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Communication

Tuning Patchy Bonds Induced by Critical Casimir Forces

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Abstract: Experimental control of patchy interactions promises new routes for the assembly of complex colloidal structures, but remains challenging. Here, we investigate the role of patch width in the assembly of patchy colloidal particles assembled by critical Casimir forces. The particles are composed of a hydrophobic dumbbell with an equatorial hydrophilic polymer shell, and are synthesized to have well-defined patch-to-shell area ratios. Patch-to-patch binding is achieved in near-critical binary solvents, in which the particle interaction strength and range are controlled by the temperature-dependent solvent correlation length. Upon decreasing the patch-to-shell area ratio, we observe a pronounced change of the bonding morphology towards directed single-bonded configurations, as clearly reflected in the formation of chain-like structures. Computer simulations using an effective critical Casimir pair potential for the patches show that the morphology change results from the geometric exclusion of the increasingly thick hydrophilic particle shells. These results highlight the experimental control of patchy interactions through the engineering of the building blocks on the way towards rationally designed colloidal superstructures.

Keywords: colloidal assembly; patchy colloid; critical Casimir effect

1. Introduction

Colloidal self-assembly, the bottom-up assembly of structures from nano- and micron-size building blocks, has recently seen major breakthroughs. Lessons learnt from colloidal self-assembly can be used to fabricate novel materials with applications in nanoscience, i.e., photonics [1,2], opto-electronics [3], and bio-materials [4]. A number of successful assembly strategies involve either controlling the interactions between the particles by use of polymer-mediated depletion interactions [5], controlled charge [6–8] and magnetic interactions [9,10], and DNA mediated interactions [11–13]; or by designing site-specific particle anisotropy with well-defined geometrical shapes and surface properties [14,15]. Indeed, recent breakthroughs in particle synthesis allow precise control of particle shape and surface affinity, offering patchy particles with specific wetting properties and symmetry [16,17]. Motivated by this experimental progress, many simulation studies have investigated the assembly behaviour of patchy particles [18–25], addressing in detail the effect of patch width, number, and geometry on the resulting topology and phase behaviour of patchy particles. Notwithstanding these simulations, experimental control in the assembly of patchy particles still remains limited. In particular, implementing the high degree of interaction control in a way similar to simulations is challenging. While model patchy interactions have been realized with DNA-coated
particles providing selective bonding [13], a generic route to patchy interactions with full control over
the interaction strength, range, and patch width remains challenging.

Particle interactions of tunable strength and range can be induced in binary solvents close to
their phase separation temperature [26,27]. In particular, close to the solvent critical point, the critical
Casimir effect allows fine control over the particle assembly determined by the temperature-dependent
solvent correlation length [28]. We recently showed that by using patchy particles with hydrophobic
patches and hydrophilic equatorial shells, we can realize selective patch-to-patch binding with tunable
strength and range [29]. The critical Casimir force results from the confinement of solvent fluctuations
between the particle surfaces in the near-critical mixture [30,31], and because it depends on the wetting
boundary conditions, it offers selective bonding of hydrophobic or hydrophilic patches. The advantage
is that the interaction range and strength is set by the solvent correlation length that adjusts with
temperature in a universal manner. This universal, in situ control of particle interactions offers new
opportunities to assemble patchy particles in and out of equilibrium, and to anneal structures once they
are formed. Moreover, by using particles with well-defined patch sizes, this allows us to experimentally
explore the effect of patch size on the assembled structures.

