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Published in:
Physical Review D. Particles, Fields, Gravitation, and Cosmology

DOI:
10.1103/PhysRevD.90.012004

Link to publication

Citation for published version (APA):

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Download date: 01 Mar 2018
Search for dark matter in events with a $Z$ boson and missing transverse momentum in $pp$ collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector

G. Aad et al.*

(ATLAS Collaboration)

(Received 31 March 2014; published 10 July 2014)

A search is presented for production of dark-matter particles recoiling against a leptonically decaying $Z$ boson in 20.3 fb$^{-1}$ of $pp$ collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector at the Large Hadron Collider. Events with large missing transverse momentum and two oppositely charged electrons or muons consistent with the decay of a $Z$ boson are analyzed. No excess above the Standard Model prediction is observed. Limits are set on the mass scale of the contact interaction as a function of the dark-matter particle mass using an effective field theory description of the interaction of dark matter with quarks or with $Z$ bosons. Limits are also set on the coupling and mediator mass of a model in which the interaction is mediated by a scalar particle.

Astrophysical measurements indicate the existence of nonbaryonic dark matter [1,2]. However, collider-based searches, nuclear scattering experiments, and searches for particles produced from dark-matter annihilation have not yet revealed its particle nature nor discovered its non-gravitational interactions, if they exist [3]. Collider-based searches for weakly interacting massive particles (WIMPs, denoted as $\chi$), specifically $pp \rightarrow \chi \bar{\chi} + X$ at the Large Hadron Collider (LHC) via some unknown intermediate state, are an important facet of the experimental program that includes initial state radiation (ISR), a photon [6,7], or a $W$ or $Z$ boson decaying hadronically [8]. Limits on dark matter produced in the decay of the Higgs boson have also been reported [9]. In this analysis, limits are set using the final state of a $Z$ boson decaying to two oppositely charged electrons or muons, plus missing transverse momentum $E_T^{\text{miss}}$.

Since the nature of the intermediate state mediating the parton-WIMP interaction is not known, a useful approach is to construct an effective field theory (EFT) [10–12]. EFTs have often been used to describe interactions between dark-matter particles and quarks or gluons, but they have recently been extended to describe direct interactions with electroweak bosons [13–15]. In the context of the EFT framework, the WIMP is considered to be the only new particle accessible at LHC energies, in addition to the SM fields. The mediator of the interaction is assumed to be heavy compared to the typical parton interaction energies involved, and the dark-matter particles are also assumed to be produced in pairs.

The EFTs considered in this analysis, depicted in Fig. 1, are expressed in terms of two parameters: $m_\chi$ and a mass scale $M_*$, described in Ref. [10]. $M_*$ parameterizes the coupling between the WIMP and SM particles, where the coupling strength is normalized, or in inverse proportion, to the heavy-mediator mass scale. The coefficients of the Lagrangian’s interaction terms appear as powers of $M_*$, e.g. for the $D1$ operator as $1/M_*^8$ and for the $D5$ and $D9$ operators as $1/M_*^3$. The definition of the $D1$, $D5$, and $D9$ operators and the region of validity of the EFT limits are discussed in Refs. [10,16].

Following the approach of Ref. [13], the coupling of dark matter to electroweak bosons is considered for dimension-5 and dimension-7 operators. The dimension-7 operator

FIG. 1. The diagrams showing different types of $pp \rightarrow \chi \bar{\chi} + Z$ production modes considered in this analysis [13]. (a) shows a diagram that includes an ISR operator, and (b) shows a diagram that includes a $ZZ\chi\bar{\chi}$ operator.

* Full author list given at the end of the article.
The production cross section is proportional to \( j \) as well as \( Z \) or \( \gamma^* \) each contribute to the matrix element. The relative contribution of the \( Z \) and \( \gamma^* \) diagrams is a parameter of the theory.

This analysis considers models of dark-matter production where a \( Z \) boson is radiated as ISR or interacts directly with WIMPs. The latter case of an interaction between a \( Z \) boson and a WIMP is a process not previously investigated in the analysis of LHC experiments.

To complement the EFT analysis, this paper also examines the results in terms of a model in which the intermediate state is specified [17]. In this model a scalar mediator \( \eta \), with mass \( m_\eta \), and a scalar-WIMP coupling strength \( f \) are responsible for the production of the dark-matter particles. The mediator \( \eta \) transforms as a color triplet and an electroweak doublet and has a hypercharge of 1/3. The production cross section is proportional to \( f^4 \). The same final state signature predicted by the EFT, \( Z\eta\eta \), is produced. This process is similar to supersymmetry processes in which the \( \chi \) is a neutralino and the \( \eta \) is a squark doublet, but without a direct gluino analogue. This corresponds to a scenario where the gluino would be too heavy to be produced at the LHC and is therefore irrelevant.

