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Measurement of Top Quark Polarization in Top-Antitop Events from Proton-Proton Collisions at $\sqrt{s} = 7$ TeV Using the ATLAS Detector

G. Aad et al.*

(ATLAS Collaboration)

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This Letter presents measurements of the polarization of the top quark in top-antitop quark pair events, using 4.7 $fb^{-1}$ of proton-proton collision data recorded with the ATLAS detector at the Large Hadron Collider at $\sqrt{s} = 7$ TeV. Final states containing one or two isolated leptons (electrons or muons) and jets are considered. Two measurements of $\alpha_t P$, the product of the leptonic spin-analyzing power and the top quark polarization, are performed assuming that the polarization is introduced by either a CP conserving or a maximally CP violating production process. The measurements obtained, $\alpha_t P_{\text{CPC}} = -0.035 \pm 0.014(\text{stat}) \pm 0.037(\text{syst})$ and $\alpha_t P_{\text{CPV}} = 0.020 \pm 0.016(\text{stat})^{+0.013}_{-0.017}(\text{syst})$, are in good agreement with the standard model prediction of negligible top quark polarization.

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The short lifetime of the top quark [1–5] implies that it decays before hadronization takes place, allowing its spin state to be studied using the angular distributions of its decay products. In the standard model (SM), parity conservation in the strong production of top-antitop quark pairs ($t\bar{t}$) in proton-proton ($pp$) collisions implies zero longitudinal polarization of the quarks. A negligible polarization (0.003) is generated by the weak interaction [6]. Physics beyond the SM can induce top quark polarization. For example, models that predict the top quark forward-backward production asymmetry to be larger than the SM prediction, as seen by the Tevatron experiments D0 [7,8] and CDF [9], can generate nonzero polarization of top quarks [10–12]. A first study of polarization in $t\bar{t}$ events has been performed by the D0 Collaboration [8], showing good agreement between the SM prediction and data.

In this Letter, measurements are presented of the polarization of the top quark in inclusive $t\bar{t}$ production in single charged lepton ($t\bar{t} \rightarrow \ell vq\bar{q}b\bar{b}$) and dilepton ($t\bar{t} \rightarrow \ell^+ \ell^- \nu\bar{\nu}bb$) events. The double differential distribution in polar angles, $\theta$, of two of the final-state decay products, with respect to a given quantization axis, is given by [13]

$$\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta_1 d\cos\theta_2} = \frac{1}{4} (1 + \alpha_t P_1 \cos\theta_1 + \alpha_t P_2 \cos\theta_2 - C \cos\theta_1 \cos\theta_2),$$

(1)

where $\theta_1$ ($\theta_2$) is the angular distribution of the decay daughter particle of the top (antitop) quark. Here, $C$ represents the $t\bar{t}$ spin correlation, $P_1$ ($P_2$) represents the degree of polarization of the top (antitop) quark along the chosen quantization axis, and $\alpha_t$ is the spin-analyzing power of the final-state object [14,15], which is a measure of the sensitivity of the daughter particle to the spin state of the parent. At leading order, charged leptons and down-type quarks from $W$-boson decays are predicted to have the largest sensitivity to the spin state of the top quark with a spin-analyzing power of $\alpha = 1$. The helicity basis is used, in which the momentum direction of the top quark in the $t\bar{t}$ center-of-mass frame is chosen as the quantization axis. The $\cos\theta$ distributions of the charged leptons are used as observables to extract a measurement of $\alpha_t P$.

The analysis is based on the full 2011 data set of $pp$ collision events, collected at a center-of-mass energy of 7 TeV by the ATLAS detector [16], corresponding to an integrated luminosity of 4.66 ± 0.08 fb$^{-1}$ [17] after data quality requirements.

