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Classical simulation of entanglement swapping with bounded communication
– Supplementary material –

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In this Supplementary material we prove that Protocols 1 and 2 generate the desired correlation
(equation (1) of the main text), or in other terms, \( P(a = b|\mathbf{x}, \mathbf{y}) = \frac{1-x}{2} \).

Correlation from Protocol 1

We first note that the definitions of \( \alpha = \text{sign}(\sin \angle A) \) and of \( \phi_A' = (\phi_A - j\frac{\pi}{4} \mod \pi) \in [0, \frac{\pi}{4}] \) on Alice’s side, and
the definitions of \( \beta = \text{sign}[\sin(\angle B - j\frac{\pi}{4})] \) (and then of \( b = \pm \beta \)) and of \( \phi_B' = (\phi_B - j\frac{\pi}{4} \mod \pi) \in [0, \pi] \) on Bob’s
side, ensure that the following relations hold, as required:

\[
\begin{align*}
P(a = b|\phi_A, \phi_B) &= P(a \neq b|\phi_A, \phi_B + \pi) = P(a \neq b|\phi_A + \pi, \phi_B) = P(a = b|\phi_A + \pi, \phi_B + \pi) \\
&= P(a = b|\phi_A + j\frac{\pi}{4}, \phi_B + j\frac{\pi}{4}) \quad \text{for any } j \in \mathbb{Z}.
\end{align*}
\]

It is therefore sufficient to check that the correct correlation is obtained for \( \phi_A \in [0, \frac{\pi}{4}] \) and \( \phi_B \in [0, \pi] \).

For such values of \( \phi_A, \phi_B \) (for which \( \phi_A' = \phi_A \) and \( \phi_B' = \phi_B \)), the probability \( P(a = b|\phi_A, \phi_B) \) obtained from
Protocol 1 can be calculated as follows:

\[
P(a = b|\phi_A < \phi_B - j\frac{\pi}{4}) = \frac{16}{\pi} \left( \int_{0}^{\phi_A} d\lambda_{AR} \int_{0}^{\lambda_{AR}} d\lambda_{RB} \psi_{010}^{j\theta}(\phi_B - \lambda_{RB}) \\
+ \int_{0}^{\phi_A} d\lambda_{AR} \int_{\frac{\phi_A - j\frac{\pi}{4}}{4}}^{\lambda_{AR}} d\lambda_{RB} \psi_{011}^{j\theta}(\phi_B - \lambda_{RB}) + \int_{0}^{\phi_A} d\lambda_{AR} \int_{\frac{\phi_A - j\frac{\pi}{4}}{4}}^{\lambda_{AR}} d\lambda_{RB} \psi_{010}^{j\theta}(\phi_B - \lambda_{RB}) \\
+ \int_{\frac{\phi_A - j\frac{\pi}{4}}{4}}^{\phi_A} d\lambda_{AR} \int_{0}^{\lambda_{AR}} d\lambda_{RB} \psi_{011}^{j\theta}(\phi_B - \lambda_{RB}) + \int_{\frac{\phi_A - j\frac{\pi}{4}}{4}}^{\phi_A} d\lambda_{AR} \int_{\frac{\phi_A - j\frac{\pi}{4}}{4}}^{\lambda_{AR}} d\lambda_{RB} \psi_{011}^{j\theta}(\phi_B - \lambda_{RB}) \\
+ \int_{\frac{\phi_A - j\frac{\pi}{4}}{4}}^{\phi_A} d\lambda_{AR} \int_{\frac{\phi_A - j\frac{\pi}{4}}{4}}^{\lambda_{AR}} d\lambda_{RB} \psi_{110}^{j\theta}(\phi_B - \lambda_{RB}) + \int_{\frac{\phi_A - j\frac{\pi}{4}}{4}}^{\phi_A} d\lambda_{AR} \int_{\frac{\phi_A - j\frac{\pi}{4}}{4}}^{\lambda_{AR}} d\lambda_{RB} \psi_{111}^{j\theta}(\phi_B - \lambda_{RB}) \\
+ \int_{\frac{\phi_A - j\frac{\pi}{4}}{4}}^{\phi_A} d\lambda_{AR} \int_{\frac{\phi_A - j\frac{\pi}{4}}{4}}^{\lambda_{AR}} d\lambda_{RB} \psi_{100}^{j\theta}(\phi_B - \lambda_{RB}) + \int_{\frac{\phi_A - j\frac{\pi}{4}}{4}}^{\phi_A} d\lambda_{AR} \int_{\frac{\phi_A - j\frac{\pi}{4}}{4}}^{\lambda_{AR}} d\lambda_{RB} \psi_{101}^{j\theta}(\phi_B - \lambda_{RB}) \right).
\]
and
\[
P(a = b | \phi_A \geq \phi_B - j_B \frac{\pi}{4}) = \frac{16}{\pi^2} \left( \int_{0}^{\phi_B - j_B \frac{\pi}{4}} d\lambda_{AR} \int_{0}^{\lambda_{AR}} \int_{\phi_B - \lambda_{RB}}^{\phi_B} \psi_{001}^{j_B} (\phi_B - \lambda_{RB}) d\lambda_{RB} \right)
\]
\[
+ \int_{\phi_B - j_B \frac{\pi}{4}}^{\phi_B} d\lambda_{AR} \int_{\lambda_{AR}}^{\phi_B} d\lambda_{RB} \psi_{001}^{j_B} (\phi_B - \lambda_{RB}) + \int_{0}^{\phi_B - j_B \frac{\pi}{4}} d\lambda_{AR} \int_{\phi_B - j_B \frac{\pi}{4}}^{\phi_B} d\lambda_{RB} \psi_{001}^{j_B} (\phi_B - \lambda_{RB})
\]
\[
+ \int_{\phi_B - j_B \frac{\pi}{4}}^{\phi_B} d\lambda_{AR} \int_{\lambda_{AR}}^{\phi_B} d\lambda_{RB} \psi_{001}^{j_B} (\phi_B - \lambda_{RB}) + \int_{\phi_B}^{\phi_B - j_B \frac{\pi}{4}} d\lambda_{AR} \int_{\phi_B - j_B \frac{\pi}{4}}^{\phi_B} d\lambda_{RB} \psi_{001}^{j_B} (\phi_B - \lambda_{RB})
\]
\[
+ \int_{\phi_B}^{\phi_B - j_B \frac{\pi}{4}} d\lambda_{AR} \int_{\lambda_{AR}}^{\phi_B} d\lambda_{RB} \psi_{001}^{j_B} (\phi_B - \lambda_{RB}) + \int_{\phi_B}^{\phi_B - j_B \frac{\pi}{4}} d\lambda_{AR} \int_{\lambda_{AR}}^{\phi_B} d\lambda_{RB} \psi_{001}^{j_B} (\phi_B - \lambda_{RB})
\]