In this analysis, a search for the production of a Z boson with subsequent decay to \( e^+e^- \) or \( \mu^+\mu^- \) in association with large \( E_T^{\text{miss}} \) from the escaping \( Z\chi \) particles is reported, based on 20.3 fb\(^{-1}\) of \( pp \) collision data collected by the ATLAS detector at a center-of-mass energy of \( \sqrt{s} = 8 \) TeV. Only data collected with stable beams and all detector subsystems fully operational are used. Several signal regions with different requirements on the \( E_T^{\text{miss}} \) are defined to best probe the variety of models tested.

The ATLAS detector [18] consists of an inner detector (ID) surrounded by a solenoid that produces a 2 T magnetic field, electromagnetic and hadronic calorimeters, and a muon spectrometer (MS) employing toroidal magnets. The ID measures charged-particle tracks over the full azimuthal angle and in a pseudorapidity [19] range of \( |\eta| < 2.5 \) using silicon pixel, silicon microstrip, and transition-radiation straw-tube detectors, the last of which also distinguishes electrons from heavier charged particles in the range \( |\eta| < 2.0 \). Liquid-argon (LAr) electromagnetic sampling calorimeters cover the range \( |\eta| < 3.2 \) with a typical granularity in \( \Delta\eta \times \Delta\phi \) of 0.025 \( \times \) 0.025. A scintillator-tile calorimeter provides hadronic calorimetry for \( |\eta| < 1.7 \). In the end caps (\( |\eta| > 1.5 \)), LAr is also used for the hadronic calorimeters matching the outer \( |\eta| \) limit of the end cap electromagnetic calorimeters. The LAr forward calorimeters extend the coverage to \( |\eta| < 4.9 \) and provide both the electromagnetic and hadronic energy measurements. The MS covers \( |\eta| < 2.7 \) and provides triggering and precision tracking for muons. A three-level trigger system is used to select interesting events to be recorded for subsequent off-line analysis.

Electrons are required to have transverse energy \( E_T \) larger than 20 GeV and \( |\eta| < 2.47 \). The \( E_T \) is measured from the energy deposited in the electromagnetic calorimeter, and the electron’s direction from the ID track. Electrons must satisfy the medium object quality requirements from Ref. [20] updated for 2012 run conditions, which are based on calorimeter shower shape, ID track quality, and the spatial match between the shower and the track. Electrons must be isolated, satisfying \( \sum_{\Delta R < 0.2} p_T^{\text{track}}/E_T < 0.1 \), where the sum is over the transverse momenta \( p_T^{\text{track}} \) of all other ID tracks with \( p_T^{\text{track}} > 1 \) GeV within a cone of radius \( \Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.2 \) around the electron direction.

Muons are required to have \( p_T > 20 \) GeV and \( |\eta| < 2.5 \). A combined fit of the ID and MS tracks is used to reconstruct the muon \( p_T \). High-quality tracks are ensured by requirements on the number of hits in the ID. Longitudinal and transverse impact parameters, \( z_0 \) and \( d_0 \), respectively, must satisfy \( |z_0| < 10 \) mm and \( |d_0| < 1 \) mm, with respect to the primary vertex, defined as the vertex with the highest \( \sum (p_T^{\text{track}})^2 \). The muon must be isolated, satisfying \( \sum_{\Delta R < 0.2} p_T^{\text{track}}/p_T < 0.1 \); here again, the muon track itself is excluded from the sum.

The anti-k\(_t\) jet algorithm [21] with radius parameter of 0.4 is used to reconstruct jets from topological clusters [22], which are three-dimensional clusters of neighboring energy deposits in the calorimeter cells. A calibration procedure is used in which the raw energy measurement from the calorimeter cluster is corrected to the jet energy scale. Jets are required to have \( p_T > 25 \) GeV and \( |\eta| < 2.5 \). Jets from secondary proton-proton collisions are removed by requiring that most of the tracks associated with the jet, weighted by \( p_T \), originate at the primary vertex.

Since muons may generate delta-ray electrons or radiate photons that produce electron-positron pairs, electrons closer than \( \Delta R = 0.2 \) to a muon that passes the analysis selection are rejected. In addition, electrons and muons closer than \( \Delta R = 0.4 \) to a jet are also rejected.