ATLAS includes an inner tracking detector, covering a pseudorapidity [18] range $|\eta| < 2.5$, surrounded by a superconducting solenoid providing a 2 T magnetic field. A liquid argon (LAr) electromagnetic sampling calorimeter ($|\eta| < 3.2$), an iron-scintillator tile hadronic calorimeter ($|\eta| < 1.7$), a LAr hadronic calorimeter ($1.4 < |\eta| < 3.2$), and a LAr forward calorimeter ($3.1 < |\eta| < 4.9$) provide the energy measurements. The muon spectrometer consists of tracking chambers covering $|\eta| < 2.7$, and trigger chambers covering $|\eta| < 2.4$, in a toroidal magnetic field. Events considered in this analysis are required to have one high-transverse-momentum ($p_T$) electron or muon that passes requirements of the three-level trigger system.

Both data-driven techniques and Monte Carlo (MC) simulations are used to estimate the sample composition of the data. For each MC sample, generated events are

*Full author list given at the end of the article.

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The data sample is enriched in $t\bar{t}$ events by applying several selection criteria based on the $t\bar{t}$ event topology. The selected $t\bar{t}$ events consist of jets, isolated leptons, and missing transverse momentum from the undetected neutrinos. Jets are reconstructed from clustered energy deposits in the calorimeters using the anti-$k_t$ algorithm [30] with a radius parameter $R = 0.4$. Their energies are corrected to correspond on average to the total energy of the stable particles emitted towards the jet using energy-dependent correction factors derived from simulation, and a residual correction derived from in situ measurements [31,32]. They are required to have $p_T > 25$ GeV and $|\eta| < 2.5$. Furthermore, at least 75% of the scalar sum of the $p_T$ of all the tracks associated with each jet must belong to tracks originating from the primary vertex, which is defined as the vertex with the highest sum of the squared $p_T$ values of the associated tracks in the event. Jets originating from $b$ quarks are selected using a neural network algorithm that combines information about the impact parameter significance of tracks with information about explicitly reconstructed secondary vertices and other variables. At the chosen working point, the algorithm identifies (“$b$ tags”) simulated $b$ jets from top quark decays with 70% efficiency and a rejection factor of about 140 for light partons [33–35]. Reconstructed electrons must have $p_T > 25$ GeV and be associated with a calorimeter cluster in the range $|\eta_{cl}| < 2.47$, excluding the transition between calorimeter sections, $1.37 < |\eta_{cl}| < 1.52$. Selected muons are required to fulfill $|\eta| < 2.5$ and $p_T > 20$ GeV. Each lepton is required to pass quality criteria, to be compatible with being produced at the primary vertex by having a longitudinal impact parameter smaller than 2 mm, and to be isolated from other calorimeter energy deposits and tracks [36]. The $E^{\text{miss}}_T$ is calculated [37] as the magnitude of the negative of the vectorial sum of all energy deposits in the calorimeters, and then corrected for the momenta of the reconstructed muons.

The details of the final event selection depend on the $W$ decay channels. This measurement uses five different channels, containing either one or two electrons or muons in the final state, including the ones coming from $\tau$ decays. The requirements for the single-lepton channels ($\ell +$ jets) include (i) exactly one electron or muon; (ii) at least four jets, at least one of which is $b$ tagged; (iii) $E^{\text{miss}}_T > 30$ GeV for the electron channel and $E^{\text{miss}}_T > 20$ GeV for the muon channel; and (iv) the transverse mass of the $W$ boson to be greater than 30 GeV for the electron channel, while $m_T + E^{\text{miss}}_T > 60$ GeV is required for the muon channel. The transverse mass is computed from the lepton $p_T$ and $\phi$ angle ($p_T^\ell$, $\phi^\ell$) and the direction of the $E^{\text{miss}}_T$ as $m_T = \sqrt{2p_T^\ell E^{\text{miss}}_T[1 - \cos(\phi^\ell - \phi(E^{\text{miss}}_T))]$.

