Search for long-lived stopped R-hadrons decaying out of time with pp collisions using the ATLAS detector


Published in: Physical Review D. Particles, Fields, Gravitation, and Cosmology

DOI: 10.1103/PhysRevD.88.112003

Link to publication

Citation for published version (APA):

General rights
It is not permitted to download or to forward/distribute the text or part of it without the consent of the author(s) and/or copyright holder(s), other than for strictly personal, individual use, unless the work is under an open content license (like Creative Commons).

Disclaimer/Complaints regulations
If you believe that digital publication of certain material infringes any of your rights or (privacy) interests, please let the Library know, stating your reasons. In case of a legitimate complaint, the Library will make the material inaccessible and/or remove it from the website. Please Ask the Library: https://uba.uva.nl/en/contact, or a letter to: Library of the University of Amsterdam, Secretariat, Singel 425, 1012 WP Amsterdam, The Netherlands. You will be contacted as soon as possible.

UvA-DARE is a service provided by the library of the University of Amsterdam (http://dare.uva.nl)
I. INTRODUCTION

Long-lived massive particles appear in many theories beyond the Standard Model [1]. They are predicted in R-parity-conserving supersymmetry (SUSY) [2–15] models, such as split SUSY [16,17] and gauge-mediated SUSY breaking [18–24], as well as other scenarios such as universal extra dimensions [25] and leptoquark extensions [26]. For instance, split SUSY addresses the hierarchy problem via the same fine-tuning mechanism that solves the cosmological constant problem; SUSY can be broken at a very high-energy scale, leading to heavy scalars, light fermions, and a light, finely tuned, Higgs boson [16]. Within this phenomenological picture, squarks would thus be much heavier than gluinos, suppressing the gluino decay. If the lifetime of the gluino is long enough, it would hadronize into R-hadrons, color-singlet states of R-mesons ($gq\bar{q}$), R-baryons ($gqqq$), and R-gluinoballs ($gg$). Other models, notably R-parity-violating SUSY, could produce a long-lived squark that would also form an R-hadron, e.g. $\tilde{t}\tilde{q}$. The phenomenology of the top squark (stop) or the bottom squark (sbottom) can be regarded as a heavy, noninteracting spectator, surrounded by a cloud of interacting quarks.

R-hadrons may change their properties through strong interactions with the detector. Most R-mesons would turn into R-baryons [29], and they could also change their electric charge through these interactions. At the Large Hadron Collider (LHC) at CERN [30], the R-hadrons would be produced in pairs and approximately back-to-back in the plane transverse to the beam direction. Some fraction of these R-hadrons would lose all of their momentum, mainly from ionization energy loss, and come to rest within the detector volume, only to decay to a neutralino ($\tilde{\chi}^0$) and hadronic jets at some later time.

A previous search for stopped gluino R-hadrons was performed by the D0 Collaboration [31], which excluded a signal for gluinos with masses up to 250 GeV. That analysis, however, could use only the filled crossings in the Tevatron bunch scheme and suppressed collision-related backgrounds by demanding that there was no nondiffractive interaction in the events. Search techniques similar to those described herein have also been employed by the CMS Collaboration [32,33] using 4 fb$^{-1}$ of 7 TeV data under the assumptions that $m_\tilde{g} - m_\tilde{\chi}^0 > 100$ GeV and $\text{BR}(g\rightarrow g\tilde{\chi}^0) = 100\%$. The resulting limit, at 95% confidence level, is $m_\tilde{g} > 640$ GeV for gluino lifetimes from 10 $\mu$s to 1000 s. ATLAS has up to now studied 31 pb$^{-1}$ of data from 2010 [34], resulting in the limit $m_\tilde{g} > 341$ GeV, under similar assumptions.

This analysis complements previous ATLAS searches for long-lived particles [35,36] that are less sensitive to particles with initial $\beta \ll 1$. By relying primarily on calorimetric measurements, this analysis is also sensitive to events where R-hadron charge flipping may make reconstruction in the inner tracker and the muon system impossible. Detection of stopped R-hadrons could also lead to a...
measurement of their lifetime and decay properties. Moreover, the search is sensitive to any new physics scenario producing large out-of-time energy deposits in the calorimeter with minimal additional detector activity.

II. THE ATLAS DETECTOR AND EVENT RECONSTRUCTION

The ATLAS detector [37] consists of an inner tracking system (ID) surrounded by a thin superconducting solenoid, electromagnetic and hadronic calorimeters, and a muon spectrometer (MS). The ID consists of silicon pixel and microstrip detectors, surrounded by a transition radiation straw-tube tracker. The calorimeter system is based on two active media for the electromagnetic and hadronic calorimeters: liquid argon in the inner barrel and end-cap forward regions, and scintillator tiles (TileCal) in the outer barrel region for $|\eta| < 1.7$ [38]. The calorimeters are segmented into cells that have typical size 0.1 by 0.1 in $\eta - \phi$ space in the TileCal section. The MS, capable of reconstructing tracks within $|\eta| < 2.7$, uses toroidal bending fields generated by three large superconducting magnet systems. There are inner, middle, and outer muon detector stations, each consisting of several precision tracking layers. Local muon track segments (abbreviated to simply “muon segments” from now on) are first found in each station, before being combined into extended muon tracks.

For this analysis, events are reconstructed using “cosmic” settings for the muon system [39], to find muon segments with high efficiency for muons that are “out of time” with respect to the expected time for a muon created from a $pp$ collision and traveling at near the speed of light. Such out-of-time muons are present in the two most important background sources. Cosmic ray muons are present at a random time compared to the bunch-crossing time. Beam-halo muons are in time with proton bunches but may appear early if they hit the muon chamber before particles created from the bunch crossing. Using cosmic settings for the muon reconstruction also loosens requirements on the segment direction and does not require the segment to point towards the interaction point.

Jets are constructed from clusters of calorimeter energy deposits [40] using the anti-$k_t$ jet algorithm [41] with the radius parameter set to $R = 0.4$, which assumes the energetic particles originate from the nominal interaction point. This assumption, while not generally valid for this analysis, has been checked and still accurately quantifies the energy released from the stopped $R$-hadrons decaying occurring in the calorimeter, as shown by comparisons of test-beam studies of calorimeter energy response with simulation [42]. Jet energy is quoted without correcting for the typical fraction of energy not deposited as ionization in the jet cone area, and the minimum jet transverse momentum ($p_T$) is 7 GeV. ATLAS jet reconstruction algorithms are described in more detail elsewhere [43]. The missing transverse momentum ($E_T^{\text{miss}}$) is calculated from the $p_T$ of all reconstructed physics objects in the event, as well as all calorimeter energy clusters not associated with jets.