The measurement of the missing transverse momentum, a vector in the transverse plane, \( \vec{E}_T^{\text{miss}} \), with magnitude \( E_T^{\text{miss}} \), is based on the measurement of the energy collected by the calorimeters and the momenta of muons. Muons and electrons with \( p_T > 10 \) GeV, jets with \( p_T > 20 \) GeV, low-\( p_T \) tracks which do not seed a topological cluster; and topological clusters not associated with a jet are included in the \( \vec{E}_T^{\text{miss}} \) calculation [23].

The candidate signal events were accepted by at least one of the several triggers that require either two leptons with low \( p_T \) or a single lepton with higher \( p_T \). An event must have at least one reconstructed vertex with at least three associated tracks with \( p_T > 400 \) MeV to remove non-collision background. In addition, events must have two oppositely charged electrons or muons with invariant mass \( m_{ee} \in [76, 106] \) GeV to form a Z boson candidate. In order
to suppress events where the \( E_{\text{T}}^{\text{miss}} \) originates from mis-measured jets, the azimuthal angle between the dilepton system and the \( \Delta\phi(E_{\text{T}}^{\text{miss}}, p_T^{\ell}) \), must be greater than 2.5, the absolute value of the pseudorapidity of the dilepton system, \( |\eta^{\ell\ell}| \), must be less than 2.5, and the ratio \( |p_T^{\ell\ell} - E_{\text{T}}^{\text{miss}}|/p_T^{\ell\ell} \) must be less than 0.5, where \( p_T^{\ell\ell} \) is the transverse momentum of the dilepton system. Events are removed if they contain one or more jets with \( p_T > 25 \text{ GeV} \) to suppress top-quark pair background. Similarly, events containing a third lepton with \( p_T > 7 \text{ GeV} \), satisfying looser identification requirements than invoked for the leptons produced in the decay of the Z boson, are removed to suppress diboson background.

The various dark-matter models considered here have different \( E_{\text{T}}^{\text{miss}} \) spectra, leading to a variety of optimal lower thresholds of \( E_{\text{T}}^{\text{miss}} \). Four inclusive signal regions are defined with lower thresholds in \( E_{\text{T}}^{\text{miss}} \) of 150, 250, 350, and 450 GeV.

The dominant background process is \( ZZ \rightarrow \ell^+\ell^-\bar{\nu}\nu \) (\( \ell = e, \mu \)), an irreducible background. The other reducible background is \( WW \rightarrow \ell^+\ell^-\bar{\nu}\nu \), which may only be reduced through the mass window requirement. Reducible backgrounds may have jets produced in association with two leptons, where the jets are misidentified or unreconstructed, such as \( t\bar{t} \rightarrow \ell^+\nu b\ell^-\bar{b} \). ZZ → \( \ell^+\ell^-\bar{q}q \), \( WZ \rightarrow q\ell^+\ell^- \), or \( Z + \text{jets} \) (including decays to \( \tau \) leptons). Additional reducible sources may produce events with more than two leptons where the additional leptons are misidentified or not reconstructed, or less than two leptons where jets are misidentified as leptons, such as \( WZ \rightarrow \ell\nu\ell^+\ell^- \) and \( W + \text{jets} \), respectively.

The WZ and ZZ backgrounds are estimated from Monte Carlo (MC) simulation. The next-to-leading-order MC generator POWHEG BOX 1.0 [24] is used, which models the production from a \( q\bar{q} \) initial state. The ZZ background sample is normalized to include the \( gg \rightarrow ZZ \) contribution using MCFM 6.2 [25]. Parton distribution functions (PDFs) for these samples are modeled using CT10 [26]. The underlying event, as well as parton showers and hadronization, are simulated with PYTHIA 8.165 [27], using the AU2 tune [28]. QED radiative corrections are calculated using PHOTOS++ 3.0 [29]. The detector response is simulated with GEANT 4 [30,31]. The simulated samples are generated with pileup conditions similar to those observed in data. Pileup refers to the multiple interactions occurring in the same, or adjacent, beam bunch crossings as the hard process. Simulated events are reweighted so that the number of pileup interactions has the same distribution as in data events.

The WW, \( t \rightarrow t, Wt \), and \( Z \rightarrow \tau\tau \) backgrounds are estimated from data using the absence of signal in the \( e\mu \) channel and the relative production rate of 1 : 1 : 2 for the \( ee, \mu\mu, \) and \( e\mu \) channels. An \( e\mu \) control region similar to the signal region is defined, and the background estimate for the \( ee \) and \( \mu\mu \) signal regions is obtained from the number of \( e\mu \) events in the control region after correcting for the slightly different electron and muon acceptances and efficiencies.