The selection of the dilepton channels ($ee$, $e\mu$, $\mu\mu$) requires (i) exactly two oppositely charged electrons or muons; (ii) at least two jets; (iii) a dilepton invariant mass larger than 15 GeV for all the channels, and more than 10 GeV away from the $Z$ boson mass for the $ee$ and $\mu\mu$ channels; (iv) $E^{\text{miss}}_T > 60$ GeV for the $ee$ and $\mu\mu$ channels; and (v) the scalar sum of the $p_T$ of all selected leptons and jets to be larger than 130 GeV for the $e\mu$ channel.

The major backgrounds are due to vector boson production with additional jets, single top quark production, and to misidentified leptons. Their contributions are estimated using data-driven methods and MC simulation. In particular, the normalization of the dominant background in the $\ell +$ jets channels, $W +$ jets production, is estimated using a measurement of the lepton charge asymmetry in data [38], while the shape of the distribution of $\cos\theta_{T}$ is taken from simulation. In the $ee$ ($\mu\mu$) channel, the normalization of the $Z/\gamma^* +$ jets background with $Z/\gamma^*$ decaying into $ee$ ($\mu\mu$) is determined from data. A $Z/\gamma^* +$ jets enriched control region is defined, where a correction factor for the simulation normalization is derived as a function of the $E^{\text{miss}}_T$ in the event, and applied to the signal region in order to account for possible $E^{\text{miss}}_T$ mismodeling.

The contributions of nonprompt leptons from semileptonic hadron decays and of jets misidentified as leptons (fakes) are determined from data using matrix methods [29,39]. For $\ell +$ jets channels this contribution comes primarily from multijet events, while for dilepton channels it originates primarily from $W +$ jets events where one charged lepton comes from $W$ decay and the other lepton is a nonprompt or fake lepton.

After selection, the expected yields for signal and background compared to data are shown in Table I.

The selected events are reconstructed under the $t\bar{t}$ event hypothesis: jets are associated with particular quarks, and the longitudinal momenta of the neutrinos in the event are determined. From the fully reconstructed decay chain we calculate the momentum of the top quark in the $t\bar{t}$ frame and from it $\cos\theta_t$. 

232002-2
TABLE I. Expected signal and background rounded yields compared to data for each of the five lepton flavor channels considered. The approximate NNLO SM prediction [25] is assumed for $t\bar{t}$ production, and the total systematic and statistical uncertainties are reported.

<table>
<thead>
<tr>
<th>Source</th>
<th>$e + \text{jets}$</th>
<th>$\mu + \text{jets}$</th>
<th>$ee$</th>
<th>$e\mu$</th>
<th>$\mu\mu$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$t\bar{t}$</td>
<td>16 200</td>
<td>26 500</td>
<td>570</td>
<td>4400</td>
<td>1660</td>
</tr>
<tr>
<td>Background</td>
<td>5100</td>
<td>9400</td>
<td>110</td>
<td>700</td>
<td>320</td>
</tr>
<tr>
<td>Total</td>
<td>21 300</td>
<td>35 900</td>
<td>690</td>
<td>5000</td>
<td>1980</td>
</tr>
<tr>
<td>Uncertainty</td>
<td>$\pm$1300</td>
<td>$\pm$1700</td>
<td>$\pm$80</td>
<td>$\pm$500</td>
<td>$\pm$180</td>
</tr>
<tr>
<td>Data</td>
<td>21 956</td>
<td>37 919</td>
<td>740</td>
<td>5328</td>
<td>2057</td>
</tr>
</tbody>
</table>

In the $\ell + \text{jets}$ channels, a kinematic likelihood fit is performed. The likelihood for the event to correspond to a $t\bar{t}$ decay topology is calculated for each possible assignment of four jets selected from the up to five jets of highest $p_T$ in the event, to the two $b$ quark jets and the two jets from the $W$ boson decay [40]. The energies of the jets and the charged lepton, as well as the $E_T^{\text{miss}}$, are allowed to vary within their respective resolutions to best meet the $W$ boson and top quark mass constraints to form the kinematic likelihood. For each assignment, the combined probability is calculated as the product of the maximum kinematic likelihood, the $b$ tagging efficiencies and light-parton rejection probabilities. The highest probability assignment is chosen as the best reconstruction and used to calculate the charged lepton $\cos\theta_{\ell}$. 

