UV flux is in the 0.5–10 keV band; hence, the slightly lower increase brighter observation means more of the neutron star atmosphere we find a factor of 2.6 increase. The higher temperature during the UVM fluxes, we get an increase by a factor of 4.0 and 0.00035324002). The XRT data have 200 s time binning.

Figure 3. Swift X-ray (black) and UVW1 (red, triangles) light curves during the first (top) and second (bottom) Swift observations (00035324001 and 00035324002). The XRT data have 200 s time binning.

neutron star surface is emitted below 0.5 keV, and thus it is important to take this flux into consideration. Furthermore, the power-law component will likely extend beyond 10 keV, though its shape there is uncertain. To provide an estimate for the total X-ray flux, we extrapolate the best-fitting neutron star atmosphere plus power-law models to cover the range 0.01–100 keV. For the faintest observation (XMM–Newton: 0654470501) we find a factor of 3 increase in the flux, as for the brightest observation (Swift: 00035324002) we find a factor of 2.6 increase. The higher temperature during the brighter observation means more of the neutron star atmosphere flux is in the 0.5–10 keV band; hence, the slightly lower increase in flux when extending the energy range. Of course, the near-UV emission is in more than just the UVW1 filter. Combining the U, UVW1 and UVM2 fluxes, we get an increase by a factor of 4.0 and 2.7 for the faintest and brightest observations, respectively, over using just the UVW1 flux. Thus, when accounting for emission over a wider wavelength range, the near-UV flux is still approximately 10 per cent of the X-ray flux. Such a fraction of the X-ray flux is reasonable for reprocessed emission.

evander Paradijs & McClintock (1994) consider X-ray reprocessing in the accretion disc and show that the reprocessed emission should be proportional to \( L_{\text{X}}^{1/2} \), where \( R \) is the outer radius of the accretion disc. We, however, find that the UV flux increases approximately linearly with the X-ray flux and a power law with index of 0.5 does give a significantly worse fit. However, with the limited number of observations and range in flux the exact relationship is uncertain, and would be significantly strengthened by further data. Considering the relevant time-scales for reprocessing, the light travel time from the neutron star to the outer disc is less than 1 s, and so we would not expect to see any time delay between the X-ray and UV light curves. This is what we observe (Fig. 3) where the X-ray and UV vary in tandem during Swift observations. The viscous time-scale from the outer disc in these objects is significantly longer (of the order of weeks), and therefore cannot be used to explain the simultaneous variability in UV and X-rays.

Another source where a similar correlation has been found is the quiescent black hole V404 Cyg, where a correlation between X-ray and H\( \alpha \) emission has been previously observed (Hynes et al. 2004), with a weaker correlation between the X-ray flux and the optical continuum. Irradiation/reprocessing can readily explain that correlation as the entire H\( \alpha \) line is seen to respond quickly to X-ray variability. Here too, we see that the X-ray and UV flux seem to respond simultaneously during the Swift observations (see Fig. 3), though we do not have the data to be able to establish any lag. This prompt response also supports irradiation/reprocessing as the origin of the correlation observed here. However, we do not observe any clear correlation between the optical continuum and the X-ray flux (unlike in V404 Cyg), which would suggest that the UV and optical emitting regions are well separated. Alternatively, it could just be the case that in Cen X-4, the optical flux is dominated by the donor star, and so any change in optical flux due to reprocessing may be too small to see. Furthermore, the longer orbital period of V404 Cyg (6.5 d compared to 15 h for Cen X-4; Cowley et al. 1988; Casares, Charles & Naylor 1992) will lead to a larger disc, which would be more luminous at optical wavelengths.

The optical emission could be dominated by the donor star, as well as some contribution from the cooler outer regions of the truncated accretion disc. The UV emission is too hot to arise from the donor star, and the X-ray/UV correlation suggests it must arise from a region that can see the central X-ray source, such as the inner region of the truncated accretion disc. The correlation between X-ray and UV flux would seem to rule out the accretion stream impact point as the location for the UV emission, which had been previously suggested by McClintock & Remillard (2000) based on the Mg\( \alpha \) emission line in an Hubble Space Telescope/STIS spectrum of Cen X-4 not showing a double-peak as would be expected for emission from the disc. Hynes & Robinson (2012) closely examine the UV Spectral Energy Distributions (SEDs) of three quiescent black hole and one quiescent neutron star and also rule out the stream impact point as the origin of the UV emission. They find that the mass accretion rate required to give the observed UV emission at the stream impact point in the black hole GU Mus and neutron star Aql X-1 is 10 times higher than realistic average rates from these sources based on their accretion histories. They therefore favour an origin for the UV emission as being located in the hot inner region of the disc as also suggested by Campana & Stella (2000) and McClintock et al. (2003). Our results also support such a picture.

It is important to consider whether there is a viable alternative to X-ray reprocessing as the origin of the X-ray/UV correlation. For instance, could the emission in Cen X-4 be jet dominated and hence the UV flux and X-ray power law both originate from the jet? While radio emission has never been detected from a quiescent neutron star, it has been detected in quiescence (Hynes et al. 2006) implies that there is a radio-emitting outflow during quiescence. Broad-band modelling of the quiescent SED in A0620–00 shows that the emission can be fitted with a maximally jet-dominated model with the jet emission dominating from radio through to the soft X-rays (Gallo et al. 2007) and implies that there are strong outflows (Froning et al. 2011). Similarly, V404 Cyg has been detected at radio wavelengths in quiescence.
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