The \( Z + \text{jets} \) background is estimated using two data-driven techniques. The first method, commonly referred to as the \( ABCD \) method [32], considers the distribution of signal and background events in a phase space defined by two uncorrelated variables, here \( E_{\text{T}}^{\text{miss}} \) and \( \eta^{\ell\ell} \), for which the signal and background have different shapes. The phase space is partitioned into four regions labeled \( A, B, C, \) and \( D \). Region \( A \) is the signal region where selection requirements on both variables are invoked, while regions \( B, C, \) and \( D \) are control regions in which one or both selections are reversed. Contamination by signal events in the control regions is found to be negligible. The number of events in one control region scaled by the ratio of background events in two other control regions estimates the background contribution in the signal region. In the second method, the contribution is measured by fitting the distributions of \( \Delta\phi(E_{\text{T}}^{\text{miss}}, p_T^{\ell\ell}) \) and \( E_{\text{T}}^{\text{miss}} \) at small values \( E_{\text{T}}^{\text{miss}} < 80 \text{ GeV} \) and \( \Delta\phi(E_{\text{T}}^{\text{miss}}, p_T^{\ell\ell}) < 2.5 \) and extrapolating them to the signal regions. The two methods invoke all the standard selection requirements and give consistent results. The \( ABCD \) method is used to provide the background estimate, and the difference between the two is taken as a systematic uncertainty on the estimate.

The \( W + \text{jets} \) background is estimated by reversing the electron isolation condition and loosening identification requirements for one electron in order to obtain a data sample enriched in jets reconstructed as electrons. The resultant \( E_{\text{T}}^{\text{miss}} \) distribution is fitted with a function of the form \( N = A \cdot (E_{\text{T}}^{\text{miss}})^b \) below 300 GeV and extrapolated to the highest 450 GeV signal region. The fitted function is integrated over \( E_{\text{T}}^{\text{miss}} \) to obtain an estimated background above a given \( E_{\text{T}}^{\text{miss}} \) threshold. A normalization factor is derived from data in the low-\( E_{\text{T}}^{\text{miss}} \) region with all the analysis selections applied. This factor is applied to the extrapolated result to obtain an estimate of this background.

Background estimates are validated in signal-depleted control regions that are determined by similar selection criteria used to define the signal region, but where a requirement may be inverted or modified. For the dominant \( ZZ \rightarrow \ell^+\ell^-\bar{\nu}\nu \) background, estimated with MC simulation, the control region probes four-lepton events and is defined by the presence of two pairs of same-flavor, oppositely charged leptons for which the invariant mass of the pairs is within the Z boson mass window. For the subdominant \( WZ \rightarrow \ell\nu\ell^+\ell^- \) background, the control region is characterized by three leptons: a pair of same-flavor, oppositely charged leptons for which the invariant mass is within the Z boson mass window, an additional electron or muon, and a reconstructed \( E_{\text{T}}^{\text{miss}} > 80 \text{ GeV} \). For the minor WW and top-quark background, derived from a data-driven technique, events containing an electron and a
muon $e^-$ with opposite charge are used. The expected signal region event yield is obtained from correction factors that account for the relative dilepton reconstruction efficiencies and the ratio of same-flavor lepton production to mixed flavor. The predicted yields from MC simulation are consistent with the data-driven estimates and in all cases are consistent with the control region yields observed in data.

Samples of $pp \to ZZ$ signals are generated using MADGRAPH 5 1.5.2 [33] with parton showering and hadronization modeled by PYTHIA 8.170 using the MSTW 2008 leading-order PDFs [34] and the AU2 tune. EFT operators $D1$ (scalar, spin independent), $D5$ (vector, spin independent), and $D9$ (tensor, spin dependent), following the definitions of Ref. [10], are representative of the full set of operators in which the $Z$ boson is emitted as ISR. Similar $E_{T}^{miss}$ distributions result from all the operators within each of the three types: scalar, vector, and tensor.

Two examples of the dimension-7 $ZZ\chi\chi$ operator mixtures are considered: one in which the $Z\gamma^*$ contribution is negligible and one in which it is maximal. Therefore the dimension-7 operators are referred to as $ZZ\chi\chi$-maximal-$\gamma^*$ and $ZZ\chi\chi$-no-$\gamma^*$ while the dimension-5 operator is referred to simply as $ZZ\chi\chi$. Cross sections for a few representative operators and for the scalar-mediator theory with representative coupling constant, $f = 6$, and $m_\eta = 1$ TeV are given in Table I.

Samples of $pp \to ZZ\chi\chi$ events are propagated through the ATLAS detector using the full simulation of the ID and muon trackers and the parameterized simulation of the calorimeter [30], tuned to full simulation and data. The shapes of the simulated $E_{T}^{miss}$ distributions for the signal operators are shown in Fig. 2 compared to the dominant SM background process $ZZ \to \ell^+\ell^-\nu\bar{\nu}$.

