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Aad, G.; et al., [Unknown]; Bentvelsen, S.; Berglund, E.; Bobbink, G.J.; Bos, K.; Boterenbrood, H.; Colijn, A.P.; de Jong, P.; de Nooij, L.; Deviveiros, P.O.; Doxiadis, A.D.; Ferrari, P.; Garitaonandia, H.; Geerts, D.A.A.; Gosselink, M.; Hartjes, F.; Hessey, N.P.; Igonkina, O.; Kayl, M.S.; Klous, S.; Kluit, P.; Koffeman, E.; Lee, H.; Lenz, T.; Linde, F.; Luijckx, G.; Massaro, G.; Mechnich, J.; Mussche, I.; Ottersbach, J.P.; Reichold, A.; Rijpstra, M.; Ruckstuhl, N.; Snuverink, J.; Ta, D.; Tsiakiris, M.; Turlay, E.; van der Graaf, H.; van der Kraaij, E.; van der Leeuw, R.; van der Poel, E.; van Kesteren, Z.; van Vulpen, I.; Verkerke, W.; Vermeulen, J.C.; Vranjes Milosavljevic, M.; Vreeswijk, M.

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Search for excited leptons in proton-proton collisions at $\sqrt{s} = 7$ TeV with the ATLAS detector

 G. Aad *et al.**

(ATLAS Collaboration)

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The ATLAS detector is used to search for excited leptons in the electromagnetic radiative decay channel $\ell^* \rightarrow \ell\gamma$. Results are presented based on the analysis of pp collisions at a center-of-mass energy of 7 TeV corresponding to an integrated luminosity of 2.05 fb^{-1} . No evidence for excited leptons is found, and limits are set on the compositeness scale Λ as a function of the excited lepton mass m_{ℓ^*} . In the special case where $\Lambda = m_{\ell^*}$, excited electron and muon masses below 1.87 TeV and 1.75 TeV are excluded at 95% C.L., respectively.

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I. INTRODUCTION

The standard model (SM) of particle physics is an extremely successful effective theory which has been extensively tested over the past 40 years. However, a number of fundamental questions are left unanswered. In particular, the SM does not provide an explanation for the source of the mass hierarchy and the generational structure of quarks and leptons. Compositeness models address these questions by proposing that quarks and leptons are composed of hypothetical constituents named preons [1]. In these models, quarks and leptons are the lowest-energy bound states of these hypothetical particles. New interactions among quarks and leptons should then be visible at the scale of the constituents' binding energies, and give rise to excited states. At the LHC, excited lepton ℓ^* production via four-fermion contact interactions can be described by the effective Lagrangian [2]

$$\mathcal{L}_{\text{contact}} = \frac{g_*^2}{2\Lambda^2} j^\mu j_\mu,$$

where g_*^2 is the coupling constant, Λ is the compositeness scale, and j_μ is the fermion current

$$j_\mu = \eta_L \bar{f}_L \gamma_\mu f_L + \eta'_L \bar{f}_L^* \gamma_\mu f_L^* + \eta''_L \bar{f}_L^* \gamma_\mu f_L + \text{H.c.} + (L \rightarrow R).$$

For simplicity and consistency with recent searches, the following prescription is used: $g_*^2 = 4\pi$, $\eta_L = \eta'_L = \eta''_L = 1$, and $\eta_R = \eta'_R = \eta''_R = 0$ such that chiral symmetry is conserved [3,4]. The above ansatz ignores underlying preon dynamics and is valid as long as the mass of the excited leptons is below the scale Λ . In the well-studied case of the homodoublet-type ℓ^* [2,5,6], the relevant

*Full author list given at the end of the article.

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gauge-mediated Lagrangian describing transitions between excited and ground-state leptons is

$$\mathcal{L}_{\text{GM}} = \frac{1}{2\Lambda} \bar{\ell}_R^* \sigma^{\mu\nu} \left[g f \frac{\tau^a}{2} W_{\mu\nu}^a + g' f' \frac{Y}{2} B_{\mu\nu} \right] \ell_L + \text{H.c.},$$

where ℓ_L is the lepton field, $W_{\mu\nu}$ and $B_{\mu\nu}$ are the $SU(2)_L$ and $U(1)_Y$ field strength tensors, g and g' are the respective electroweak couplings, and f and f' are phenomenological constants chosen to be equal to 1. The \mathcal{L}_{GM} term allows the decay of excited leptons via the electromagnetic radiative mode $\ell^{*\pm} \rightarrow \ell^\pm \gamma$, a very clean signature which is exploited in this search. For a fixed value of Λ , the branching ratio $B(\ell^{*\pm} \rightarrow \ell^\pm \gamma)$ decreases rapidly with increasing ℓ^* mass. For $\Lambda = 2$ TeV, $B(\ell^{*\pm} \rightarrow \ell^\pm \gamma)$ is 30% for $m_{\ell^*} = 0.2$ TeV and decreases exponentially to about 2.3% for $m_{\ell^*} = 2$ TeV.

Previous searches at LEP [7], HERA [8], and the Tevatron [9] have found no evidence for such excited leptons. For the case where $\Lambda = m_{\ell^*}$, the CMS experiment has excluded masses below 1.07 TeV for e^* and 1.09 TeV for μ^* at the 95% credibility level (C.L.) [10].

II. ANALYSIS STRATEGY

This article reports on searches for excited electrons and muons in the $\ell^* \rightarrow \ell\gamma$ channel based on 2.05 fb^{-1} of 7 TeV pp collision data recorded in 2011 with the ATLAS detector [11]. The benchmark signal model considered is based upon theoretical calculations from Ref. [2]. In this model, excited leptons may be produced singly via $q\bar{q} \rightarrow \ell^* \bar{\ell}$ or in pairs via $q\bar{q} \rightarrow \ell^* \bar{\ell}^*$, due to contact interactions. As the cross section for pair production is much less than for single production, the search for excited leptons is based on the search for events with $\ell\bar{\ell}\gamma$ in the final state: three very energetic particles, isolated, and well separated from one another.

For both the e^* and μ^* searches, the dominant background arises from Drell-Yan (DY) processes accompanied either by a prompt photon from initial- or final-state radiation ($Z + \gamma$) or by a jet misidentified as a photon

($Z + \text{jets}$). The dominant irreducible $Z + \gamma$ background results in the same final state as the signal, whereas $Z + \text{jets}$ background can be highly suppressed by imposing stringent requirements on the quality of the reconstructed photon candidate. Small contributions from $t\bar{t}$ and diboson (WW , WZ , and ZZ) production are also present in both channels. $W + \text{jets}$ events, as well as semileptonic decays of heavy flavor hadrons, and multi-jet events can be heavily suppressed by requiring the leptons and photons to be isolated and thus have a negligible contribution to the total background.

The signature for excited leptons can present itself as a peak in the invariant mass of the $\ell + \gamma$ system because the width of the ℓ^* is predicted to be narrower than the detector mass resolution for excited lepton masses $m_{\ell^*} \lesssim 0.5\Lambda$. This peak could be easily resolved from the $Z + \gamma$ background. However, it is difficult to identify which of the two leading leptons in the event comes from the ℓ^* decay. To avoid this ambiguity, one can search for an excess in the $\ell\bar{\ell}\gamma$ invariant mass ($m_{\ell\bar{\ell}\gamma}$) spectrum. This approach is effective for the whole $m_{\ell^*} - \Lambda$ parameter space probed, as one can search for an excess of events with $m_{\ell\bar{\ell}\gamma} > 350$ GeV, which defines a nearly background-free signal region. Optimization studies demonstrate that the observable $m_{\ell\bar{\ell}\gamma}$ provides better signal sensitivity than $m_{\ell\gamma}$, particularly for lower ℓ^* masses. The analysis strategy therefore relies on $m_{\ell\bar{\ell}\gamma}$ for the statistical interpretation of the results.

III. ATLAS DETECTOR

ATLAS is a multipurpose detector with a forward-backward symmetric cylindrical geometry and nearly 4π coverage in solid angle. It consists of an inner tracking detector immersed in a 2 T solenoidal field, electromagnetic and hadronic calorimeters, and a muon spectrometer. Charged particle tracks and vertices are reconstructed in silicon-based pixel and microstrip tracking detectors that cover $|\eta| < 2.5$ and transition radiation detectors extending to $|\eta| < 2.0$ [12]. A hermetic calorimetry system, which covers $|\eta| < 4.9$, surrounds the superconducting solenoid. The liquid-argon electromagnetic calorimeter, which plays an important role in electron and photon identification and measurement, is finely segmented. It has a readout granularity varying by layer and cells as small as 0.025×0.025 in $\eta \times \phi$, and extends to $|\eta| < 2.5$ to provide excellent energy and position resolution. Hadron calorimetry is provided by an iron-scintillator tile calorimeter in the central rapidity range $|\eta| < 1.7$ and a liquid-argon calorimeter with copper and tungsten as absorber material in the rapidity range $1.5 < |\eta| < 4.9$. Outside the calorimeter, there is a muon spectrometer which is designed to identify muons and measure their momenta with high precision. The muon spectrometer comprises three toroidal air-core magnet systems: one for the barrel and one per endcap, each composed of eight

coils. Three layers of drift tube chambers and/or cathode strip chambers provide precision (η) coordinates for momentum measurement in the region $|\eta| < 2.7$. A muon trigger system consisting of resistive plate chambers in the barrel and thin-gap chambers for $|\eta| > 1$ provides triggering capability up to $|\eta| = 2.4$ and measurements of the ϕ coordinate.

IV. SIMULATED SAMPLES

The excited lepton signal samples are generated based on calculations from Ref. [2] at LO with COMPHEP 4.5.1 [13] interfaced with PYTHIA 6.421 to handle parton showers and hadronization [14,15], using MRST2007 LO* [16] parton distribution functions (PDFs). Only single production of excited leptons is simulated, with the ℓ^* decaying exclusively via the electromagnetic channel. The $Z + \gamma$ sample is generated with SHERPA 1.2.3 [17] using CTEQ6.6 [18] PDFs, requiring the dilepton mass to be above 40 GeV. To avoid phase-space regions where matrix elements diverge, the angular separation between the photon and leptons is required to be $R(\ell, \gamma) = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} > 0.5$ and the transverse momentum (p_T) of the photon is required to be $p_T^\gamma > 10$ GeV. To ensure adequate statistics at large $m_{\ell\bar{\ell}\gamma}$, an additional $Z + \gamma$ sample is generated with $p_T^\gamma > 40$ GeV, and is equivalent to ~ 300 fb $^{-1}$ of data. The $Z + \text{jets}$ background is generated with ALPGEN 2.13 [19], while the $t\bar{t}$ background is produced with MC@NLO 3.41 [20]. In both cases, JIMMY 4.31 [21] is used to describe multiple parton interactions and HERWIG 6.510 [22] is used to simulate the remaining underlying event and parton showers and hadronization. CTEQ6.6 PDFs are used for both backgrounds. To remove overlaps between the $Z + \text{jets}$ and the $Z + \gamma$ samples, $Z + \text{jets}$ events with prompt energetic photons are rejected if the photon-lepton separation is such that $R(\ell, \gamma) > 0.5$. The diboson processes are generated with HERWIG using MRST2007 LO* PDFs. For all samples, final-state photon radiation is handled via PHOTOS [23]. The generated samples are then processed through a detailed detector simulation [24] based on GEANT4 [25] to propagate the particles and account for the detector response. A large sample of MC minimum bias events is then mixed with the signal and background MC events to simulate pileup from additional pp collisions. Simulations are normalized on an event-by-event basis such that the distribution of the number of interactions per event agrees with the spectrum observed in data.

Although SHERPA includes higher-order QCD contributions beyond the $Z + \gamma$ Born amplitude, such as the real emission of partons in the initial state, it omits virtual corrections. For this reason, the $Z + \gamma$ cross section is calculated at next-to-leading order (σ_{NLO}) using MCFM [26] with MSTW2008 NLO PDFs [27]. The theoretical precision of the σ_{NLO} estimate is $\sim 6\%$, and the ratio $\sigma_{\text{NLO}}/\sigma_{\text{SHERPA}}$ is used to determine a correction factor as

a function of $m_{\ell\ell\gamma}$. The $Z + \text{jets}$ cross section is initially normalized to predictions calculated at next-to-next-to-leading order (NNLO) in perturbative QCD as determined by the FEWZ [28] program using MSTW2008 NNLO PDFs. Since the misidentification of jets as photons is not well modeled, the $Z + \text{jets}$ prediction is adjusted at the analysis level using data-driven techniques described below. Cross sections for diboson processes are known at NLO with an uncertainty of 5%, while the $t\bar{t}$ cross section is predicted at approximately NNLO, with better than 10% uncertainty [29,30].

V. DATA AND PRESELECTION

The data, which correspond to a total integrated luminosity of 2.05 fb^{-1} , were collected in 2011 during stable beam periods of 7 TeV pp collisions. For the e^* search, events are required to pass the lowest unrescaled single electron trigger available. For the first half of the data this corresponds to a p_T^e threshold of 20 GeV, and a p_T^e threshold of 22 GeV for the later runs. For the μ^* search, a single muon trigger with matching tracks in the muon spectrometer and inner detector with combined $p_T^\mu > 22 \text{ GeV}$ is used to select events. In addition, events with a muon with $p_T^\mu > 40 \text{ GeV}$ in the muon spectrometer are also kept. Collision candidates are then identified by requiring a primary vertex with a z position along the beam line of $|z| < 200 \text{ mm}$ and at least three associated charged particle tracks with $p_T > 0.4 \text{ GeV}$.

The lepton selection consists of the same requirements used in the ATLAS search for new heavy resonances decaying to dileptons [31]. Electron candidates are formed from clusters of cells in the electromagnetic calorimeter associated with a charged particle track in the inner detector. For the e^* search, two electron candidates with $p_T^e > 25 \text{ GeV}$ and $|\eta| < 2.47$ are required. Electrons within the transition region $1.37 < |\eta| < 1.52$ between the barrel and the endcap calorimeters are rejected. The *medium* electron identification criteria [32] on the transverse shower shape, the longitudinal leakage into the hadronic calorimeter, and the association with an inner detector track are applied to the cluster. The electron's reconstructed energy is obtained from the calorimeter measurement and its direction from the associated track. A hit in the first active pixel layer is required to suppress the background from photon conversions. To further suppress background from jets, the leading electron is required to be isolated by demanding that the sum of the transverse energies in the cells around the electron direction in a cone of radius $R < 0.2$ be less than 7 GeV. The core of the electron energy deposition is excluded, and the sum is corrected for transverse shower leakage and pileup from additional pp collisions to make the isolation variable essentially independent of p_T^e [33]. In cases where more than two electrons are found to satisfy the above requirements, the pair with the largest invariant mass is chosen. To minimize the impact of possible charge

misidentification, the electrons are not required to have opposite electric charges.

Muon tracks are reconstructed independently in both the inner detector and the muon spectrometer, and their momenta are determined from a combined fit to these two measurements. For the μ^* search, two muons with $p_T^\mu > 25 \text{ GeV}$ are required. To optimize the momentum resolution, each muon candidate is required to have a minimum number of hits in the inner detector and to have at least three hits in each of the inner, middle, and outer layers of the muon spectrometer. This requirement results in a muon fiducial acceptance of $|\eta| < 2.5$. Muons with hits in the barrel-endcap overlap regions of the muon spectrometer are discarded because of large residual misalignments. The effects of misalignments and intrinsic position resolution are otherwise included in the simulation. The p_T^μ resolution at 1 TeV ranges from 13% to 20%. To suppress background from cosmic rays, the muon tracks are required to have transverse and longitudinal impact parameters $|d_0| < 0.2 \text{ mm}$ and $|z_0| < 1 \text{ mm}$ with respect to the primary vertex. To reduce background from heavy flavor hadrons, each muon is required to be isolated such that $\sum p_T(R < 0.3)/p_T^\mu < 0.05$, where only inner detector tracks with $p_T > 1 \text{ GeV}$ enter the sum. Muons are required to have opposite electric charges. In cases where more than two muons are found to satisfy the above requirements, the pair of muons with the largest invariant mass is considered.

The dielectron and dimuon distributions are inspected for consistency with background predictions to ensure that the resolution and efficiency corrections were adjusted properly in the simulation. Excellent agreement is found around the mass of the Z , in terms of both the peak position and width of the dilepton invariant mass distributions. For the mass range $70 < m_{\ell\ell} < 110 \text{ GeV}$, the number of events observed in data agrees to within 1% of the background predictions for both the electron and muon channels. Furthermore, the tails of the p_T^e and p_T^μ distributions in the simulation are found to closely match the data.

The presence of at least one photon candidate with $p_T^\gamma > 20 \text{ GeV}$ and pseudorapidity $|\eta| < 2.37$ is then necessary for the events to be kept. Photons within the transition region between the barrel and the endcap calorimeters are excluded. Photon candidates are formed from clusters of cells in the electromagnetic calorimeter. They include unconverted photons, with no associated track, and photons that converted to electron-positron pairs, associated with one or two tracks. All photon candidates are required to satisfy the *tight* photon definition [34]. This selection includes constraints on the energy leakage into the hadronic calorimeter as well as stringent requirements on the energy distribution in the first sampling layer of the electromagnetic calorimeter, and on the shower width in the second sampling layer. The *tight* photon definition is designed to increase the purity of the photon selection sample by rejecting most of the jet background, including jets with

a leading neutral hadron (usually a π^0) that decays to a pair of collimated photons. To further reduce background from misidentified jets, photon candidates are required to be isolated by demanding that the sum of the transverse energies of the cells within a cone $R < 0.4$ of the photon be less than 10 GeV. As for the electron isolation, the core of the photon energy deposition is excluded and the sum is corrected for transverse shower leakage and pileup. Because no background predictions are simulated for $R(\ell, \gamma) < 0.5$, photons are required to be well separated from the leptons with $R(\ell, \gamma) > 0.7$. This requirement has a negligible impact on signal efficiency. Finally, if more than one photon in an event satisfies the above requirements, the one with the largest p_T is used in the search.

For the above selection criteria, the total signal acceptance times efficiency ($A \times \epsilon$) is $\sim 56\%$ in the e^* channel for masses $m_{e^*} > 600$ GeV. This value includes the acceptance of all selection cuts and the reconstruction efficiencies, and reflects the lepton and photon angular distributions. In comparison, $A \times \epsilon$ is $\sim 32\%$ for $m_{\mu^*} > 600$ GeV. The lower acceptance in the μ^* channel is due to the stringent selection on the muon spectrometer hits used to maximize the p_T^μ resolution, in particular, the limited geometrical coverage of the muon spectrometer with three layers of precision chambers.

VI. BACKGROUND DETERMINATION

All background predictions are evaluated with simulated samples. These include the dominant and irreducible $Z + \gamma$ background, as well as $Z + \text{jets}$ events where a jet is misidentified as a photon. The rate of jet misidentification is overestimated in the simulation so the $Z + \text{jets}$ predictions are adjusted to data as described below. Small contributions from $t\bar{t}$ and diboson production are also present at low $m_{\ell\ell\gamma}$. Background from multijet events and semileptonic decays of heavy flavor hadrons are heavily suppressed by the isolation requirements and are negligible in the signal region.

The $Z + \text{jets}$ estimates are adjusted to data in a control region defined by $m_{\ell\ell\gamma} < 300$ GeV. This region represents less than 1% of the signal parameter space for $m_{e^*} \geq 200$ GeV. The nominal strategy consists of counting the number of events in data in this control region and comparing it to the MC background predictions. The excess of background events found in the simulation is attributed to the mismodeling of the rate of jets misidentified as photons, and the number of $Z + \text{jets}$ events is scaled down accordingly. As a result, the number of events in the control region is the same in the MC simulations as in data, as shown in Table I. The $Z + \text{jets}$ estimates are validated using various data-driven methods, notably by using misidentification rates evaluated in jet-enriched samples, and applying these rates to $Z + \text{jets}$ data samples using an approach similar to the one described in Ref. [34]. The main reason for the overestimation of

TABLE I. Data yields and background expectations inside ($m_{\ell\ell\gamma} < 300$ GeV) and outside the $m_{\ell\ell\gamma}$ control region after adjusting the $Z + \text{jets}$ background. The uncertainties shown are purely statistical, except for the $Z + \text{jets}$ background for which the total uncertainty is dominated by systematic uncertainties.

Region (GeV)	$Z + \gamma$	$Z + \text{jets}$	Diboson	$t\bar{t}$	Data
$m_{e\ell\gamma} < 300$	306 ± 8	138 ± 38	8.3 ± 0.8	2.4 ± 0.5	455
$m_{e\ell\gamma} > 300$	25 ± 2	8.1 ± 1.6	0.8 ± 0.2	0.5 ± 0.2	29
$m_{\mu\mu\gamma} < 300$	255 ± 8	89 ± 31	4.9 ± 0.6	0.9 ± 0.3	350
$m_{\mu\mu\gamma} > 300$	14 ± 1	5.4 ± 1.4	0.9 ± 0.3	0.1 ± 0.1	19

the jet misidentification rate in the simulation is due to the mismodeling of the jet shower shapes. A $Z + \text{jets}$ enriched sample was used to correct the shower shapes of jets in the simulations, such that the efficiency for jets to pass the *tight* photon requirement in the MC simulation is comparable to the rate measured in data. This correction depends strongly on the generator used (e.g. PYTHIA vs ALPGEN) and results in a 15% uncertainty in the $Z + \text{jets}$ background estimate.