III. LHC BUNCH STRUCTURE AND TRIGGER STRATEGY

The LHC accelerates two counterrotating proton beams, each divided into 3564 slots for proton bunches separated by 25 ns. When protons are injected into the LHC, not every bunch slot (BCID) is filled. During 2011 and 2012, alternate BCIDs within a “bunch train” were filled, leading to collisions every 50 ns, but there were also many gaps of various lengths between the bunch trains containing adjacent unfilled BCIDs. Filled BCIDs typically had $>10^{11}$ protons. Unfilled BCIDs could contain protons due to diffusion from filled BCIDs, but typically $<10^9$ protons per BCID [44,45]. The filled and unfilled BCIDs from the two beams can combine to make three different “bunch-crossing” scenarios. A paired crossing consists of a filled BCID from each beam colliding in ATLAS and is when $R$-hadrons would be produced. An unpaired crossing has a filled BCID from one beam and an unfilled BCID from the other. Finally, in an empty crossing the BCIDs from both beams are unfilled.

Standard ATLAS analyses use data collected from the paired crossings, while this analysis searches for physics signatures of metastable $R$-hadrons produced in paired crossings and decaying during selected empty crossings. This is accomplished with a set of dedicated low-threshold calorimeter triggers that can fire only in the selected empty or unpaired crossings where the background to this search is much lower. The type of each bunch crossing is defined at the start of each LHC “fill” using the ATLAS beam monitors [46]. Crossings at least six BCIDs after a filled crossing are selected, to reduce background in the muon system.

ATLAS has a three-level trigger system consisting of one hardware and two software levels [47]. Signal candidates are collected using a hardware trigger requiring localized calorimeter activity, a so-called jet trigger, with a 30 GeV transverse energy threshold. This trigger could fire only during an empty crossing at least 125 ns after the most recent paired crossing, such that the detector is mostly free of background from previous interactions. The highest-level software trigger then requires a jet with $p_T > 50$ GeV, $|\eta| < 1.3$, and $E_T^{\text{miss}} > 50$ GeV. The software trigger is more robust against detector noise, keeping the final trigger rate to $<1$ Hz. After offline reconstruction, only 5% of events with more than two muon segments are saved, and no events with more than 20 muon segments are saved, to reduce the data storage needs since events with muon segments are vetoed in the analysis. A data sample enriched with beam-halo muons is also accepted with a lower-threshold jet trigger that fires in the unpaired crossings, and a sample is collected using a trigger that accepts random events from the empty crossings to study background conditions.
IV. DATA SAMPLES

The data used are summarized in Table I, where the corresponding delivered integrated luminosity and recorded live time in the selected empty BCIDs are provided. The early periods of data taking in 2011 are selected as a “cosmic background region” to estimate the rate of background events (mostly from cosmic ray muons, as discussed below). This is motivated by the low integrated luminosity and small number of paired crossings during these initial periods. For a typical signal model that this analysis excludes, less than 3% of events in the cosmic background region are expected to arise from signal processes. As discussed in detail in Sec. VII B, the cosmic ray muon background rate is constant, but the signal production rate scales with luminosity. The later data of 2011 and all of 2012 are used as the “search region,” where an excess of events from R-hadron decays is sought. ATLAS data are taken in runs, which typically span one LHC fill, lasting approximately 1 day.

V. SIMULATION OF R-HADRONS

Monte Carlo simulations are used primarily to determine the reconstruction efficiency and stopping fraction of the R-hadrons, and to study associated systematic uncertainties on the quantities used in the selections. The simulated samples have gluino or squark masses in the range 300–1000 GeV, to which the present analysis is sensitive. The PYTHIA program [48], version 6.427, with CTEQ6L1 parton distribution functions (PDF) [49], is used to simulate pair production of gluinos, stops, or sbottoms. The string hadronization model [50], incorporating specialized hadronization routines [1] for R-hadrons, is used to produce final states containing two R-hadrons.

To compensate for the fact that R-hadron scattering is not strongly constrained, the simulation of R-hadron interactions with matter is handled by a special detector response simulation [29] using GEANT4 [51,52] routines based on several scattering and spectrum models with different sets of assumptions: the generic [29,53], regge [54,55], and intermediate [56] models. Each model makes different assumptions about the R-hadron nuclear interactions and the spectrum of R-hadron states.

![FIG. 1 (color online). The kinetic energies of simulated gluino R-hadrons with a mass of 800 GeV in the generic model are shown at initial production (black line). The energy lost through hadronic interactions with the detector (red, dotted line), electromagnetic ionization (red, dashed line), and total (red, dashed-dotted line) are also shown, for those R-hadrons that have stopped.](112003-3)
The stopping locations are used as input for a second step of PYTHIA where the decays of the $R$-hadrons are simulated. A uniform random time translation is applied in a 25 ns time window, from $-15$ to $+10$ ns, relative to the bunch-crossing time, since the $R$-hadron would decay at a random time relative to the bunch structure of the LHC, but would be triggered by the ATLAS detector during the corresponding empty BCID. These simulated events then proceed through the standard ATLAS digitization simulation [52] and event reconstruction (but with cosmic ray muon settings). The effects of cavern background, a long-lived background component made up of low-energy $\gamma$-hadrons are studied. The planar background in this region is used to generate the corresponding empty BCID. These simulated events then proceed through the standard ATLAS digitization simulation [52] and event reconstruction (but with cosmic ray muon settings). The effects of cavern background, a long-lived background component made up of low-energy $\gamma$-hadrons are studied.

### Table II

The selection efficiency after all selection criteria have been applied, its systematic uncertainty, and the stopping fraction, for all signal samples.