Contributions to the systematic uncertainty of the expected SM backgrounds are due largely to experimental sources affecting the $E_{T}^{miss}$ measurement and to the efficiencies for the reconstruction and identification of electrons and muons. For example, when $E_{T}^{miss} > 120$ GeV, the experimental systematic uncertainty for the ZZ background is dominated by the jet-energy scale (1.7% and 2.3% for electron and muon final states, respectively) and the electron and muon momentum scale (2.3% and 0.8%, respectively). Smaller systematic uncertainties are associated with the $E_{T}^{miss}$ measurement and with the efficiencies for the reconstruction and identification of electrons and muons.

For the dominant background, $ZZ \to \ell^+\ell^-\nu\bar{\nu}$, determined from simulated samples, systematic theoretical uncertainties are derived from the generator differences, QCD factorization and renormalization scales, and PDF modeling. The largest theoretical uncertainty, the generator difference, is evaluated as the difference in yields calculated from samples simulated with SHERPA 1.4.1 [35] and POWHEG BOX. The systematic uncertainties associated with the ZZ background are summarized in Table II for each signal region. The luminosity uncertainty is 2.8% and

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**Table I.** The power dependence of $1/M_\chi$ for the EFT and the cross sections of WIMP production in association with an on-shell $Z$ boson for various EFT operators and the scalar-mediator theory are shown. For the calculation of the production cross section, $M_\chi$ is taken to be 1 TeV for the EFT operators. The coupling constant of the scalar-mediator theory, $f$, is taken to be 6 and the mass of the mediator, $m_\eta$, is 1 TeV.

<table>
<thead>
<tr>
<th>$m_\eta$ [GeV]</th>
<th>$D5$</th>
<th>$D9$</th>
<th>$ZZ\chi\chi$ max. $\gamma^*$</th>
<th>Scalar mediator</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>7.2</td>
<td>130</td>
<td>3.1</td>
<td>830</td>
</tr>
<tr>
<td>10</td>
<td>7.1</td>
<td>120</td>
<td>3.1</td>
<td>810</td>
</tr>
<tr>
<td>200</td>
<td>5.6</td>
<td>89</td>
<td>2.0</td>
<td>300</td>
</tr>
<tr>
<td>400</td>
<td>3.1</td>
<td>47</td>
<td>0.83</td>
<td>70</td>
</tr>
<tr>
<td>1000</td>
<td>0.25</td>
<td>3.4</td>
<td>0.023</td>
<td>...</td>
</tr>
<tr>
<td>$M_\chi^{-1}$ power</td>
<td>2</td>
<td>2</td>
<td>3</td>
<td>...</td>
</tr>
</tbody>
</table>

**Table II.** Summary of the systematic uncertainties for the largest background process: ($ZZ \to \ell^+\ell^-\nu\bar{\nu}$). Statistical uncertainties are from MC simulation sample size.

<table>
<thead>
<tr>
<th>Uncertainty source</th>
<th>$E_{T}^{miss}$ threshold [GeV]</th>
</tr>
</thead>
<tbody>
<tr>
<td>Statistical [%]</td>
<td>2 6 13 24</td>
</tr>
<tr>
<td>Experimental [%]</td>
<td>3 6 9 8</td>
</tr>
<tr>
<td>Theoretical [%]</td>
<td>36 37 37 38</td>
</tr>
<tr>
<td>Luminosity [%]</td>
<td>3 3 3 3</td>
</tr>
<tr>
<td>Total [%]</td>
<td>36 38 40 46</td>
</tr>
</tbody>
</table>
TABLE III. Observed yields and expected SM backgrounds in each signal region. Statistical, systematic, and luminosity uncertainties are added in quadrature to give the total background estimate and uncertainties.

<table>
<thead>
<tr>
<th>Process</th>
<th>( E_{\text{T}}^{\text{miss}} ) threshold [GeV]</th>
</tr>
</thead>
<tbody>
<tr>
<td>( ZZ )</td>
<td>150</td>
</tr>
<tr>
<td>( WZ )</td>
<td>8.0 ± 3.1</td>
</tr>
<tr>
<td>( WW, \bar{t}, \bar{t} )</td>
<td>1.9 ± 1.4</td>
</tr>
<tr>
<td>( Z \rightarrow \tau^+ \tau^- )</td>
<td>0.1 ± 0.1</td>
</tr>
<tr>
<td>( Z + \text{jets} )</td>
<td>0.1 ± 0.1</td>
</tr>
<tr>
<td>( W + \text{jets} )</td>
<td>0.1 ± 0.1</td>
</tr>
<tr>
<td>Total</td>
<td>52 ± 18</td>
</tr>
<tr>
<td>Data</td>
<td>45</td>
</tr>
</tbody>
</table>

The dark-matter particle mass is \( m_{\chi} = 200 \text{ GeV} \). The last bin contains the events with \( E_{\text{T}}^{\text{miss}} > 450 \text{ GeV} \). The ratio of data to simulated backgrounds is also shown. The band shows the experimental systematic uncertainties on the ratio.