One can then check that with the choice of functions \( \psi_{001}^{j_B} \) indicated in Table I (see main text), this leads (for all values of \( j_B \)) to
\[
P(a = b | \phi_A, \phi_B) = \frac{1 - \cos(\phi_A - \phi_B)}{2},
\]
as desired.

**Correlation from Protocol 2**

After running Protocol 2 for inputs \( x = (\sin \theta_A \cos \phi_A, \sin \theta_A \sin \phi_A, \cos \theta_A) \) and \( y = (\sin \theta_B \cos \phi_B, \sin \theta_B \sin \phi_B, \cos \theta_B) \), the probability that Alice and Bob’s outputs are the same is
\[
P(a = b | x, y) = \sum_{a_0, b_0 = \pm 1} P(a_0, b_0 | \phi_A, \phi_B) P(a = b | \theta_A, \theta_B, a_0, b_0) = \sum_{a_0, b_0 = \pm 1} P(a_0, b_0 | \phi_A, \phi_B) \frac{1 - \cos(a_0 \theta_A + b_0 \theta_B)}{2}
\]
\[
= P(a_0 = b_0 | \phi_A, \phi_B) \frac{1 - \cos(\theta_A + \theta_B)}{2} + P(a_0 \neq b_0 | \phi_A, \phi_B) \frac{1 - \cos(\theta_A - \theta_B)}{2}
\]
\[
= \frac{1}{2} \cos(\theta_A - \theta_B) \frac{1 - \cos(\theta_A + \theta_B)}{2} + \frac{1}{2} \cos(\theta_A + \theta_B) \frac{1 - \cos(\theta_A - \theta_B)}{2}
\]
\[
= \frac{1}{2} \cos(\theta_A \cos \theta_B - \sin \theta_A \sin \theta_B \cos(\phi_A - \phi_B)) \frac{1 - \cos(\theta_A - \theta_B)}{2} = \frac{1 - x \cdot y}{2}.
\]

Protocol 2 thus reproduces the desired entanglement swapping correlation (equation (1) of the main text).