In this paper, we experimentally investigate the effect of patch width on the bonding topology of
patchy particles assembled by critical Casimir forces. Using particles synthesized with well-defined
patch-to-shell area ratios in near-critical binary solvents, we achieve control over both the width
and strength of the patchy particle interaction. We show that decreasing the patch area fraction
towards narrow patches tailors the bonding to change from multi-particle bonding, where more than
two patches bind together, to single bonding, where only one patch binds to another one. Using
dimer particles as an example, we demonstrate that as a result, the topology of colloidal aggregates
becomes much more distinct, clearly favoring directed chain-like structures. Computer simulations
on patchy particles with shape and area ratios taken from the experiment and interacting with an
effective critical Casimir pair potential model are in good agreement with the experimentally observed
topologies, and show that this morphology change is due to geometric exclusion of the bonding
particles. These qualitative results demonstrate the experimental control in tuning the assembly
behaviour via the patchy particle design, offering a new generic approach to the rational design of
colloidal superstructures.

2. Experiment

Our multivalent particles are made by swelling and polymerizing clusters of polymethylmethacrylate
(PMMA) spheres of radius $R = 1.15 \mu m$ with a methacrylate/methacrylic acid shell, resulting in
geometrically well-defined patches with rotational symmetry [14], as shown schematically in Figure 1a.
We use fluorescently labelled 4-methylaminoethylmethacrylate-7-nitrobenzo-2-oxa-1,3-diazol (NBD-MAEM)
PMMA spheres with grafted poly(12-hydroxystearic acid) copolymer (PHS-g-PMMA) as steric stabilizer
as suggested by Elsesser et al. [32]. Colloidal clusters are fabricated by encapsulating the spherical
PMMA particles in toluene droplets in an aqueous phase and selectively evaporating the toluene
(anhydrous, 99.8%, Sigma-Aldrich, St. Louis, MO, USA); see Figure 1a, steps 1 and 2 [14]. To obtain
patchy particles, we grew PMMA shells around the colloidal clusters; see Figure 1a, step 3. By varying the
amount of monomer during the growth of the shells, we achieved particles with different, well-defined
patch-to-matrix area ratios. Specifically, the shells were grown using a suspension of 0.5 mL PMMA clusters
(2% w/w), 0.2 mL 1% w/w sodium dodecyl sulfate, 0.3 mL water, and varying amounts of a monomer
mixture consisting of 97:2:1 w/w methylmethacrylate:methacrylic acid:ethylene glycol dimethacrylate:
0.1, 0.2, and 0.4 mL. The suspension was stirred in a 5 mL round bottom flask with a Teflon-coated
magnetic stir bar for 30 min. We then polymerized the swollen particles in an oil bath at 80 °C for 6 h by
the addition of 350 \mu L 7.4 mM aqueous potassium persulfate (99.99%, Sigma-Aldrich). The potassium
persulfate imparted sulfate charges to the PMMA shells, rendering them hydrophilic with a low surface
charge density. The surface of the patches not covered by shells was equivalent to that of the PMMA
spheres, whose steric stabilization by PHS-g-PMMA makes them hydrophobic. The sample was purified
The angle \( \alpha \) defines the respective orientation of the bonded particles as the angle between the long axis vectors of the dimers. The angle \( \beta \) defines the position of the bond as the angle between the connecting vector of the bonded spheres and the long axis vector of one of the bonded dimers. With this definition, linear, bonded chains have angles \( \alpha \sim 180^\circ \) (or \( 0^\circ \)) and \( \beta \sim 180^\circ \) (or \( 0^\circ \)), while side-by-side bonded particles have angles \( \alpha \sim 0^\circ \) (or \( 180^\circ \)) and \( 60^\circ < \beta < 120^\circ \).
3. Simulation