The largest difference between the nominal $Z + \text{jets}$ background determination and the alternative estimates is

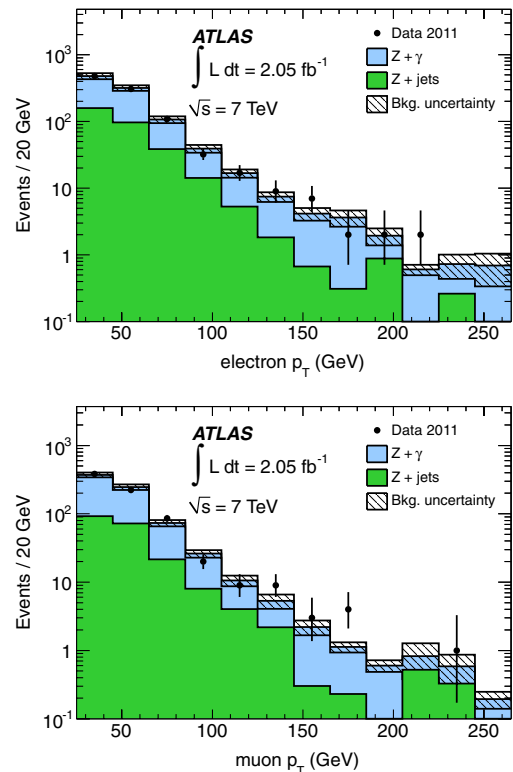


FIG. 1 (color online). Lepton p_T distributions for the e^* (top panel) and μ^* (bottom panel) channels. The expected background uncertainties shown correspond to the sum in quadrature of the statistical uncertainties as well as the uncertainty in the $Z + \text{jets}$ normalization measured in the control region.

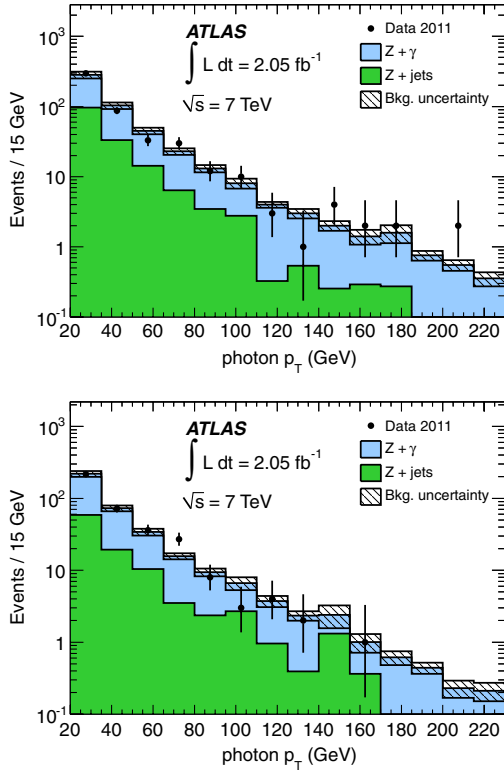


FIG. 2 (color online). Photon p_T distributions for the e^* (top panel) and μ^* (bottom panel) channels. The expected background uncertainties shown correspond to the sum in quadrature of the statistical uncertainties as well as the uncertainty in the $Z + \text{jets}$ normalization measured in the control region.

assigned as a systematic uncertainty and dominates the total error in the $Z + \text{jets}$ estimates presented in Table I. The corresponding scaling factors applied to the $Z + \text{jets}$ simulation are 0.51 ± 0.14 and 0.61 ± 0.21 for the e^* and μ^* channels, respectively, i.e. within uncertainties of one another. Furthermore, the ratio of the number of $Z + \text{jets}$ events outside the control region to the number of events inside is found to be the same in the MC simulations as in the data-driven techniques: 0.06 for both the e^* and μ^* channels. This finding indicates that the jet p_T misidentification rate as a function of the jet is modeled properly.

Comparisons between data and the resulting background expectations for the p_T^ℓ , p_T^γ , $m_{\ell\gamma}$, and $m_{\ell\ell\gamma}$ distributions are shown in Figs. 1–4. No significant discrepancies are observed between data and the simulations. In particular, the background prediction for the photon p_T shape matches the data for both the e^* and μ^* searches, which suggests that the tuning of the jet misidentification rate for the $Z + \text{jets}$ background is adequate.

VII. SIGNAL REGION OPTIMIZATION

The signal search region is optimized as a function of m_{ℓ^*} using simulated events by determining the lower bound on $m_{\ell\ell\gamma}$ that maximizes the significance defined as

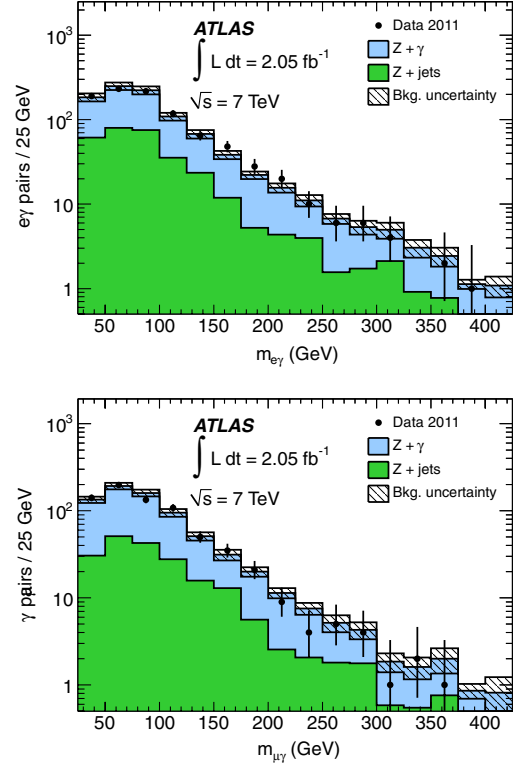


FIG. 3 (color online). Distributions of the invariant mass of the $\ell\gamma$ systems for the e^* (top panel) and μ^* (bottom panel) channels. Combinations with both the leading and subleading leptons are shown. The expected background uncertainties shown correspond to the sum in quadrature of the statistical uncertainties as well as the uncertainty in the $Z + \text{jets}$ normalization measured in the control region. For both channels, one event lies outside the mass range shown.

$$S_L = \sqrt{2 \ln[(1 + S/B)^{S+B} e^{-S}]},$$

where S and B are the number of signal and background events, respectively. The optimum threshold value is found to be $m_{\ell\ell\gamma} = m_{\ell^*} + 150$ GeV. Additionally, to improve the sensitivity, particularly at low m_{ℓ^*} , background contributions from DY processes are suppressed further by requiring events to satisfy $m_{\ell\ell} > 110$ GeV. The signal efficiency for these two additional requirements is $>99\%$ for $m_{\ell^*} \geq 200$ GeV.

Because few events survive the complete set of requirements, the shapes of the $Z + \gamma$ and $Z + \text{jets}$ backgrounds are individually fitted using an exponential function $\exp(P_0 + P_1 \times m_{\ell\ell\gamma})$ over the mass range $250 \text{ GeV} < m_{\ell\ell\gamma} < 950 \text{ GeV}$. The sum of these two fits is then used to obtain the total background prediction for $m_{\ell\ell\gamma} > 350$ GeV. The resulting background estimates and data yields are shown in Table II for the e^* and μ^* searches, as well as in Figs. 5 and 6.

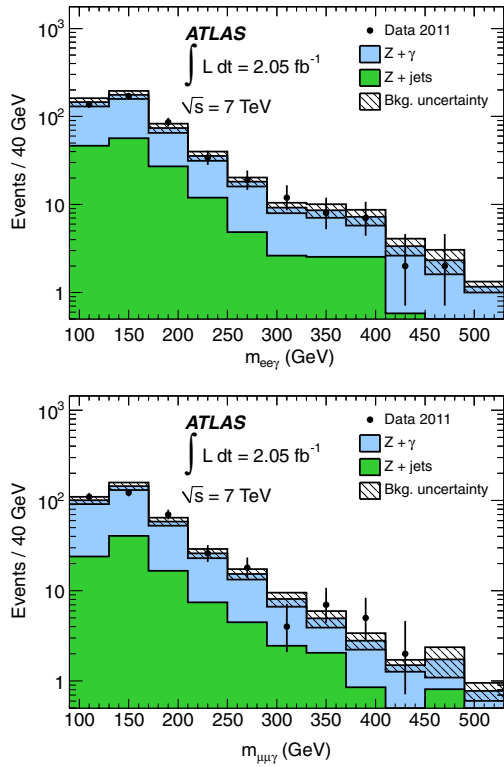


FIG. 4 (color online). Distributions of the invariant mass for the $l\ell\gamma$ system for the e^* (top panel) and μ^* (bottom panel) channels. The expected background uncertainties shown correspond to the sum in quadrature of the statistical uncertainties as well as the uncertainty in the $Z + \gamma$ and $Z + \text{jets}$ normalization measured in the control region. For both channels, one event lies outside the mass range shown.

VIII. SYSTEMATIC UNCERTAINTIES

In this section, the dominant systematic uncertainties in the $Z + \gamma$ and $Z + \text{jets}$ background predictions are first described, followed by a description of the experimental systematic uncertainties that affect both the background and signal yields, and by a discussion of the theoretical uncertainties which affect both the e^* and μ^* .

The dominant systematic uncertainty in the irreducible $Z + \gamma$ background comes from the fit of its background shape and normalization due to the limited number of events with $m_{\ell\ell} > 110$ GeV. This uncertainty increases with m_{ℓ^*} from about 20% at 200 GeV to 100% for $m_{\ell^*} > 800$ GeV. The second largest uncertainty in the $Z + \gamma$ background is of theoretical nature and arises from the NLO computations. This uncertainty is obtained by varying the renormalization and factorization scales by factors of 2 around their nominal values and combining with uncertainties arising from the PDFs and values of the strong coupling constant α_s . For $m_{\ell^*} = 200$ GeV ($m_{\ell^*} > 800$ GeV), the resulting theoretical uncertainty in the number of $Z + \gamma$ background events in our signal region is 7% (10%) for both channels.

The uncertainty in the $Z + \text{jets}$ normalization is determined to be 38% (35%) for the e^* (μ^*) channel, which covers the range of values obtained by the different estimates as well as their uncertainties in the $m_{\ell\ell\gamma} < 300$ GeV control region. Uncertainties in the $Z + \text{jets}$ prediction from the shape of the fitted distribution are added in quadrature to the normalization uncertainty.

Experimental systematic uncertainties that affect both signal and background yields include the uncertainty from the luminosity measurement of 3.7% [35] and uncertainties in particle reconstruction and identification as described below.

A 3% systematic uncertainty is assigned to the photon efficiency. This value is obtained by comparing the signal efficiency with and without photon shower shape corrections (2%), by studying the impact of material mismodeling in the inner detector (1%), and by determining the reconstruction efficiency for various pileup conditions (1%) [36].

The electron trigger and reconstruction efficiency is evaluated in data and in MC simulations in several $\eta \times \phi$ bins to high precision. Correction factors are applied to the simulations accordingly and have negligible uncertainties. A 1% systematic uncertainty in the electron efficiency at high p_T is assigned. This uncertainty is

TABLE II. Data yields and background expectation as a function of a lower bound on $m_{\ell\ell\gamma} = m_{\ell^*} + 150$ GeV. The uncertainties represent the sum in quadrature of the statistical and systematic uncertainties. The probability for the background-only hypothesis (p value) is also provided.

$m_{\ell\ell\gamma}$ region (TeV)	e^* search				μ^* search			
	$Z + \gamma$	Total bkg	Data	p value	$Z + \gamma$	Total bkg	Data	p value
>0.35	10.1 ± 1.9	11.5 ± 2.2	8	0.92	5.2 ± 1.4	6.0 ± 1.6	6	0.40
>0.45	4.6 ± 1.0	5.1 ± 1.2	2	0.83	3.1 ± 0.8	3.4 ± 0.9	3	0.42
>0.55	2.1 ± 0.7	2.3 ± 0.8	1	0.80	1.8 ± 0.6	2.0 ± 0.7	1	0.72
>0.65	0.98 ± 0.47	1.02 ± 0.49	1	0.32	1.09 ± 0.49	1.14 ± 0.51	1	0.72
>0.75	0.45 ± 0.29	0.46 ± 0.30	1	0.16	0.65 ± 0.39	0.67 ± 0.39	1	0.28
>0.85	0.20 ± 0.16	0.21 ± 0.17	1	0.11	0.39 ± 0.29	0.39 ± 0.29	1	0.17
>0.95	0.09 ± 0.09	0.10 ± 0.09	1	0.03	0.23 ± 0.21	0.23 ± 0.21	0	0.78
>1.05	0.05 ± 0.05	0.05 ± 0.05	0	0.81	0.14 ± 0.14	0.14 ± 0.14	0	0.92

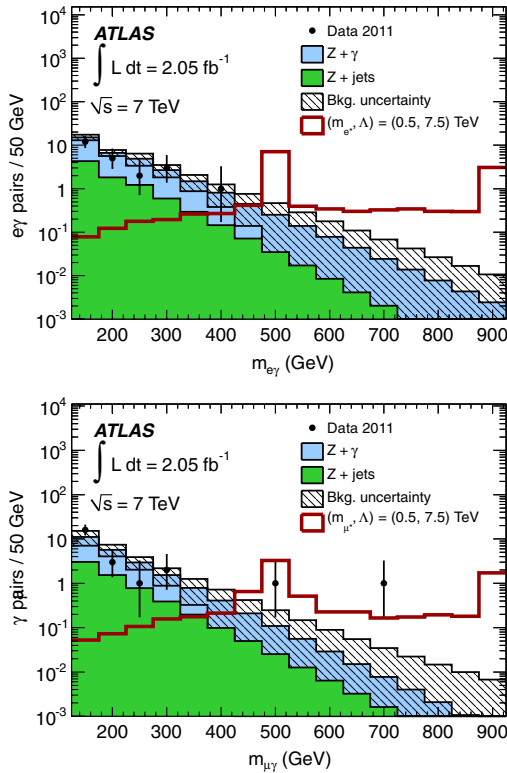


FIG. 5 (color online). Distributions of the invariant mass of the $\ell\gamma$ systems for the e^* (top panel) and μ^* (bottom panel) channels after requiring $m_{\ell\ell} > 110$ GeV. Combinations with both the leading and subleading leptons are shown. The expected background uncertainties shown correspond to the sum in quadrature of the statistical uncertainties as well as the uncertainty in the $Z + \text{jets}$ normalization measured in the control region. Note that the last bin contains the sum of all entries with $m_{\ell\gamma} > 950$ GeV.

estimated by studying the electron efficiency as a function of the calorimeter isolation criteria.

The calorimeter energy resolution is dominated at high p_T by a constant term which is 1.1% in the barrel and 1.8% in the endcaps. The simulation is adjusted to reproduce this resolution at high energy, and the uncertainty in this correction has a negligible effect on p_T^e and p_T^γ . The calorimeter energy scale is corrected by studying $J/\psi \rightarrow ee$ and $Z \rightarrow ee$ events. Calibration constants are obtained for several η regions and deviate at most by 1.5% of unity, and have small uncertainties. Thus, uncertainties on the calorimeter energy scale and resolution result in negligible uncertainties in the background and signal yields.

The combined uncertainty in yields arising from the trigger and reconstruction efficiency for muons is estimated to increase linearly as a function of p_T^μ to about 1.5% at 1 TeV. This uncertainty is dominated by a conservative estimate of the impact of large energy loss from muon bremsstrahlung in the calorimeter, which can affect reconstruction in the muon spectrometer. The uncertainty from the resolution due to residual misalignments in the

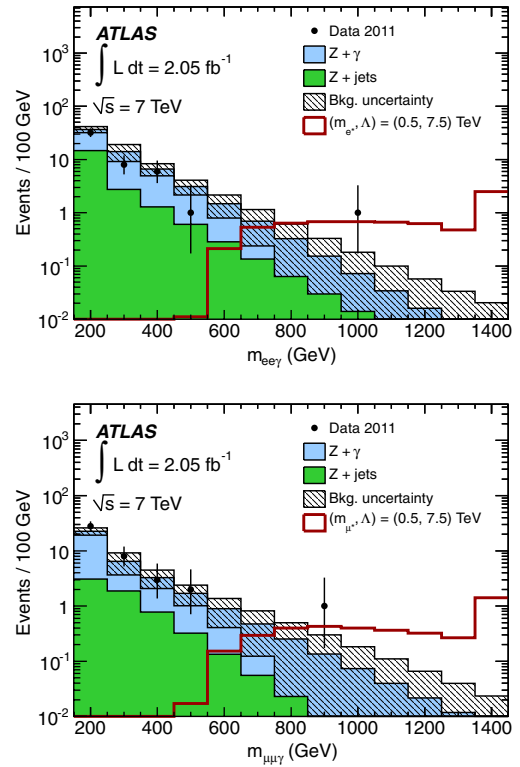


FIG. 6 (color online). Distributions of the invariant mass for the $\ell\ell\gamma$ system for the e^* (top panel) and μ^* (bottom panel) searches after requiring $m_{\ell\ell} > 110$ GeV. The $Z + \text{jets}$ and $Z + \gamma$ backgrounds were fitted, and the total uncertainties from the fit as well as the uncertainty in the $Z + \text{jets}$ normalization measured in the control region are displayed as the shaded area. Note that the last bin contains the sum of all events with $m_{\ell\ell\gamma} > 1450$ GeV.

muon spectrometer propagates to a change in the number of events passing the $m_{\mu\mu\gamma}$ cut and affects the sensitivity of the search. The muon momentum scale is calibrated with a statistical precision of 0.1% using the $Z \rightarrow \mu\mu$ mass peak. Thus, uncertainties on the muon momentum scale and resolution result in negligible uncertainties in the background and signal yields.

An additional 1% systematic uncertainty is assigned to the e^* and μ^* signal efficiencies to account for the fact that the dependence on Λ is neglected in this analysis. This uncertainty is obtained by studying the signal $A \times \epsilon$ for various excited lepton masses and compositeness scales. Theoretical uncertainties from renormalization and factorization scales and PDFs have a negligible impact on the signal efficiency and are not included in the results presented below.

IX. RESULTS

A summary of the data yields and background expectations as a function of a lower bound on $m_{\ell\ell\gamma}$ is shown in Table II for the e^* and μ^* searches. The uncertainties

displayed correspond to the sum in quadrature of the statistical and systematic uncertainties. The significance for an excited lepton signal is estimated by means of a p value, the probability of observing an outcome at least as signal-like as the one observed in data, assuming that a signal is absent. The lowest p values obtained are 3% in the e^* channel (for $m_{ee\gamma} > 950$ GeV) and 17% in the μ^* channel (for $m_{\mu\mu\gamma} > 850$ GeV), which indicates that the data are consistent with the background hypothesis.

Given the absence of a signal, an upper limit on the ℓ^* cross section times branching ratio σB is determined at the 95% C.L. using a Bayesian approach [37] with a flat, positive prior on σB . Systematic uncertainties are incorporated in the limit calculation as nuisance parameters. The limits are translated into bounds on the compositeness scale as a function of the mass of the excited leptons by comparing them with theoretical predictions of σB for various values of Λ .

The expected exclusion limits are determined using simulated pseudoexperiments (PE) containing only SM

processes, by evaluating the 95% C.L. upper limits for each PE for each fixed value of m_{ℓ^*} . The median of the distribution of limits represents the expected limit. The ensemble of limits is used to find the 1σ and 2σ envelopes of the expected limits as a function of m_{ℓ^*} .

Figure 7 shows the 95% C.L. expected and observed limits on $\sigma B(\ell^* \rightarrow \ell\gamma)$ for the e^* and μ^* searches. For $m_{\ell^*} > 0.9$ TeV, the observed and expected limits on σB are 2.3 fb and 4.5 fb for the e^* and μ^* , respectively. The green and yellow bands show the expected 1σ and 2σ contours of the expected limits. When the expected number of background events is zero, there is an effective quantization of the expected limits obtained from the PE, and no downward fluctuation of the background is possible. These effects explain the behavior of the 1σ and 2σ contours of the expected limits for large ℓ^* masses. Theoretical predictions of σB for three different values of Λ are also displayed in Fig. 7, as well as the theoretical uncertainties from renormalization and factorization scales and PDFs for $\Lambda = 2$ TeV. These uncertainties are shown for illustrative

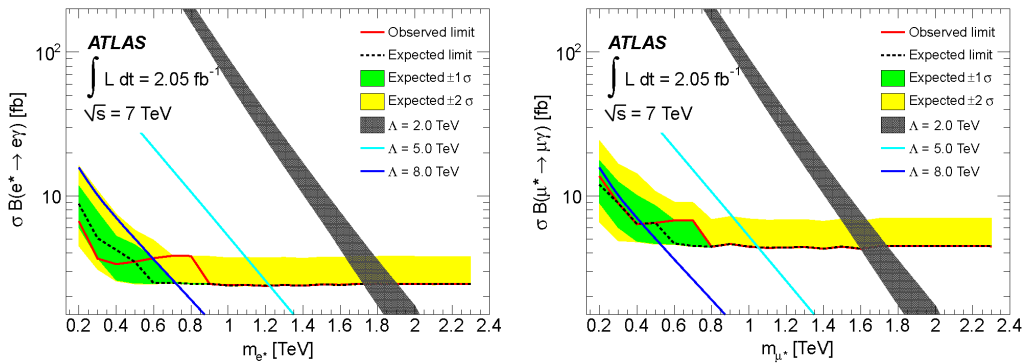


FIG. 7 (color online). Limits at 95% C.L. on the cross section times branching ratio as a function of e^* and of μ^* mass. Theoretical predictions for excited leptons produced for three different compositeness scales are shown, as well as the theoretical uncertainties from renormalization and factorization scales and PDFs for $\Lambda = 2$ TeV. For $m_{\ell^*} > 0.9$ TeV, the observed limit on σB is 2.3 fb (4.5 fb) for e^* (μ^*).

