Search for new phenomena in the dijet mass distribution using pp collision data at $\sqrt{s} = 8$ TeV with the ATLAS detector


DOI
10.1103/PhysRevD.91.052007

Publication date
2015

Document Version
Final published version

Published in
Physical Review D. Particles and Fields

Citation for published version (APA):

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I. INTRODUCTION

This paper describes the search for new phenomena (NP) in the two-jet (dijet) invariant mass spectrum in the full 2012 data set delivered by the Large Hadron Collider (LHC) at CERN, and collected with the ATLAS detector. The studies reported here select events containing two or more jets. The two highest-

p_T (“leading” and “subleading”) jets are combined to determine the dijet invariant mass, \( m_{jj} \). High-transverse-momentum (high-

p_T) dijet events are produced copiously by QCD processes, and can reach the highest mass scales accessible in LHC \( pp \) collisions. The QCD processes, along with a subpercent admixture of additional Standard Model (SM) processes, create a smooth rapidly falling spectrum in \( m_{jj} \). Many NP models describe new particles or excitations created as s-channel resonances with appreciable branching ratios to final states involving quarks and gluons (\( q \) and \( g \)), and can produce dijet final states. If the resonance width is sufficiently narrow, these NP signals would appear as local excesses (bumps) in the dijet mass spectrum over the smooth SM background.

Studies searching for excesses in dijet mass spectra have been performed in all previous collider experiments, including CDF and D0 at the Tevatron [1–4]. At the LHC, the CMS and ATLAS experiments have continued this program, starting from the first 2010 data [5–14]. The most recent published ATLAS results used the 2011 data set of 4.8 fb\(^{-1}\) at a center-of-mass energy of \( \sqrt{s} = 7 \) TeV [13]. In 2012, the LHC delivered an integrated luminosity of 20.3 fb\(^{-1}\) [15] in proton-proton (\( pp \)) collisions, roughly a factor of 4 more than used for previous dijet studies. In addition, the increase in center-of-mass energy to \( \sqrt{s} = 8 \) TeV in 2012 increased the sensitivity in searches for new phenomena, and pushed exclusion limits to higher masses and energy scales. This increased kinematic reach, combined with a new online event selection strategy employed in the present analysis, provides the largest dijet invariant mass range coverage to date, from 250 GeV to 4.5 TeV.

No significant excess is observed above the background. A number of NP models are compared to data to derive limits. The models chosen span a range of characteristic masses (or energy scales) and cross sections, and are complementary in terms of the flavor of their final-state partons. The benchmark models under consideration include excited quarks (\( q^* \)) decaying to \( gg \) [16,17], color-octet scalars (\( s8 \)) decaying to \( gg \) [18–21], heavy W’ gauge bosons decaying to \( q\bar{q}^{*} \) [22–29], two forms of chiral W’ gauge bosons [30–33] and quantum black holes [34–37] decaying to a mixture of quarks and gluons. The current dijet search also sets limits on new generic resonances whose intrinsic width is convolved with effects due to parton distribution functions (PDFs), parton shower, nonperturbative effects and detector resolution.

II. THE ATLAS DETECTOR

A detailed description of the ATLAS detector has been published elsewhere [38]. The detector is instrumented over almost the entire solid angle around the \( pp \) collision point with layers of tracking detectors, calorimeters, and muon chambers.

In ATLAS, high-

p_T hadronic jets are measured using a finely segmented calorimeter system, designed to achieve high reconstruction efficiency and excellent energy resolution. Electromagnetic calorimetry is provided by high-granularity liquid-argon (LAr) sampling calorimeters,
using lead as an absorber. The calorimeters are split into a barrel region ($|\eta| < 1.475$) and two end-cap ($1.375 < |\eta| < 3.2$) regions. The hadronic calorimeter is divided into barrel, extended barrel ($|\eta| < 1.7$) and end-cap ($1.5 < |\eta| < 3.2$) regions. The barrel and extended barrel are instrumented with scintillator tiles and steel absorbers, while the end-caps use copper with LAr modules. The forward calorimeter ($3.1 < |\eta| < 4.9$) is instrumented with modules using LAr as the active medium and copper or tungsten as absorbers to provide electromagnetic and hadronic energy measurements, respectively.

### III. JET AND EVENT SELECTION

Jets are reconstructed from contiguous groups of calorimeter cells that contain significant energy above noise (topological clusters) [39]. The anti-$k_t$ jet algorithm [40, 41] is applied to these clusters using a distance parameter of $R = 0.6$. Effects from additional $pp$ interactions in the same and neighboring bunch crossings are corrected for using the calibration procedure described in Ref. [42]. Simulated QCD multijet events are used for the derivation of the jet calibration to the hadronic scale. They are produced with the event generator PYTHIA [43] 8.160, using the CT10 PDF [44] and the AU2 set of parameters to describe the nonperturbative effects tuned to ATLAS data (the AU2 tune) [45]. Detector effects are simulated using GEANT4 within the ATLAS software infrastructure [46, 47]. Additional simulated minimum-bias events are overlaid onto the hard scattering, both within the same bunch crossings and within trains of consecutive bunches. The same software used to reconstruct data is used for the Monte Carlo (MC) samples. The level of agreement between data and MC simulation is further improved by the application of calibration constants obtained with in situ techniques based on momentum balancing between central and forward jets, between photons or $Z$ bosons and jets, and between high-momentum jets and a recoil system of low-momentum jets [48].