In the dilepton channels, the neutrino weighting method is used [41]. Because of the presence of two neutrinos from $W$ boson decays, the final system is underconstrained and assumptions must be made to calculate all particle momenta. Making a hypothesis for the pseudorapidities of the two neutrinos ($\eta_1, \eta_2$), a weight is assigned for each permutation of jets, based on the compatibility of the total neutrino transverse momentum vector and the measured $E_T^{\text{miss}}$, accounting for $E_T^{\text{miss}}$ resolution [37]. For each event, 10 000 different hypotheses for ($\eta_1, \eta_2$) are scanned, drawn from the observed probability distribution in the signal MC sample. The configuration with the maximal weight is selected and used to reconstruct the values of $\cos\theta_{\ell}$ for both charged leptons. Events for which no physical solution can be found with this method are discarded, corresponding to 15% of the selected events in the simulated dilepton $t\bar{t}$ sample. The assumed $\eta$ distributions of the neutrinos are insensitive to top quark polarization.

To extract the value of $\alpha_t P$ from the data, a fit using templates for partially polarized top quarks is performed. The signal templates are obtained by reweighting the top and antitop quark decay products in the simulated $t\bar{t}$ sample according to Eq. (1) using the helicity basis and setting $C$ to the SM value of $t\bar{t}$ spin correlation, $C = 0.31$ [6]. Two different assumptions about the top quark polarization are made to produce two template fits. In one case, the polarization is assumed to be induced by a charge-parity ($CP$) conserving process, labeled CPC, which leads to top and antitop quarks having equal values of $\alpha_t P$ and therefore the same angular distribution for the daughter particles. In the other, maximal $CP$ violation, labeled CPV, is assumed, leading to opposite values of $\alpha_t P$ for the top and antitop quarks. In this case, when a value of $\alpha_t P$ is quoted its sign refers to the sign of the coefficient for positively charged leptons. The positive and negative templates used in the fit are built assuming a value of $\alpha_t P = \pm 0.3$, to guarantee that the differential decay distribution is positive for all values of $\cos\theta_{\ell}$ given the degree of spin correlation. The fraction, $f$, of the positive template component and the $t\bar{t}$ production cross section are fitted simultaneously, in order to reduce the influence of normalization uncertainties on the measured polarization. The polarization is computed as $\alpha_t P = 0.6 f - 0.3$.

For all the considered channels, a template fit is performed with a binned maximum likelihood method on the positive lepton and on the negative lepton distributions. Channels are combined by multiplying their respective likelihood functions. The fitting method is unbiased, which was shown using pseudoeperiments.

For each source of systematic uncertainty, new templates corresponding to the respective 1 standard deviation up and down variation are considered. When an uncertainty is evaluated as the difference between two points, it is symmetrized around the central value. The mean of the distribution of the respective differences between the central fit values and the up and down results from 1000 pseudodata sets is taken as the systematic uncertainty on that source. Systematic uncertainties arising from the same source are treated as being correlated between the different lepton charge and flavor samples.

Detector systematic uncertainties, related to the determination of the energy or momentum scales, resolutions, and efficiencies for jets, electrons, and muons, as well as the $E_T^{\text{miss}}$, are considered [32,37,42–46]. Simulated samples are corrected in order to match the reconstructed object properties observed in data, and the correction factors are varied depending on the uncertainties of their values, in order to estimate the uncertainty on the final measurement. The largest uncertainty in this measurement comes from the jet energy scale.