Similar to the experimental dimer particles, we constructed anisotropic patchy particles from two fused (tangent) spherical particles of radius $R$ (forming a dumbbell) and one hard-core shell of varying radii located at the center of mass of the dumbbell. The spherical particles of the dumbbell interacted with those of a neighboring dumbbell via the potential $u(r) = u_{\text{rep}}(r) + u_{\text{attr}}(r)$, determined by the balance of electrostatic repulsion and critical Casimir attraction [35–37]. Here, $u_{\text{rep}} = A_{\text{rep}} \exp[-(r - 2R)/\xi]$ and $u_{\text{attr}} = A_{\text{Cas}} \exp[-(r - 2R)/\xi]$, with $r$ being the separation of the centers of two spheres not belonging to the same dumbbell, $A_{\text{rep}} = (2\pi R \sigma_{\text{rep}}^2 \xi^2)/(|\varepsilon| \varepsilon_0)$ the strength of the electrostatic repulsion, $A_{\text{Cas}}$ the amplitude of the critical Casimir force and the solvent correlation length $\xi = \xi(\Delta T, \Delta T)$ that depends on both solvent composition and temperature. The central sphere representing the hydrophilic shell was modeled by a hard-core repulsive particle interacting via $u_{\text{rep}}(r)$ with all other spheres. For our solvents with compositions close to the critical composition, we applied the simplified scaling $\xi = \xi_{0}(\Delta T/T_c)^{-0.63}$ of the critical composition, which for the temperatures applied here were within 10% (at $T_a$) and 20% (at 0.6$T_a$) of the actual off-critical correlation length [34]. We measured the surface charge $\sigma_C = -0.188 \mu C/cm^2$ by electrophoresis, and the Debye screening length $l_D = 24$ nm by conductivity measurements; the only free parameters were then the amplitude $A_{\text{Cas}}$ and correlation length $\xi_{0}$ far away from the critical point. As shown in Reference [37], we obtained good fits of experimentally measured pair correlation functions with $\xi_{0} = 1.6$ nm, larger than the typical literature values of around 0.3 nm [38], but in agreement with a recent value of 1.5 nm used in the modeling of critical Casimir potentials [39], and $A_{\text{Cas}} = 2\pi R / \xi$, the amplitude at the critical composition. The potential was thus optimized to reproduce the right temperature dependence of the experimental pair correlation function [37]. While the patchy particle as a whole is anisotropic and interacts via an orientation-dependent potential, the simple model was constructed as a colloidal molecule in which the different parts of the composite particle (molecule) interact isotropically. With this model, we performed Monte Carlo simulations using translation, orientation, and cluster moves. The particles were put in a cubic box and simulated under gravity, with a gravitational length based on the experimental density mismatch, to obtain a quasi-2D system. For consistency, the simulation data were analyzed in exactly the same way as the experimental data.

4. Results and Discussion

Suspended in the binary solvent, the particles exhibited specific binding with temperature-dependent strength and range when the suspension was heated close to $T_c$ [34]. Specifically, we found that in
3MP-poor solvents, the hydrophobic patches attract, while in 3MP-rich solvents the shells attract \cite{29}, consistent with the fact that aggregation occurs between surfaces favoring the minority component of the solvent. Here we investigated the effect of patch width and temperature (interaction strength) on the bonding topology. In particular, one of the major challenges is to achieve single-bonded structures with a narrow bond-angle range, to create the most distinct structures. To investigate the influence of particle patch width on the bonding morphology, we focused on particles with small and large shells (Figure 1a,c). We observed that already at temperatures below $T_a$, where particles interact weakly and no permanent bonds form, the interaction of thin-shell and thick-shell particles is very different, as seen in the representative microscope images in Figure 2b,c. Particles with thin shells showed more ambiguous configurations with many side-by-side bonded configurations (Figure 2b), while particles with thick shells showed predominantly linear, chain-like configurations (Figure 2c). This was confirmed when we plotted the distribution of bond angles $\alpha$ and $\beta$ as shown in Figure 2d,e. In these figures, the probability of a configuration with angle $\alpha$ and $\beta$ is indicated in a shade of red. Frequently observed configurations are demarcated in dark red. Thin-shell particles (Figure 2d) show an abundance of bond angles $\alpha \sim 0^\circ$ and $\beta \sim 60^\circ$ or $120^\circ$, indicating side-by-side bonded configurations. In contrast, for thick shells, the distribution is clearly shifted towards higher values of $\alpha$, reflecting more linear, patch-to-patch bonded configurations. Yet, due to the low attraction, the distribution is still spread and no clear peak is observed at $\alpha, \beta \sim 180^\circ$, which would indicate linear bonded chains.