The energy scale of central jets relevant to this search is known to within 4% [48]. The jet energy resolution is estimated both in data and in simulation using transverse momentum balance studies in dijet events [49], and are found to be in agreement within uncertainties. The dijet mass resolution is approximately 8% at 200 GeV, and improves to less than 4% above 2 TeV. The measured dijet mass distribution is not corrected for detector resolution effects.

The data were collected using single-jet triggers [50]. These triggers are designed to select events that have at least one large transverse energy deposition in the calorimeter. To match the data rate to the processing and storage capacity available to ATLAS, the triggers with low-$p_T$ thresholds are prescaled; only a preselected fraction of all events passing the threshold is recorded. Combinations of prescaled single-jet triggers are used to reach lower dijet invariant masses. The highest prescale for the trigger combinations used in this analysis is $1/460000$. This trigger combination is used for jets with $p_T$ between 59 and 99 GeV.

For a given leading-jet $p_T$, a predetermined combination of triggers with efficiencies exceeding 99.5% is used to select the event. Each event is weighted according to the average effective integrated luminosity recorded by the given trigger combination [51].

During the 2012 data taking, ATLAS recorded data at a rate that was higher than the rate of the offline reconstruction: 400 Hz of recorded data were promptly reconstructed, while 100 Hz of data from hadronic triggers were recorded and reconstructed later (the “delayed stream”). Dijet events from the delayed data stream fall primarily into the region between 750 GeV and 1 TeV. They are used to further increase the size of the analysis data set as follows. First, two independent data sets are built from the delayed and normal trigger streams. The $m_{jj}$ distributions from the two data sets have checked to be in agreement with a shape-only Kolmogorov-Smirnov test, leading to a probability of 86% for their compatibility. The two dijet mass distributions measured from these data sets are then averaged, using the effective integrated luminosity as a weight. The delayed stream increases the luminosity recorded in this region of phase space by up to an order of magnitude. The effective integrated luminosity attained in this analysis (using the 2012 normal stream with the added delayed stream) is compared to previous ATLAS analyses in Fig. 1.

Events used for the search are required to have at least one collision vertex defined by two or more charged-particle tracks. In the presence of multiple $pp$ interactions, the collision vertex with the largest scalar sum of $p_T$ of its associated tracks is chosen as the primary vertex.

Events are rejected if either the leading jets, or any of the other jets with $p_T$ greater than 30% of the $p_T$ of the subleading jet, are poorly measured or have a topology characteristic of noncollision background or calorimeter noise [52]. Poorly measured jets correspond to energy depositions in regions where the energy measurement is known to be inaccurate. Events are also rejected if one of the jets relevant to this analysis falls into regions of the calorimeter that were nonoperational during data taking. An inefficiency of roughly 10% due to this veto is emulated in MC signal samples following the same conditions as data. Central values and statistical errors of the dijet mass
spectra of both the data and MC signal samples are scaled, in order to correct for this inefficiency.

Additional kinematic selection criteria are used to enrich the dijet sample with events in the hard-scatter region of phase space. The rapidity \( y \) of the two leading jets must be within \( |y| < 2.8 \). The leading and subleading jets are required to have \( p_T > 50 \text{ GeV} \), ensuring a jet reconstruction efficiency of 100% [53] both for QCD background and for all benchmark models under consideration. Events must satisfy \( |y^*| = \frac{1}{2} |y_{\text{lead}} - y_{\text{sublead}}| < 0.6 \) and \( m_{jj} > 250 \text{ GeV} \). The invariant mass cut of \( m_{jj} > 250 \text{ GeV} \) is chosen such that the dijet mass spectrum is unbiased by the kinematic selection on \( p_T \).

IV. COMPARISON OF THE DIJET MASS SPECTRUM TO A SMOOTH BACKGROUND

The observed dijet mass distribution in data, after all selection requirements, is shown in Fig. 2. The bin width varies with mass and is chosen to approximately equal the dijet mass resolution derived from simulation of QCD processes. The predictions for an excited quark \( q^* \) with three different mass hypotheses are also shown.

The search for resonances in \( m_{jj} \) uses a data-driven background estimate derived by fitting a smooth functional form to the spectrum. An important feature of this functional form is that it allows for smooth background variations, but does not accommodate localized excesses that could indicate the presence of NP signals. In previous studies, ATLAS and other experiments [54] have found that the following function provides a satisfactory fit to the QCD prediction of dijet production:

\[
f(x) = p_1 (1 - x)^{p_2 x^{p_3} + p_4 \ln x},
\]

where the \( p_i \) are fit parameters, and \( x \equiv m_{jj}/\sqrt{s} \). The uncertainty associated with the stability of the fit is carried forward as a nuisance parameter in the statistical analysis.