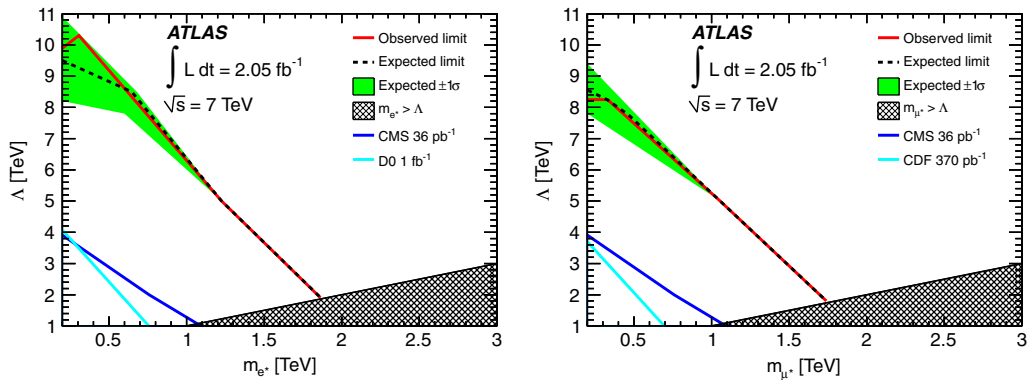


FIG. 8 (color online). Exclusion limits in the $m_{\ell^*} - \Lambda$ parameter space for e^* and μ^* . Regions to the left of the experimental limits are excluded at 95% C.L. No limits are set for the hashed region, as the approximations made in the effective contact interaction model do not hold for $m_{\ell^*} > \Lambda$. The best limits from the Tevatron experiments as well as from the CMS experiment based on 36 pb^{-1} are also shown.

purposes only and are not included in determining mass limits. The mass limits obtained for various Λ values are used to produce exclusion limits on the $m_{\ell^*} - \Lambda$ plane as shown in Fig. 8. In the special case where $\Lambda = m_{\ell^*}$, masses below 1.87 TeV and 1.75 TeV are excluded for excited electrons and muons, respectively.

X. CONCLUSIONS

The results of a search for excited electrons and muons with the ATLAS detector are reported, using a sample of $\sqrt{s} = 7$ TeV pp collisions corresponding to an integrated luminosity of 2.05 fb^{-1} . The observed invariant mass spectra are consistent with SM background expectations. Limits are set on the cross section times branching ratio $\sigma B(\ell^* \rightarrow \ell \gamma)$ at 95% C.L. For $m_{\ell^*} > 0.9$ TeV, the observed upper limits on σB are 2.3 fb and 4.5 fb for the e^* and μ^* channels, respectively. The limits are translated into bounds on the compositeness scale Λ as a function of the mass of the excited leptons. In the special case where $\Lambda = m_{\ell^*}$, masses below 1.87 TeV and 1.75 TeV are excluded for e^* and μ^* , respectively. These limits are the most stringent bounds to date on excited leptons for the parameter-space region with $m_{\ell^*} \geq 200$ GeV.

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G. Aad,⁴⁷ B. Abbott,¹¹⁰ J. Abdallah,¹¹ A. A. Abdelalim,⁴⁸ A. Abdesselam,¹¹⁷ O. Abdinov,¹⁰ B. Abi,¹¹¹ M. Abolins,⁸⁷ O. S. AbouZeid,¹⁵⁷ H. Abramowicz,¹⁵² H. Abreu,¹¹⁴ E. Acerbi,^{88a,88b} B. S. Acharya,^{163a,163b} L. Adamczyk,³⁷ D. L. Adams,²⁴ T. N. Addy,⁵⁵ J. Adelman,¹⁷⁴ M. Aderholz,⁹⁸ S. Adomeit,⁹⁷ P. Adragna,⁷⁴ T. Adye,¹²⁸ S. Aefsky,²² J. A. Aguilar-Saavedra,^{123b,b} M. Aharrouche,⁸⁰ S. P. Ahlen,²¹ F. Ahles,⁴⁷ A. Ahmad,¹⁴⁷ M. Ahsan,⁴⁰ G. Aielli,^{132a,132b} T. Akdogan,^{18a} T. P. A. Åkesson,⁷⁸ G. Akimoto,¹⁵⁴ A. V. Akimov,⁹³ A. Akiyama,⁶⁶ M. S. Alam,¹ M. A. Alam,⁷⁵ J. Albert,¹⁶⁸ S. Albrand,⁵⁴ M. Aleksa,²⁹ I. N. Aleksandrov,⁶⁴ F. Alessandria,^{88a} C. Alexa,^{25a} G. Alexander,¹⁵² G. Alexandre,⁴⁸ T. Alexopoulos,⁹ M. Alhroob,²⁰ M. Aliev,¹⁵ G. Alimonti,^{88a} J. Alison,¹¹⁹ M. Aliyev,¹⁰ B. M. M. Allbrooke,¹⁷ P. P. Allport,⁷² S. E. Allwood-Spiers,⁵² J. Almond,⁸¹ A. Aloisio,^{101a,101b} R. Alon,¹⁷⁰ A. Alonso,⁷⁸ B. Alvarez Gonzalez,⁸⁷ M. G. Alviggi,^{101a,101b} K. Amako,⁶⁵ P. Amaral,²⁹ C. Amelung,²² V. V. Ammosov,¹²⁷ A. Amorim,^{123a,c} G. Amorós,¹⁶⁶ N. Amram,¹⁵² C. Anastopoulos,²⁹ L. S. Ancu,¹⁶ N. Andari,¹¹⁴ T. Andeen,³⁴ C. F. Anders,²⁰ G. Anders,^{57a} K. J. Anderson,³⁰ A. Andreazza,^{88a,88b} V. Andrei,^{57a} M-L. Andrieux,⁵⁴ X. S. Anduaga,⁶⁹ A. Angerami,³⁴ F. Anghinolfi,²⁹ A. Anisenkov,¹⁰⁶ N. Anjos,^{123a} A. Annovi,⁴⁶ A. Antonaki,⁸ M. Antonelli,⁴⁶ A. Antonov,⁹⁵ J. Antos,^{143b} F. Anulli,^{131a} S. Aoun,⁸² L. Aperio Bella,⁴ R. Apolle,^{117,d} G. Arabidze,⁸⁷ I. Aracena,¹⁴² Y. Arai,⁶⁵ A. T. H. Arce,⁴⁴ S. Arfaoui,¹⁴⁷ J-F. Arguin,¹⁴ E. Arik,^{18a,a} M. Arik,^{18a} A. J. Armbruster,⁸⁶ O. Arnaez,⁸⁰ C. Arnault,¹¹⁴ A. Artamonov,⁹⁴ G. Artoni,^{131a,131b} D. Arutinov,²⁰ S. Asai,¹⁵⁴ R. Asfandiyarov,¹⁷¹ S. Ask,²⁷ B. Åsman,^{145a,145b} L. Asquith,⁵ K. Assamagan,²⁴ A. Astbury,¹⁶⁸ A. Astvatsatourov,⁵¹ B. Aubert,⁴ E. Auge,¹¹⁴ K. Augsten,¹²⁶ M. Aourousseau,^{144a} G. Avolio,¹⁶² R. Avramidou,⁹ D. Axen,¹⁶⁷ C. Ay,⁵³ G. Azuelos,^{92,e} Y. Azuma,¹⁵⁴ M. A. Baak,²⁹ G. Baccaglioni,^{88a} C. Bacci,^{133a,133b} A. M. Bach,¹⁴ H. Bachacou,¹³⁵ K. Bachas,²⁹ M. Backes,⁴⁸ M. Backhaus,²⁰ E. Badescu,^{25a} P. Bagnaia,^{131a,131b} S. Bahinipati,² Y. Bai,^{32a} D. C. Bailey,¹⁵⁷ T. Bain,¹⁵⁷ J. T. Baines,¹²⁸ O. K. Baker,¹⁷⁴ M. D. Baker,²⁴ S. Baker,⁷⁶ E. Banas,³⁸ P. Banerjee,⁹² Sw. Banerjee,¹⁷¹ D. Banfi,²⁹ A. Bangert,¹⁴⁹ V. Bansal,¹⁶⁸ H. S. Bansil,¹⁷ L. Barak,¹⁷⁰ S. P. Baranov,⁹³ A. Barashkou,⁶⁴ A. Barbaro Galtieri,¹⁴ T. Barber,⁴⁷ E. L. Barberio,⁸⁵ D. Barberis,^{49a,49b} M. Barbero,²⁰ D. Y. Bardin,⁶⁴ T. Barillari,⁹⁸ M. Barisonzi,¹⁷³ T. Barklow,¹⁴² N. Barlow,²⁷ B. M. Barnett,¹²⁸ R. M. Barnett,¹⁴ A. Baroncelli,^{133a} G. Barone,⁴⁸ A. J. Barr,¹¹⁷ F. Barreiro,⁷⁹ J. Barreiro Guimarães da Costa,⁵⁶ P. Barrillon,¹¹⁴ R. Bartoldus,¹⁴² A. E. Barton,⁷⁰ V. Bartsch,¹⁴⁸ R. L. Bates,⁵² L. Batkova,^{143a} J. R. Batley,²⁷ A. Battaglia,¹⁶ M. Battistin,²⁹ F. Bauer,¹³⁵ H. S. Bawa,^{142,f} S. Beale,⁹⁷ T. Beau,⁷⁷ P. H. Beauchemin,¹⁶⁰ R. Beccherle,^{49a} P. Bechtel,²⁰ H. P. Beck,¹⁶ S. Becker,⁹⁷ M. Beckingham,¹³⁷ K. H. Becks,¹⁷³ A. J. Beddall,^{18c} A. Beddall,^{18c} S. Bedikian,¹⁷⁴ V. A. Bednyakov,⁶⁴ C. P. Bee,⁸² M. Begel,²⁴ S. Behar Harpaz,¹⁵¹ P. K. Behera,⁶² M. Beimforde,⁹⁸ C. Belanger-Champagne,⁸⁴ P. J. Bell,⁴⁸ W. H. Bell,⁴⁸ G. Bella,¹⁵² L. Bellagamba,^{19a} F. Bellina,²⁹ M. Bellomo,²⁹ A. Belloni,⁵⁶ O. Beloborodova,^{106,g} K. Belotskiy,⁹⁵ O. Beltramello,²⁹ S. Ben Ami,¹⁵¹ O. Benary,¹⁵² D. Benchekroun,^{134a} C. Benchouk,⁸² M. Bendel,⁸⁰ N. Benekos,¹⁶⁴

- Y. Benhammou,¹⁵² E. Benhar Noccioli,⁴⁸ J. A. Benitez Garcia,^{158b} D. P. Benjamin,⁴⁴ M. Benoit,¹¹⁴ J. R. Bensinger,²² K. Benslama,¹²⁹ S. Bentvelsen,¹⁰⁴ D. Berge,²⁹ E. Bergeaas Kuutmann,⁴¹ N. Berger,⁴ F. Berghaus,¹⁶⁸ E. Berglund,¹⁰⁴ J. Beringer,¹⁴ P. Bernat,⁷⁶ R. Bernhard,⁴⁷ C. Bernius,²⁴ T. Berry,⁷⁵ C. Bertella,⁸² A. Bertin,^{19a,19b} F. Bertinelli,²⁹ F. Bertolucci,^{121a,121b} M. I. Besana,^{88a,88b} N. Besson,¹³⁵ S. Bethke,⁹⁸ W. Bhimji,⁴⁵ R. M. Bianchi,²⁹ M. Bianco,^{71a,71b} O. Biebel,⁹⁷ S. P. Bieniek,⁷⁶ K. Bierwagen,⁵³ J. Biesiada,¹⁴ M. Biglietti,^{133a} H. Bilokon,⁴⁶ M. Bindi,^{19a,19b} S. Binet,¹¹⁴ A. Bingul,^{18c} C. Bini,^{131a,131b} C. Biscarat,¹⁷⁶ U. Bitenc,⁴⁷ K. M. Black,²¹ R. E. Blair,⁵ J.-B. Blanchard,¹³⁵ G. Blanchot,²⁹ T. Blazek,^{143a} C. Blocker,²² J. Blocki,³⁸ A. Blondel,⁴⁸ W. Blum,⁸⁰ U. Blumenschein,⁵³ G. J. Bobbink,¹⁰⁴ V. B. Bobrovnikov,¹⁰⁶ S. S. Bocchetta,⁷⁸ A. Bocci,⁴⁴ C. R. Boddy,¹¹⁷ M. Boehler,⁴¹ J. Boek,¹⁷³ N. Boelaert,³⁵ J. A. Bogaerts,²⁹ A. Bogdanchikov,¹⁰⁶ A. Bogouch,^{89a} C. Bohm,^{145a} V. Boisvert,⁷⁵ T. Bold,³⁷ V. Boldea,^{25a} N. M. Bolnet,¹³⁵ M. Bona,⁷⁴ V. G. Bondarenko,⁹⁵ M. Bondioli,¹⁶² M. Boonekamp,¹³⁵ C. N. Booth,¹³⁸ S. Bordini,⁷⁷ C. Borer,¹⁶ A. Borisov,¹²⁷ G. Borissov,⁷⁰ I. Borjanovic,^{12a} M. Borri,⁸¹ S. Borroni,⁸⁶ V. Bortolotto,^{133a,133b} K. Bos,¹⁰⁴ D. Boscherini,^{19a} M. Bosman,¹¹ H. Boterenbrood,¹⁰⁴ D. Botterill,¹²⁸ J. Bouchami,⁹² J. Boudreau,¹²² E. V. Bouhova-Thacker,⁷⁰ D. Boumediene,³³ C. Bourdarios,¹¹⁴ N. Bousson,⁸² A. Boveia,³⁰ J. Boyd,²⁹ I. R. Boyko,⁶⁴ N. I. Bozhko,¹²⁷ I. Bozovic-Jelisavcic,^{12b} J. Bracinek,¹⁷ A. Braem,²⁹ P. Branchini,^{133a} G. W. Brandenburg,⁵⁶ A. Brandt,⁷ G. Brandt,¹¹⁷ O. Brandt,⁵³ U. Bratzler,¹⁵⁵ B. Brau,⁸³ J. E. Brau,¹¹³ H. M. Braun,¹⁷³ B. Brelier,¹⁵⁷ J. Bremer,²⁹ R. Brenner,¹⁶⁵ S. Bressler,¹⁷⁰ D. Breton,¹¹⁴ D. Britton,⁵² F. M. Brochu,²⁷ I. Brock,²⁰ R. Brock,⁸⁷ T. J. Brodbeck,⁷⁰ E. Brodet,¹⁵² F. Broggi,^{88a} C. Bromberg,⁸⁷ J. Bronner,⁹⁸ G. Brooijmans,³⁴ W. K. Brooks,^{31b} G. Brown,⁸¹ H. Brown,⁷ P. A. Bruckman de Renstrom,³⁸ D. Bruncko,^{143b} R. Bruneliere,⁴⁷ S. Brunet,⁶⁰ A. Bruni,^{19a} G. Bruni,^{19a} M. Bruschi,^{19a} T. Buanes,¹³ Q. Buat,⁵⁴ F. Bucci,⁴⁸ J. Buchanan,¹¹⁷ N. J. Buchanan,² P. Buchholz,¹⁴⁰ R. M. Buckingham,¹¹⁷ A. G. Buckley,⁴⁵ S. I. Buda,^{25a} I. A. Budagov,⁶⁴ B. Budick,¹⁰⁷ V. Büscher,⁸⁰ L. Bugge,¹¹⁶ O. Bulekov,⁹⁵ M. Bunse,⁴² T. Buran,¹¹⁶ H. Burckhart,²⁹ S. Burdin,⁷² T. Burgess,¹³ S. Burke,¹²⁸ E. Busato,³³ P. Bussey,⁵² C. P. Buszello,¹⁶⁵ F. Butin,²⁹ B. Butler,¹⁴² J. M. Butler,²¹ C. M. Buttar,⁵² J. M. Butterworth,⁷⁶ W. Buttinger,²⁷ S. Cabrera Urbán,¹⁶⁶ D. Caforio,^{19a,19b} O. Cakir,^{3a} P. Calafiura,¹⁴ G. Calderini,⁷⁷ P. Calfayan,⁹⁷ R. Calkins,¹⁰⁵ L. P. Caloba,^{23a} R. Caloi,^{131a,131b} D. Calvet,³³ S. Calvet,³³ R. Camacho Toro,³³ P. Camarri,^{132a,132b} M. Cambiaghi,^{118a,118b} D. Cameron,¹¹⁶ L. M. Caminada,¹⁴ S. Campana,²⁹ M. Campanelli,⁷⁶ V. Canale,^{101a,101b} F. Canelli,^{30,h} A. Canepa,^{158a} J. Cantero,⁷⁹ L. Capasso,^{101a,101b} M. D. M. Capeans Garrido,²⁹ I. Caprini,^{25a} M. Caprini,^{25a} D. Capriotti,⁹⁸ M. Capua,^{36a,36b} R. Caputo,⁸⁰ C. Caramarcu,²⁴ R. Cardarelli,^{132a} T. Carli,²⁹ G. Carlino,^{101a} L. Carminati,^{88a,88b} B. Caron,⁸⁴ S. Caron,¹⁰³ G. D. Carrillo Montoya,¹⁷¹ A. A. Carter,⁷⁴ J. R. Carter,²⁷ J. Carvalho,^{123a,i} D. Casadei,¹⁰⁷ M. P. Casado,¹¹ M. Cascella,^{121a,121b} C. Caso,^{49a,49b,a} A. M. Castaneda Hernandez,¹⁷¹ E. Castaneda-Miranda,¹⁷¹ V. Castillo Gimenez,¹⁶⁶ N. F. Castro,^{123a} G. Cataldi,^{71a} F. Cataneo,²⁹ A. Catinaccio,²⁹ J. R. Catmore,²⁹ A. Cattai,²⁹ G. Cattani,^{132a,132b} S. Caughron,⁸⁷ D. Cauz,^{163a,163c} P. Cavalleri,⁷⁷ D. Cavalli,^{88a} M. Cavalli-Sforza,¹¹ V. Cavasinni,^{121a,121b} F. Ceradini,^{133a,133b} A. S. Cerqueira,^{23b} A. Cerri,²⁹ L. Cerrito,⁷⁴ F. Cerutti,⁴⁶ S. A. Cetin,^{18b} F. Cevenini,^{101a,101b} A. Chafaq,^{134a} D. Chakraborty,¹⁰⁵ K. Chan,² B. Chapleau,⁸⁴ J. D. Chapman,²⁷ J. W. Chapman,⁸⁶ E. Chareyre,⁷⁷ D. G. Charlton,¹⁷ V. Chavda,⁸¹ C. A. Chavez Barajas,²⁹ S. Cheatham,⁸⁴ S. Chekanov,⁵ S. V. Chekulaev,^{158a} G. A. Chelkov,⁶⁴ M. A. Chelstowska,¹⁰³ C. Chen,⁶³ H. Chen,²⁴ S. Chen,^{32c} T. Chen,^{32c} X. Chen,¹⁷¹ S. Cheng,^{32a} A. Cheplakov,⁶⁴ V. F. Chepurinov,⁶⁴ R. Cherkaoui El Moursli,^{134e} V. Chernyatin,²⁴ E. Cheu,⁶ S. L. Cheung,¹⁵⁷ L. Chevalier,¹³⁵ G. Chiefari,^{101a,101b} L. Chikovani,^{50a} J. T. Childers,²⁹ A. Chilingarov,⁷⁰ G. Chiodini,^{71a} A. S. Chisholm,¹⁷ M. V. Chizhov,⁶⁴ G. Choudalakis,³⁰ S. Chouridou,¹³⁶ I. A. Christidi,⁷⁶ A. Christov,⁴⁷ D. Chromek-Burckhart,²⁹ M. L. Chu,¹⁵⁰ J. Chudoba,¹²⁴ G. Ciapetti,^{131a,131b} K. Ciba,³⁷ A. K. Ciftci,^{3a} R. Ciftci,^{3a} D. Cinca,³³ V. Cindro,⁷³ M. D. Ciobotaru,¹⁶² C. Ciocca,^{19a} A. Ciocio,¹⁴ M. Cirilli,⁸⁶ M. Citterio,^{88a} M. Ciubancan,^{25a} A. Clark,⁴⁸ P. J. Clark,⁴⁵ W. Cleland,¹²² J. C. Clemens,⁸² B. Clement,⁵⁴ C. Clement,^{145a,145b} R. W. Clift,¹²⁸ Y. Coadou,⁸² M. Cobal,^{163a,163c} A. Coccaro,¹⁷¹ J. Cochran,⁶³ P. Coe,¹¹⁷ J. G. Cogan,¹⁴² J. Coggeshall,¹⁶⁴ E. Cogneras,¹⁷⁶ J. Colas,⁴ A. P. Colijn,¹⁰⁴ N. J. Collins,¹⁷ C. Collins-Tooth,⁵² J. Collot,⁵⁴ G. Colon,⁸³ P. Conde Muiño,^{123a} E. Coniavitis,¹¹⁷ M. C. Conidi,¹¹ M. Consonni,¹⁰³ V. Consorti,⁴⁷ S. Constantinescu,^{25a} C. Conta,^{118a,118b} F. Conventi,^{101a,j} J. Cook,²⁹ M. Cooke,¹⁴ B. D. Cooper,⁷⁶ A. M. Cooper-Sarkar,¹¹⁷ K. Copic,¹⁴ T. Cornelissen,¹⁷³ M. Corradi,^{19a} F. Corriveau,^{84,k} A. Cortes-Gonzalez,¹⁶⁴ G. Cortiana,⁹⁸ G. Costa,^{88a} M. J. Costa,¹⁶⁶ D. Costanzo,¹³⁸ T. Costin,³⁰ D. Côté,²⁹ R. Coura Torres,^{23a} L. Courneyea,¹⁶⁸ G. Cowan,⁷⁵ C. Cowden,²⁷ B. E. Cox,⁸¹ K. Cranmer,¹⁰⁷ F. Crescioli,^{121a,121b} M. Cristinziani,²⁰ G. Crosetti,^{36a,36b} R. Crupi,^{71a,71b} S. Crépe-Renaudin,⁵⁴ C.-M. Cuciuc,^{25a} C. Cuenca Almenar,¹⁷⁴ T. Cuhadar Donszelmann,¹³⁸ M. Curatolo,⁴⁶ C. J. Curtis,¹⁷ C. Cuthbert,¹⁴⁹ P. Cwetanski,⁶⁰ H. Czirr,¹⁴⁰