**Figure 2.** Effect of patch width on weak critical Casimir bonding. (a) Definition of bond angles $\alpha$ and $\beta$; (b,c) Confocal microscope images of weakly bonded dimer particles with thin (b) and thick shells (c). The temperature was set to just below $T_a$, where bonds are weak and break frequently, with an attractive strength of the patches (total potential minimum) of $\sim 3k_B T_a$ as estimated from combined modeling and pair correlation measurements \cite{37}. (d,e) Map showing $\alpha$–$\beta$ angle probability distributions of weakly bonded thin (d) and thick-shell particles (e). Frequently observed configurations are labelled and illustrated below. Bottom: Illustration of typical bonding configurations of dimer particles. The bond angle $\beta$ is visualized by arrows.

This situation changes for temperatures above $T_a$ where due to a stronger Casimir attraction the topologies of bonded structures become more pronounced. Representative images of the aggregate structures above $T_a$ are shown in Figure 3a,b. Clearly, particles with small shells exhibit significant bonding ambiguity: the aggregates show more compact structures, branch points, and a variety of
particles with thick shells showed a strong tendency towards linear bonded structures, as shown by Figure 4. The characteristic change in the bonding morphology is clearly seen, supported by the angle above and Reference [37]. A comparison of simulated thin and thick-shell particles is shown in Figure 3c. The smaller patches resulting from the thicker shells, in contrast, leave less room to bond, narrowing down bond angles, and favoring single-bonded configurations. Yet, the pronounced trend was the shift of the distribution to the upper right corner of the $\alpha$, $\beta$ plane, indicating linear, single-bonded configurations. We thus concluded that the thinner shells and concomitant larger patches offer more room to bond and thus allow a larger range of bond angles and multiple patches to bind with each other, resulting in both sideways and inclined head-to-head bonded configurations, which ultimately yield the close-packed structures abundant in the microscope image (Figure 3c). The smaller patches resulting from the thicker shells, in contrast, leave less room to bond, narrowing down bond angles, and favoring single-bonded configurations.

Figure 3. Effect of patch width on aggregate morphology for strong bonding. Confocal microscope images (a,b) and corresponding distribution of bond angles (c,d) for strongly bonded thin-shell (left) and thick-shell particles (right). The temperatures $0.6\Delta T_u$ and $0.85\Delta T_u$ correspond to solvent correlation length $\xi = 15$ nm and $\xi = 10$ nm, respectively, resulting in a total potential minimum $u_{\text{min}} = -6.5k_B T$ with attractive critical Casimir contribution $u_{\text{att}}(r_{\text{min}}) = -9.2k_B T$ (a,c), and $u_{\text{min}} = -4.1k_B T$ with $u_{\text{att}}(r_{\text{min}}) = -6.3k_B T$ (b,d) [37]. Clear change in the bonding morphology from close-packed to head-to-head bonding was observed. The most frequently observed bonding configurations are labelled; see bottom of Figure 2 for illustrations.

These conclusions were supported by Monte Carlo simulations of dimer particles interacting via effective pair potentials, where the area ratio between the particle patches and shell was taken into account by tuning the hard-core shell to that of the experimental particles (see Simulation section above and Reference [37]). A comparison of simulated thin and thick-shell particles is shown in Figure 4. The characteristic change in the bonding morphology is clearly seen, supported by the angle
distributions, which qualitatively resemble the experimental ones well. In particular, the qualitative change from the bonding ambiguity of several distinct bonding configurations for thin shells to straight head-to-head bonded configurations for thick shells was well reproduced. Again, thin-shell particles showed an abundance of side-by-side bonded (red blobs 1 and 1’) and inclined head-to-head bonded configurations (red blob 2), as well as configurations interpolating in between (red connecting lines). For thick shells, the distribution shifts to the upper right corner: side-by-side bonded configurations vanished, and those with \( \alpha, \beta \) close to 180° emerged (red blob 3). As the only difference in the simulations was the presence of a differently sized hard-core, we concluded that a simple geometric exclusion is enough to cause the change in bonding topology, and for thick enough shells, it is sufficient to achieve single bonded configurations.