The functional form was selected using a data set consisting of a quarter of the full data, a quantity known to be insensitive to resonant new physics at dijet masses above 1.5 TeV after the previous public result on 13 fb\(^{-1}\) of data [55]. A range of parametrizations were tested on the blinded data set using a k-fold cross-validation and there was found to be no substantial difference between the standard function of Eq. (1) and higher-order parametrizations, so the function with a simpler form and a published precedent was selected. The \( \chi^2 \) value of the fit to the blinded data set was 37 for 56 degrees of freedom using the parametrization of Eq. (1). The fit function showed good agreement to both the fully simulated dijet mass spectrum obtained from the simulated PYTHIA 8.160 QCD multijet events mentioned in Sec. III, corrected for next-to-leading-order effects using the NLOJET ++v4.1.3 program [56,57] as described in Ref. [11], and from a large-statistics sample of generator-level events, for which the \( \chi^2 \) of the fit was 58 for 55 degrees of freedom. While the number of data events is matched or surpassed by the number of fully simulated events starting from dijet masses of roughly 2 TeV, the generator-level statistics is sufficient to reproduce that of data. The \( \chi^2 \) value of the fit to data shown in Fig. 2 is 79 for 56 degrees of freedom.
The center panel of Fig. 2 shows the relative difference between the data and the background fit, and overlays the shapes that would be expected in the presence of three sample $q^+$ signals. The bottom panel of Fig. 2 shows the significance of the difference between the data and the fit in each bin. The significance is calculated taking only statistical uncertainties into account, and assuming that the data follow a Poisson distribution with the expected value given by the fit function.

For each bin a $p$-value is determined by assessing the probability of a background fluctuation leading to a number of events higher than or equal to the observed excess, or lower than or equal to the observed deficit. This $p$-value is converted to a significance in terms of an equivalent number of standard deviations (the $z$-value) [58]. Where there is an excess (deficit) in data in a given bin, the significance is plotted as positive (negative). To test the degree of consistency between the data and the fitted background, the $p$-value of the fit is determined by calculating the $\chi^2$-value from the data and comparing this result to the $\chi^2$ distribution obtained from pseudoexperiments drawn from the background fit, as described in a previous publication [11]. The resulting $p$-value is 0.027.

The BUMPHUNTER algorithm [59,60] is used to establish the presence or absence of a narrow resonance in the dijet mass spectrum, as described in greater detail in previous publications [11,12]. Starting with a two-bin window, the algorithm increases the signal window and shifts its location until all possible bin ranges, up to the widest window corresponding to half the mass range spanned by the data, are tested. The most significant excess of data above the smooth spectrum (“bump”) is defined by the set of bins over which the integrated excess of data over the fit prediction has the smallest probability of arising from a background fluctuation, assuming Poisson statistics.

The BUMPHUNTER algorithm accounts for the so-called “look-elsewhere effect” [61], by performing a series of pseudoexperiments drawn from the background estimate to determine the probability that random fluctuations in the background-only hypothesis would create an excess anywhere in the spectrum at least as significant as the one observed. Furthermore, to prevent any NP signal from biasing the background estimate, if the most significant local excess from the background fit has a $p$-value smaller than 0.01, the corresponding region is excluded and a new background fit is performed. The exclusion is then progressively widened bin by bin until the $p$-value of the remaining fitted region is acceptable. No such exclusion is needed for the current data set.

The most significant discrepancy identified by the BUMPHUNTER algorithm in the observed dijet mass distribution in Fig. 2 is a seven-bin excess in the interval 390–599 GeV. The probability of observing an excess at least as large somewhere in the mass spectrum for a background-only hypothesis is 0.075, corresponding to a $z$-value of 1.44σ. To conclude, this test shows no evidence for a resonant signal in the observed $m_{jj}$ spectrum.

V. SIMULATION OF HYPOTHETICAL NEW PHENOMENA

In the absence of any significant signals indicating the presence of phenomena beyond the SM, Bayesian 95% credibility level (C.L.) limits are determined for a number of NP hypotheses that would produce localized excesses.

Samples for NP models are produced by a variety of event generators. The partons originating from the initial $2 \rightarrow 2$ matrix elements are passed to $\text{PYTHIA}$ 8.160 with the AU2 tune [45]. $\text{PYTHIA}$ uses $p_T$-ordered parton showers to model additional radiation in the leading-logarithm approximation [62]. Multiple parton interactions [63], as well as fragmentation and hadronization based on the Lund string model [64], are also simulated. Renormalization and factorization scales for the NP models are set to the mean $p_T$ of the two leading jets.

Excited $u$ and $d$ quarks ($q^*$), one possible manifestation of quark compositeness, are simulated in all decay modes using the $\text{PYTHIA}$ 8.162 generator, using the CT10 PDF. Excited quarks are assumed to decay to quarks via gauge couplings set to unity, leading to a $qq$ final state approximately 83% of the time (the remaining generated decays involve $W/Z$ or $\gamma$ emission). The acceptance, defined as the fraction of generated events passing all reconstruction steps and the analysis selection described in Sec. III using jets reconstructed from stable particles excluding muons and neutrinos, is approximately 58%. The largest reduction in acceptance arises from the rapidity selection criteria.

The color-octet scalar model describes the production of exotic colored resonances ($s_8$). $\text{MADGRAPH}$ 5 (v1.5.5) [65] with the MSTW2008LO PDF [66] is employed to generate parton-level events at leading-order approximation. Parton showering and nonperturbative effects are simulated with $\text{PYTHIA}$ 8.170. Color-octet scalars can decay into two gluons and can then have a broader mass distribution and larger tails than resonances decaying to quarks. For resonances produced by this model, the acceptance ranges from 61% to 63%.