- P. Czodrowski,⁴³ Z. Czyczula,¹⁷⁴ S. D'Auria,⁵² M. D'Onofrio,⁷² A. D'Orazio,^{131a,131b} P. V. M. Da Silva,^{23a}
 C. Da Via,⁸¹ W. Dabrowski,³⁷ T. Dai,⁸⁶ C. Dallapiccola,⁸³ M. Dam,³⁵ M. Dameri,^{49a,49b} D. S. Damiani,¹³⁶
 H. O. Danielsson,²⁹ D. Dannheim,⁹⁸ V. Dao,⁴⁸ G. Darbo,^{49a} G. L. Darlea,^{25b} W. Davey,²⁰ T. Davidek,¹²⁵
 N. Davidson,⁸⁵ R. Davidson,⁷⁰ E. Davies,^{117,d} M. Davies,⁹² A. R. Davison,⁷⁶ Y. Davygora,^{57a} E. Dawe,¹⁴¹
 I. Dawson,¹³⁸ J. W. Dawson,^{5,a} R. K. Daya-Ishmukhametova,²² K. De,⁷ R. de Asmundis,^{101a} S. De Castro,^{19a,19b}
 P. E. De Castro Faria Salgado,²⁴ S. De Cecco,⁷⁷ J. de Graat,⁹⁷ N. De Groot,¹⁰³ P. de Jong,¹⁰⁴ C. De La Taille,¹¹⁴
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 D. della Volpe,^{101a,101b} M. Delmastro,⁴ N. Delruelle,²⁹ P. A. Delsart,⁵⁴ C. Deluca,¹⁴⁷ S. Demers,¹⁷⁴ M. Demichev,⁶⁴
 B. Demirkoz,^{11,l} J. Deng,¹⁶² S. P. Denisov,¹²⁷ D. Derendarz,³⁸ J. E. Derkaoui,^{134d} F. Derue,⁷⁷ P. Dervan,⁷² K. Desch,²⁰
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 A. Di Ciaccio,^{132a,132b} L. Di Ciaccio,⁴ A. Di Girolamo,²⁹ B. Di Girolamo,²⁹ S. Di Luise,^{133a,133b} A. Di Mattia,¹⁷¹
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 J. Dietrich,⁴¹ T. A. Dietzsch,^{57a} S. Diglio,⁸⁵ K. Dindar Yagci,³⁹ J. Dingfelder,²⁰ C. Dionisi,^{131a,131b} P. Dita,^{25a}
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 M. Dobbs,⁸⁴ R. Dobinson,^{29,a} D. Dobos,²⁹ E. Dobson,^{29,n} J. Dodd,³⁴ C. Doglioni,⁴⁸ T. Doherty,⁵² Y. Doi,^{65,a}
 J. Dolejsi,¹²⁵ I. Dolenc,⁷³ Z. Dolezal,¹²⁵ B. A. Dolgoshein,^{95,a} T. Dohmae,¹⁵⁴ M. Donadelli,^{23d} M. Donega,¹¹⁹
 J. Donini,³³ J. Dopke,²⁹ A. Doria,^{101a} A. Dos Anjos,¹⁷¹ M. Dosil,¹¹ A. Dotti,^{121a,121b} M. T. Dova,⁶⁹ J. D. Dowell,¹⁷
 A. D. Doxiadis,¹⁰⁴ A. T. Doyle,⁵² Z. Drasal,¹²⁵ J. Drees,¹⁷³ N. Dressnandt,¹¹⁹ H. Drevermann,²⁹ C. Driouichi,³⁵
 M. Dris,⁹ J. Dubbert,⁹⁸ S. Dube,¹⁴ E. Duchovni,¹⁷⁰ G. Duckeck,⁹⁷ A. Dudarev,²⁹ F. Dudziak,⁶³ M. Dührssen,²⁹
 I. P. Duerdoth,⁸¹ L. Dufлот,¹¹⁴ M-A. Dufour,⁸⁴ M. Dunford,²⁹ H. Duran Yildiz,^{3a} R. Duxfield,¹³⁸ M. Dwuznik,³⁷
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 N. C. Edwards,⁵² W. Ehrenfeld,⁴¹ T. Ehrich,⁹⁸ T. Eifert,¹⁴² G. Eigen,¹³ K. Einsweiler,¹⁴ E. Eisenhandler,⁷⁴
 T. Ekelof,¹⁶⁵ M. El Kacimi,^{134c} M. Ellert,¹⁶⁵ S. Elles,⁴ F. Ellinghaus,⁸⁰ K. Ellis,⁷⁴ N. Ellis,²⁹ J. Elmsheuser,⁹⁷
 M. Elsing,²⁹ D. Emelianov,¹²⁸ R. Engelmann,¹⁴⁷ A. Engl,⁹⁷ B. Epp,⁶¹ A. Eppig,⁸⁶ J. Erdmann,⁵³ A. Ereditato,¹⁶
 D. Eriksson,^{145a} J. Ernst,¹ M. Ernst,²⁴ J. Ernwein,¹³⁵ D. Errede,¹⁶⁴ S. Errede,¹⁶⁴ E. Ertel,⁸⁰ M. Escalier,¹¹⁴
 C. Escobar,¹²² X. Espinal Curull,¹¹ B. Esposito,⁴⁶ F. Etienne,⁸² A. I. Etienvre,¹³⁵ E. Etzion,¹⁵² D. Evangelakou,⁵³
 H. Evans,⁶⁰ L. Fabbri,^{19a,19b} C. Fabre,²⁹ R. M. Fakhruddinov,¹²⁷ S. Falciano,^{131a} Y. Fang,¹⁷¹ M. Fanti,^{88a,88b}
 A. Farbin,⁷ A. Farilla,^{133a} J. Farley,¹⁴⁷ T. Farooque,¹⁵⁷ S. M. Farrington,¹¹⁷ P. Farthouat,²⁹ P. Fassnacht,²⁹
 D. Fassouliotis,⁸ B. Fatholahzadeh,¹⁵⁷ A. Favareto,^{88a,88b} L. Fayard,¹¹⁴ S. Fazio,^{36a,36b} R. Febbraro,³³ P. Federic,^{143a}
 O. L. Fedin,¹²⁰ W. Fedorko,⁸⁷ M. Fehling-Kaschek,⁴⁷ L. Feligioni,⁸² D. Fellmann,⁵ C. Feng,^{32d} E. J. Feng,³⁰
 A. B. Fenyuk,¹²⁷ J. Ferencei,^{143b} J. Ferland,⁹² W. Fernando,¹⁰⁸ S. Ferrag,⁵² J. Ferrando,⁵² V. Ferrara,⁴¹ A. Ferrari,¹⁶⁵
 P. Ferrari,¹⁰⁴ R. Ferrari,^{118a} D. E. Ferreira de Lima,⁵² A. Ferrer,¹⁶⁶ M. L. Ferrer,⁴⁶ D. Ferrere,⁴⁸ C. Ferretti,⁸⁶
 A. Ferretto Parodi,^{49a,49b} M. Fiascaris,³⁰ F. Fiedler,⁸⁰ A. Filipčič,⁷³ A. Filippas,⁹ F. Filthaut,¹⁰³ M. Fincke-Keeler,¹⁶⁸
 M. C. N. Fiolhais,^{123a,i} L. Fiorini,¹⁶⁶ A. Firan,³⁹ G. Fischer,⁴¹ P. Fischer,²⁰ M. J. Fisher,¹⁰⁸ M. Flechl,⁴⁷ I. Fleck,¹⁴⁰
 J. Fleckner,⁸⁰ P. Fleischmann,¹⁷² S. Fleischmann,¹⁷³ T. Flick,¹⁷³ A. Floderus,⁷⁸ L. R. Flores Castillo,¹⁷¹
 M. J. Flowerdew,⁹⁸ M. Fokitis,⁹ T. Fonseca Martin,¹⁶ D. A. Forbush,¹³⁷ A. Formica,¹³⁵ A. Forti,⁸¹ D. Fortin,^{158a}
 J. M. Foster,⁸¹ D. Fournier,¹¹⁴ A. Foussat,²⁹ A. J. Fowler,⁴⁴ K. Fowler,¹³⁶ H. Fox,⁷⁰ P. Francavilla,¹¹
 S. Franchino,^{118a,118b} D. Francis,²⁹ T. Frank,¹⁷⁰ M. Franklin,⁵⁶ S. Franz,²⁹ M. Fraternali,^{118a,118b} S. Fratina,¹¹⁹
 S. T. French,²⁷ F. Friedrich,⁴³ R. Froeschl,²⁹ D. Froidevaux,²⁹ J. A. Frost,²⁷ C. Fukunaga,¹⁵⁵ E. Fullana Torregrosa,²⁹
 J. Fuster,¹⁶⁶ C. Gabaldon,²⁹ O. Gabizon,¹⁷⁰ T. Gadfort,²⁴ S. Gadomski,⁴⁸ G. Gagliardi,^{49a,49b} P. Gagnon,⁶⁰ C. Galea,⁹⁷
 E. J. Gallas,¹¹⁷ V. Gallo,¹⁶ B. J. Gallop,¹²⁸ P. Gallus,¹²⁴ K. K. Gan,¹⁰⁸ Y. S. Gao,^{142,f} V. A. Gapienko,¹²⁷
 A. Gaponenko,¹⁴ F. Garberon,¹⁷⁴ M. Garcia-Sciveres,¹⁴ C. García,¹⁶⁶ J. E. García Navarro,¹⁶⁶ R. W. Gardner,³⁰
 N. Garelli,²⁹ H. Garitaonandia,¹⁰⁴ V. Garonne,²⁹ J. Garvey,¹⁷ C. Gatti,⁴⁶ G. Gaudio,^{118a} B. Gaur,¹⁴⁰ L. Gauthier,¹³⁵
 I. L. Gavrilenko,⁹³ C. Gay,¹⁶⁷ G. Gaycken,²⁰ J-C. Gayde,²⁹ E. N. Gazis,⁹ P. Ge,^{32d} C. N. P. Gee,¹²⁸ D. A. A. Geerts,¹⁰⁴
 Ch. Geich-Gimbel,²⁰ K. Gellerstedt,^{145a,145b} C. Gemme,^{49a} A. Gemmell,⁵² M. H. Genest,⁵⁴ S. Gentile,^{131a,131b}
 M. George,⁵³ S. George,⁷⁵ P. Gerlach,¹⁷³ A. Gershon,¹⁵² C. Geweniger,^{57a} H. Ghazlane,^{134b} N. Ghodbane,³³
 B. Giacobbe,^{19a} S. Giagu,^{131a,131b} V. Giakoumopoulou,⁸ V. Giangiobbe,¹¹ F. Gianotti,²⁹ B. Gibbard,²⁴ A. Gibson,¹⁵⁷
 S. M. Gibson,²⁹ L. M. Gilbert,¹¹⁷ V. Gilevsky,⁹⁰ D. Gillberg,²⁸ A. R. Gillman,¹²⁸ D. M. Gingrich,^{2,e} J. Ginzburg,¹⁵²

- N. Giokaris,⁸ M. P. Giordani,^{163c} R. Giordano,^{101a,101b} F. M. Giorgi,¹⁵ P. Giovannini,⁹⁸ P. F. Giraud,¹³⁵ D. Giugni,^{88a} M. Giunta,⁹² P. Giusti,^{19a} B. K. Gjelsten,¹¹⁶ L. K. Gladilin,⁹⁶ C. Glasman,⁷⁹ J. Glatzer,⁴⁷ A. Glazov,⁴¹ K. W. Glitza,¹⁷³ G. L. Glonti,⁶⁴ J. R. Goddard,⁷⁴ J. Godfrey,¹⁴¹ J. Godlewski,²⁹ M. Goebel,⁴¹ T. Göpfert,⁴³ C. Goeringer,⁸⁰ C. Gössling,⁴² T. Göttfert,⁹⁸ S. Goldfarb,⁸⁶ T. Golling,¹⁷⁴ A. Gomes,^{123a,c} L. S. Gomez Fajardo,⁴¹ R. Gonçalves,⁷⁵ J. Goncalves Pinto Firmino Da Costa,⁴¹ L. Gonella,²⁰ A. Gonidec,²⁹ S. Gonzalez,¹⁷¹ S. González de la Hoz,¹⁶⁶ G. Gonzalez Parra,¹¹ M. L. Gonzalez Silva,²⁶ S. Gonzalez-Sevilla,⁴⁸ J. J. Goodson,¹⁴⁷ L. Goossens,²⁹ P. A. Gorbounov,⁹⁴ H. A. Gordon,²⁴ I. Gorelov,¹⁰² G. Gorfine,¹⁷³ B. Gorini,²⁹ E. Gorini,^{71a,71b} A. Gorišek,⁷³ E. Gornicki,³⁸ S. A. Gorokhov,¹²⁷ V. N. Goryachev,¹²⁷ B. Gosdzik,⁴¹ M. Gosselink,¹⁰⁴ M. I. Gostkin,⁶⁴ I. Gough Eschrich,¹⁶² M. Gouighri,^{134a} D. Goujdami,^{134c} M. P. Goulette,⁴⁸ A. G. Goussiou,¹³⁷ C. Goy,⁴ S. Gozpinar,²² I. Grabowska-Bold,³⁷ P. Grafström,²⁹ K.-J. Grahm,⁴¹ F. Grancagnolo,^{71a} S. Grancagnolo,¹⁵ V. Grassi,¹⁴⁷ V. Gratchev,¹²⁰ N. Grau,³⁴ H. M. Gray,²⁹ J. A. Gray,¹⁴⁷ E. Graziani,^{133a} O. G. Grebenyuk,¹²⁰ T. Greenshaw,⁷² Z. D. Greenwood,^{24,m} K. Gregersen,³⁵ I. M. Gregor,⁴¹ P. Grenier,¹⁴² J. Griffiths,¹³⁷ N. Grigalashvili,⁶⁴ A. A. Grillo,¹³⁶ S. Grinstein,¹¹ Y. V. Grishkevich,⁹⁶ J.-F. Grivaz,¹¹⁴ M. Groh,⁹⁸ E. Gross,¹⁷⁰ J. Grosse-Knetter,⁵³ J. Groth-Jensen,¹⁷⁰ K. Grybel,¹⁴⁰ V. J. Guarino,⁵ D. Guest,¹⁷⁴ C. Guicheney,³³ A. Guida,^{71a,71b} S. Guindon,⁵³ H. Guler,^{84,o} J. Gunther,¹²⁴ B. Guo,¹⁵⁷ J. Guo,³⁴ A. Gupta,³⁰ Y. Gusakov,⁶⁴ V. N. Gushchin,¹²⁷ P. Gutierrez,¹¹⁰ N. Guttman,¹⁵² O. Gutzwiller,¹⁷¹ C. Guyot,¹³⁵ C. Gwenlan,¹¹⁷ C. B. Gwilliam,⁷² A. Haas,¹⁴² S. Haas,²⁹ C. Haber,¹⁴ H. K. Hadavand,³⁹ D. R. Hadley,¹⁷ P. Haefner,⁹⁸ F. Hahn,²⁹ S. Haider,²⁹ Z. Hajduk,³⁸ H. Hakobyan,¹⁷⁵ D. Hall,¹¹⁷ J. Haller,⁵³ K. Hamacher,¹⁷³ P. Hamal,¹¹² M. Hamer,⁵³ A. Hamilton,^{144b,p} S. Hamilton,¹⁶⁰ H. Han,^{32a} L. Han,^{32b} K. Hanagaki,¹¹⁵ K. Hanawa,¹⁵⁹ M. Hance,¹⁴ C. Handel,⁸⁰ P. Hanke,^{57a} J. R. Hansen,³⁵ J. B. Hansen,³⁵ J. D. Hansen,³⁵ P. H. Hansen,³⁵ P. Hansson,¹⁴² K. Hara,¹⁵⁹ G. A. Hare,¹³⁶ T. Harenberg,¹⁷³ S. Harkusha,⁸⁹ D. Harper,⁸⁶ R. D. Harrington,⁴⁵ O. M. Harris,¹³⁷ K. Harrison,¹⁷ J. Hartert,⁴⁷ F. Hartjes,¹⁰⁴ T. Haruyama,⁶⁵ A. Harvey,⁵⁵ S. Hasegawa,¹⁰⁰ Y. Hasegawa,¹³⁹ S. Hassani,¹³⁵ M. Hatch,²⁹ D. Hauff,⁹⁸ S. Haug,¹⁶ M. Hauschild,²⁹ R. Hauser,⁸⁷ M. Havranek,²⁰ B. M. Hawes,¹¹⁷ C. M. Hawkes,¹⁷ R. J. Hawkings,²⁹ A. D. Hawkins,⁷⁸ D. Hawkins,¹⁶² T. Hayakawa,⁶⁶ T. Hayashi,¹⁵⁹ D. Hayden,⁷⁵ H. S. Hayward,⁷² S. J. Haywood,¹²⁸ E. Hazen,²¹ M. He,^{32d} S. J. Head,¹⁷ V. Hedberg,⁷⁸ L. Heelan,⁷ S. Heim,⁸⁷ B. Heinemann,¹⁴ S. Heisterkamp,³⁵ L. Helary,⁴ C. Heller,⁹⁷ M. Heller,²⁹ S. Hellman,^{145a,145b} D. Hellmich,²⁰ C. Helsen,¹¹ R. C. W. Henderson,⁷⁰ M. Henke,^{57a} A. Henrichs,⁵³ A. M. Henriques Correia,²⁹ S. Henrot-Versille,¹¹⁴ F. Henry-Couannier,⁸² C. Hensel,¹⁵³ T. Henß,¹⁷³ C. M. Hernandez,⁷ Y. Hernández Jiménez,¹⁶⁶ R. Herrberg,¹⁵ A. D. Hershendorff,¹⁵¹ G. Herten,⁴⁷ R. Hertenberger,⁹⁷ L. Hervas,²⁹ G. G. Hesketh,⁷⁶ N. P. Hessey,¹⁰⁴ E. Higón-Rodríguez,¹⁶⁶ D. Hill,^{5,a} J. C. Hill,²⁷ N. Hill,⁵ K. H. Hiller,⁴¹ S. Hillert,²⁰ S. J. Hillier,¹⁷ I. Hinchliffe,¹⁴ E. Hines,¹¹⁹ M. Hirose,¹¹⁵ F. Hirsch,⁴² D. Hirschbuehl,¹⁷³ J. Hobbs,¹⁴⁷ N. Hod,¹⁵² M. C. Hodgkinson,¹³⁸ P. Hodgson,¹³⁸ A. Hoecker,²⁹ M. R. Hoferkamp,¹⁰² J. Hoffman,³⁹ D. Hoffmann,⁸² M. Hohlfeld,⁸⁰ M. Holder,¹⁴⁰ S. O. Holmgren,^{145a} T. Holy,¹²⁶ J. L. Holzbauer,⁸⁷ Y. Homma,⁶⁶ T. M. Hong,¹¹⁹ L. Hooft van Huysduynen,¹⁰⁷ T. Horazdovsky,¹²⁶ C. Horn,¹⁴² S. Horner,⁴⁷ J.-Y. Hostachy,⁵⁴ S. Hou,¹⁵⁰ M. A. Houlden,⁷² A. Hoummada,^{134a} J. Howarth,⁸¹ D. F. Howell,¹¹⁷ I. Hristova,¹⁵ J. Hrivnac,¹¹⁴ I. Hruska,¹²⁴ T. Hryn'ova,⁴ P. J. Hsu,⁸⁰ S.-C. Hsu,¹⁴ G. S. Huang,¹¹⁰ Z. Hubacek,¹²⁶ F. Hubaut,⁸² F. Huegging,²⁰ A. Huettmann,⁴¹ T. B. Huffman,¹¹⁷ E. W. Hughes,³⁴ G. Hughes,⁷⁰ R. E. Hughes-Jones,⁸¹ M. Huhtinen,²⁹ P. Hurst,⁵⁶ M. Hurwitz,¹⁴ U. Husemann,⁴¹ N. Huseynov,^{64,q} J. Huston,⁸⁷ J. Huth,⁵⁶ G. Iacobucci,⁴⁸ G. Iakovidis,⁹ M. Ibbotson,⁸¹ I. Ibragimov,¹⁴⁰ R. Ichimiya,⁶⁶ L. Iconomidou-Fayard,¹¹⁴ J. Idarraga,¹¹⁴ P. Iengo,^{101a} O. Igonkina,¹⁰⁴ Y. Ikegami,⁶⁵ M. Ikeno,⁶⁵ Y. Ilchenko,³⁹ D. Iliadis,¹⁵³ N. Ilic,¹⁵⁷ M. Imori,¹⁵⁴ T. Ince,²⁰ J. Inigo-Golfín,²⁹ P. Ioannou,⁸ M. Iodice,^{133a} V. Ippolito,^{131a,131b} A. Irlés Quiles,¹⁶⁶ C. Isaksson,¹⁶⁵ A. Ishikawa,⁶⁶ M. Ishino,⁶⁷ R. Ishmukhametov,³⁹ C. Issever,¹¹⁷ S. Istin,^{18a} A. V. Ivashin,¹²⁷ W. Iwanski,³⁸ H. Iwasaki,⁶⁵ J. M. Izen,⁴⁰ V. Izzo,^{101a} B. Jackson,¹¹⁹ J. N. Jackson,⁷² P. Jackson,¹⁴² M. R. Jaekel,²⁹ V. Jain,⁶⁰ K. Jakobs,⁴⁷ S. Jakobsen,³⁵ J. Jakubek,¹²⁶ D. K. Jana,¹¹⁰ E. Jankowski,¹⁵⁷ E. Jansen,⁷⁶ H. Jansen,²⁹ A. Jantsch,⁹⁸ M. Janus,²⁰ G. Jarlskog,⁷⁸ L. Jeanty,⁵⁶ K. Jelen,³⁷ I. Jen-La Plante,³⁰ P. Jenni,²⁹ A. Jeremie,⁴ P. Jež,³⁵ S. Jézéquel,⁴ M. K. Jha,^{19a} H. Ji,¹⁷¹ W. Ji,⁸⁰ J. Jia,¹⁴⁷ Y. Jiang,^{32b} M. Jimenez Belenguer,⁴¹ G. Jin,^{32b} S. Jin,^{32a} O. Jinnouchi,¹⁵⁶ M. D. Joergensen,³⁵ D. Joffe,³⁹ L. G. Johansen,¹³ M. Johansen,^{145a,145b} K. E. Johansson,^{145a} P. Johansson,¹³⁸ S. Johnert,⁴¹ K. A. Johns,⁶ K. Jon-And,^{145a,145b} G. Jones,¹¹⁷ R. W. L. Jones,⁷⁰ T. W. Jones,⁷⁶ T. J. Jones,⁷² O. Jonsson,²⁹ C. Joram,²⁹ P. M. Jorge,^{123a} J. Joseph,¹⁴ J. Jovicevic,¹⁴⁶ T. Jovin,^{12b} X. Ju,¹⁷¹ C. A. Jung,⁴² R. M. Jungst,²⁹ V. Juranek,¹²⁴ P. Jussel,⁶¹ A. Juste Rozas,¹¹ V. V. Kabachenko,¹²⁷ S. Kabana,¹⁶ M. Kaci,¹⁶⁶ A. Kaczmarska,³⁸ P. Kadlecik,³⁵ M. Kado,¹¹⁴ H. Kagan,¹⁰⁸ M. Kagan,⁵⁶ S. Kaiser,⁹⁸ E. Kajomovitz,¹⁵¹ S. Kalinin,¹⁷³ L. V. Kalinovskaya,⁶⁴ S. Kama,³⁹ N. Kanaya,¹⁵⁴ M. Kaneda,²⁹ S. Kaneti,²⁷ T. Kanno,¹⁵⁶ V. A. Kantserov,⁹⁵ J. Kanzaki,⁶⁵ B. Kaplan,¹⁷⁴