Systematic uncertainties from the modeling of $t\bar{t}$ production are accounted for using alternative signal templates. These templates are produced by varying the MC event generator, initial- and final- state radiation, color reconnection, fragmentation modeling, and the PDF sets, as detailed in Ref. [47]. The estimation of the uncertainty due to the top quark mass is performed by repeating the fitting procedure using seven samples with different mass settings in the simulation, and interpolating the change in the parameter $f$ corresponding to a variation of the top quark mass of $\pm 1.4$ GeV [48] around the nominal value. Because an assumption on the degree of spin correlation is
made when constructing the template, an additional uncertainty is applied based on the difference in the parton-level spin correlation in simulated \( t\bar{t} \) events between the MC@NLO and POWHEG [49] generators.

For the W + jets background in the \( \ell + \) jets final state, the overall normalization is varied according to the residual uncertainty after the rescaling based on the measured charge asymmetry [38]. In addition, the \( W + \) jets template is varied in shape and normalization by reweighting events according to both the uncertainty in the associated heavy quark production fraction and the parameters of the simulation of extra jets [39]. For the estimate of the systematic uncertainty due to events with nonprompt or fake leptons, the templates are varied according to the uncertainties in the matrix method inputs [29,39]. The MC statistical uncertainty is taken into account by performing pseudoexperiments, where the bin content of each template is varied independently according to the uncertainty. Table II summarizes the sources of systematic uncertainty and their effect on \( \alpha_tP \) for the combined fit. The two largest uncertainties come from jet reconstruction and MC modeling, both affecting the shape of the \( \cos\theta_{t\ell} \) distribution. For sources of systematic uncertainty that do not depend on the lepton charge in the event, the uncertainty in the \( CP \) violating scenario is greatly reduced. These uncertainties push the fit parameters in opposite directions for the samples with different lepton charge, leading to a smaller total uncertainty in the combination.

The results of the fit to the data in the single-lepton and dilepton channels are summarized in Table III. Figure 1 shows the fitted observable in the single-lepton and dilepton final states with the \( CP \) conserving hypothesis, and Fig. 2 shows the same observable in the \( CP \) violating hypothesis. The deviation from the expected linear behavior of the \( \cos\theta_{t\ell} \) distributions is primarily a result of the detector acceptance.

The combined results are

\[
\alpha_tP_{\text{CPC}} = -0.035 \pm 0.014(\text{stat}) \pm 0.037(\text{syst})
\]

in the \( CP \) conserving scenario, and

\[
\alpha_tP_{\text{CPV}} = 0.020 \pm 0.016(\text{stat}) \pm 0.013(\text{syst})
\]

in the \( CP \) violating scenario. The polarization in both scenarios agrees with the SM prediction of negligible polarization. The fitted \( \sigma_P \) is in good agreement with the SM prediction as obtained from NNLO QCD calculations [50,51].

In conclusion, the first measurement of top quark polarization in \( t\bar{t} \) events has been performed for two different scenarios with 4.7 fb\(^{-1}\) of proton-proton collision data at 7 TeV center-of-mass energy with the ATLAS detector at

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**TABLE II.** Summary of the systematic uncertainties on \( \alpha_tP \) for the \( CP \) conserving and \( CP \) violating fits in the combined channels. The systematic uncertainties have been added in quadrature to obtain the total uncertainty.

<table>
<thead>
<tr>
<th>Source</th>
<th>( \Delta \alpha_tP_{\text{CPC}} )</th>
<th>( \Delta \alpha_tP_{\text{CPV}} )</th>
</tr>
</thead>
<tbody>
<tr>
<td>Jet reconstruction</td>
<td>+0.031 – 0.031</td>
<td>+0.009 – 0.005</td>
</tr>
<tr>
<td>Lepton reconstruction</td>
<td>+0.006 – 0.007</td>
<td>+0.002 – 0.001</td>
</tr>
<tr>
<td>( E_T \text{miss} )</td>
<td>+0.008 – 0.007</td>
<td>+0.004 – 0.001</td>
</tr>
<tr>
<td>( t\bar{t} ) modeling</td>
<td>+0.015 – 0.016</td>
<td>+0.005 – 0.013</td>
</tr>
<tr>
<td>Background modeling</td>
<td>+0.011 – 0.010</td>
<td>+0.005 – 0.007</td>
</tr>
<tr>
<td>Template statistics</td>
<td>+0.005 – 0.005</td>
<td>+0.006 – 0.006</td>
</tr>
<tr>
<td>Total systematic</td>
<td>+0.037 – 0.037</td>
<td>+0.013 – 0.017</td>
</tr>
</tbody>
</table>