The production of heavy charged gauge bosons, $W'$, has been sought through decays to $qq'$ via $s_8$. The specific model...
A new excited $W^*$ boson [31,32] is generated through a simplified model [30] in the CalcHEP 3.4.2 generator, in combination with the MSTW2008LO PDF and PYTHIA 8.165 for the simulation of nonperturbative effects. The sine of the mixing angle in this model ($\sin \phi_X$) is set to zero, producing leptophobic decays of the $W^*$ that are limited to quarks. With $\sin \phi_X = 1$, a leptophilic $W^*$ would instead be produced with branching ratios divided equally between quarks (3 families $\times$ 3 flavors $\times$ 8.3% BR = 75%) and leptons (3 families $\times$ 8.3% BR = 25%). The angular distribution of the $W^*$ differs from that of the other signals under study, preferring decays with a wider separation in $y$. The acceptance for both leptophobic and leptophilic $W^*$ spans 25% to 27%.

A model for quantum black holes (QBH) that decay to two jets is simulated using both the BLACKMAX [68] and the QBH [69] generators, to produce a simple two-body final-state scenario of quantum gravitational effects at the fundamental Planck scale $M_D$, with $n = 6$ extra spatial dimensions in the context of the ADD model [70]. In this model, the Planck scale is set equal to the threshold mass for the quantum black hole production $m_{th}$. These QBH models are used as benchmarks to represent any quantum gravitational effect that produces events containing dijets. The PDF used for the generation and parton shower of BLACKMAX is CT10, while the QBH samples employ the MSTW2008LO PDF. In the mass range considered, the branching ratio of QBH to dijets is above 85% and the acceptance is between 52% and 55% for BLACKMAX and 54%–56% for QBH. Nonperturbative effects for events coming from both event generators are simulated with PYTHIA 8.170.

Further information on cross sections, branching ratio to dijets and acceptances for the benchmark models under consideration can be found in HepData [71].

All MC signal samples except for the excited quark signals are passed through a fast detector simulation [72] with the jet calibration appropriately corrected to full simulation. The excited quark signals are simulated using GEANT4 within the ATLAS simulation infrastructure [47].
estimated using \textsc{herwig} ++ 2.6.3 [74,75] and convoluted with the signal shape.

The detector resolution is accounted for by convolving the signal shape with a Gaussian function of width equal to the detector resolution at each signal mass. The result is then truncated below 250 GeV due to the dijet mass cut. This produces a signal template shape that is still generic but more likely to match the forms visible in actual physical processes. The effect on the shape of the signal template originating from the $y^*$ cut is not simulated in the signal templates used for these limits, due to the possible model dependence of the angular distribution of the considered NP process.\textsuperscript{6} Tests of the benchmark model-specific templates indicate that the combined effect of the $y$ and $y^*$ acceptance requirement is constant within 20% as a function of the dijet mass, with the largest discrepancies outside the mass peak.

A. Systematic uncertainties in limit setting

The effects of several systematic uncertainties are considered when setting limits on new phenomena. These are incorporated into the Bayesian marginalization limit setting procedure using Gaussian priors, with one nuisance parameter for each uncertainty. They are listed below.

(i) Choice of fitting function: a tenfold cross validation [76] using the full data set shows that the background is also well described when introducing an additional degree of freedom to Eq. (1),

$$f(x) = p_1 (1 - x)^{p_2} x^{p_3} + p_4 \ln x + p_5 (\ln x)^2.$$  \hspace{1cm} (2)

The $\chi^2$-value of this five-parameter function fit to data is 45 for 57 degrees of freedom. Since the two fit functions provide background estimates that differ beyond statistical uncertainties, an additional uncertainty is introduced due to the choice of fitting function. The difference between the two background estimates is treated as a one-sided nuisance parameter, with a Gaussian prior centered at zero corresponding to the background estimate from Eq. (1) and truncated to one-$\sigma$ corresponding to the background estimate from Eq. (2). The marginalized posterior indicates a preference for the alternative function.

(ii) Background fit quality: the uncertainty on the background parametrization from the fit is estimated by refitting bin-by-bin Poisson fluctuations of the data, as described in Ref. [77]. The resulting uncertainty is calculated by refitting a large number of pseudoexperiments drawn from the data, and defining the fit error from the variation in fit results in each bin: $\pm 1\sigma$ in the uncertainty corresponds to the central 68% of pseudoexperiment fit values in the bin.

(3) Jet energy scale: shifts to the jet energy due to the various jet energy scale (JES) uncertainty components are propagated separately through the analysis of the signal templates. Changes in both shape and acceptance due to the JES uncertainty in the simulated signal templates are considered in the limit setting. Combined, the JES uncertainty shifts the resonance mass peaks by less than 3%: this is the JES shift used for Gaussian and Breit-Wigner limits.

(4) Luminosity: a 2.8% uncertainty [15] is applied to the overall normalization of the signal templates.

(5) Theoretical uncertainties: the uncertainty on the signal acceptance for the model-dependent limits due to the choice of PDF is derived employing the PDF4LHC recommendation [78] using the envelope of the error sets of the NNPDF 2.1 [79] and MSTW2008LO. Renormalization and factorization scale uncertainties on the signal acceptance are considered for the $W'$ and $s8$ signals but found negligible. Since the $W'$ cross section estimation used in this analysis includes NNLO corrections, the uncertainties on cross section due to variations of the renormalization and factorization scales, the choice of PDF, and PDF + $\alpha_s$ variations on the theoretical cross section are considered as well.