A. Kapliy,³⁰ J. Kaplon,²⁹ D. Kar,⁴³ M. Karagounis,²⁰ M. Karagoz,¹¹⁷ M. Karnevskiy,⁴¹ K. Karr,⁵ V. Kartvelishvili,⁷⁰
A. N. Karyukhin,¹²⁷ L. Kashif,¹⁷¹ G. Kasieczka,^{57b} R. D. Kass,¹⁰⁸ A. Kastanas,¹³ M. Kataoka,⁴ Y. Kataoka,¹⁵⁴
E. Katsoufis,⁹ J. Katzy,⁴¹ V. Kaushik,⁶ K. Kawagoe,⁶⁶ T. Kawamoto,¹⁵⁴ G. Kawamura,⁸⁰ M. S. Kayl,¹⁰⁴
V. A. Kazanin,¹⁰⁶ M. Y. Kazarinov,⁶⁴ R. Keeler,¹⁶⁸ R. Kehoe,³⁹ M. Keil,⁵³ G. D. Kekelidze,⁶⁴ J. Kennedy,⁹⁷
C. J. Kenney,¹⁴² M. Kenyon,⁵² O. Kepka,¹²⁴ N. Kerschen,²⁹ B. P. Kerševan,⁷³ S. Kersten,¹⁷³ K. Kessoku,¹⁵⁴
J. Keung,¹⁵⁷ F. Khalil-zada,¹⁰ H. Khandanyan,¹⁶⁴ A. Khanov,¹¹¹ D. Kharchenko,⁶⁴ A. Khodinov,⁹⁵
A. G. Kholodenko,¹²⁷ A. Khomich,^{57a} T. J. Khoo,²⁷ G. Khoriali,²⁰ A. Khoroshilov,¹⁷³ N. Khovanskiy,⁶⁴
V. Khovanskiy,⁹⁴ E. Khrarov,⁶⁴ J. Khubua,^{50b} H. Kim,^{145a,145b} M. S. Kim,² S. H. Kim,¹⁵⁹ N. Kimura,¹⁶⁹ O. Kind,¹⁵
B. T. King,⁷² M. King,⁶⁶ R. S. B. King,¹¹⁷ J. Kirk,¹²⁸ L. E. Kirsch,²² A. E. Kiryunin,⁹⁸ T. Kishimoto,⁶⁶
D. Kisielewska,³⁷ T. Kittelmann,¹²² A. M. Kiver,¹²⁷ E. Kladiva,^{143b} J. Klaiber-Lodewigs,⁴² M. Klein,⁷² U. Klein,⁷²
K. Kleinknecht,⁸⁰ M. Klemetti,⁸⁴ A. Klier,¹⁷⁰ P. Klimek,^{145a,145b} A. Klimentov,²⁴ R. Klingenberg,⁴² J. A. Klinger,⁸¹
E. B. Klinkby,³⁵ T. Klioutchnikova,²⁹ P. F. Klok,¹⁰³ S. Klous,¹⁰⁴ E.-E. Kluge,^{57a} T. Kluge,⁷² P. Kluit,¹⁰⁴ S. Kluth,⁹⁸
N. S. Knecht,¹⁵⁷ E. Kneringer,⁶¹ J. Knobloch,²⁹ E.B.F.G. Knoops,⁸² A. Knue,⁵³ B. R. Ko,⁴⁴ T. Kobayashi,¹⁵⁴
M. Kobel,⁴³ M. Kocian,¹⁴² P. Kodys,¹²⁵ K. Köneke,²⁹ A. C. König,¹⁰³ S. Koenig,⁸⁰ L. Köpke,⁸⁰ F. Koetsveld,¹⁰³
P. Koevesarki,²⁰ T. Koffas,²⁸ E. Koffeman,¹⁰⁴ L. A. Kogan,¹¹⁷ F. Kohn,⁵³ Z. Kohout,¹²⁶ T. Kohriki,⁶⁵ T. Koi,¹⁴²
T. Kokott,²⁰ G. M. Kolachev,¹⁰⁶ H. Kolanoski,¹⁵ V. Kolesnikov,⁶⁴ I. Koletsou,^{88a} J. Koll,⁸⁷ M. Kollfrath,⁴⁷
S. D. Kolya,⁸¹ A. A. Komar,⁹³ Y. Komori,¹⁵⁴ T. Kondo,⁶⁵ T. Kono,^{41,r} A. I. Kononov,⁴⁷ R. Konoplich,^{107,s}
N. Konstantinidis,⁷⁶ A. Kootz,¹⁷³ S. Koperny,³⁷ K. Korcyl,³⁸ K. Kordas,¹⁵³ V. Koreshev,¹²⁷ A. Korn,¹¹⁷ A. Korol,¹⁰⁶
I. Korolkov,¹¹ E. V. Korolkova,¹³⁸ V. A. Korotkov,¹²⁷ O. Kortner,⁹⁸ S. Kortner,⁹⁸ V. V. Kostyukhin,²⁰
M. J. Kotamäki,²⁹ S. Kotov,⁹⁸ V. M. Kotov,⁶⁴ A. Kotwal,⁴⁴ C. Kourkoumelis,⁸ V. Kouskoura,¹⁵³ A. Koutsman,^{158a}
R. Kowalewski,¹⁶⁸ T. Z. Kowalski,³⁷ W. Kozanecki,¹³⁵ A. S. Kozhin,¹²⁷ V. Kral,¹²⁶ V. A. Kramarenko,⁹⁶
G. Kramberger,⁷³ M. W. Krasny,⁷⁷ A. Krasznahorkay,¹⁰⁷ J. Kraus,⁸⁷ J. K. Kraus,²⁰ A. Kreisel,¹⁵² F. Krejci,¹²⁶
J. Kretzschmar,⁷² N. Krieger,⁵³ P. Krieger,¹⁵⁷ K. Kroeninger,⁵³ H. Kroha,⁹⁸ J. Kroll,¹¹⁹ J. Kroseberg,²⁰ J. Krstic,^{12a}
U. Kruchonak,⁶⁴ H. Krüger,²⁰ T. Kruker,¹⁶ N. Krumnack,⁶³ Z. V. Krumshyteyn,⁶⁴ A. Kruth,²⁰ T. Kubota,⁸⁵ S. Kuday,^{3a}
S. Kuehn,⁴⁷ A. Kugel,^{57c} T. Kuhl,⁴¹ D. Kuhn,⁶¹ V. Kukhtin,⁶⁴ Y. Kulchitsky,⁸⁹ S. Kuleshov,^{31b} C. Kummer,⁹⁷
M. Kuna,⁷⁷ N. Kundu,¹¹⁷ J. Kunkle,¹¹⁹ A. Kupco,¹²⁴ H. Kurashige,⁶⁶ M. Kurata,¹⁵⁹ Y. A. Kurochkin,⁸⁹ V. Kus,¹²⁴
E. S. Kuwertz,¹⁴⁶ M. Kuze,¹⁵⁶ J. Kvita,¹⁴¹ R. Kwee,¹⁵ A. La Rosa,⁴⁸ L. La Rotonda,^{36a,36b} L. Labarga,⁷⁹ J. Labbe,⁴
S. Lablak,^{134a} C. Lacasta,¹⁶⁶ F. Lacava,^{131a,131b} H. Lacker,¹⁵ D. Lacour,⁷⁷ V. R. Lacuesta,¹⁶⁶ E. Ladygin,⁶⁴
R. Lafaye,⁴ B. Laforge,⁷⁷ T. Lagouri,⁷⁹ S. Lai,⁴⁷ E. Laisne,⁵⁴ M. Lamanna,²⁹ C. L. Lampen,⁶ W. Lampl,⁶
E. Lancon,¹³⁵ U. Landgraf,⁴⁷ M. P. J. Landon,⁷⁴ J. L. Lane,⁸¹ C. Lange,⁴¹ A. J. Lankford,¹⁶² F. Lanni,²⁴
K. Lantzsch,¹⁷³ S. Laplace,⁷⁷ C. Lapoire,²⁰ J. F. Laporte,¹³⁵ T. Lari,^{88a} A. V. Larionov,¹²⁷ A. Larner,¹¹⁷ C. Lasseur,²⁹
M. Lassnig,²⁹ P. Laurelli,⁴⁶ V. Lavorini,^{36a,36b} W. Lavrijsen,¹⁴ P. Laycock,⁷² A. B. Lazarev,⁶⁴ O. Le Dortz,⁷⁷
E. Le Guirrec,⁸² C. Le Maner,¹⁵⁷ E. Le Menedeu,⁹ C. Lebel,⁹² T. LeCompte,⁵ F. Ledroit-Guillon,⁵⁴ H. Lee,¹⁰⁴
J. S. H. Lee,¹¹⁵ S. C. Lee,¹⁵⁰ L. Lee,¹⁷⁴ M. Lefebvre,¹⁶⁸ M. Legendre,¹³⁵ A. Leger,⁴⁸ B. C. LeGeyt,¹¹⁹ F. Legger,⁹⁷
C. Leggett,¹⁴ M. Lehmacher,²⁰ G. Lehmann Miotto,²⁹ X. Lei,⁶ M. A. L. Leite,^{23d} R. Leitner,¹²⁵ D. Lellouch,¹⁷⁰
M. Leltchouk,³⁴ B. Lemmer,⁵³ V. Lendermann,^{57a} K. J. C. Leney,^{144b} T. Lenz,¹⁰⁴ G. Lenzen,¹⁷³ B. Lenzi,²⁹
K. Leonhardt,⁴³ S. Leontsinis,⁹ C. Leroy,⁹² J.-R. Lessard,¹⁶⁸ J. Lesser,^{145a} C. G. Lester,²⁷ A. Leung Fook Cheong,¹⁷¹
J. Levêque,⁴ D. Levin,⁸⁶ L. J. Levinson,¹⁷⁰ M. S. Levitski,¹²⁷ A. Lewis,¹¹⁷ G. H. Lewis,¹⁰⁷ A. M. Leyko,²⁰
M. Leyton,¹⁵ B. Li,⁸² H. Li,^{171,t} S. Li,^{32b,u} X. Li,⁸⁶ Z. Liang,^{117,v} H. Liao,³³ B. Liberti,^{132a} P. Lichard,²⁹
M. Lichtnecker,⁹⁷ K. Lie,¹⁶⁴ W. Liebig,¹³ R. Lifshitz,¹⁵¹ C. Limbach,²⁰ A. Limosani,⁸⁵ M. Limper,⁶² S. C. Lin,^{150,w}
F. Linde,¹⁰⁴ J. T. Linnemann,⁸⁷ E. Lipeles,¹¹⁹ L. Lipinsky,¹²⁴ A. Lipniacka,¹³ T. M. Liss,¹⁶⁴ D. Lissauer,²⁴ A. Lister,⁴⁸
A. M. Litke,¹³⁶ C. Liu,²⁸ D. Liu,¹⁵⁰ H. Liu,⁸⁶ J. B. Liu,⁸⁶ M. Liu,^{32b} Y. Liu,^{32b} M. Livan,^{118a,118b}
S. S. A. Livermore,¹¹⁷ A. Lleres,⁵⁴ J. Llorente Merino,⁷⁹ S. L. Lloyd,⁷⁴ E. Lobodzinska,⁴¹ P. Loch,⁶
W. S. Lockman,¹³⁶ T. Loddenkoetter,²⁰ F. K. Loebinger,⁸¹ A. Loginov,¹⁷⁴ C. W. Loh,¹⁶⁷ T. Lohse,¹⁵ K. Lohwasser,⁴⁷
M. Lokajicek,¹²⁴ J. Loken,¹¹⁷ V. P. Lombardo,⁴ R. E. Long,⁷⁰ L. Lopes,^{123a} D. Lopez Mateos,⁵⁶ J. Lorenz,⁹⁷
N. Lorenzo Martinez,¹¹⁴ M. Losada,¹⁶¹ P. Loscutoff,¹⁴ F. Lo Sterzo,^{131a,131b} M. J. Losty,^{158a} X. Lou,⁴⁰ A. Lounis,¹¹⁴
K. F. Loureiro,¹⁶¹ J. Love,²¹ P. A. Love,⁷⁰ A. J. Lowe,^{142,f} F. Lu,^{32a} H. J. Lubatti,¹³⁷ C. Luci,^{131a,131b} A. Lucotte,⁵⁴
A. Ludwig,⁴³ D. Ludwig,⁴¹ I. Ludwig,⁴⁷ J. Ludwig,⁴⁷ F. Luehring,⁶⁰ G. Luijckx,¹⁰⁴ D. Lumb,⁴⁷ L. Luminari,^{131a}
E. Lund,¹¹⁶ B. Lund-Jensen,¹⁴⁶ B. Lundberg,⁷⁸ J. Lundberg,^{145a,145b} J. Lundquist,³⁵ M. Lungwitz,⁸⁰ G. Lutz,⁹⁸
D. Lynn,²⁴ J. Lys,¹⁴ E. Lytken,⁷⁸ H. Ma,²⁴ L. L. Ma,¹⁷¹ J. A. Macana Goia,⁹² G. Maccarrone,⁴⁶ A. Macchiolo,⁹⁸
B. Maček,⁷³ J. Machado Miguens,^{123a} R. Mackeprang,³⁵ R. J. Madaras,¹⁴ W. F. Mader,⁴³ R. Maenner,^{57c} T. Maeno,²⁴