These results are also in line with simulations based on other dimer patchy particle models, consisting of a single sphere with opposing patchy caps [23]. Varying the width of the patchy caps, the authors observed that for patches sufficiently small so that a patch can only bind with one other patch (surface coverage below ~6% of the total sphere surface), the simulated particles formed chain-like structures, while for larger patches allowing two or more patchy bonds, the particles formed close-packed configurations, assembling into close-packed planes. We note, however, that for critical Casimir forces, many-body interactions [42] and synchronization phenomena [43] can lead to additional effects not accounted for in the simulations based on pairwise additive potentials; such effects are expected to be particularly pronounced close to \( T_c \), where they can influence the transition from multi-particle bonding to chain-like structures.

![Figure 4. Computer simulations of bonding morphology. Snapshots of computer simulations (a,b) and corresponding bond angle distributions (c,d) of thin (left) and thick-shell particles (right). Temperatures correspond to the experimental temperatures in Figure 3a,b. The experimentally observed narrowing of bond angles towards head-to-head single bonding is clearly reproduced. This narrowing of bond angles is a direct consequence of the geometric exclusion due to the thicker shells.](image)

The dimer patch-to-shell area ratio therefore drastically affects the bonding configuration. We summarize our results in Figure 5. A large patch-to-shell ratio leads to a large range of bond angles and bond positions, resulting in a wide variety of bonding configurations (left). A small
patch-to-shell ratio leads to a much narrower range of bond angles and binding sites, resulting in much more distinct, chain-like structures (right). These results provide guideline for particle design in the experiments. The thick-shell particles employed in this work provide a good step in the direction towards distinct single-bonded structures, and it will be interesting to see which structures result from higher-symmetries, such as trimers, tetramers, and higher-order particles and mixtures thereof.

**Figure 5.** Tuning patchy critical Casimir bonds with patch size. Schematic of the bonding of dimer particles with large (left) and small attractive patches (right). Top: Artist’s impression of critical Casimir bonding of patchy particles. Bottom: Representative particle configurations reconstructed from confocal microscope images of bonded particles. The hydrophobic patches attract in solvents with composition $c_{3MP} < c_c$. Decreasing the patch-to-shell area ratio results in increasingly chain-like, linear structures.

5. Conclusions

By employing critical Casimir forces between patchy particles in near-critical binary solvents, we qualitatively investigated the influence of patch size on the bonding topology of dimer particles and the structure of the resulting aggregates. Selective bonding of the particle patches was achieved by grafting poly-hydroxystearic acid onto the patch surfaces, rendering them hydrophobic, and suspending these particles in 3MP-poor solvents to induce critical Casimir interactions between the patches only. Using dimer particles with varying patch-to-shell ratios, we clearly observed that a large patch-to-shell ratio leads to a large variety of bond angles and bond configurations that, in turn, yield a wide variety of aggregate structures from chains to closely packed particles. A small patch-to-shell ratio, in contrast, narrows down bond angles towards distinct linear configurations, yielding chain-like aggregates. Computer simulations indicate that this morphology change is due to the simple geometric exclusion of the hard shells. These qualitative results provide a proof-of-principle demonstration of the structural control resulting from the combination of critical Casimir forces and well-designed patchy particle building blocks. As the critical Casimir interaction is based on a universal effect, many material and solvent systems can be used, and assembly in more environmental-friendly solvent systems has been recently demonstrated [44].

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**Conflicts of Interest:** The authors declare no conflict of interest.
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