---

**TABLE III.** Summary of fitted \( \alpha_tP \) in the individual channels for the \( CP \) conserving and \( CP \) violating fits. The uncertainties quoted are first statistical and then systematic.

<table>
<thead>
<tr>
<th>Channel</th>
<th>( \alpha_tP_{\text{CPC}} )</th>
<th>( \alpha_tP_{\text{CPV}} )</th>
</tr>
</thead>
<tbody>
<tr>
<td>ee</td>
<td>0.12 ( \pm 0.10 ) ( ^{+0.09}_{-0.12} )</td>
<td>-0.04 ( \pm 0.12 ) ( ^{+0.18}_{-0.12} )</td>
</tr>
<tr>
<td>( e\mu )</td>
<td>-0.07 ( \pm 0.04 ) ( ^{+0.05}_{-0.06} )</td>
<td>0.00 ( \pm 0.04 ) ( ^{+0.05}_{-0.04} )</td>
</tr>
<tr>
<td>( \mu\mu )</td>
<td>-0.04 ( \pm 0.06 ) ( ^{+0.07}_{-0.07} )</td>
<td>0.04 ( \pm 0.07 ) ( ^{+0.06}_{-0.06} )</td>
</tr>
<tr>
<td>Dilepton</td>
<td>-0.04 ( \pm 0.03 ) ( ^{+0.05}_{-0.05} )</td>
<td>0.01 ( \pm 0.03 ) ( ^{+0.04}_{-0.04} )</td>
</tr>
<tr>
<td>( e+) jets</td>
<td>-0.031 ( \pm 0.028 ) ( ^{+0.043}_{-0.040} )</td>
<td>0.001 ( \pm 0.031 ) ( ^{+0.019}_{-0.019} )</td>
</tr>
<tr>
<td>( \mu+) jets</td>
<td>-0.033 ( \pm 0.021 ) ( ^{+0.039}_{-0.039} )</td>
<td>0.036 ( \pm 0.023 ) ( ^{+0.018}_{-0.017} )</td>
</tr>
<tr>
<td>( \ell+) jets</td>
<td>-0.034 ( \pm 0.017 ) ( ^{+0.038}_{-0.037} )</td>
<td>0.023 ( \pm 0.019 ) ( ^{+0.012}_{-0.011} )</td>
</tr>
<tr>
<td>Combined</td>
<td>-0.035 ( \pm 0.014 ) ( ^{+0.037}_{-0.037} )</td>
<td>0.020 ( \pm 0.016 ) ( ^{+0.013}_{-0.017} )</td>
</tr>
</tbody>
</table>
the LHC. Single-lepton and dilepton final states have been analyzed and no deviation from the SM prediction of negligible polarization is observed for either the CP conserving or maximally CP violating scenario.

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FIG. 2 (color online). The result of the full combined fit to the data with the CP violating hypothesis in (a) the single-lepton channel and (b) the dilepton channel, adding together electrons and muons. It is compared to the polarization templates used and the SM prediction of no polarization. Positively charged leptons are on the left, and negatively charged leptons are on the right.

[18] ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the center of the detector and the z axis along the beam pipe. The x axis points from the IP to the center of the LHC ring, and the y axis points upward. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle θ as η = −ln[tan(θ/2)]. The transverse energy E_T is defined as E sinθ, where E is the energy associated with the calorimeter cell or energy cluster. Similarly, p_T is the momentum component transverse to the beam line.