- P. Mättig,¹⁷³ S. Mättig,⁴¹ L. Magnoni,²⁹ E. Magradze,⁵³ Y. Mahalalel,¹⁵² K. Mahboubi,⁴⁷ G. Mahout,¹⁷
 C. Maiani,^{131a,131b} C. Maidantchik,^{23a} A. Maio,^{123a,c} S. Majewski,²⁴ Y. Makida,⁶⁵ N. Makovec,¹¹⁴ P. Mal,¹³⁵
 B. Malaescu,²⁹ Pa. Malecki,³⁸ P. Malecki,³⁸ V. P. Maleev,¹²⁰ F. Malek,⁵⁴ U. Mallik,⁶² D. Malon,⁵ C. Malone,¹⁴²
 S. Maltezos,⁹ V. Malyshev,¹⁰⁶ S. Malyukov,²⁹ R. Mameghani,⁹⁷ J. Mamuzic,^{12b} A. Manabe,⁶⁵ L. Mandelli,^{88a}
 I. Mandić,⁷³ R. Mandrysch,¹⁵ J. Maneira,^{123a} P. S. Mangeard,⁸⁷ L. Manhaes de Andrade Filho,^{23a} I. D. Manjavidze,⁶⁴
 A. Mann,⁵³ P. M. Manning,¹³⁶ A. Manousakis-Katsikakis,⁸ B. Mansoulie,¹³⁵ A. Manz,⁹⁸ A. Mapelli,²⁹ L. Mapelli,²⁹
 L. March,⁷⁹ J. F. Marchand,²⁸ F. Marchese,^{132a,132b} G. Marchiori,⁷⁷ M. Marcisovsky,¹²⁴ A. Marin,^{21a} C. P. Marino,¹⁶⁸
 F. Marroquim,^{23a} R. Marshall,⁸¹ Z. Marshall,²⁹ F. K. Martens,¹⁵⁷ S. Marti-Garcia,¹⁶⁶ A. J. Martin,¹⁷⁴ B. Martin,²⁹
 B. Martin,⁸⁷ F. F. Martin,¹¹⁹ J. P. Martin,⁹² Ph. Martin,⁵⁴ T. A. Martin,¹⁷ V. J. Martin,⁴⁵ B. Martin dit Latour,⁴⁸
 S. Martin-Haugh,¹⁴⁸ M. Martinez,¹¹ V. Martinez Outschoorn,⁵⁶ A. C. Martyniuk,¹⁶⁸ M. Marx,⁸¹ F. Marzano,^{131a}
 A. Marzin,¹¹⁰ L. Masetti,⁸⁰ T. Mashimo,¹⁵⁴ R. Mashinistov,⁹³ J. Masik,⁸¹ A. L. Maslennikov,¹⁰⁶ I. Massa,^{19a,19b}
 G. Massaro,¹⁰⁴ N. Massol,⁴ P. Mastrandrea,^{131a,131b} A. Mastroberardino,^{36a,36b} T. Masubuchi,¹⁵⁴ M. Mathes,²⁰
 P. Matricon,¹¹⁴ H. Matsumoto,¹⁵⁴ H. Matsunaga,¹⁵⁴ T. Matsushita,⁶⁶ C. Mattraversi,^{117,d} J. M. Maugain,²⁹
 J. Maurer,⁸² S. J. Maxfield,⁷² D. A. Maximov,^{106,g} E. N. May,⁵ A. Mayne,¹³⁸ R. Mazini,¹⁵⁰ M. Mazur,²⁰
 M. Mazzanti,^{88a} E. Mazzeni,^{121a,121b} S. P. Mc Kee,⁸⁶ A. McCarn,¹⁶⁴ R. L. McCarthy,¹⁴⁷ T. G. McCarthy,²⁸
 N. A. McCubbin,¹²⁸ K. W. McFarlane,⁵⁵ J. A. McFayden,¹³⁸ H. McGlone,⁵² G. Mchedlidze,^{50b} R. A. McLaren,²⁹
 T. McLaughlan,¹⁷ S. J. McMahon,¹²⁸ R. A. McPherson,^{168,k} A. Meade,⁸³ J. Mechnich,¹⁰⁴ M. Mechtel,¹⁷³
 M. Medinnis,⁴¹ R. Meera-Lebbai,¹¹⁰ T. Meguro,¹¹⁵ R. Mehdiyev,⁹² S. Mehlhase,³⁵ A. Mehta,⁷² K. Meier,^{57a}
 B. Meirose,⁷⁸ C. Melachrinou,³⁰ B. R. Mellado Garcia,¹⁷¹ L. Mendoza Navas,¹⁶¹ Z. Meng,^{150,t} A. Mengarelli,^{19a,19b}
 S. Menke,⁹⁸ C. Menot,²⁹ E. Meoni,¹¹ K. M. Mercurio,⁵⁶ P. Mermod,⁴⁸ L. Merola,^{101a,101b} C. Meroni,^{88a}
 F. S. Merritt,³⁰ H. Merritt,¹⁰⁸ A. Messina,²⁹ J. Metcalfe,¹⁰² A. S. Mete,⁶³ C. Meyer,⁸⁰ C. Meyer,³⁰ J-P. Meyer,¹³⁵
 J. Meyer,¹⁷² J. Meyer,⁵³ T. C. Meyer,²⁹ W. T. Meyer,⁶³ J. Miao,^{32d} S. Michal,²⁹ L. Micu,^{25a} R. P. Middleton,¹²⁸
 S. Migas,⁷² L. Mijović,⁴¹ G. Mikenberg,¹⁷⁰ M. Mikestikova,¹²⁴ M. Mikuž,⁷³ D. W. Miller,³⁰ R. J. Miller,⁸⁷
 W. J. Mills,¹⁶⁷ C. Mills,⁵⁶ A. Milov,¹⁷⁰ D. A. Milstead,^{145a,145b} D. Milstein,¹⁷⁰ A. A. Minaenko,¹²⁷
 M. Miñano Moya,¹⁶⁶ I. A. Minashvili,⁶⁴ A. I. Mincer,¹⁰⁷ B. Mindur,³⁷ M. Mineev,⁶⁴ Y. Ming,¹⁷¹ L. M. Mir,¹¹
 G. Mirabelli,^{131a} L. Miralles Verge,¹¹ A. Misiejuk,⁷⁵ J. Mitrevski,¹³⁶ G. Y. Mitrofanov,¹²⁷ V. A. Mitsou,¹⁶⁶
 S. Mitsui,⁶⁵ P. S. Miyagawa,¹³⁸ K. Miyazaki,⁶⁶ J. U. Mjörnmark,⁷⁸ T. Moa,^{145a,145b} P. Mockett,¹³⁷ S. Moed,⁵⁶
 V. Moeller,²⁷ K. Mönig,⁴¹ N. Möser,²⁰ S. Mohapatra,¹⁴⁷ W. Mohr,⁴⁷ S. Mohrdieck-Möck,⁹⁸ A. M. Moiseev,^{127,a}
 R. Moles-Valls,¹⁶⁶ J. Molina-Perez,²⁹ J. Monk,⁷⁶ E. Monnier,⁸² S. Montesano,^{88a,88b} F. Monticelli,⁶⁹
 S. Monzani,^{19a,19b} R. W. Moore,² G. F. Moorhead,⁸⁵ C. Mora Herrera,⁴⁸ A. Moraes,⁵² N. Morange,¹³⁵ J. Morel,⁵³
 G. Morello,^{36a,36b} D. Moreno,⁸⁰ M. Moreno Llácer,¹⁶⁶ P. Morettini,^{49a} M. Morgenstern,⁴³ M. Morii,⁵⁶ J. Morin,⁷⁴
 A. K. Morley,²⁹ G. Mornacchi,²⁹ S. V. Morozov,⁹⁵ J. D. Morris,⁷⁴ L. Morvaj,¹⁰⁰ H. G. Moser,⁹⁸ M. Mosidze,^{50b}
 J. Moss,¹⁰⁸ R. Mount,¹⁴² E. Mountricha,^{9,x} S. V. Mouraviev,⁹³ E. J. W. Moyse,⁸³ M. Mudrinic,^{12b} F. Mueller,^{57a}
 J. Mueller,¹²² K. Mueller,²⁰ T. A. Müller,⁹⁷ T. Mueller,⁸⁰ D. Muenstermann,²⁹ A. Muir,¹⁶⁷ Y. Munwes,¹⁵²
 W. J. Murray,¹²⁸ I. Mussche,¹⁰⁴ E. Musto,^{101a,101b} A. G. Myagkov,¹²⁷ M. Myska,¹²⁴ J. Nadal,¹¹ K. Nagai,¹⁵⁹
 K. Nagano,⁶⁵ A. Nagarkar,¹⁰⁸ Y. Nagasaka,⁵⁹ M. Nagel,⁹⁸ A. M. Nairz,²⁹ Y. Nakahama,²⁹ K. Nakamura,¹⁵⁴
 T. Nakamura,¹⁵⁴ I. Nakano,¹⁰⁹ G. Nanava,²⁰ A. Napier,¹⁶⁰ R. Narayan,^{57b} M. Nash,^{76,d} N. R. Nation,²¹
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 A. Negri,^{118a,118b} G. Negri,²⁹ S. Nektarijevic,⁴⁸ A. Nelson,¹⁶² S. Nelson,¹⁴² T. K. Nelson,¹⁴² S. Nemecek,¹²⁴
 P. Nemethy,¹⁰⁷ A. A. Nepomuceno,^{23a} M. Nessi,^{29,y} M. S. Neubauer,¹⁶⁴ A. Neusiedl,⁸⁰ R. M. Neves,¹⁰⁷ P. Nevski,²⁴
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 F. Niedercorn,¹¹⁴ J. Nielsen,¹³⁶ T. Niinikoski,²⁹ N. Nikiforou,³⁴ A. Nikiforov,¹⁵ V. Nikolaenko,¹²⁷ K. Nikolaev,⁶⁴
 I. Nikolic-Audit,⁷⁷ K. Nikolics,⁴⁸ K. Nikolopoulos,²⁴ H. Nilsen,⁴⁷ P. Nilsson,⁷ Y. Ninomiya,¹⁵⁴ A. Nisati,^{131a}
 T. Nishiyama,⁶⁶ R. Nisius,⁹⁸ L. Nodulman,⁵ M. Nomachi,¹¹⁵ I. Nomidis,¹⁵³ M. Nordberg,²⁹ B. Nordkvist,^{145a,145b}
 P. R. Norton,¹²⁸ J. Novakova,¹²⁵ M. Nozaki,⁶⁵ L. Nozka,¹¹² I. M. Nugent,^{158a} A.-E. Nuncio-Quiroz,²⁰
 G. Nunes Hanninger,⁸⁵ T. Nunnemann,⁹⁷ E. Nurse,⁷⁶ B. J. O'Brien,⁴⁵ S. W. O'Neale,^{17,a} D. C. O'Neil,¹⁴¹
 V. O'Shea,⁵² L. B. Oakes,⁹⁷ F. G. Oakham,^{28,e} H. Oberlack,⁹⁸ J. Ocariz,⁷⁷ A. Ochi,⁶⁶ S. Oda,¹⁵⁴ S. Odaka,⁶⁵
 J. Odier,⁸² H. Ogren,⁶⁰ A. Oh,⁸¹ S. H. Oh,⁴⁴ C. C. Ohm,^{145a,145b} T. Ohshima,¹⁰⁰ H. Ohshita,¹³⁹ T. Ohsugi,⁵⁸
 S. Okada,⁶⁶ H. Okawa,¹⁶² Y. Okumura,¹⁰⁰ T. Okuyama,¹⁵⁴ A. Olariu,^{25a} M. Olcese,^{49a} A. G. Olchevski,⁶⁴
 S. A. Olivares Pino,^{31a} M. Oliveira,^{123a,i} D. Oliveira Damazio,²⁴ E. Oliver Garcia,¹⁶⁶ D. Olivito,¹¹⁹ A. Olszewski,³⁸
 J. Olszowska,³⁸ C. Omachi,⁶⁶ A. Onofre,^{123a,z} P. U. E. Onyisi,³⁰ C. J. Oram,^{158a} M. J. Oreglia,³⁰ Y. Oren,¹⁵²

- D. Orestano,^{133a,133b} I. Orlov,¹⁰⁶ C. Oropeza Barrera,⁵² R. S. Orr,¹⁵⁷ B. Osculati,^{49a,49b} R. Ospanov,¹¹⁹ C. Osuna,¹¹ G. Otero y Garzon,²⁶ J. P. Ottersbach,¹⁰⁴ M. Ouchrif,^{134d} E. A. Ouellette,¹⁶⁸ F. Ould-Saada,¹¹⁶ A. Ouraou,¹³⁵ Q. Ouyang,^{32a} A. Ovcharova,¹⁴ M. Owen,⁸¹ S. Owen,¹³⁸ V. E. Ozcan,^{18a} N. Ozturk,⁷ A. Pacheco Pages,¹¹ C. Padilla Aranda,¹¹ S. Pagan Griso,¹⁴ E. Paganis,¹³⁸ F. Paige,²⁴ P. Pais,⁸³ K. Pajchel,¹¹⁶ G. Palacino,^{158b} C. P. Paleari,⁶ S. Palestini,²⁹ D. Pallin,³³ A. Palma,^{123a} J. D. Palmer,¹⁷ Y. B. Pan,¹⁷¹ E. Panagiotopoulou,⁹ B. Panes,^{31a} N. Panikashvili,⁸⁶ S. Panitkin,²⁴ D. Pantea,^{25a} M. Panuskova,¹²⁴ V. Paolone,¹²² A. Papadelis,^{145a} Th.D. Papadopoulou,⁹ A. Paramonov,⁵ D. Paredes Hernandez,³³ W. Park,^{24,aa} M. A. Parker,²⁷ F. Parodi,^{49a,49b} J. A. Parsons,³⁴ U. Parzefall,⁴⁷ E. Pasqualucci,^{131a} S. Passaggio,^{49a} A. Passeri,^{133a} F. Pastore,^{133a,133b} Fr. Pastore,⁷⁵ G. Pásztor,^{48,bb} S. Pataraiia,¹⁷³ N. Patel,¹⁴⁹ J. R. Pater,⁸¹ S. Patricelli,^{101a,101b} T. Pauly,²⁹ M. Pecsý,^{143a} M. I. Pedraza Morales,¹⁷¹ S. V. Peleganchuk,¹⁰⁶ H. Peng,^{32b} R. Pengo,²⁹ A. Penson,³⁴ J. Penwell,⁶⁰ M. Perantoni,^{23a} K. Perez,^{34,cc} T. Perez Cavalcanti,⁴¹ E. Perez Codina,¹¹ M. T. Pérez García-Estañ,¹⁶⁶ V. Perez Reale,³⁴ L. Perini,^{88a,88b} H. Pernegger,²⁹ R. Perrino,^{71a} P. Perrodo,⁴ S. Persema,^{3a} A. Perus,¹¹⁴ V. D. Peshekhonov,⁶⁴ K. Peters,²⁹ B. A. Petersen,²⁹ J. Petersen,²⁹ T. C. Petersen,³⁵ E. Petit,⁴ A. Petridis,¹⁵³ C. Petridou,¹⁵³ E. Petrolu,^{131a} F. Petrucci,^{133a,133b} D. Petschull,⁴¹ M. Petteni,¹⁴¹ R. Pezoa,^{31b} A. Phan,⁸⁵ P. W. Phillips,¹²⁸ G. Piacquadio,²⁹ E. Piccaro,⁷⁴ M. Piccinini,^{19a,19b} S. M. Piec,⁴¹ R. Piegaia,²⁶ D. T. Pignotti,¹⁰⁸ J. E. Pilcher,³⁰ A. D. Pilkington,⁸¹ J. Pina,^{123a,c} M. Pinamonti,^{163a,163c} A. Pinder,¹¹⁷ J. L. Pinfeld,² J. Ping,^{32c} B. Pinto,^{123a} O. Pirote,²⁹ C. Pizio,^{88a,88b} M. Plamondon,¹⁶⁸ M.-A. Pleier,²⁴ A. V. Pleskach,¹²⁷ A. Poblaguev,²⁴ S. Poddar,^{57a} F. Podlyski,³³ L. Poggioli,¹¹⁴ T. Poghosyan,²⁰ M. Pohl,⁴⁸ F. Polci,⁵⁴ G. Polesello,^{118a} A. Policicchio,^{36a,36b} A. Polini,^{19a} J. Poll,⁷⁴ V. Polychronakos,²⁴ D. M. Pomarede,¹³⁵ D. Pomeroy,²² K. Pommès,²⁹ L. Pontecorvo,^{131a} B. G. Pope,⁸⁷ G. A. Popeneciu,^{25a} D. S. Popovic,^{12a} A. Poppleton,²⁹ X. Portell Bueso,²⁹ C. Posch,²¹ G. E. Pospelov,⁹⁸ S. Pospisil,¹²⁶ I. N. Potrap,⁹⁸ C. J. Potter,¹⁴⁸ C. T. Potter,¹¹³ G. Poulard,²⁹ J. Poveda,¹⁷¹ R. Prabhu,⁷⁶ P. Pralavorio,⁸² A. Pranko,¹⁴ S. Prasad,⁵⁶ R. Pravahan,⁷ S. Prell,⁶³ K. Pretzl,¹⁶ L. Pribyl,²⁹ D. Price,⁶⁰ J. Price,⁷² L. E. Price,⁵ M. J. Price,²⁹ D. Prieur,¹²² M. Primavera,^{71a} K. Prokofiev,¹⁰⁷ F. Prokoshin,^{31b} S. Protopopescu,²⁴ J. Proudfoot,⁵ X. Prudent,⁴³ M. Przybycien,³⁷ H. Przysiezniak,⁴ S. Psoroulas,²⁰ E. Ptacek,¹¹³ E. Pueschel,⁸³ J. Purdham,⁸⁶ M. Purohit,^{24,aa} P. Puzo,¹¹⁴ Y. Pylypchenko,⁶² J. Qian,⁸⁶ Z. Qian,⁸² Z. Qin,⁴¹ A. Quadt,⁵³ D. R. Quarrie,¹⁴ W. B. Quayle,¹⁷¹ F. Quinonez,^{31a} M. Raas,¹⁰³ V. Radescu,^{57b} B. Radics,²⁰ P. Radloff,¹¹³ T. Rador,^{18a} F. Ragusa,^{88a,88b} G. Rahal,¹⁷⁶ A. M. Rahimi,¹⁰⁸ D. Rahm,²⁴ S. Rajagopalan,²⁴ M. Rammensee,⁴⁷ M. Rammes,¹⁴⁰ A. S. Randle-Conde,³⁹ K. Randrianarivony,²⁸ P. N. Ratoff,⁷⁰ F. Rauscher,⁹⁷ T. C. Rave,⁴⁷ M. Raymond,²⁹ A. L. Read,¹¹⁶ D. M. Rebuffi,^{118a,118b} A. Redelbach,¹⁷² G. Redlinger,²⁴ R. Reece,¹¹⁹ K. Reeves,⁴⁰ A. Reichold,¹⁰⁴ E. Reinherz-Aronis,¹⁵² A. Reinsch,¹¹³ I. Reisinger,⁴² C. Rembser,²⁹ Z. L. Ren,¹⁵⁰ A. Renaud,¹¹⁴ P. Renkel,³⁹ M. Rescigno,^{131a} S. Resconi,^{88a} B. Resende,¹³⁵ P. Reznicek,⁹⁷ R. Rezvani,¹⁵⁷ A. Richards,⁷⁶ R. Richter,⁹⁸ E. Richter-Was,⁴ M. Ridel,⁷⁷ M. Rijpstra,¹⁰⁴ M. Rijssenbeek,¹⁴⁷ A. Rimoldi,^{118a,118b} L. Rinaldi,^{19a} R. R. Rios,³⁹ I. Riu,¹¹ G. Rivoltella,^{88a,88b} F. Rizatdinova,¹¹¹ E. Rizvi,⁷⁴ S. H. Robertson,^{84,k} A. Robichaud-Veronneau,¹¹⁷ D. Robinson,²⁷ J. E. M. Robinson,⁷⁶ M. Robinson,¹¹³ A. Robson,⁵² J. G. Rocha de Lima,¹⁰⁵ C. Roda,^{121a,121b} D. Roda Dos Santos,²⁹ D. Rodriguez,¹⁶¹ A. Roe,⁵³ S. Roe,²⁹ O. Røhne,¹¹⁶ V. Rojo,¹ S. Rolli,¹⁶⁰ A. Romaniouk,⁹⁵ M. Romano,^{19a,19b} V. M. Romanov,⁶⁴ G. Romeo,²⁶ E. Romero Adam,¹⁶⁶ L. Roos,⁷⁷ E. Ros,¹⁶⁶ S. Rosati,^{131a} K. Rosbach,⁴⁸ A. Rose,¹⁴⁸ M. Rose,⁷⁵ G. A. Rosenbaum,¹⁵⁷ E. I. Rosenberg,⁶³ P. L. Rosendahl,¹³ O. Rosenthal,¹⁴⁰ L. Rosselet,⁴⁸ V. Rossetti,¹¹ E. Rossi,^{131a,131b} L. P. Rossi,^{49a} M. Rotaru,^{25a} I. Roth,¹⁷⁰ J. Rothberg,¹³⁷ D. Rousseau,¹¹⁴ C. R. Royon,¹³⁵ A. Rozanov,⁸² Y. Rozen,¹⁵¹ X. Ruan,^{32a,dd} I. Rubinskiy,⁴¹ B. Ruckert,⁹⁷ N. Ruckstuhl,¹⁰⁴ V. I. Rud,⁹⁶ C. Rudolph,⁴³ G. Rudolph,⁶¹ F. Rühr,⁶ F. Ruggieri,^{133a,133b} A. Ruiz-Martinez,⁶³ V. Rumiantsev,^{90,a} L. Rummyantsev,⁶⁴ K. Runge,⁴⁷ Z. Rurikova,⁴⁷ N. A. Rusakovich,⁶⁴ D. R. Rust,⁶⁰ J. P. Rutherford,⁶ C. Ruwiedel,¹⁴ P. Ruzicka,¹²⁴ Y. F. Ryabov,¹²⁰ V. Ryadovikov,¹²⁷ P. Ryan,⁸⁷ M. Rybar,¹²⁵ G. Rybkin,¹¹⁴ N. C. Ryder,¹¹⁷ S. Rzaeva,¹⁰ A. F. Saavedra,¹⁴⁹ I. Sadeh,¹⁵² H. F. W. Sadrozinski,¹³⁶ R. Sadykov,⁶⁴ F. Safai Tehrani,^{131a} H. Sakamoto,¹⁵⁴ G. Salamanna,⁷⁴ A. Salamon,^{132a} M. Saleem,¹¹⁰ D. Salihagic,⁹⁸ A. Salnikov,¹⁴² J. Salt,¹⁶⁶ B. M. Salvachua Ferrando,⁵ D. Salvatore,^{36a,36b} F. Salvatore,¹⁴⁸ A. Salvucci,¹⁰³ A. Salzburger,²⁹ D. Sampsonidis,¹⁵³ B. H. Samset,¹¹⁶ A. Sanchez,^{101a,101b} V. Sanchez Martinez,¹⁶⁶ H. Sandaker,¹³ H. G. Sander,⁸⁰ M. P. Sanders,⁹⁷ M. Sandhoff,¹⁷³ T. Sandoval,²⁷ C. Sandoval,¹⁶¹ R. Sandstroem,⁹⁸ S. Sandvoss,¹⁷³ D. P. C. Sankey,¹²⁸ A. Sansoni,⁴⁶ C. Santamarina Rios,⁸⁴ C. Santoni,³³ R. Santonic,^{132a,132b} H. Santos,^{123a} J. G. Saraiva,^{123a} T. Sarangi,¹⁷¹ E. Sarkisyan-Grinbaum,⁷ F. Sarri,^{121a,121b} G. Sartisohn,¹⁷³ O. Sasaki,⁶⁵ N. Sasao,⁶⁷ I. Satsounkevitch,⁸⁹ G. Sauvage,⁴ E. Sauvan,⁴ J. B. Sauvan,¹¹⁴ P. Savard,^{157,e} V. Savinov,¹²² D. O. Savu,²⁹ L. Sawyer,^{24,m} D. H. Saxon,⁵² L. P. SAYS,³³ C. Sbarra,^{19a} A. Sbrizzi,^{19a,19b} O. Scallan,⁹² D. A. Scannicchio,¹⁶² M. Scarcella,¹⁴⁹ J. Schaarschmidt,¹¹⁴ P. Schacht,⁹⁸