The effect of the jet energy resolution uncertainty is found to be negligible. Similarly, effects due to jet reconstruction efficiency and jet angular resolution lead to negligible uncertainties.

B. Constraints on NP benchmark models

The resulting limits for excited quarks are shown in Fig. 3. The expected lower mass limit at 95% C.L. for $q'$ is 3.98 TeV, and the observed limit is 4.06 TeV. The limits for color-octet scalars are shown in Fig. 4. The expected mass limit at 95% C.L. is 2.808 TeV, and the observed limit is 2.708 TeV.

The limits for heavy charged gauge bosons, $W'$, are shown in Fig. 5. The expected mass limit at 95% C.L. is 2.51 TeV, and the observed limit is 2.45 TeV.

The limits for the excited $W'$ boson are shown in Fig. 6. The plot shows the observed and expected limits calculated for a leptophobic $W'$ but includes theory curves for both leptophobic and nonleptophobic $W'$ given that the acceptances for the two samples are the same to within 1%. The expected mass limit for the leptophobic model at 95% C.L. is 1.95 TeV and the observed limit is 1.75 TeV. The expected mass limit for the nonleptophobic model at 95% C.L. is 1.66 TeV and the observed limit is 1.65 TeV.

\textsuperscript{6}If a flat distribution in the jet polar angle in the center-of-mass rest frame for the new particle is assumed for the NP model, the acceptance can be calculated analytically and it amounts to roughly 56% for all considered dijet masses.
The limits for black holes generated using QBH and BLACKMAX are shown in Fig. 7. The observed limit is consistent between the two generators, but the cross sections differ, hence the difference in the mass limit. The observed limits for the two models have visually matching shapes and normalizations, so only one (BLACKMAX) is selected for display. The limits for both models are, however, computed separately and recorded. The expected mass limit for QBH black holes at 95% C.L. is 5.66 TeV, and the observed limit is 5.66 TeV. For BLACKMAX black holes, the expected limit at 95% C.L. is 5.62 TeV and the observed limit is 5.62 TeV. Above \( \sim 4.5 \text{ TeV} \) the observed and expected limits are driven by the absence of any observed data events, leading to identical observed and expected mass limits.

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Although the search phase of the analysis starts at 250 GeV, \( \sigma \times A \) exclusion limits on benchmark NP models are set starting at 800 GeV for the \( q^* \), \( s^8 \), and \( W^* \) models. For BLACKMAX black holes, the expected limit at 95% C.L. is 5.62 TeV and the observed limit is 5.62 TeV. Above \( \sim 4.5 \text{ TeV} \) the observed and expected limits are driven by the absence of any observed data events, leading to identical observed and expected mass limits.
Exclusion limits on quantum black holes are set starting from 1 TeV in light of the large cross section and of previous exclusion limits [77,80].

C. Generic resonance limits on dijet production

The resulting limits on \( \sigma \times A \) for the Gaussian template shape are shown in Fig. 8. Limits resulting from the convolution of Breit-Wigner signals of different intrinsic widths (\( \Gamma_{BW} \)) with the appropriate parton distribution function, parton shower, nonperturbative effects and detector resolution are shown in Figs. 9 and 10. For the initial Breit-Wigner signal the following nonrelativistic function was chosen:

\[
    f(x, \mu, \Gamma) = \frac{1}{2\pi} \frac{\Gamma (x - \mu)}{(x - \mu)^2 + (\Gamma^2/4)},
\]

where \( \mu \) and \( \Gamma \) are the mass and the width of the resonance. The use of a relativistic Breit-Wigner signal for the resonance line shape may lead to different limits than the ones derived using the nonrelativistic approximation above. Parton showers and nonperturbative effects have been simulated using HERWIG ++ 2.6.3, which gives a more conservative limit with respect to what is obtained from PYTHIA.

The difference in shapes between the two Breit-Wigner limits is a result of the much larger low-mass tails resulting from the.gg parton luminosity, which becomes especially pronounced at high masses. The convolution with parton shower and nonperturbative effects enhances this effect further.
For sufficiently narrow resonances, these results may be used to set limits on NP models beyond those considered in the current studies, as described in detail in Appendix A. It should be noted that these limits will be conservative at high masses with respect to the limits obtained with full benchmark templates. This is due to the simplifying assumptions made in their derivation, in particular from the use of a nonrelativistic and mass-independent Breit-Wigner shape.

Gaussian limits should be used when tails from PDF and nonperturbative effects can be safely truncated or neglected. Otherwise, convolved Breit-Wigner signals would be more reliable.

In the case of the Gaussian limits, the signal distribution after applying the kinematic selection criteria on $y^*$, $m_{jj}$ and $\eta$ of the leading jets (Sec. III) should approach a Gaussian distribution. The acceptance should include the jet reconstruction efficiency (100% for the current analysis and detector conditions, since inefficiencies due to calorimeter problems are corrected for in data) and the efficiency with respect to the kinematic selection above. NP models with a width smaller than 5% should be compared to the results with width equal to the experimental resolution only (see Appendix B). For models with a larger width after detector effects, the limit that best matches their width should be used.

