

# Tunable long range forces mediated by self-propelled colloidal hard spheres

## Supplementary Information

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In the 2D system of active hard spheres, the interaction between particle  $i$  and  $j$  is modeled with a steep Weeks-Chandler-Andersen (WCA) potential given by

$$\frac{U_{WCA}(r_{ij})}{k_B T} = \begin{cases} 4\epsilon \left[ \left( \frac{\sigma}{r_{ij}} \right)^{12} - \left( \frac{\sigma}{r_{ij}} \right)^6 + \frac{1}{4} \right] & r_{ij}/\sigma \leq 2^{1/6}, \\ 0 & r_{ij}/\sigma > 2^{1/6}, \end{cases} \quad (1)$$

where  $r_{ij}$  is the center-to-center distance between the two spheres, and  $\sigma$  is the diameter of the particles, with  $k_B$  and  $T$  the Boltzmann constant and the temperature of the system, respectively. Here, the strength of the interaction is set  $\epsilon = 40$ . Moreover, the interaction between particle  $i$  and wall  $j$  is given by  $U_{wall}(i, j) = U_{WCA}(r_{ij})$ , in which  $r_{ij}$  is minimal distance between the center of particle  $i$  and points on the line segment  $j$ .

In the quasi-2D system of active hard spheres confined between two horizontal walls, the interaction between particle  $i$  at  $\mathbf{r}_i = (x_i, y_i, z_i)$  and the two horizontal walls is

$$\frac{U_{conf}(\mathbf{r}_i)}{k_B T} = \begin{cases} 0 & 0.5\sigma \leq z_i \leq H - 0.5\sigma, \\ \infty & \text{otherwise,} \end{cases} \quad (2)$$

where  $H$  is the height of the confinement.

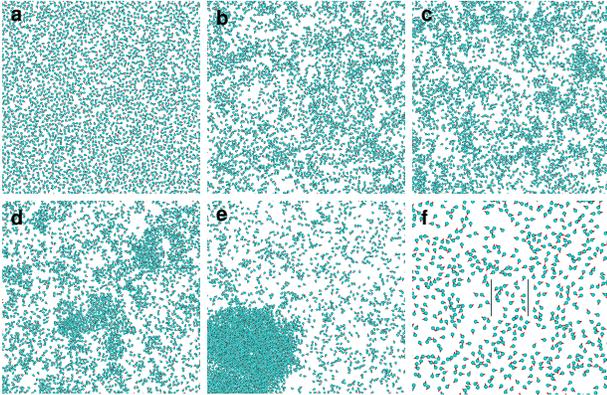


FIG. S1: (a-e) Typical snapshots of 2D systems of self-propelled colloidal hard spheres at density  $\rho_{bulk}\sigma^2 = 0.4$  with self-propulsion  $f\sigma/k_B T = 0, 20, 40, 60$  and  $80$ , respectively. (f) Illustration of the system to study the interaction mediated by the dynamic clusters of self-propelled colloidal hard spheres, and the two vertical lines indicate the two parallel hard walls. The red arrows indicate the direction of random self-propulsion.

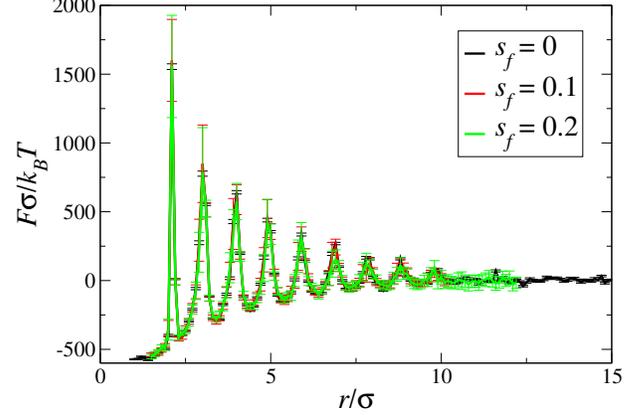


FIG. S2: Effective interaction between two parallel walls  $F\sigma/k_B T$  as a function of wall-wall distance  $r$  for systems of active hard spheres with different polydisperse active propulsions  $s_f$ . The density of the system is fixed at  $\rho_{bulk}\sigma^2 = 0.4$ , and the average self-propulsion is  $f\sigma/k_B T = 40$  with  $\sigma$  the particle diameter.

The effect of polydispersity of active force on the dynamic wetting force is shown in Fig. S2. The self-propulsion on particle  $i$  is given by

$$f_i = f \cdot [1 + \xi(s_f)], \quad (3)$$

where  $f$  is average active force in the system, and  $\xi(s_f)$  is a Gaussian random number with zero mean and the standard deviation  $s_f$ .

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