<table>
<thead>
<tr>
<th>$R$-hadron model</th>
<th>Gluino/squark decay</th>
<th>Mass (GeV)</th>
<th>Selection efficiency</th>
<th>Relative systematic uncert.</th>
<th>Stopping fraction (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>$g \to g/gq + \tilde{\chi}^0$</td>
<td>$g$</td>
<td>$\tilde{\chi}^0$</td>
<td>$E &gt; 100$ GeV</td>
<td>$E &gt; 300$ GeV</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to g/gq + \tilde{\chi}^0$</td>
<td>400</td>
<td>100</td>
<td>14.1%</td>
<td>0.5%</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to g/gq + \tilde{\chi}^0$</td>
<td>600</td>
<td>100</td>
<td>15.0%</td>
<td>10.6%</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to g/gq + \tilde{\chi}^0$</td>
<td>800</td>
<td>100</td>
<td>15.5%</td>
<td>13.9%</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to g/gq + \tilde{\chi}^0$</td>
<td>1000</td>
<td>100</td>
<td>14.8%</td>
<td>14.1%</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to g/gq + \tilde{\chi}^0$</td>
<td>400</td>
<td>300</td>
<td>3.4%</td>
<td>&lt;0.1%</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to g/gq + \tilde{\chi}^0$</td>
<td>600</td>
<td>500</td>
<td>4.2%</td>
<td>&lt;0.1%</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to g/gq + \tilde{\chi}^0$</td>
<td>800</td>
<td>700</td>
<td>4.5%</td>
<td>&lt;0.1%</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to g/gq + \tilde{\chi}^0$</td>
<td>1000</td>
<td>900</td>
<td>5.7%</td>
<td>&lt;0.1%</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to t + \tilde{\chi}^0$</td>
<td>600</td>
<td>100</td>
<td>9.9%</td>
<td>7.2%</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to t + \tilde{\chi}^0$</td>
<td>800</td>
<td>100</td>
<td>10.1%</td>
<td>8.9%</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to t + \tilde{\chi}^0$</td>
<td>1000</td>
<td>100</td>
<td>9.5%</td>
<td>8.9%</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to t + \tilde{\chi}^0$</td>
<td>400</td>
<td>20</td>
<td>8.7%</td>
<td>4.3%</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to t + \tilde{\chi}^0$</td>
<td>600</td>
<td>220</td>
<td>9.8%</td>
<td>5.4%</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to t + \tilde{\chi}^0$</td>
<td>800</td>
<td>420</td>
<td>8.3%</td>
<td>4.5%</td>
</tr>
<tr>
<td>Generic</td>
<td>$g \to t + \tilde{\chi}^0$</td>
<td>1000</td>
<td>620</td>
<td>8.7%</td>
<td>4.7%</td>
</tr>
</tbody>
</table>

The effects of cavern background, a long-lived background component made up of low-energy $\gamma$-hadrons are studied.
rays and x rays from low-energy neutrons in the cavern, are not included in the simulation directly, but they are accounted for by measuring the muon activity in the randomly triggered empty data (see Sec. IX). Using the randomly triggered data, the calorimeter activity due to preceding interactions is found to be negligible compared to the jet energy uncertainty and is ignored.

Different models allow the gluinos to decay via the radiative process, $\tilde{g} \rightarrow g \tilde{\chi}_1^0$, or via $\tilde{g} \rightarrow q\bar{q} \tilde{\chi}_1^0$. The results are interpreted assuming either a 50% branching ratio to $g \tilde{\chi}_1^0$ and 50% to $q\bar{q} \tilde{\chi}_1^0$, or 100% to $t\bar{t} \tilde{\chi}_1^0$ as would be the case if the top squark was significantly lighter than the other squarks. Reconstruction efficiencies are typically higher for $q\bar{q} \tilde{\chi}_1^0$ compared to $g \tilde{\chi}_1^0$ decays. The stop (bottom) is assumed to always decay to a top (bottom)
quark and a neutralino. The neutralino mass, $m_{\tilde{\chi}}$, is fixed either to 100 GeV, or such that there is only 100 GeV of free energy left in the decay (a compressed scenario).

VI. CANDIDATE SELECTION

First, events are required to pass tight data quality constraints that verify that all parts of the detector were operating normally. Events with calorimeter noise bursts are rejected; this has negligible impact on signal efficiency. The basic selection criteria, imposed to isolate signal-like events from background events, demand at least one high-energy jet and no muon segments reconstructed in the muon system passing selections. Since most of the $R$-hadrons are produced centrally in $\eta$, the analysis uses only the central barrel of the calorimeter and requires that the leading jet satisfies $|\eta| < 1.2$. In order to reduce the background, the analysis demands the leading jet energy $> 50$ GeV. Up to five additional jets are allowed, more than expected for the signal models considered.

![Graphs showing event yields and selection criteria](image)

FIG. 3 (color online). The event yields in the signal region for candidates with all selection criteria applied (in Table III) including the muon segment veto, but omitting the jet energy $> 100$ GeV requirement. All samples are scaled to represent their anticipated yields in the search region. The top hashed band shows the total statistical uncertainty on the background estimate.
Muon segment veto

The fractional missing $E_T$ is the $E_T^{\text{miss}}$ divided by the leading jet $p_T$ and is required to be $>0.5$. This eliminates background from beam-gas and residual $pp$ events, and has minimal impact on the signal efficiencies. To remove events with a single, narrow energy spike in the calorimeter, due to noise in the electronics or data corruption, events are vetoed where the smallest number of cells containing 90% of the energy deposit of the leading jet (n90) is fewer than four. This n90 > 3 requirement also reduces other background significantly since most large energy deposits from muons in the calorimeter result from hard bremsstrahlung photons, which create short, narrow electromagnetic showers. Large, broad, hadronic showers from deep-inelastic scattering of the muons off nuclei are far rarer. To further exploit the difference between calorimeter energy deposits from muons and the expected signal, the jet width is required to be $>0.04$, where jet width is the $p_T$-weighted $\Delta R$ average of each constituent from the jet axis and $\Delta R = \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2}$. The fraction of the leading jet energy deposited in the TileCal must be $>0.5$, to reduce background from beam halo where the incoming muon cannot be detected due to lack of MS coverage at low radius in the forward region.

The analysis then requires that no muon segment with more than four associated hits in a muon station be reconstructed in the event. Muon segments with a small number of measurements are often present from cavern background, noise, and pileup, as studied in the randomly triggered data. The events before the muon segment veto, only requiring the leading jet energy $>50$ GeV, are studied as a control sample, since the expected signal-to-background ratio is small. Comparisons of the distributions of several jet variables between backgrounds and data can be seen in Fig. 2 for events in this control sample. The backgrounds shown in these figures are estimated using the techniques described in Sec. VII. To remove overlap between the cosmic ray and beam-halo backgrounds in these plots, an event is not considered “cosmic” if it has a muon segment with more than four hits and an angle within 0.2 rad from parallel to the beam line. The same distributions are shown for events after the muon segment veto in Fig. 3. Finally, a leading jet energy requirement of $>100$ or $>300$ GeV defines two signal regions, sensitive to either a small or large mass difference between the $R$-hadron and the neutralino in the signal model, respectively. Table II shows the efficiencies of these selections on the signal simulations, and Table III presents the number of data events surviving each of the imposed selection criteria.