VII. CONCLUSIONS

In the 2012 running of the ATLAS experiment at the LHC, the collision energy was raised from 7 TeV to 8 TeV, accompanied by a more than fourfold increase in integrated luminosity. The higher energy, and the associated rise in parton luminosity for high masses, have increased the sensitivity of the search and its mass reach for various model hypotheses. In addition, novel trigger techniques have been employed to extend the search to low dijet masses. The data sample used in the current analysis consists of $20.3 \, \text{fb}^{-1}$ of $pp$ collision data at $\sqrt{s} = 8$ TeV, and the resulting dijet mass distribution extends from 250 GeV to approximately 4.5 TeV.

No resonancelike features are observed in the dijet mass spectrum. This analysis places limits on the cross section times acceptance at the 95% credibility level on the mass or

<table>
<thead>
<tr>
<th>Model and final state</th>
<th>95% C.L. limits [TeV]</th>
</tr>
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<tbody>
<tr>
<td>$q^* \rightarrow q\bar{q}$</td>
<td>3.98</td>
</tr>
<tr>
<td>$s8 \rightarrow gg$</td>
<td>2.80</td>
</tr>
<tr>
<td>$W' \rightarrow q\bar{q}'$</td>
<td>2.51</td>
</tr>
<tr>
<td>Leptophobic $W^* \rightarrow q\bar{q}'$</td>
<td>1.95</td>
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<tr>
<td>Leptophilic $W^* \rightarrow q\bar{q}'$</td>
<td>1.66</td>
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<tr>
<td>QBH black holes</td>
<td>5.66</td>
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<tr>
<td>BLACKMAX black holes (all decays)</td>
<td>5.62</td>
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(q and $g$ decays only)
energy scale of a variety of hypotheses for physics phenomena beyond the Standard Model.

To illustrate the typical increases in sensitivity to new phenomena at the LHC up to the end of 2012 running, Table II shows the history of expected limits from ATLAS studies using dijet resonance analysis of two benchmark models, excited quarks and color-octet scalars. The limits set by this analysis on excited quarks, color-octet scalars, heavy $W'$ bosons, chiral $W'$ bosons, and quantum black holes, are summarized in Table I.

**ACKNOWLEDGMENTS**

We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently. We acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWFW and FWF, Austria; ANAS, Azerbaijan; SSTC, Belarus; CNPq and FAPESP, Brazil; NSERC, NRC and CFI, Canada; CERN; CONICYT, Chile; CAS, MOST and NSFC, China; COLCIENCIAS, Colombia; MSMT CR, MPO CR and VSC CR, Czech Republic; DNRF, DNSRC and Lundbeck Foundation, Denmark; EPLANET, ERC and NSRF, European Union; IN2P3-CNRS, CEA-DSM/IRFU, France; GNSF, Georgia; BMBF, DFG, HGF, MPG and AvH Foundation, Germany; GSRT and NSRF, Greece; ISF, MINERVA, GIF, I-CORE and Benoziyo Center, Israel; INFN, Italy; MEXT and JSPS, Japan; CNRST, Morocco; FOM and NWO, Netherlands; BRF and RCN, Norway; MNiSW and NCN, Poland; GRICES and FCT, Portugal; MINE/IFA, Romania; MES of Russia and ROSATOM, Russian Federation; JINR; MSTD, Serbia; MSSR, Slovakia; ARRS and MIZŠ, Slovenia; DST/NRF, South Africa; MINECO, Spain; SRC and Wallenberg Foundation, Sweden; SER, SNSF and Cantons of Bern and Geneva, Switzerland; NSC, Taiwan; TAEK, Turkey; STFC, the Royal Society and Leverhulme Trust, United Kingdom; DOE and NSF, United States of America. The crucial computing support from all WLCG partners is acknowledged gratefully, in particular from CERN and the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CC-IN2P3 (France), KIT/GridKA (Germany), INFN-CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (UK) and BNL (USA) and in the Tier-2 facilities worldwide.

**APPENDIX A: SUGGESTED PROCEDURE FOR SETTING LIMITS ON GENERIC NP MODELS**

1. Setting limits for NP models with a Gaussian shape

The following detailed procedure is appropriate for setting limits involving resonances that are approximately Gaussian near the core, and with tails that are much smaller than the background. The results of Fig. 8 are provided in tables on HepData.

1. Generate an MC sample of a hypothetical new particle with mass set to $M$. Nonperturbative effects should be included in the event generation. Apply the kinematic selection on the parton $\eta$, $p_T$, and $|y^*|$ used in this analysis, as in Sec. III.
2. Smear the signal mass distribution to reflect the detector resolution. The smearing factors derived from full ATLAS simulation of QCD dijet events can be taken from Fig. 11.
3. Since a Gaussian signal shape has been assumed in determining the limits, any long tails in the reconstructed $m_{jj}$ should be removed in the sample under study. The recommendation (based on optimization using $q^*$ templates) is to retain events with $m_{jj}$ between $0.8M$ and $1.2M$. The mean mass, $m$, should be recalculated for this truncated signal.
4. The fraction of MC events surviving the first four steps determines the modified acceptance, $A$.
5. From the table in HepData, select $m_G$ so that $m_G = m$. If the exact value of $m$ is not among the listed values of $m_G$, check the limit for the two values of $m_G$ that are directly above and below $m$, and use the larger of the two limits to be conservative.
6. To retain enough of the information in the full signal template, and at the same time reject tails that would invalidate the Gaussian approximation, the following truncation procedure is recommended. For