- U. Schäfer,⁸⁰ S. Schaepe,²⁰ S. Schaetzel,^{57b} A. C. Schaffer,¹¹⁴ D. Schaile,⁹⁷ R. D. Schamberger,¹⁴⁷ A. G. Schamov,¹⁰⁶ V. Scharf,^{57a} V. A. Schegelsky,¹²⁰ D. Scheirich,⁸⁶ M. Schernau,¹⁶² M. I. Scherzer,³⁴ C. Schiavi,^{49a,49b} J. Schieck,⁹⁷ M. Schioppa,^{36a,36b} S. Schlenker,²⁹ J. L. Schlereth,⁵ E. Schmidt,⁴⁷ K. Schmieden,²⁰ C. Schmitt,⁸⁰ S. Schmitt,^{57b} M. Schmitz,²⁰ A. Schöning,^{57b} M. Schott,²⁹ D. Schouten,^{158a} J. Schovancova,¹²⁴ M. Schram,⁸⁴ C. Schroeder,⁸⁰ N. Schroer,^{57c} S. Schuh,²⁹ G. Schuler,²⁹ M. J. Schultens,²⁰ J. Schultes,¹⁷³ H.-C. Schultz-Coulon,^{57a} H. Schulz,¹⁵ J. W. Schumacher,²⁰ M. Schumacher,⁴⁷ B. A. Schumm,¹³⁶ Ph. Schune,¹³⁵ C. Schwanenberger,⁸¹ A. Schwartzman,¹⁴² Ph. Schwemling,⁷⁷ R. Schwienhorst,⁸⁷ R. Schwierz,⁴³ J. Schwindling,¹³⁵ T. Schwindt,²⁰ M. Schwoerer,⁴ W. G. Scott,¹²⁸ J. Searcy,¹¹³ G. Sedov,⁴¹ E. Sedykh,¹²⁰ E. Segura,¹¹ S. C. Seidel,¹⁰² A. Seiden,¹³⁶ F. Seifert,⁴³ J. M. Seixas,^{23a} G. Sekhniaidze,^{101a} K. E. Selbach,⁴⁵ D. M. Seliverstov,¹²⁰ B. Sellden,^{145a} G. Sellers,⁷² M. Seman,^{143b} N. Semprini-Cesari,^{19a,19b} C. Serfon,⁹⁷ L. Serin,¹¹⁴ L. Serkin,⁵³ R. Seuster,⁹⁸ H. Severini,¹¹⁰ M. E. Sevier,⁸⁵ A. Sfyrla,²⁹ E. Shabalina,⁵³ M. Shamim,¹¹³ L. Y. Shan,^{32a} J. T. Shank,²¹ Q. T. Shao,⁸⁵ M. Shapiro,¹⁴ P. B. Shatalov,⁹⁴ L. Shaver,⁶ K. Shaw,^{163a,163c} D. Sherman,¹⁷⁴ P. Sherwood,⁷⁶ A. Shibata,¹⁰⁷ H. Shichi,¹⁰⁰ S. Shimizu,²⁹ M. Shimojima,⁹⁹ T. Shin,⁵⁵ M. Shiyakova,⁶⁴ A. Shmeleva,⁹³ M. J. Shochet,³⁰ D. Short,¹¹⁷ S. Shrestha,⁶³ E. Shulga,⁹⁵ M. A. Shupe,⁶ P. Sicho,¹²⁴ A. Sidoti,^{131a} F. Siegert,⁴⁷ Dj. Sijacki,^{12a} O. Silbert,¹⁷⁰ J. Silva,^{123a,c} Y. Silver,¹⁵² D. Silverstein,¹⁴² S. B. Silverstein,^{145a} V. Simak,¹²⁶ O. Simard,¹³⁵ Lj. Simic,^{12a} S. Simion,¹¹⁴ B. Simmons,⁷⁶ M. Simonyan,³⁵ P. Sinervo,¹⁵⁷ N. B. Sinev,¹¹³ V. Sipica,¹⁴⁰ G. Siragusa,¹⁷² A. Sircar,²⁴ A. N. Sisakyan,⁶⁴ S. Yu. Sivoklokov,⁹⁶ J. Sjölin,^{145a,145b} T. B. Sjursen,¹³ L. A. Skinnari,¹⁴ H. P. Skottowe,⁵⁶ K. Skovpen,¹⁰⁶ P. Skubic,¹¹⁰ N. Skvorodnev,²² M. Slater,¹⁷ T. Slavicek,¹²⁶ K. Sliwa,¹⁶⁰ J. Sloper,²⁹ V. Smakhtin,¹⁷⁰ B. H. Smart,⁴⁵ S. Yu. Smirnov,⁹⁵ Y. Smirnov,⁹⁵ L. N. Smirnova,⁹⁶ O. Smirnova,⁷⁸ B. C. Smith,⁵⁶ D. Smith,¹⁴² K. M. Smith,⁵² M. Smizanska,⁷⁰ K. Smolek,¹²⁶ A. A. Snesarev,⁹³ S. W. Snow,⁸¹ J. Snow,¹¹⁰ J. Snuverink,¹⁰⁴ S. Snyder,²⁴ M. Soares,^{123a} R. Sobie,^{168,k} J. Sodomka,¹²⁶ A. Soffer,¹⁵² C. A. Solans,¹⁶⁶ M. Solar,¹²⁶ J. Solc,¹²⁶ E. Soldatov,⁹⁵ U. Soldevila,¹⁶⁶ E. Solfaroli Camillocci,^{131a,131b} A. A. Solodkov,¹²⁷ O. V. Solovyanov,¹²⁷ N. Soni,² V. Sopko,¹²⁶ B. Sopko,¹²⁶ M. Sosebee,⁷ R. Soualah,^{163a,163c} A. Soukharev,¹⁰⁶ S. Spagnolo,^{71a,71b} F. Spanò,⁷⁵ R. Spighi,^{19a} G. Spigo,²⁹ F. Spila,^{131a,131b} R. Spiwoks,²⁹ M. Spousta,¹²⁵ T. Spreitzer,¹⁵⁷ B. Spurlock,⁷ R. D. St. Denis,⁵² J. Stahlman,¹¹⁹ R. Stamen,^{57a} E. Stanecka,³⁸ R. W. Stanek,⁵ C. Stanescu,^{133a} S. Stapnes,¹¹⁶ E. A. Starchenko,¹²⁷ J. Stark,⁵⁴ P. Staroba,¹²⁴ P. Starovoitov,⁹⁰ A. Staude,⁹⁷ P. Stavina,^{143a} G. Stavropoulos,¹⁴ G. Steele,⁵² P. Steinbach,⁴³ P. Steinberg,²⁴ I. Stekl,¹²⁶ B. Stelzer,¹⁴¹ H. J. Stelzer,⁸⁷ O. Stelzer-Chilton,^{158a} H. Stenzel,⁵¹ S. Stern,⁹⁸ K. Stevenson,⁷⁴ G. A. Stewart,²⁹ J. A. Stillings,²⁰ M. C. Stockton,⁸⁴ K. Stoerig,⁴⁷ G. Stoicea,^{25a} S. Stonjek,⁹⁸ P. Strachota,¹²⁵ A. R. Stradling,⁷ A. Straessner,⁴³ J. Strandberg,¹⁴⁶ S. Strandberg,^{145a,145b} A. Strandlie,¹¹⁶ M. Strang,¹⁰⁸ E. Strauss,¹⁴² M. Strauss,¹¹⁰ P. Strizenec,^{143b} R. Ströhmer,¹⁷² D. M. Strom,¹¹³ J. A. Strong,^{75,a} R. Stroynowski,³⁹ J. Strube,¹²⁸ B. Stugu,¹³ I. Stumer,^{24,a} J. Stupak,¹⁴⁷ P. Sturm,¹⁷³ N. A. Styles,⁴¹ D. A. Soh,^{150,v} D. Su,¹⁴² HS. Subramania,² A. Succurro,¹¹ Y. Sugaya,¹¹⁵ T. Sugimoto,¹⁰⁰ C. Suhr,¹⁰⁵ K. Suita,⁶⁶ M. Suk,¹²⁵ V. V. Sulin,⁹³ S. Sultansoy,^{3d} T. Sumida,⁶⁷ X. Sun,⁵⁴ J. E. Sundermann,⁴⁷ K. Suruliz,¹³⁸ S. Sushkov,¹¹ G. Susinno,^{36a,36b} M. R. Sutton,¹⁴⁸ Y. Suzuki,⁶⁵ Y. Suzuki,⁶⁶ M. Svatos,¹²⁴ Yu. M. Sviridov,¹²⁷ S. Swedish,¹⁶⁷ I. Sykora,^{143a} T. Sykora,¹²⁵ B. Szeless,²⁹ J. Sánchez,¹⁶⁶ D. Ta,¹⁰⁴ K. Tackmann,⁴¹ A. Taffard,¹⁶² R. Tafirout,^{158a} N. Taiblum,¹⁵² Y. Takahashi,¹⁰⁰ H. Takai,²⁴ R. Takashima,⁶⁸ H. Takeda,⁶⁶ T. Takeshita,¹³⁹ Y. Takubo,⁶⁵ M. Talby,⁸² A. Talyshev,^{106,g} M. C. Tamsett,²⁴ J. Tanaka,¹⁵⁴ R. Tanaka,¹¹⁴ S. Tanaka,¹³⁰ S. Tanaka,⁶⁵ Y. Tanaka,⁹⁹ A. J. Tanasijczuk,¹⁴¹ K. Tani,⁶⁶ N. Tannoury,⁸² G. P. Tappern,²⁹ S. Tapprogge,⁸⁰ D. Tardif,¹⁵⁷ S. Tarem,¹⁵¹ F. Tarrade,²⁸ G. F. Tartarelli,^{88a} P. Tas,¹²⁵ M. Tasevsky,¹²⁴ E. Tassi,^{36a,36b} M. Tatarkhanov,¹⁴ Y. Tayalati,^{134d} C. Taylor,⁷⁶ F. E. Taylor,⁹¹ G. N. Taylor,⁸⁵ W. Taylor,^{158b} M. Teinturier,¹¹⁴ M. Teixeira Dias Castanheira,⁷⁴ P. Teixeira-Dias,⁷⁵ K. K. Temming,⁴⁷ H. Ten Kate,²⁹ P. K. Teng,¹⁵⁰ S. Terada,⁶⁵ K. Terashi,¹⁵⁴ J. Terron,⁷⁹ M. Testa,⁴⁶ R. J. Teuscher,^{157,k} J. Thadome,¹⁷³ J. Therhaag,²⁰ T. Theveneaux-Pelzer,⁷⁷ M. Thioye,¹⁷⁴ S. Thoma,⁴⁷ J. P. Thomas,¹⁷ E. N. Thompson,³⁴ P. D. Thompson,¹⁷ P. D. Thompson,¹⁵⁷ A. S. Thompson,⁵² L. A. Thomsen,³⁵ E. Thomson,¹¹⁹ M. Thomson,²⁷ R. P. Thun,⁸⁶ F. Tian,³⁴ M. J. Tibbetts,¹⁴ T. Tic,¹²⁴ V. O. Tikhomirov,⁹³ Y. A. Tikhonov,^{106,g} S. Timoshenko,⁹⁵ P. Tipton,¹⁷⁴ F. J. Tique Aires Viegas,²⁹ S. Tisserant,⁸² B. Toczek,³⁷ T. Todorov,⁴ S. Todorova-Nova,¹⁶⁰ B. Toggerson,¹⁶² J. Tojo,⁶⁵ S. Tokár,^{143a} K. Tokunaga,⁶⁶ K. Tokushuku,⁶⁵ K. Tollefson,⁸⁷ M. Tomoto,¹⁰⁰ L. Tompkins,³⁰ K. Toms,¹⁰² G. Tong,^{32a} A. Tonoyan,¹³ C. Topfel,¹⁶ N. D. Topilin,⁶⁴ I. Torchiani,²⁹ E. Torrence,¹¹³ H. Torres,⁷⁷ E. Torró Pastor,¹⁶⁶ J. Toth,^{82,bb} F. Touchard,⁸² D. R. Tovey,¹³⁸ T. Trefzger,¹⁷² L. Tremblet,²⁹ A. Tricoli,²⁹ I. M. Trigger,^{158a} S. Trincaz-Duvoid,⁷⁷ T. N. Trinh,⁷⁷ M. F. Tripijana,⁶⁹ W. Trischuk,¹⁵⁷ A. Trivedi,^{24,aa} B. Trocmé,⁵⁴ C. Troncon,^{88a} M. Trottier-McDonald,¹⁴¹ M. Trzebinski,³⁸ A. Trzupek,³⁸ C. Tsarouchas,²⁹ J. C.-L. Tseng,¹¹⁷ M. Tsiakiris,¹⁰⁴ P. V. Tsiarehshka,⁸⁹ D. Tsonou,^{4,cc} G. Tsipolitis,⁹ V. Tsiskaridze,⁴⁷

E. G. Tskhadadze,^{50a} I. I. Tsukerman,⁹⁴ V. Tsulaia,¹⁴ J.-W. Tsung,²⁰ S. Tsuno,⁶⁵ D. Tsybychev,¹⁴⁷ A. Tua,¹³⁸ A. Tudorache,^{25a} V. Tudorache,^{25a} J. M. Tuggle,³⁰ M. Turala,³⁸ D. Turecek,¹²⁶ I. Turk Cakir,^{3e} E. Turlay,¹⁰⁴ R. Turra,^{88a,88b} P. M. Tuts,³⁴ A. Tykhonov,⁷³ M. Tylmad,^{145a,145b} M. Tyndel,¹²⁸ G. Tzanakos,⁸ K. Uchida,²⁰ I. Ueda,¹⁵⁴ R. Ueno,²⁸ M. Uglan,¹³ M. Uhlenbrock,²⁰ M. Uhrmacher,⁵³ F. Ukegawa,¹⁵⁹ G. Unal,²⁹ D. G. Underwood,⁵ A. Undrus,²⁴ G. Unel,¹⁶² Y. Unno,⁶⁵ D. Urbaniec,³⁴ G. Usai,⁷ M. Uslenghi,^{118a,118b} L. Vacavant,⁸² V. Vacek,¹²⁶ B. Vachon,⁸⁴ S. Vahsen,¹⁴ J. Valenta,¹²⁴ P. Valente,^{131a} S. Valentineti,^{19a,19b} S. Valkar,¹²⁵ E. Valladolid Gallego,¹⁶⁶ S. Vallecorsa,¹⁵¹ J. A. Valls Ferrer,¹⁶⁶ H. van der Graaf,¹⁰⁴ E. van der Kraaij,¹⁰⁴ R. Van Der Leeuw,¹⁰⁴ E. van der Poel,¹⁰⁴ D. van der Ster,²⁹ N. van Eldik,⁸³ P. van Gemmeren,⁵ Z. van Kesteren,¹⁰⁴ I. van Vulpen,¹⁰⁴ M. Vanadia,⁹⁸ W. Vandelli,²⁹ G. Vandoni,²⁹ A. Vaniachine,⁵ P. Vankov,⁴¹ F. Vannucci,⁷⁷ F. Varela Rodriguez,²⁹ R. Vari,^{131a} E. W. Varnes,⁶ D. Varouchas,¹⁴ A. Vartapetian,⁷ K. E. Varvell,¹⁴⁹ V. I. Vassilakopoulos,⁵⁵ F. Vazeille,³³ G. Vegni,^{88a,88b} J. J. Veillet,¹¹⁴ C. Vellidis,⁸ F. Veloso,^{123a} R. Veness,²⁹ S. Veneziano,^{131a} A. Ventura,^{71a,71b} D. Ventura,¹³⁷ M. Venturi,⁴⁷ N. Venturi,¹⁵⁷ V. Vercesi,^{118a} M. Verducci,¹³⁷ W. Verkerke,¹⁰⁴ J. C. Vermeulen,¹⁰⁴ A. Vest,⁴³ M. C. Vetterli,^{141,e} I. Vichou,¹⁶⁴ T. Vickey,^{144b,ff} O. E. Vickey Boeriu,^{144b} G. H. A. Viehhauser,¹¹⁷ S. Viel,¹⁶⁷ M. Villa,^{19a,19b} M. Villaplana Perez,¹⁶⁶ E. Vilucchi,⁴⁶ M. G. Vincker,²⁸ E. Vinek,²⁹ V. B. Vinogradov,⁶⁴ M. Virchaux,^{135,a} J. Virzi,¹⁴ O. Vitells,¹⁷⁰ M. Viti,⁴¹ I. Vivarelli,⁴⁷ F. Vives Vaque,² S. Vlachos,⁹ D. Vladoiu,⁹⁷ M. Vlasak,¹²⁶ N. Vlasov,²⁰ A. Vogel,²⁰ P. Vokac,¹²⁶ G. Volpi,⁴⁶ M. Volpi,⁸⁵ G. Volpini,^{88a} H. von der Schmitt,⁹⁸ J. von Loeben,⁹⁸ H. von Radziewski,⁴⁷ E. von Toerne,²⁰ V. Vorobel,¹²⁵ A. P. Vorobiev,¹²⁷ V. Vorwerk,¹¹ M. Vos,¹⁶⁶ R. Voss,²⁹ T. T. Voss,¹⁷³ J. H. Vosseveld,⁷² N. Vranjes,¹³⁵ M. Vranjes Milosavljevic,¹⁰⁴ V. Vrba,¹²⁴ M. Vreeswijk,¹⁰⁴ T. Vu Anh,⁴⁷ R. Vuillermet,²⁹ I. Vukotic,¹¹⁴ W. Wagner,¹⁷³ P. Wagner,¹¹⁹ H. Wahlen,¹⁷³ J. Wakabayashi,¹⁰⁰ J. Walbersloh,⁴² S. Walch,⁸⁶ J. Walder,⁷⁰ R. Walker,⁹⁷ W. Walkowiak,¹⁴⁰ R. Wall,¹⁷⁴ P. Waller,⁷² C. Wang,⁴⁴ H. Wang,¹⁷¹ H. Wang,^{32b,gg} J. Wang,¹⁵⁰ J. Wang,⁵⁴ J. C. Wang,¹³⁷ R. Wang,¹⁰² S. M. Wang,¹⁵⁰ A. Warburton,⁸⁴ C. P. Ward,²⁷ M. Warsinsky,⁴⁷ P. M. Watkins,¹⁷ A. T. Watson,¹⁷ I. J. Watson,¹⁴⁹ M. F. Watson,¹⁷ G. Watts,¹³⁷ S. Watts,⁸¹ A. T. Waugh,¹⁴⁹ B. M. Waugh,⁷⁶ M. Weber,¹²⁸ M. S. Weber,¹⁶ P. Weber,⁵³ A. R. Weidberg,¹¹⁷ P. Weigell,⁹⁸ J. Weingarten,⁵³ C. Weiser,⁴⁷ H. Wellenstein,²² P. S. Wells,²⁹ M. Wen,⁴⁶ T. Wenaus,²⁴ S. Wendler,¹²² Z. Weng,^{150,v} T. Wengler,²⁹ S. Wenig,²⁹ N. Wermes,²⁰ M. Werner,⁴⁷ P. Werner,²⁹ M. Werth,¹⁶² M. Wessels,^{57a} C. Weydert,⁵⁴ K. Whalen,²⁸ S. J. Wheeler-Ellis,¹⁶² S. P. Whitaker,²¹ A. White,⁷ M. J. White,⁸⁵ S. R. Whitehead,¹¹⁷ D. Whiteson,¹⁶² D. Whittington,⁶⁰ F. Wicek,¹¹⁴ D. Wicke,¹⁷³ F. J. Wickens,¹²⁸ W. Wiedenmann,¹⁷¹ M. Wielers,¹²⁸ P. Wienemann,²⁰ C. Wiglesworth,⁷⁴ L. A. M. Wiik-Fuchs,⁴⁷ P. A. Wijeratne,⁷⁶ A. Wildauer,¹⁶⁶ M. A. Wildt,^{41,r} I. Wilhelm,¹²⁵ H. G. Wilkens,²⁹ J. Z. Will,⁹⁷ E. Williams,³⁴ H. H. Williams,¹¹⁹ W. Willis,³⁴ S. Willocq,⁸³ J. A. Wilson,¹⁷ M. G. Wilson,¹⁴² A. Wilson,⁸⁶ I. Wingerter-Seez,⁴ S. Winkelmann,⁴⁷ F. Winklmeier,²⁹ M. Wittgen,¹⁴² M. W. Wolter,³⁸ H. Wolters,^{123a,i} W. C. Wong,⁴⁰ G. Wooden,⁸⁶ B. K. Wosiek,³⁸ J. Wotschack,²⁹ M. J. Woudstra,⁸³ K. W. Wozniak,³⁸ K. Wraight,⁵² C. Wright,⁵² M. Wright,⁵² B. Wrona,⁷² S. L. Wu,¹⁷¹ X. Wu,⁴⁸ Y. Wu,^{32b,hh} E. Wulf,³⁴ R. Wunstorf,⁴² B. M. Wynne,⁴⁵ S. Xella,³⁵ M. Xiao,¹³⁵ S. Xie,⁴⁷ Y. Xie,^{32a} C. Xu,^{32b,x} D. Xu,¹³⁸ G. Xu,^{32a} B. Yabsley,¹⁴⁹ S. Yacoob,^{144b} M. Yamada,⁶⁵ H. Yamaguchi,¹⁵⁴ A. Yamamoto,⁶⁵ K. Yamamoto,⁶³ S. Yamamoto,¹⁵⁴ T. Yamamura,¹⁵⁴ T. Yamanaka,¹⁵⁴ J. Yamaoka,⁴⁴ T. Yamazaki,¹⁵⁴ Y. Yamazaki,⁶⁶ Z. Yan,²¹ H. Yang,⁸⁶ U. K. Yang,⁸¹ Y. Yang,⁶⁰ Y. Yang,^{32a} Z. Yang,^{145a,145b} S. Yanush,⁹⁰ Y. Yao,¹⁴ Y. Yasu,⁶⁵ G. V. Ybeles Smit,¹²⁹ J. Ye,³⁹ S. Ye,²⁴ M. Yilmaz,^{3c} R. Yoosofmiya,¹²² K. Yorita,¹⁶⁹ R. Yoshida,⁵ C. Young,¹⁴² S. Youssef,²¹ D. Yu,²⁴ J. Yu,⁷ J. Yu,¹¹¹ L. Yuan,^{32a,ii} A. Yurkewicz,¹⁰⁵ B. Zabinski,³⁸ V. G. Zaets,¹²⁷ R. Zaidan,⁶² A. M. Zaitsev,¹²⁷ Z. Zajacova,²⁹ L. Zanello,^{131a,131b} P. Zarzhitsky,³⁹ A. Zaytsev,¹⁰⁶ C. Zeitnitz,¹⁷³ M. Zeller,¹⁷⁴ M. Zeman,¹²⁴ A. Zemla,³⁸ C. Zender,²⁰ O. Zenin,¹²⁷ T. Ženiš,^{143a} Z. Zinonos,^{121a,121b} S. Zenz,¹⁴ D. Zerwas,¹¹⁴ G. Zevi della Porta,⁵⁶ Z. Zhan,^{32d} D. Zhang,^{32b,gg} H. Zhang,⁸⁷ J. Zhang,⁵ X. Zhang,^{32d} Z. Zhang,¹¹⁴ L. Zhao,¹⁰⁷ T. Zhao,¹³⁷ Z. Zhao,^{32b} A. Zhemchugov,⁶⁴ S. Zheng,^{32a} J. Zhong,¹¹⁷ B. Zhou,⁸⁶ N. Zhou,¹⁶² Y. Zhou,¹⁵⁰ C. G. Zhu,^{32d} H. Zhu,⁴¹ J. Zhu,⁸⁶ Y. Zhu,^{32b} X. Zhuang,⁹⁷ V. Zhuravlov,⁹⁸ D. Ziemska,⁶⁰ R. Zimmermann,²⁰ S. Zimmermann,²⁰ S. Zimmermann,⁴⁷ M. Ziolkowski,¹⁴⁰ R. Zitoun,⁴ L. Živković,³⁴ V. V. Zmouchko,^{127,a} G. Zobernig,¹⁷¹ A. Zoccoli,^{19a,19b} Y. Zolnierowski,⁴ A. Zsenei,²⁹ M. zur Nedden,¹⁵ V. Zutshi,¹⁰⁵ and L. Zwalinski²⁹