Poisson distribution, for either signal and background or background alone, includes the effect of varying the nuisance parameters.

The \( E_{\text{T}}^{\text{miss}} \) region with the best expected limit is used to calculate the observed limit for each operator and mass point. Limits on the cross section for \( pp \rightarrow Z\chi\bar{\chi} \) production are translated into limits on the mass scale of the effective operators mediating the interaction of the dark-matter particles with the initial state quarks or the \( Z/\gamma' \) intermediate state. This is done using the relation \( M_{\chi}^{\text{limit}} = M_{\chi}^{\text{generator}} \times (\sigma^{\text{generator}}/\sigma^{\text{limit}})^{1/2} \), where the superscript indicates whether the parameter is a measured limit or calculated using MC simulation and \( p \) indicates the power of \((1/M_{\chi})\) appearing in the EFT Lagrangian. These limits are shown in Fig. 4. They are also translated into limits on the \( \chi \)-nucleon scattering cross section using the method in Ref. [10] for several effective operators mediating the interaction of the dark-matter particles with the initial state quarks or the \( Z/\gamma' \) intermediate state.

A lower limit on the coupling, \( f \), of the scalar-mediator model is also calculated based on the WIMP relic abundance in Ref. [47] and the expression for the freeze-out temperature from Ref. [48]. If the relic abundance lower limit calculated at some mass point \((m_{\psi}, m_{\chi})\) is greater than the upper limit measured in this analysis, that mass point is excluded. Limits on the cross section times branching ratio in the scalar-mediator model are shown in Fig. 7, and limits on \( f \) as a function of mediator mass \( m_{\psi} \) and \( m_{\chi} \), as well as the exclusion region, are shown in Fig. 8.

Fiducial cross-section limits are calculated in each signal region to complement the limits on specific models. The
XENON100 [43], CDMS [44,45], and LUX [46]. These limits are characterized by selection requirements at particle level identical to all the requirements on the reconstructed dilepton + $E_T^{miss}$ system, where the $E_T^{miss}$ is calculated summing over all neutrinos and dark-matter particles. The acceptance is the ratio of the number of generated events within the fiducial region to the total number of generated events. In addition, the generated leptons are required to be separated by at least $\Delta R = 0.2$ to match the isolation requirement. The reconstruction efficiency ranges from (56.9 ± 0.9)% for $ZZ\chi\chi$-max. $\gamma^*$ at $m_\chi = 1000$ GeV to (77.9 ± 3.1)% for $D5$ at $m_\chi = 400$ GeV. The lowest

FIG. 7 (color online). Observed 95% C.L. upper limits on the cross section multiplied by the branching ratio of $Z \rightarrow \ell^+\ell^-$ of the scalar-mediator theory as a function of $m_\chi$. The observed cross-section limit for a scalar-mediator mass, $m_\eta$, of 1000 GeV is shown. Production cross sections predicted from theory are shown for $m_\eta = 1$ TeV and for different values of the coupling strength $f$.

FIG. 8. Observed 95% C.L. upper limits on the coupling constant $f$ of the scalar-mediator theory as a function of $m_\chi$ and the mediator mass $m_\eta$. The cross-hatching shows the theoretically accessible region outside the range covered by this analysis. The white region is phase space beyond the model’s validity. In the excluded region in the upper left-hand corner, demarcated by the black line, the lower limit on $f$ from the relic abundance calculations based on [47,48] is greater than the upper limit measured in this analysis.

FIG. 5 (color online). Observed 90% C.L. upper limits on the $\chi$-nucleon scattering cross section as a function of $m_\chi$ for the spin-dependent $D9$ effective operators mediating the interaction of the dark-matter particles with the $q\bar{q}$ initial state. The limits are compared with results from the published ATLAS hadronically decaying $W/Z$ [8] and $j + \chi\chi$ [4] searches, COUPP [38], SIMPLE [39], PICASSO [40], and IceCube [41]. These limits are shown as they are given in the corresponding publications and are only shown for comparison with the results from this analysis, since they are obtained assuming the interactions are mediated by operators different from those used for the ATLAS limits.