VII. BACKGROUND ESTIMATION

A. Beam-halo background

Protons in either beam can interact with residual gas in the beam pipe, or with the beam pipe itself if they stray off orbit, leading to a hadronic shower. If the interaction takes place several hundred meters from ATLAS, most of the shower is absorbed in shielding or surrounding material before reaching ATLAS. The muons from the shower can survive and enter the detector, traveling parallel to the beam line and in time with the (filled) proton BCIDs [58,59]. The unpaired-crossing data with a jet passing the selection criteria are dominantly beam-halo background. Figure 4 (left) shows an event display of an example beam-halo background candidate event.

To estimate the number of beam-halo events in the empty crossings of the search region, an orthogonal sample of events from the unpaired crossings is used. The ratio of the number of beam-halo events that pass the jet criteria but fail to have a muon segment identified to those that do have a muon segment identified is measured. This ratio is then

<table>
<thead>
<tr>
<th>Selection criteria</th>
<th>Cosmic region</th>
<th>Scaled cosmic region</th>
<th>Search region</th>
<th>$\tilde{g} \rightarrow g/qq\chi^0$</th>
<th>$\tilde{g} \rightarrow g/qq\chi^0$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Trigger</td>
<td>49390 ± 920</td>
<td>152800 ± 2800</td>
<td>218076</td>
<td>71.1 ± 0.7%</td>
<td></td>
</tr>
<tr>
<td>Leading jet $</td>
<td>\eta</td>
<td>&lt; 1.2$</td>
<td>44760 ± 870</td>
<td>138500 ± 2700</td>
<td>202015</td>
</tr>
<tr>
<td>Number of jets &lt; 6</td>
<td>44690 ± 870</td>
<td>138300 ± 2700</td>
<td>201628</td>
<td>60.4 ± 0.7%</td>
<td></td>
</tr>
<tr>
<td>Fractional $E_T^{\text{miss}} &gt; 0.5$</td>
<td>44680 ± 870</td>
<td>138200 ± 2700</td>
<td>201618</td>
<td>60.2 ± 0.7%</td>
<td></td>
</tr>
<tr>
<td>Leading jet n90 &gt; 3</td>
<td>12680 ± 470</td>
<td>39200 ± 1500</td>
<td>85866</td>
<td>54.2 ± 0.7%</td>
<td></td>
</tr>
<tr>
<td>Leading jet width $&gt; 0.04$</td>
<td>4130 ± 260</td>
<td>12770 ± 810</td>
<td>34445</td>
<td>29.9 ± 0.7%</td>
<td></td>
</tr>
<tr>
<td>Leading jet tile $E_T$ frac &gt; 0.5</td>
<td>1640 ± 180</td>
<td>5070 ± 560</td>
<td>5396</td>
<td>22.3 ± 0.6%</td>
<td></td>
</tr>
<tr>
<td>Leading jet energy $&gt;50$ GeV</td>
<td>1640 ± 180</td>
<td>5070 ± 560</td>
<td>5396</td>
<td>22.3 ± 0.6%</td>
<td></td>
</tr>
<tr>
<td>Muon segment veto</td>
<td>2.0 ± 1.4</td>
<td>4.7 ± 3.4</td>
<td>10</td>
<td>15.5 ± 0.6%</td>
<td></td>
</tr>
<tr>
<td>Leading jet energy $&gt;100$ GeV</td>
<td>1.0 ± 1.0</td>
<td>2.4 ± 2.4</td>
<td>5</td>
<td>15.5 ± 0.6%</td>
<td></td>
</tr>
<tr>
<td>Leading jet energy $&gt;300$ GeV</td>
<td>1.0 ± 1.0</td>
<td>2.4 ± 2.4</td>
<td>0</td>
<td>13.9 ± 0.5%</td>
<td></td>
</tr>
</tbody>
</table>
multiplied by the number of beam-halo events observed in the signal region that do have an identified muon segment to give the estimate of the number that do not have a muon segment and thus contribute to background in the signal region. Beam-halo events in the unpaired-crossing data are identified by applying a modified version of the search selection criteria. The muon segment veto is removed, events with leading jet energy \( > 50 \text{ GeV} \) are used, and the \( n_{90} > 3 \) requirement is not applied. Studies show that the muon efficiency is not significantly correlated with the energy or shape of the jet in the calorimeter for beam-halo events. A muon segment is required to be nearly parallel to the beam pipe, \( \theta < 0.2 \) or \( \theta > (\pi - 0.2) \), and have more than four muon station measurements. Next, beam-halo events that failed to leave a muon segment are counted, allowing the ratio of beam-halo events with no muon segment identified to be calculated. Then the number of beam-halo muons in the search region (the empty crossings) that did leave a muon segment is counted. The same selection criteria as listed in Table III are used, omitting the 100 GeV requirement. However, instead of a muon segment veto, a parallel muon segment is required. If no events are present, the uncertainty is taken as \( \pm 1 \) event. Findings are summarized in Table IV.

### B. Cosmic ray muon background

The background from cosmic ray muons is estimated using the cosmic background region (described in Sec. IV). The beam-halo background is estimated for this data sample as described above, and this estimate is subtracted from the observed events passing all selections. Finally, this number of cosmic ray events in the cosmic region is scaled by the ratio of the signal region to cosmic region live times to estimate the cosmic ray background in the signal region. Additionally, the cosmic background estimate is multiplied by the muon-veto efficiency (see Sec. IX) to account for the rejection of background caused by the muon veto. An example cosmic ray muon background event candidate is shown in Fig. 4 (right).