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**TABLE II. ATLAS previous and current expected 95% C.L. upper limits [TeV] on excited quarks and color-octet scalars.**

<table>
<thead>
<tr>
<th>$\sqrt{s}$ (TeV)</th>
<th>$\mathcal{L}$ (fb$^{-1}$)</th>
<th>Citation</th>
<th>$q^*$ (TeV)</th>
<th>$s_8$ (TeV)</th>
</tr>
</thead>
<tbody>
<tr>
<td>7</td>
<td>36</td>
<td>[11]</td>
<td>2.07</td>
<td>-</td>
</tr>
<tr>
<td>7</td>
<td>1.0</td>
<td>[12]</td>
<td>2.81</td>
<td>1.77</td>
</tr>
<tr>
<td>7</td>
<td>4.8</td>
<td>[13]</td>
<td>2.94</td>
<td>1.97</td>
</tr>
<tr>
<td>8</td>
<td>20.3</td>
<td>Current</td>
<td>3.98</td>
<td>2.80</td>
</tr>
</tbody>
</table>

---

**FIG. 11. Dijet mass resolution obtained from fully simulated Pythia QCD Monte Carlo Pythia 8.175 [43], with the AU2 tune obtained from ATLAS data [45]. The dijet mass resolution is interpolated linearly between the bin centers.**
this mass point, choose a value of $\sigma_G/m_G$ such that the region within $\pm 2\sigma_G$ is well contained in the (truncated) mass range. For the $q^*$ case a good choice is $\sigma_G=(1.2M-0.8M)/5$ so that 95% of the Gaussian spans $4 \times (0.4/5)M$. Use this value to pick the closest $\sigma_G/m_G$ value, rounded up to be conservative.

(7) Compare the tabulated 95% C.L. upper limit corresponding to the chosen $m_G$ and $\sigma_G/m_G$ values to the $\sigma \times A$ obtained from the theoretical cross section of the model multiplied by the acceptance defined in step (4) above and taking into account its branching ratio into dijets.

2. Setting limits for NP models with a Breit-Wigner shape, accounting for PDF effects

The following detailed procedure is appropriate for setting limits involving resonances that approximate a Breit-Wigner (BW) shape and extend with a low-mass tail due to effects of parton luminosity. For signals that are very narrow or whose tails deviate significantly from a BW, a truncation of the signal template suggested in the Gaussian narrow or whose tails deviate significantly from a BW, a truncation of the signal template suggested in the Gaussian shape, accounting for PDF effects should be simulated after the hard scattering.

(2) Smear the signal mass distribution to reflect the detector resolution. The smearing factors for the dijet mass are derived from full ATLAS simulation and can be taken from Fig. 11.

(3) The kinematic selection detailed in Sec. III should be applied to the simulated events. It should be checked at this point that the shape of the template after the $y^* < 0.6$ cut does not change significantly. For example, in a simple model with a flat distribution for the cosine of the polar angle of the jets in the rest frame of the resonance ($\cos \theta'$) that decays into two back-to-back jets ($y_{lead} \approx -y_{sublead}$), a cut on $|y^*| = 0.5 \times |y_{lead} - y_{sublead}| < 0.6$ imposes a cut on the unboosted rapidity distribution $(|y_{lead,sublead}| < 0.6)$. In the mass ranges investigated, this corresponds to a more stringent constraint than the $m_{BW} < 2.8$ acceptance correction, leading to an acceptance of $\sim 0.5$ that does not depend on dijet mass. Deviations from a flat acceptance of up to 20% can be observed in the tails of models with different angular distributions ($q^*, s8, W'$).

FIG. 12 (color online). Dijet invariant mass for models corresponding to a resonance mass of 2.5 TeV. All distributions are normalized to the same area.

(4) The fraction of generated events surviving the first three steps determines the signal acceptance, $A$.
(5) From the tables available in HepData, select the one corresponding to the production mode for the new resonance (gluon-gluon, gluon-quark, quark-quark or quark-antiquark) as the parton luminosities and hence the signal shapes differ.

(6) Compare the tabulated 95% C.L. upper limit corresponding to the chosen $M$ and $\Gamma/M$ values to the $\sigma \times A$ obtained from the theoretical cross section of the model, multiplied by the acceptance defined in step (3) above and taking into account its branching ratio into dijets.

If the exact values of $M$ and $\Gamma/M$ are not among the listed values of $m_{BW}$ and $\Gamma_{BW}/m_{BW}$, check the limit for the two values of $m_{BW}$ that are directly above and below $M$, and use the more conservative of the two limits.

APPENDIX B: DIJET MASS RESOLUTION

The dijet mass resolution in Fig. 11 is derived from fully simulated QCD Monte Carlo, generated with PYTHIA 8.175 [43], using the AU2 tune obtained from ATLAS data [45] using the analysis selection detailed in Sec. III. The dijet mass resolution $\sigma_m/m_j$ is 8% at $m_j \approx 250$ GeV falls to 4% at 2 TeV and approaches 4% at $m_j$ of 3 TeV and above, and it is interpolated linearly between the bin centers.