FIG. 6 (color online). Observed 90% C.L. upper limits on the $\chi$-nucleon scattering cross section as a function of $m_\chi$ for spin-independent effective operators mediating the interaction of the dark-matter particles with the $q\bar{q}$ initial state. The limits are compared with results from the published ATLAS hadronically decaying $W/Z$ [8] and $j + \chi\chi$ [4] searches, COUPP [38], SIMPLE [39], PICASSO [40], and IceCube [41]. These limits are shown as they are given in the corresponding publications and are only shown for comparison with the results from this analysis, since they are obtained assuming the interactions are mediated by operators different from those used for the ATLAS limits.

FIG. 8. Observed 95% C.L. upper limits on the coupling constant $f$ of the scalar-mediator theory as a function of $m_\chi$ and the mediator mass $m_\eta$. The cross-hatching shows the theoretically accessible region outside the range covered by this analysis. The white region is phase space beyond the model’s validity. In the excluded region in the upper left-hand corner, demarcated by the black line, the lower limit on $f$ from the relic abundance calculations based on [47,48] is greater than the upper limit measured in this analysis.

FIG. 7 (color online). Observed 95% C.L. upper limits on the cross section multiplied by the branching ratio of $Z \rightarrow \ell^+\ell^-$ of the scalar-mediator theory as a function of $m_\chi$. The observed cross-section limit for a scalar-mediator mass, $m_\eta$, of 1000 GeV is shown. Production cross sections predicted from theory are shown for $m_\eta = 1$ TeV and for different values of the coupling strength $f$. 

FIG. 8. Observed 95% C.L. upper limits on the coupling constant $f$ of the scalar-mediator theory as a function of $m_\chi$ and the mediator mass $m_\eta$. The cross-hatching shows the theoretically accessible region outside the range covered by this analysis. The white region is phase space beyond the model’s validity. In the excluded region in the upper left-hand corner, demarcated by the black line, the lower limit on $f$ from the relic abundance calculations based on [47,48] is greater than the upper limit measured in this analysis.
value of the reconstruction efficiency is used to calculate the fiducial cross-section limits in order to be conservative. The corresponding acceptances for the previous operators are (30.3 ± 0.5)% and (2.6 ± 0.2)%), respectively, where the uncertainties are purely statistical and the variation in the acceptance arises primarily from the different $E_T^{\text{miss}}$ spectra of the operators. The observed and expected upper limits on the fiducial cross section are given in Table IV.

In conclusion, a search for the production of dark-matter particles in association with a Z boson that decays leptonically in 20.3 fb$^{-1}$ of $pp$ collisions at $\sqrt{s} = 8$ TeV is presented for three EFT operators where the dark matter interacts directly with quarks: D1, D5, and D9. The new limits complement the limits reported in other LHC analyses. The results are also interpreted using EFT models where the dark matter interacts directly with pairs of electroweak bosons. Initial limits are set on the mass scale of the $ZZ\chi\chi$ EFT operators describing the interaction between dark matter and a Z or $\gamma^*$ intermediate state. Upper limits are also set on the scattering cross section of dark-matter particles with nucleons for effective operators mediating the interaction of dark-matter particles with a $q\bar{q}$ initial state, and on a model in which the interaction between the dark matter and $Z/\gamma^*$ is mediated by a scalar particle.

<table>
<thead>
<tr>
<th>$E_T^{\text{miss}}$ threshold [GeV]</th>
<th>Fiducial cross section [fb]</th>
</tr>
</thead>
<tbody>
<tr>
<td>150</td>
<td>3.0 0.73 0.36 0.27</td>
</tr>
<tr>
<td>250</td>
<td>2.7 0.57 0.27 0.26</td>
</tr>
<tr>
<td>350</td>
<td></td>
</tr>
<tr>
<td>450</td>
<td></td>
</tr>
</tbody>
</table>

TABLE IV. The observed and expected upper limits on the fiducial cross section at 95% C.L. for each signal region.