### VIII. EVENT YIELDS

Some candidate event displays are shown in Fig. 5. Distributions of jet variables are plotted for events in the

<table>
<thead>
<tr>
<th>Data region</th>
<th>Leading jet energy (GeV)</th>
<th>Unpaired (all data combined)</th>
<th>Empty</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td>Parallel ( \mu )</td>
<td>No ( \mu )</td>
</tr>
<tr>
<td>Cosmic</td>
<td>50</td>
<td>1634</td>
<td>22</td>
</tr>
<tr>
<td>Search</td>
<td>50</td>
<td>1634</td>
<td>22</td>
</tr>
<tr>
<td>Cosmic</td>
<td>100</td>
<td>1634</td>
<td>22</td>
</tr>
<tr>
<td>Search</td>
<td>100</td>
<td>1634</td>
<td>22</td>
</tr>
<tr>
<td>Cosmic</td>
<td>300</td>
<td>1634</td>
<td>22</td>
</tr>
<tr>
<td>Search</td>
<td>300</td>
<td>1634</td>
<td>22</td>
</tr>
</tbody>
</table>

FIG. 4 (color online). Left: a beam-halo candidate event during an unpaired crossing in data. This event passed all the selection criteria except for the muon segment veto. Right: a cosmic ray muon candidate event during an empty crossing in data. This event passed all the selection criteria except for the muon segment veto. In both plots, white squares filled with red squares show reconstructed energy deposits in TileCal cells above noise threshold (the fraction of red area indicates the amount of energy in the cell), purple bars show a histogram of total energy in projective TileCal towers, jets are shown by open red trapezoids, muon segments by green line segments in each muon station, and muon tracks by continuous thin yellow lines. \( E_T^{\text{miss}} \) is shown as an orange arrow.
jet energy $> 100$ GeV signal region after applying all selection criteria and are compared to the estimated backgrounds in Fig. 6. The shapes of these distributions and event yields are consistent. Table V shows the signal region event yields and background estimates. There is no evidence of an excess of events over the background estimate.

IX. CONTRIBUTIONS TO SIGNAL EFFICIENCY

Quantifying the signal efficiency for the stopped $R$-hadron search presents several unique challenges due to the nonprompt nature of their decays. Specifically there are four sources of inefficiency: stopping fraction (Sec. V), reconstruction efficiency (Table II), accidental muon veto, and probability to have the decay occur in an empty crossing (timing acceptance). Since the first two have been discussed above, only the accidental muon veto and timing acceptance are described here.

A. Accidental muon veto

Operating in the empty crossings has the significant advantage of eliminating collision backgrounds. However because such a stringent muon activity veto is employed, a significant number of events is rejected in the offline analysis due to spurious track segments in the muon system, which are not modeled in the signal simulation. Both $\beta$-decays from activated nuclei and $\delta$-electrons could produce segments with more than four muon station measurements. This effect is separate from a signal decay producing a muon segment that then vetoes the event. To study the rate of muon segments, events from the empty random trigger data in 2011 and 2012 are examined as a function of run number, since the effect can depend strongly on beam conditions. The rate of these events that have a muon segment from noise or other background is calculated. The efficiency per run to pass the muon segment veto is applied on a live-time weighted basis to the cosmic background estimate and varies from 98% at the start of 2011 to 70% at the end of 2012. It is still applied to the cosmic background estimate after the muon veto, since the probability to have the cosmic background event pass the analysis selections and contribute to the signal region events depends on it passing the muon veto. The beam-halo background estimate already implicitly accounts for this effect across run periods. For the signal, this effect is accounted for inside the timing acceptance calculation, on a per-run basis.

B. Timing acceptance

The expected signal decay rate does not scale with instantaneous luminosity. Rather, at any moment in time,
the decay rate is a function of the hypothetical R-hadron lifetime and the entire history of delivered luminosity. For example, for longer R-hadron lifetimes the decay rate anticipated in today’s run is boosted by luminosity delivered yesterday. To address the complicated time behavior of the R-hadron decays, a timing acceptance is defined for each R-hadron lifetime hypothesis, \(\epsilon_{\text{stop}} \times \epsilon_{\text{recon}} \times \epsilon_{\tau} \), as the number of R-hadrons decaying in an empty crossing divided by the total number that stopped. The \(\epsilon_{\tau} \) factor thus accounts for the full history of the delivered luminosity and live time recorded in empty crossings. This means the number of R-hadrons expected to be reconstructed is \(L \times \sigma \times \epsilon_{\text{stop}} \times \epsilon_{\text{recon}} \times \epsilon_{\tau} \), where \(L\) is the integrated luminosity, \(\sigma\) is the R-hadron production cross section weighted by integrated luminosity at 7 and 8 TeV, \(\epsilon_{\text{stop}}\) is the stopping fraction, and \(\epsilon_{\text{recon}}\) is the reconstruction efficiency.

To calculate the timing acceptance for the actual 2011 and 2012 LHC and ATLAS run schedule, measurements are combined of the delivered luminosity in each BCID, the

FIG. 6 (color online). The event yields in the signal region for candidates with all selections (in Table III) except jet energy > 300 GeV. All samples are scaled to represent their anticipated yields in the search region. The top hashed band shows the total statistical uncertainty on the background estimate.
bunch structure of each LHC fill, and the live time recorded in empty crossings during each fill, all kept in the ATLAS online conditions database. The efficiency calculation is split into short and long R-hadron lifetimes, to simplify the calculation. For R-hadron lifetimes less than 10 s, the bunch structure is taken into account, but not the possibility that an R-hadron produced in one run could decay in a later one. For longer R-hadron lifetimes, the bunch structure is averaged over, but the chance that stopped R-hadrons from one run decay in a later one is considered. The resulting timing acceptance is presented in Fig. 7.

### X. SYSTEMATIC UNCERTAINTIES

Three sources of systematic uncertainty on the signal efficiency are studied: the R-hadron interaction with matter, the out-of-time decays in the calorimeters, and the effect of the selection criteria. The total uncertainties, added in quadrature, are shown in Table II. In addition to these, a 2.6% uncertainty is assigned to the luminosity measurement [60], fully correlated between the 2011 and 2012 data. To account for occasional dead-time due to high trigger rates, a 5% uncertainty is assigned to the timing acceptance; this accounts for any mismodeling of the accidental muon veto as well. The gluino, stop, or sbottom pair production cross section uncertainty is not included as a systematic uncertainty but is used when extracting limits on their mass by finding the intersection with the cross section $-1\sigma$ of its uncertainty.