(ATLAS Collaboration)

¹University at Albany, Albany, New York, USA²Department of Physics, University of Alberta, Edmonton, Alberta, Canada^{3a}Department of Physics, Ankara University, Ankara, Turkey

- ^{3b}*Department of Physics, Dumlupinar University, Kutahya, Turkey*
- ^{3c}*Department of Physics, Gazi University, Ankara, Turkey*
- ^{3d}*Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey*
- ^{3e}*Turkish Atomic Energy Authority, Ankara, Turkey*
- ⁴*LAPP, CNRS/IN2P3 and Université de Savoie, Annecy-le-Vieux, France*
- ⁵*High Energy Physics Division, Argonne National Laboratory, Argonne, Illinois, USA*
- ⁶*Department of Physics, University of Arizona, Tucson, Arizona, USA*
- ⁷*Department of Physics, The University of Texas at Arlington, Arlington, Texas, USA*
- ⁸*Physics Department, University of Athens, Athens, Greece*
- ⁹*Physics Department, National Technical University of Athens, Zografou, Greece*
- ¹⁰*Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan*
- ¹¹*Institut de Física d'Altes Energies and Departament de Física de la Universitat Autònoma de Barcelona and ICREA, Barcelona, Spain*
- ^{12a}*Institute of Physics, University of Belgrade, Belgrade, Serbia*
- ^{12b}*Vinca Institute of Nuclear Sciences, Belgrade, Serbia*
- ¹³*Department for Physics and Technology, University of Bergen, Bergen, Norway*
- ¹⁴*Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, California, USA*
- ¹⁵*Department of Physics, Humboldt University, Berlin, Germany*
- ¹⁶*Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland*
- ¹⁷*School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom*
- ^{18a}*Department of Physics, Bogazici University, Istanbul, Turkey*
- ^{18b}*Division of Physics, Dogus University, Istanbul, Turkey*
- ^{18c}*Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey*
- ^{18d}*Department of Physics, Istanbul Technical University, Istanbul, Turkey*
- ^{19a}*INFN Sezione di Bologna, Italy*
- ^{19b}*Dipartimento di Fisica, Università di Bologna, Bologna, Italy*
- ²⁰*Physikalisches Institut, University of Bonn, Bonn, Germany*
- ²¹*Department of Physics, Boston University, Boston, Massachusetts, USA*
- ²²*Department of Physics, Brandeis University, Waltham, Massachusetts, USA*
- ^{23a}*Universidade Federal do Rio de Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil*
- ^{23b}*Federal University of Juiz de Fora (UFJF), Juiz de Fora, Brazil*
- ^{23c}*Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei, Brazil*
- ^{23d}*Instituto de Física, Universidade de Sao Paulo, Sao Paulo, Brazil*
- ²⁴*Physics Department, Brookhaven National Laboratory, Upton New York, United States of America, USA*
- ^{25a}*National Institute of Physics and Nuclear Engineering, Bucharest, Romania*
- ^{25b}*University Politehnica Bucharest, Bucharest, Romania*
- ^{25c}*West University in Timisoara, Timisoara, Romania*
- ²⁶*Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina*
- ²⁷*Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom*
- ²⁸*Department of Physics, Carleton University, Ottawa, Ontario, Canada*
- ²⁹*CERN, Geneva, Switzerland*
- ³⁰*Enrico Fermi Institute, University of Chicago, Chicago, Illinois, USA*
- ^{31a}*Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile*
- ^{31b}*Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile*
- ^{32a}*Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China*
- ^{32b}*Department of Modern Physics, University of Science and Technology of China, Anhui, China*
- ^{32c}*Department of Physics, Nanjing University, Jiangsu, China*
- ^{32d}*School of Physics, Shandong University, Shandong, China*
- ³³*Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Aubiere Cedex, France*
- ³⁴*Nevis Laboratory, Columbia University, Irvington, New York, USA*
- ³⁵*Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark*
- ^{36a}*INFN Gruppo Collegato di Cosenza, Italy*
- ^{36b}*Dipartimento di Fisica, Università della Calabria, Arcavata di Rende, Italy*
- ³⁷*AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland*
- ³⁸*The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland*
- ³⁹*Physics Department, Southern Methodist University, Dallas, Texas, USA*
- ⁴⁰*Physics Department, University of Texas at Dallas, Richardson, Texas, USA*
- ⁴¹*DESY, Hamburg and Zeuthen, Germany*
- ⁴²*Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany*

- ⁴³*Institut für Kern- und Teilchenphysik, Technical University Dresden, Dresden, Germany*
- ⁴⁴*Department of Physics, Duke University, Durham, North Carolina, USA*
- ⁴⁵*SUPA—School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom*
- ⁴⁶*INFN Laboratori Nazionali di Frascati, Frascati, Italy*
- ⁴⁷*Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg i.Br., Germany*
- ⁴⁸*Section de Physique, Université de Genève, Geneva, Switzerland*
- ^{49a}*INFN Sezione di Genova, Italy*
- ^{49b}*Dipartimento di Fisica, Università di Genova, Genova, Italy*
- ^{50a}*E.Andronikashvili Institute of Physics, Georgian Academy of Sciences, Tbilisi, Georgia*
- ^{50b}*High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia*
- ⁵¹*II Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany*
- ⁵²*SUPA—School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom*
- ⁵³*II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany*
- ⁵⁴*Laboratoire de Physique Subatomique et de Cosmologie, Université Joseph Fourier and CNRS/IN2P3 and Institut National Polytechnique de Grenoble, Grenoble, France*
- ⁵⁵*Department of Physics, Hampton University, Hampton, Virginia, USA*
- ⁵⁶*Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, Massachusetts, USA*
- ^{57a}*Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany*
- ^{57b}*Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany*
- ^{57c}*ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany*
- ⁵⁸*Faculty of Science, Hiroshima University, Hiroshima, Japan*
- ⁵⁹*Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan*
- ⁶⁰*Department of Physics, Indiana University, Bloomington, Indiana, USA*
- ⁶¹*Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria*
- ⁶²*University of Iowa, Iowa City, Iowa, USA*
- ⁶³*Department of Physics and Astronomy, Iowa State University, Ames, Iowa, USA*
- ⁶⁴*Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia*
- ⁶⁵*KEK, High Energy Accelerator Research Organization, Tsukuba, Japan*
- ⁶⁶*Graduate School of Science, Kobe University, Kobe, Japan*
- ⁶⁷*Faculty of Science, Kyoto University, Kyoto, Japan*
- ⁶⁸*Kyoto University of Education, Kyoto, Japan*
- ⁶⁹*Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina*
- ⁷⁰*Physics Department, Lancaster University, Lancaster, United Kingdom*
- ^{71a}*INFN, Sezione di Lecce, Italy*
- ^{71b}*Dipartimento di Fisica, Università del Salento, Lecce, Italy*
- ⁷²*Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom*
- ⁷³*Department of Physics, Jožef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia*
- ⁷⁴*School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom*
- ⁷⁵*Department of Physics, Royal Holloway University of London, Surrey, United Kingdom*
- ⁷⁶*Department of Physics and Astronomy, University College London, London, United Kingdom*
- ⁷⁷*Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France*
- ⁷⁸*Fysiska Institutionen, Lunds Universitet, Lund, Sweden*
- ⁷⁹*Departamento de Física Teórica C-15, Universidad Autónoma de Madrid, Madrid, Spain*
- ⁸⁰*Institut für Physik, Universität Mainz, Mainz, Germany*
- ⁸¹*School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom*
- ⁸²*CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France*
- ⁸³*Department of Physics, University of Massachusetts, Amherst, Massachusetts, USA*
- ⁸⁴*Department of Physics, McGill University, Montreal, Quebec, Canada*
- ⁸⁵*School of Physics, University of Melbourne, Victoria, Australia*
- ⁸⁶*Department of Physics, The University of Michigan, Ann Arbor, Michigan, USA*
- ⁸⁷*Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan, USA*
- ^{88a}*INFN, Sezione di Milano, Italy*
- ^{88b}*Dipartimento di Fisica, Università di Milano, Milano, Italy*
- ⁸⁹*B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus*
- ⁹⁰*National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Republic of Belarus*
- ⁹¹*Department of Physics, Massachusetts Institute of Technology, Cambridge, Massachusetts, USA*
- ⁹²*Group of Particle Physics, University of Montreal, Montreal, Quebec, Canada*
- ⁹³*P.N. Lebedev Institute of Physics, Academy of Sciences, Moscow, Russia*
- ⁹⁴*Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia*
- ⁹⁵*Moscow Engineering and Physics Institute (MEPhI), Moscow, Russia*
- ⁹⁶*Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia*

- ⁹⁷*Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany*
- ⁹⁸*Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany*
- ⁹⁹*Nagasaki Institute of Applied Science, Nagasaki, Japan*
- ¹⁰⁰*Graduate School of Science, Nagoya University, Nagoya, Japan*
- ^{101a}*INFN, Sezione di Napoli, Italy*
- ^{101b}*Dipartimento di Scienze Fisiche, Università di Napoli, Napoli, Italy*
- ¹⁰²*Department of Physics and Astronomy, University of New Mexico, Albuquerque, New Mexico, USA*
- ¹⁰³*Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, The Netherlands*
- ¹⁰⁴*Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, The Netherlands*
- ¹⁰⁵*Department of Physics, Northern Illinois University, DeKalb, Illinois, USA*
- ¹⁰⁶*Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia*
- ¹⁰⁷*Department of Physics, New York University, New York, New York, USA*
- ¹⁰⁸*Ohio State University, Columbus, Ohio, USA*
- ¹⁰⁹*Faculty of Science, Okayama University, Okayama, Japan*
- ¹¹⁰*Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, Oklahoma, USA*
- ¹¹¹*Department of Physics, Oklahoma State University, Stillwater, Oklahoma, USA*
- ¹¹²*Palacký University, RCPTM, Olomouc, Czech Republic*
- ¹¹³*Center for High Energy Physics, University of Oregon, Eugene, Oregon, USA*
- ¹¹⁴*LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France*
- ¹¹⁵*Graduate School of Science, Osaka University, Osaka, Japan*
- ¹¹⁶*Department of Physics, University of Oslo, Oslo, Norway*
- ¹¹⁷*Department of Physics, Oxford University, Oxford, United Kingdom*
- ^{118a}*INFN, Sezione di Pavia, Italy*
- ^{118b}*Dipartimento di Fisica Nucleare e Teorica, Università di Pavia, Pavia, Italy*
- ¹¹⁹*Department of Physics, University of Pennsylvania, Philadelphia, Pennsylvania, USA*
- ¹²⁰*Petersburg Nuclear Physics Institute, Gatchina, Russia*
- ^{121a}*INFN, Sezione di Pisa, Italy*
- ^{121b}*Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy*
- ¹²²*Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, Pennsylvania, USA*
- ^{123a}*Laboratorio de Instrumentacao e Fisica Experimental de Particulas, Lisboa, Portugal*
- ^{123b}*Departamento de Fisica Teorica y del Cosmos and CAFPE, Universidad de Granada, Granada, Portugal*
- ¹²⁴*Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic*
- ¹²⁵*Faculty of Mathematics and Physics, Charles University in Prague, Praha, Czech Republic*
- ¹²⁶*Czech Technical University in Prague, Praha, Czech Republic*
- ¹²⁷*State Research Center Institute for High Energy Physics, Protvino, Russia*
- ¹²⁸*Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom*
- ¹²⁹*Physics Department, University of Regina, Regina, Saskatchewan, Canada*
- ¹³⁰*Ritsumeikan University, Kusatsu, Shiga, Japan*
- ^{131a}*INFN, Sezione di Roma I, Italy*
- ^{131b}*Dipartimento di Fisica, Università La Sapienza, Roma, Italy*
- ^{132a}*INFN, Sezione di Roma Tor Vergata, Italy*
- ^{132b}*Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy*
- ^{133a}*INFN, Sezione di Roma Tre, Italy*
- ^{133b}*Dipartimento di Fisica, Università Roma Tre, Roma, Italy*
- ^{134a}*Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies—Université Hassan II, Casablanca, Morocco*
- ^{134b}*Centre National de l'Energie des Sciences Techniques Nucleaires, Rabat, Morocco*
- ^{134c}*Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech, Morocco*
- ^{134d}*Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda, Morocco*
- ^{134e}*Faculté des Sciences, Université Mohammed V-Agdal, Rabat, Morocco*
- ¹³⁵*DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l'Univers), CEA Saclay (Commissariat a l'Energie Atomique), Gif-sur-Yvette, France*
- ¹³⁶*Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, California, USA*
- ¹³⁷*Department of Physics, University of Washington, Seattle, Washington, USA*
- ¹³⁸*Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom*
- ¹³⁹*Department of Physics, Shinshu University, Nagano, Japan*
- ¹⁴⁰*Fachbereich Physik, Universität Siegen, Siegen, Germany*
- ¹⁴¹*Department of Physics, Simon Fraser University, Burnaby, British Columbia, Canada*
- ¹⁴²*SLAC National Accelerator Laboratory, Stanford, California, USA*
- ^{143a}*Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava, Slovak Republic*
- ^{143b}*Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic*

- ^{144a}*Department of Physics, University of Johannesburg, Johannesburg, South Africa*
^{144b}*School of Physics, University of the Witwatersrand, Johannesburg, South Africa*
^{145a}*Department of Physics, Stockholm University, Sweden*
^{145b}*The Oskar Klein Centre, Stockholm, Sweden*
¹⁴⁶*Physics Department, Royal Institute of Technology, Stockholm, Sweden*
¹⁴⁷*Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook, New York, USA*
¹⁴⁸*Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom*
¹⁴⁹*School of Physics, University of Sydney, Sydney, Australia*
¹⁵⁰*Institute of Physics, Academia Sinica, Taipei, Taiwan*
¹⁵¹*Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel*
¹⁵²*Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel*
¹⁵³*Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece*
¹⁵⁴*International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan*
¹⁵⁵*Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan*
¹⁵⁶*Department of Physics, Tokyo Institute of Technology, Tokyo, Japan*
¹⁵⁷*Department of Physics, University of Toronto, Toronto, Ontario, Canada*
^{158a}*TRIUMF, Vancouver, British Columbia, Canada*
^{158b}*Department of Physics and Astronomy, York University, Toronto, Ontario, Canada*
¹⁵⁹*Institute of Pure and Applied Sciences, University of Tsukuba, 1-1-1 Tennodai, Tsukuba, Ibaraki 305-8571, Japan*
¹⁶⁰*Science and Technology Center, Tufts University, Medford, Massachusetts, USA*
¹⁶¹*Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia*
¹⁶²*Department of Physics and Astronomy, University of California Irvine, Irvine, California, USA*
^{163a}*INFN Gruppo Collegato di Udine, Italy*
^{163b}*ICTP, Trieste, Italy*
^{163c}*Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy*
¹⁶⁴*Department of Physics, University of Illinois, Urbana, Illinois, USA*
¹⁶⁵*Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden*
¹⁶⁶*Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain*
¹⁶⁷*Department of Physics, University of British Columbia, Vancouver, British Columbia, Canada*
¹⁶⁸*Department of Physics and Astronomy, University of Victoria, Victoria, British Columbia, Canada*
¹⁶⁹*Waseda University, Tokyo, Japan*
¹⁷⁰*Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel*
¹⁷¹*Department of Physics, University of Wisconsin, Madison, Wisconsin, USA*
¹⁷²*Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany*
¹⁷³*Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany*
¹⁷⁴*Department of Physics, Yale University, New Haven, Connecticut, USA*
¹⁷⁵*Yerevan Physics Institute, Yerevan, Armenia*
¹⁷⁶*Domaine Scientifique de la Doua, Centre de Calcul CNRS/IN2P3, Villeurbanne Cedex, France*

^aDeceased.

^bAlso at Laboratório de Instrumentação e Física Experimental de Partículas—LIP, Lisboa, Portugal.

^cAlso at Faculdade de Ciências and CFNUL, Universidade de Lisboa, Lisboa, Portugal.

^dAlso at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom.

^eAlso at TRIUMF, Vancouver BC, Canada.

^fAlso at Department of Physics, California State University, Fresno, CA, USA.

^gAlso at Novosibirsk State University, Novosibirsk, Russia.

^hAlso at Fermilab, Batavia, IL, USA.

ⁱAlso at Department of Physics, University of Coimbra, Coimbra, Portugal.

^jAlso at Università di Napoli Parthenope, Napoli, Italy.

^kAlso at Institute of Particle Physics (IPP), Canada.

^lAlso at Department of Physics, Middle East Technical University, Ankara, Turkey.

^mAlso at Louisiana Tech University, Ruston, LA, USA.

ⁿAlso at Department of Physics and Astronomy, University College London, London, United Kingdom.

^oAlso at Group of Particle Physics, University of Montreal, Montreal QC, Canada.

^pAlso at Department of Physics, University of Cape Town, Cape Town, South Africa.

^qAlso at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.

^rAlso at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany.

^sAlso at Manhattan College, New York, NY, USA.

^tAlso at School of Physics, Shandong University, Shandong, China.

^uAlso at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France.

^vAlso at School of Physics and Engineering, Sun Yat-sen University, Guanzhou, China.

^wAlso at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan.

^xAlso at DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l'Univers), CEA Saclay (Commissariat à l'Énergie Atomique), Gif-sur-Yvette, France.

^yAlso at Section de Physique, Université de Genève, Geneva, Switzerland.

^zAlso at Departamento de Física, Universidade de Minho, Braga, Portugal.

^{aa}Also at Department of Physics and Astronomy, University of South Carolina, Columbia, SC, USA.

^{bb}Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.

^{cc}Also at California Institute of Technology, Pasadena, CA, USA.

^{dd}Also at LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France.

^{ee}Also at Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom.

^{ff}Also at Department of Physics, Oxford University, Oxford, United Kingdom.

^{gg}Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.

^{hh}Also at Department of Physics, The University of Michigan, Ann Arbor, MI, USA.

ⁱⁱAlso at Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France.