ACKNOWLEDGMENTS

We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently. We acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWF and FWF, Austria; ANAS, Azerbaijan; SSTC, Belarus; CNPq and FAPERJ, Brazil; NSERC, NRC and CFI, Canada; CERN; CONICYT, Chile; CAS, MOST and NSFC, China; COLCIENCIAS, Colombia; MSMT CR, MPO CR and VSC CR, Czech Republic; DNRF, DNSRC and Lundbeck Foundation, Denmark; EPLANET, ERC and NSRF, European Union; IN2P3-CNRS, CEA-DSM/IRFU, France; GNSF, Georgia; BMBF, DFG, HGF, MPG and AvH Foundation, Germany; GSRT and NSRF, Greece; ISF, MINERVA, GIF, I-CORE and Benoziyo Center, Israel; INFN, Italy; MEXT and JSPS, Japan; CNRST, Morocco; FOM and NWO, Netherlands; BRF and RCN, Norway; MNiSW and NCN, Poland; GRICES and FCT, Portugal; MNE/IFA, Romania; MES of Russia and ROSATOM, Russian Federation; JINR; MSTD, Serbia; MSSR, Slovakia; ARRS and MIZŠ, Slovenia; DST/NRF, South Africa; MINECO, Spain; SRC and Wallenberg Foundation, Sweden; SER, SNSF and Cantons of Bern and Geneva, Switzerland; NSE, Taiwan; TAEK, Turkey; STFC, the Royal Society and Leverhulme Trust, United Kingdom; DOE and NSF, United States of America. The crucial computing support from all WLCG partners is acknowledged gratefully, in particular from CERN and the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CC-IN2P3 (France), KIT/GridKA (Germany), INFN-CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (United Kingdom) and BNL (USA) and in the Tier-2 facilities worldwide.

[19] ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the center of the
the detector and the $z$ axis along the beam pipe. The $x$ axis points from the IP to the center of the LHC ring, and the $y$ axis points upward. Cylindrical coordinates $(r, \phi)$ are used in the transverse plane, $\phi$ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle $\theta$ as $\eta = -\ln \tan(\theta/2)$.

SEARCH FOR DARK MATTER IN EVENTS WITH A Z ...

(ATLAS Collaboration)
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PHYSICAL REVIEW D 90, 012004 (2014)

21 Physikalisches Institut, University of Bonn, Bonn, Germany
22 Department of Physics, Boston University, Boston, Massachusetts, United States of America
23 Department of Physics, Brandeis University, Waltham, Massachusetts, United States of America
24a Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil
24b Federal University of Juiz de Fora (UFJF), Juiz de Fora, Brazil
24c Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei, Brazil
24d Instituto de Fisica, Universidade de Sao Paulo, Sao Paulo, Brazil
25 Physics Department, Brookhaven National Laboratory, Upton New York, United States of America
26a National Institute of Physics and Nuclear Engineering, Bucharest, Romania
26b National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj Napoca, Romania
26c University Politehnica Bucharest, Bucharest, Romania
26d West University in Timisoara, Timisoara, Romania
27 Departamento de Fisica, Universidad de Buenos Aires, Buenos Aires, Argentina
28 Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
29 Department of Physics, Carleton University, Ottawa, Ontario, Canada
30 CERN, Geneva, Switzerland
31 Enrico Fermi Institute, University of Chicago, Chicago, Illinois, United States of America
32a Departamento de Fisica, Pontificia Universidad Católica de Chile, Santiago, Chile
32b Departamento de Fisica, Universidad Técnica Federico Santa Maria, Valparaíso, Chile
33a Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China
33b Department of Modern Physics, University of Science and Technology of China, Anhui, China
33c Department of Physics, Nanjing University, Jiangsu, China
33d School of Physics, Shandong University, Shandong, China
33e Physics Department, Shanghai Jiao Tong University, Shanghai, China
34 Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France
35 Nevis Laboratory, Columbia University, Irvington, New York, United States of America
36 Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark
36a INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Italy
36b Dipartimento di Fisica, Università della Calabria, Rende, Italy
37a AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland
37b Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland
39 The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland
40 Physics Department, Southern Methodist University, Dallas, Texas, United States of America
41 Physics Department, University of Texas at Dallas, Richardson, Texas, United States of America
42 DESY, Hamburg and Zeuthen, Germany
43 Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
44 Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
45 Department of Physics, Duke University, Durham, North Carolina, United States of America
46 SUPA–School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
46a INFN Laboratori Nazionali di Frascati, Frascati, Italy
48 Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany
49 Section de Physique, Université de Genève, Geneva, Switzerland
50a INFN Sezione di Genova, Italy
50b Dipartimento di Fisica, Università di Genova, Genova, Italy
51a E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi, Georgia
51b High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
52a II Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
52b II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany
52c Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
53 SUPA–School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
54 Laboratoire de Physique Subatomique et de Cosmologie, Université Joseph Fourier and CNRS/IN2P3 and Institut National Polytechnique de Grenoble, Grenoble, France
55 Department of Physics, Hampton University, Hampton, Virginia, United States of America
56 Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, Massachusetts, United States of America
57 Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
58 Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
58c ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany

012004-17