#### A. R-hadron–matter interactions

The various simulated signal samples are used to estimate the systematic uncertainty on the stopping fraction due to the scattering model. There are two sources of theoretical uncertainty: the spectrum of R-hadrons and nuclear interactions. To estimate the effect from different allowed R-hadron states, three different scattering models are employed: generic, Regge, and intermediate (see Sec. V). Each allows a different set of charged states that affect the R-hadron’s electromagnetic interaction with the calorimeters. Since these models have large differences for the R-hadron stopping fraction, limits are quoted separately for each model, rather than including the differences as a systematic uncertainty on the signal efficiency. There is also uncertainty from the modeling of nuclear interactions of the R-hadron with the calorimeter since these can affect the stopping fraction. The effect is estimated by recalculating the stopping fraction after doubling and halving the nuclear cross section. The difference gives a relative uncertainty of 11%, which is used as the systematic uncertainty in limit setting.

#### B. Timing in the calorimeters

Since the R-hadron decay is not synchronized with a BCID it is possible that the calorimeters respond differently to the energy deposits in the simulated signals than in data. The simulation considers only a single BCID for each event; it does not simulate the trigger in multiple BCIDs and the firing of the trigger for the first BCID that passes the trigger. In reality, a decay at $-15$ ns relative to a given BCID might fire the trigger for that BCID, or it may fire the trigger for the following BCID. The reconstructed energy response of the calorimeter can vary between these two
FIG. 8 (color online). Bayesian upper limits on gluino events produced versus gluino mass for the various signal models considered, with gluino lifetimes in the plateau acceptance region between $10^{-5}$ and $10^3$ s, compared to the theoretical expectations.
cases by up to 10% since the reconstruction is optimized for in-time energy deposits. To estimate the systematic uncertainty, the total number of simulated signal events passing the offline selections is studied when varying the timing offset by 5 ns in each direction (keeping the 25 ns range). This variation conservatively covers the timing difference observed between simulated signal jets and cosmic ray muon showers. The minimum and maximum efficiency for each mass point is calculated, and the difference is used as the uncertainty, which is always less than 3% across all mass points.

C. Selection criteria

The systematic uncertainty on the signal efficiency due to selection criteria is evaluated by varying each criterion up and down by its known uncertainty. The uncertainties from each criterion are combined in quadrature and the results are shown in Table II. Varying only the jet energy scale produces most of the total uncertainty from the selection criteria. The jet energy scale uncertainty is taken to be ±10% to allow for nonpointing R-hadron decays and is significantly larger than is used in standard ATLAS analyses. Although test-beam studies showed the energy response agrees between data and simulation for hadronic showers to within a few percent [42], even for nonprojective showers, a larger uncertainty is conservatively assigned to cover possible differences between single pions and full jets, and between the test-beam detectors studied and the final ATLAS calorimeter.

D. Systematic uncertainties on background yield

The systematic uncertainty on the estimated cosmic background arises from the small number of events in the cosmic background region. This statistical uncertainty
is scaled by the same factor used to propagate the cosmic background region data yield into expectations of background events in the search regions. Similarly, for the beam-halo background, a systematic uncertainty is assigned based on the statistical uncertainty of the estimates in the search regions.

**XI. RESULTS**

The predicted number of background events agrees well with the observed number of events in the search region, as shown in Table V. Using these yields, upper limits on the number of pair-produced gluino, stop, or sbottom signal events are calculated with a simple event-counting method and then interpreted as a function of their masses for a given range of lifetimes.

**A. Limit setting**

A Bayesian method is used to set 95% credibility-level upper limits on the number of signal events that could have been produced. For each limit extraction, pseudoexperiments are run. The number of events is sampled from a Poisson distribution, with mean equal to the signal plus background expectation. The systematic uncertainties are taken into account by varying the Poisson mean according to the effect of variations of the sources of the systematic uncertainties [61]. The latter variations are assumed to follow a Gaussian distribution, which is convolved with the Poisson function. A flat prior is used for the signal strength, to be consistent with previous analyses. A Poisson prior gives less conservative limits that are within 10% of those obtained with the flat prior. Since little background is expected and no pseudoexperiment may produce fewer than zero observed events, the distribution of upper limits is bounded from below at $-1.15\sigma$. The input data for the limit-setting algorithm can be seen in Table V. The leading jet energy $>300$ GeV region is used to set the limits, except for the compressed models with a small difference between the gluino or squark mass and $m_{\tilde{g}}$, where the leading jet energy $>100$ GeV signal region is used.

Signal cross sections are calculated to next-to-leading order in the strong coupling constant, adding the resummation of soft gluon emission at next-to-leading-logarithmic accuracy (NLO + NLL) [62–66]. The nominal cross section and the uncertainty are taken from an envelope of cross section predictions using different PDF sets and factorization and renormalization scales, as described in Ref. [67]. The number of expected signal events is given by the signal cross sections at 7 and 8 TeV, weighted by the integrated luminosities in the 2011 and 2012 data. Figures 8 and 9 show the limits on the number of produced signal events for the various signal models considered, for $R$-hadron lifetimes in the plateau acceptance region between $10^{-3}$ and $10^{3}$ seconds.

To provide limits in terms of the gluino, stop, or sbottom masses, the mass is found where the theoretically predicted number of signal events, using the signal cross sections at $-1\sigma$ of their uncertainties, intersects the experimental limit on the number of signal events produced. The gluino, stop, and sbottom mass limits for each of the signal models, for lifetimes in the plateau acceptance region between $10^{-5}$ and $10^{3}$ s, can be seen in Table VI. Figure 10 shows the mass limits for various signal models as a function of lifetime. Limits are also studied as a function of the mass splitting between the gluino or squark and the neutralino. Figure 11 shows the total reconstruction efficiency for a gluino mass of 600 GeV and mass limits for gluino $R$-hadrons, in the generic $R$-hadron model with lifetimes in the plateau acceptance region between $10^{-5}$ and $10^{3}$ s, as a function of this mass splitting between the gluino and the neutralino. A similar dependence exists for stop and sbottom efficiencies and mass limits.