(ATLAS Collaboration)

1School of Chemistry and Physics, University of Adelaide, Adelaide, Australia
2Physics Department, SUNY Albany, Albany, New York, USA
3Department of Physics, University of Alberta, Edmonton, Alberta, Canada
4aDepartment of Physics, Ankara University, Ankara, Turkey
4bDepartment of Physics, Gazi University, Ankara, Turkey
4cDivision of Physics, TOBB University of Economics and Technology, Ankara, Turkey
4dTurkish Atomic Energy Authority, Ankara, Turkey
5LAPP, CNRS/IN2P3 and Université de Savoie, Annecy-le-Vieux, France
6High Energy Physics Division, Argonne National Laboratory, Argonne Illinois, USA
7School of Physics, University of Arizona, Tucson, Arizona, USA
8Physics Department, The University of Texas at Arlington, Arlington, Texas, USA
9Physics Department, National Technical University of Athens, Athens, Greece
10Physics Department, National Technical University of Athens, Zografou, Greece
11Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
12Institut de Física d’Altes Energies and Departamento de Física de la Universitat Autònoma de Barcelona, Spain
13aInstitute of Physics, University of Belgrade, Belgrade, Serbia

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Vinca Institute of Nuclear Sciences, University of Belgrade, Belgrade, Serbia

Department for Physics and Technology, University of Bergen, Bergen, Norway

Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, California, USA

Department of Physics, Humboldt University, Berlin, Germany

Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland

School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom

Department of Physics, Bogazici University, Istanbul, Turkey

Department of Physics, Dogus University, Istanbul, Turkey

Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey

INFN Sezione di Bologna, Italy

Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy

Physikalisches Institut, University of Bonn, Bonn, Germany

Department of Physics, Boston University, Boston, Massachusetts, USA

Department of Physics, Brandeis University, Waltham, Massachusetts, USA

Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil

Federal University of Juiz de Fora (UFJF), Juiz de Fora, Brazil

Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei, Brazil

Instituto de Física, Universidade de Sao Paulo, Sao Paulo, Brazil

Physics Department, Brookhaven National Laboratory, Upton, New York, USA

National Institute of Physics and Nuclear Engineering, Bucharest, Romania

National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj Napoca, Romania

University Politehnica Bucharest, Bucharest, Romania

West University in Timisoara, Timisoara, Romania

Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina

Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom

Department of Physics, Carleton University, Ottawa, Ontario, Canada

CERN, Geneva, Switzerland

Enrico Fermi Institute, University of Chicago, Chicago, Illinois, USA

Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile

Departamento de Física, Universidad Técnica Federico Santa María, Valparaiso, Chile

Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China

Department of Modern Physics, University of Science and Technology of China, Anhui, China

Department of Physics, Nanjing University, Jiangsu, China

School of Physics, Shandong University, Shandong, China

Physics Department, Shanghai Jiao Tong University, Shanghai, China

Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France

Nevis Laboratory, Columbia University, Irvington, New York, USA

Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark

INFN Gruppo Collegato di Cosenza, Italy

Dipartimento di Física, Università della Calabria, Rende, Italy

AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland

Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland

The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland

Physics Department, Southern Methodist University, Dallas, Texas, USA

Physics Department, University of Texas at Dallas, Richardson, Texas, USA

DESY, Hamburg and Zeuthen, Germany

Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany

Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany

Department of Physics, Duke University, Durham, North Carolina, USA

SUPA-School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom

INFN Laboratori Nazionali di Frascati, Frascati, Italy

Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany

Section de Physique, Université de Genève, Geneva, Switzerland

INFN Sezione di Genova, Italy

Dipartimento di Física, Università di Genova, Genova, Italy

E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi, Georgia

High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia

If Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany

SUPA-School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel

Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece

International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan

Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan

Department of Physics, Tokyo Institute of Technology, Tokyo, Japan

Department of Physics, University of Toronto, Toronto, Ontario, Canada

TRIUMF, Vancouver, British Columbia, Canada

Department of Physics and Astronomy, York University, Toronto, Ontario, Canada

Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan

Department of Physics and Astronomy, Tufts University, Medford, Massachusetts, USA

Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia

Department of Physics and Astronomy, University of California Irvine, Irvine, California, USA

INFN Gruppo Collegato di Udine, Italy

ICTP, Trieste, Italy

Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy

Department of Physics, University of Illinois, Urbana, Illinois, USA

Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden

Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNMT, University of Valencia and CSIC, Valencia, Spain

Department of Physics, University of British Columbia, Vancouver, British Columbia, Canada

Department of Physics and Astronomy, University of Victoria, Victoria, British Columbia, Canada

Department of Physics, University of Warwick, Coventry, United Kingdom

Waseda University, Tokyo, Japan

Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel

Department of Physics, University of Wisconsin, Madison, Wisconsin, USA

Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany

Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany

Department of Physics, Yale University, New Haven, Connecticut, USA

Yerevan Physics Institute, Yerevan, Armenia

Centre de Calcul de l’Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France

\( ^a \)Deceased.

\( ^b \)Also at Department of Physics, King’s College London, London, United Kingdom.

\( ^c \)Also at Laboratorio de Instrumentacao e Fisica Experimental de Particulas-LIP, Lisboa, Portugal.

\( ^d \)Also at Facultade de Ciencias and CFNUL, Universidade de Lisboa, Lisboa, Portugal.

\( ^e \)Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom.

\( ^f \) Also at TRIUMF, Vancouver, British Columbia, Canada.

\( ^g \)Also at Department of Physics, California State University, Fresno, CA, USA.

\( ^h \)Also at Novosibirsk State University, Novosibirsk, Russia.

\( ^i \)Also at Department of Physics, University of Coimbra, Coimbra, Portugal.

\( ^j \)Also at Università di Napoli Parthenope, Napoli, Italy.

\( ^k \)Also at Institute of Particle Physics (IPP), Canada.

\( ^l \)Also at Department of Physics, Middle East Technical University, Ankara, Turkey.

\( ^m \)Also at Louisiana Tech University, Ruston, LA, USA.

\( ^n \)Also at Dep Fisica and CEFITEC of Faculdade de Ciencias e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal.

\( ^o \)Also at Department of Physics and Astronomy, Michigan State University, East Lansing MI, USA.

\( ^p \)Also at Department of Financial and Management Engineering, University of the Aegean, Chios, Greece.

\( ^q \)Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain.

\( ^r \)Also at Department of Physics, University of Cape Town, Cape Town, South Africa.

\( ^s \)Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.

\( ^t \)Also at CERN, Geneva, Switzerland.

\( ^u \)Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany.

\( ^v \)Also at Manhattan College, New York, NY, USA.

\( ^w \)Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.

\( ^x \)Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China.

\( ^y \)Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan.
Also at Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France.

Also at School of Physical Sciences, National Institute of Science Education and Research, Bhubaneswar, India.

Also at Dipartimento di Fisica, Università La Sapienza, Roma, Italy.

Also at DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat à l’Energie Atomique et aux Energies Alternatives), Gif-sur-Yvette, France.

Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia.

Also at Section de Physique, Université de Genève, Geneva, Switzerland.

Also at Department of Physics, The University of Texas at Austin, Austin, TX, USA.

Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.

Also at DESY, Hamburg and Zeuthen, Germany.

Also at International School for Advanced Studies (SISSA), Trieste, Italy.

Also at Department of Physics and Astronomy, University of South Carolina, Columbia, SC, USA.

Also at Faculty of Physics, M.V. Lomonosov Moscow State University, Moscow, Russia.

Also at Nevis Laboratory, Columbia University, Irvington, NY, USA.

Also at Physics Department, Brookhaven National Laboratory, Upton, NY, USA.

Also at Department of Physics, Oxford University, Oxford, United Kingdom.

Also at Department of Physics, The University of Michigan, Ann Arbor, MI, USA.

Also at Discipline of Physics, University of KwaZulu-Natal, Durban, South Africa.