<table>
<thead>
<tr>
<th>Leading jet energy (GeV)</th>
<th>$R$-hadron model</th>
<th>Gluino/squark decay</th>
<th>Neutralino mass (GeV)</th>
<th>Gluino/squark mass limit (GeV)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td></td>
<td></td>
<td>Expected</td>
</tr>
<tr>
<td>100</td>
<td>Generic</td>
<td>$g \to q\bar{q} + \tilde{\chi}^0$</td>
<td>$m_{\tilde{g}} - 100$</td>
<td>526</td>
</tr>
<tr>
<td>100</td>
<td>Generic</td>
<td>$g \to \tilde{t} + \tilde{\chi}^0$</td>
<td>$m_{\tilde{g}} - 380$</td>
<td>694</td>
</tr>
<tr>
<td>300</td>
<td>Generic</td>
<td>$g \to q\bar{q} + \tilde{\chi}^0$</td>
<td>100</td>
<td>731</td>
</tr>
<tr>
<td>300</td>
<td>Generic</td>
<td>$g \to \tilde{t} + \tilde{\chi}^0$</td>
<td>100</td>
<td>700</td>
</tr>
<tr>
<td>300</td>
<td>Intermediate</td>
<td>$g \to q\bar{q} + \tilde{\chi}^0$</td>
<td>100</td>
<td>615</td>
</tr>
<tr>
<td>300</td>
<td>Regge</td>
<td>$g \to q\bar{q} + \tilde{\chi}^0$</td>
<td>100</td>
<td>664</td>
</tr>
<tr>
<td>100</td>
<td>Generic</td>
<td>$\tilde{t} \to t + \tilde{\chi}^0$</td>
<td>$m_{\tilde{t}} - 200$</td>
<td>389</td>
</tr>
<tr>
<td>100</td>
<td>Generic</td>
<td>$\tilde{t} \to t + \tilde{\chi}^0$</td>
<td>100</td>
<td>384</td>
</tr>
<tr>
<td>100</td>
<td>Regge</td>
<td>$\tilde{t} \to t + \tilde{\chi}^0$</td>
<td>100</td>
<td>371</td>
</tr>
<tr>
<td>100</td>
<td>Regge</td>
<td>$\tilde{b} \to b + \tilde{\chi}^0$</td>
<td>100</td>
<td>334</td>
</tr>
</tbody>
</table>
FIG. 10 (color online). Bayesian lower limits on gluino, stop, or sbottom mass versus its lifetime, for the two signal regions, with $R$-hadron lifetimes in the plateau acceptance region between $10^{-3}$ and $10^3$ s. An 800 GeV gluino (stop or sbottom), in the generic (Regge) $R$-hadron model is used as a reference for the stopping fraction and reconstruction efficiency.

FIG. 11 (color online). Total reconstruction efficiency and Bayesian lower limits on gluino mass, as a function of the mass splitting between the gluino and the neutralino, in the generic $R$-hadron model with gluino lifetimes in the plateau acceptance region between $10^{-5}$ and $10^3$ s. A 600 GeV gluino is used as a reference for the reconstruction efficiencies.
XII. SUMMARY

An updated search is presented for stopped long-lived gluino, stop, and sbottom R-hadrons decaying in the calorimeter, using a jet trigger operating in empty crossings of the LHC. Data from the ATLAS experiment recorded in 2011 and 2012 are used, from 5.0 and 22.9 fb⁻¹ of pp collisions at 7 and 8 TeV, respectively. The remaining events after all selections are compatible with the expected rate from backgrounds, predominantly cosmic ray and beam-halo muons where no muon track segment is identified. Limits are set on the gluino, stop, and sbottom masses, for different decays, lifetimes, and neutralino masses. With a neutralino of mass 100 GeV, the analysis excludes \( m_\tilde{q} < 832 \text{ GeV} \) (with an expected lower limit of 731 GeV), for a gluino lifetime between \( 10 \, \mu s \) and 1000 s in the generic R-hadron model with equal branching ratios for decays to \( \tilde{q}\tilde{q}' \) and \( \tilde{g}\tilde{g}' \). For the same \( m_\tilde{q} \) and squark lifetime assumptions, stop and sbottom are excluded with masses below 379 and 344 GeV, respectively, in the Regge R-hadron model.

ACKNOWLEDGMENTS

We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently. We acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWF and FWF, Austria; ANAS, Azerbaijan; SSTC, Belarus; CNPq and FAPESP, Brazil; NSERC, NRC, and CFI, Canada; CERN; CONICYT, Chile; CAS, MOST, and NSFC, China; COLCIENCIAS, Colombia; MSMT CR, MPO CR, and VSC CR, Czech Republic; DNRF, DFG, and MPG, and AvH Foundation, Germany; GRSF and NSRF, Greece; ISF, MINERVA, GIF, DIP, and Benoziyo Center, Israel; INFN, Italy; MEXT and JSPS, Japan; CNRST, Morocco; FOM and NWO, Netherlands; BRF and RCN, Norway; MNiSW, Poland; GRICES and FCT, Portugal; MNE/IFA, Romania; MES of Russia and ROSATOM, Russian Federation; JINR; MSTD, Serbia; MSSR, Slovakia; ARRS and MIZŠ, Slovenia; DST/NRF, South Africa; MINECO, Spain; SRC and Wallenberg Foundation, Sweden; SER, SNSF, and Cantons of Bern and Geneva, Switzerland; NSC, Taiwan; TAEK, Turkey; STFC, the Royal Society and Leverhulme Trust, United Kingdom; DOE and NSF, United States of America. The crucial computing support from all WLCG partners is acknowledged gratefully, in particular, from CERN and the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CC-IN2P3 (France), KIT/GridKA (Germany), INFN-CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (UK), and BNL (USA) and in the Tier-2 facilities worldwide.

SEARCH FOR LONG-LIVED STOPPED R-HADRONS ...


[38] ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the center of the detector and the z axis along the beam pipe. The x axis points from the IP to the center of the LHC ring, and the y axis points upward. Cylindrical coordinates (r, φ) are used, φ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle θ as η = −ln tan(θ/2).


SEARCH FOR LONG-LIVED STOPPED $R$-HADRONS

...
G. AAD et al.  

PHYSICAL REVIEW D 88, 112003 (2013)

LAPP, CNRS/IN2P3 and Université de Savoie, Annecy-le-Vieux, France
6 High Energy Physics Division, Argonne National Laboratory, Argonne, Illinois, USA
7 Department of Physics, University of Arizona, Tucson, Arizona, USA
8 Department of Physics, The University of Texas at Arlington, Arlington, Texas, USA
9 Physics Department, University of Athens, Athens, Greece
10 Physics Department, National Technical University of Athens, Zografou, Greece
11 Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
12 Institut de Física d’Altes Energies and Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain
13 Institute of Physics, University of Belgrade, Belgrade, Serbia
13b Vinca Institute of Nuclear Sciences, University of Belgrade, Belgrade, Serbia
14 Department for Physics and Technology, University of Bergen, Bergen, Norway
15 Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, Berkeley, California, USA
16 Department of Physics, Humboldt University, Berlin, Germany
17 Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
18 School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
19a Department of Physics, Bogazici University, Istanbul, Turkey
19b Department of Physics, Dogus University, Istanbul, Turkey
19c Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey
20a INFN Sezione di Bologna, Italy
20b Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy
21 Physikalisches Institut, University of Bonn, Bonn, Germany
22 Department of Physics, Boston University, Boston, Massachusetts, USA
23 Department of Physics, Brandeis University, Waltham, Massachusetts, USA
24a Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio De Janeiro, Brazil
24b Federal University of Juiz de Fora (UFJF), Juiz de Fora, Brazil
24c Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei, Brazil
24d Instituto de Fisica, Universidade de Sao Paulo, Sao Paulo, Brazil
25 Physics Department, Brookhaven National Laboratory, Upton, New York, USA
26a National Institute of Physics and Nuclear Engineering, Bucharest, Romania
26b National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj Napoca, Romania
26c University Politehnica Bucharest, Bucharest, Romania
26d West University in Timisoara, Timisoara, Romania
27 Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
28 Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
29 Department of Physics, Carleton University, Ottawa, Ontario, Canada
30 CERN, Geneva, Switzerland
31 Enrico Fermi Institute, University of Chicago, Chicago, Illinois, USA
32a Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile
32b Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
33a Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China
33b Department of Modern Physics, University of Science and Technology of China, Anhui, China
33c Department of Physics, Nanjing University, Jiangsu, China
33d School of Physics, Shandong University, Shandong, China
34 Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France
35 Nevis Laboratory, Columbia University, Irvington, New York, USA
36 Niels Bohr Institute, University of Copenhagen, København, Denmark
37a INFN Gruppo Collegato di Cosenza, Italy
37b Dipartimento di Fisica, Università della Calabria, Rende, Italy
38a AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland
39 Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland
40 The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland
41 Physics Department, Southern Methodist University, Dallas, Texas, USA
42 DESY, Hamburg and Zeuthen, Germany
43 Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
44 Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
45 Department of Physics, Duke University, Durham, North Carolina, USA
46 SUPA-School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
SEARCH FOR LONG-LIVED STOPPED R-HADRONS . . .

PHYSICAL REVIEW D 88, 112003 (2013)
100 Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
101 Nagasaki Institute of Applied Science, Nagasaki, Japan
102 Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan
103 INFN Sezione di Napoli, Italy
104 Department of Physics and Astronomy, University of New Mexico, Albuquerque, New Mexico, USA
105 Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
106 Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
107 Department of Physics, Northern Illinois University, DeKalb, Illinois, USA
108 Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia
109 Department of Physics, New York University, New York, New York, USA
110 Ohio State University, Columbus, Ohio, USA
111 Faculty of Science, Okayama University, Okayama, Japan
112 Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, Oklahoma, USA
113 Department of Physics, Oklahoma State University, Stillwater, Oklahoma, USA
114 Palacký University, RCPTM, Olomouc, Czech Republic
115 Center for High Energy Physics, University of Oregon, Eugene, Oregon, USA
116 LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France
117 Graduate School of Science, Osaka University, Osaka, Japan
118 Department of Physics, University of Oslo, Oslo, Norway
119 Department of Physics, Oxford University, Oxford, United Kingdom
120 INFN Sezione di Pavia, Italy
121 Department of Physics, University of Pennsylvania, Philadelphia, Pennsylvania, USA
122 Petersburg Nuclear Physics Institute, Gatchina, Russia
123 INFN Sezione di Pisa, Italy
124 Dipartimento di Fisica, Università di Pisa, Pisa, Italy
125 Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, Pennsylvania, USA
126 Laboratorio de Instrumentaciono y del Cosmos y del CAFPE, Universidad de Granada, Granada, Spain
127 Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic
128 Czech Technical University in Prague, Praha, Czech Republic
129 Faculty of Mathematics and Physics, Charles University in Prague, Prague, Czech Republic
130 State Research Center Institute for High Energy Physics, Protvino, Russia
131 Department of Physics, Rutherford Appleton Laboratory, Didcot, United Kingdom
132 Physics Department, University of Regina, Regina, Saskatchewan, Canada
133 Ritsumeikan University, Kusatsu, Shiga, Japan
134 INFN Sezione di Roma I, Italy
135 Dipartimento di Fisica, Università La Sapienza, Roma, Italy
136 INFN Sezione di Roma Tor Vergata, Italy
137 DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat à l’Energie Atomique et aux Energies Alternatives), Gif-sur-Yvette, France
138 Santa Cruz Institute for Particle Physics, University of California, Santa Cruz, California
139 Department of Physics, University of Washington, Seattle, Washington, USA
140 Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
141 Department of Physics, Shinshu University, Nagano, Japan
142 Fachbereich Physik, Universität Siegen, Siegen, Germany
143 Department of Physics, Simon Fraser University, Burnaby, British Columbia, Canada
144 SLAC National Accelerator Laboratory, Stanford, California, USA
145 Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava, Slovak Republic
146 Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
147 Department of Physics, University of Cape Town, Cape Town, South Africa
SEARCH FOR LONG-LIVED STOPPED R-HADRONS... PHYSICAL REVIEW D 88, 112003 (2013)

aDeceased.
bAlso at Department of Physics, King’s College London, London, United Kingdom.
cAlso at Laboratorio de Instrumentacao e Fisica Experimental de Particulas-LIP, Lisboa, Portugal.
dAlso at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.
eAlso at Faculdade de Ciencias and CFNUL, Universidade de Lisboa, Lisboa, Portugal.
fAlso at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom.
gAlso at TRIUMF, Vancouver, British Columbia, Canada.
hAlso at Department of Physics, California State University, Fresno, CA, USA.
iAlso at Novosibirsk State University, Novosibirsk, Russia.
jAlso at Department of Physics, University of Coimbra, Coimbra, Portugal.
kAlso at Università di Napoli Parthenope, Napoli, Italy.
lAlso at Institute of Particle Physics (IPP), Canada.
mAlso at Department of Physics, Middle East Technical University, Ankara, Turkey.
nAlso at Louisiana Tech University, Ruston, LA, USA.
oAlso at Dep Fisica and CEFITEC of Faculdade de Ciencias e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal.
pAlso at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France.
qAlso at Department of Physics and Astronomy, Michigan State University, East Lansing, MI, USA.
rAlso at Department of Financial and Management Engineering, University of the Aegean, Chios, Greece.