APPENDIX C: SIGNAL TEMPLATE SHAPES

For ease of comparison of the shapes of different signals used in this paper, the various signal template shapes are overlaid in Fig. 12 for the mass point of 2.5 TeV, after normalizing to the same area.
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(ATLAS Collaboration)

1Department of Physics, University of Adelaide, Adelaide, Australia
2Physics Department, SUNY Albany, Albany, New York, USA
3Department of Physics, University of Alberta, Edmonton, Alberta, Canada
4aDepartment of Physics, Ankara University, Ankara, Turkey
4bDepartment of Physics, Gazi University, Ankara, Turkey
4cDivision of Physics, TOBB University of Economics and Technology, Ankara, Turkey
4dTurkish Atomic Energy Authority, Ankara, Turkey
5LAPP, CNRS/IN2P3 and Université de Savoie, Annecy-le-Vieux, France
6High Energy Physics Division, Argonne National Laboratory, Argonne IL, USA
7Department of Physics, University of Arizona, Tucson, Arizona, USA
8Department of Physics, The University of Texas at Arlington, Arlington, Texas, USA
9Physics Department, University of Athens, Athens, Greece
10Physics Department, National Technical University of Athens, Zografou, Greece
11Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
12Institut de Física d’Altes Energies and Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain
13aInstitute of Physics, University of Belgrade, Belgrade, Serbia
13bVinca Institute of Nuclear Sciences, University of Belgrade, Belgrade, Serbia
14Department for Physics and Technology, University of Bergen, Bergen, Norway
15Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, California, USA
16Department of Physics, Humboldt University, Berlin, Germany
17Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
18School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
19aDepartment of Physics, Bogazici University, Istanbul, Turkey
19bDepartment of Physics, Dogus University, Istanbul, Turkey
19cDepartment of Physics Engineering, Gaziantep University, Gaziantep, Turkey
20aINFN Sezione di Bologna, Italy
20bDipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy
21Physikalisches Institut, University of Bonn, Bonn, Germany
22Department of Physics, Boston University, Boston, Massachusetts, USA
23Department of Physics, Brandeis University, Waltham, Massachusetts, USA
24aUniversidade Federal do Rio de Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil
24bFederal University of Juiz de Fora (UFJF), Juiz de Fora, Brazil
24cFederal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei, Brazil
24dInstituto de Física, Universidade de Sao Paulo, Sao Paulo, Brazil
25Physics Department, Brookhaven National Laboratory, Upton, New York, USA
26aNational Institute of Physics and Nuclear Engineering, Bucharest, Romania
26bNational Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj Napoca, Romania
26cUniversity Politehnica Bucharest, Bucharest, Romania
26dWest University in Timisoara, Timisoara, Romania
27Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
28Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
29Department of Physics, Carleton University, Ottawa, Ontario, Canada
30CERN, Geneva, Switzerland
31Enrico Fermi Institute, University of Chicago, Chicago, Illinois, USA
32aDepartamento de Física, Pontifícia Universidade Católica de Chile, Santiago, Chile
32bDepartamento de Física, Universidad Técnica Federico Santa María, Valparaiso, Chile
33aInstitute of High Energy Physics, Chinese Academy of Sciences, Beijing, China
33bDepartment of Modern Physics, University of Science and Technology of China, Anhui, China
33cDepartment of Physics, Nanjing University, Jiangsu, China
School of Physics, Shandong University, Shandong, China

Physics Department, Shanghai Jiao Tong University, Shanghai, China

Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France

Nevis Laboratory, Columbia University, Irvington, New York, USA

Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark

INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Italy

Dipartimento di Fisica, Università della Calabria, Rende, Italy

AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland

Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland

The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland

Physics Department, Southern Methodist University, Dallas, Texas, USA

Physics Department, University of Texas at Dallas, Richardson, Texas, USA

DESY, Hamburg and Zeuthen, Germany

Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany

Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany

Department of Physics, Duke University, Durham, North Carolina, USA

SUPA-School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom

INFN Laboratori Nazionali di Frascati, Frascati, Italy

Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany

Section de Physique, Université de Genève, Geneva, Switzerland

INFN Sezione di Genova, Italy

Dipartimento di Fisica, Università di Genova, Genova, Italy

E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi, Georgia

High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia

II Physikalisches Institut, Justus-Liebig-Universität Gießen, Gießen, Germany

SUPA-School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom

II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany

Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS/IN2P3, Grenoble, France

Department of Physics, Hampton University, Hampton, Virginia, USA

Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, Massachusetts, USA

Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany

Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany

ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany

Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan

Department of Physics, Indiana University, Bloomington, Indiana, USA

Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria

University of Iowa, Iowa City, Iowa, USA

University of Iowa, Iowa State University, Ames, Iowa, USA

Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia

KEK, High Energy Accelerator Research Organization, Tsukuba, Japan

Graduate School of Science, Kobe University, Kobe, Japan

Faculty of Science, Kyoto University, Kyoto, Japan

Kyoto University of Education, Kyoto, Japan

Department of Physics, Kyushu University, Fukuoka, Japan

Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina

Physics Department, Lancaster University, Lancaster, United Kingdom

INFN Sezione di Lecce, Italy

 Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy

 Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom

 Department of Physics, Jožef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia

 School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom

 Department of Physics, Royal Holloway University of London, Surrey, United Kingdom

 Department of Physics and Astronomy, University College London, London, United Kingdom

 Louisiana Tech University, Ruston LA, USA

 Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France

 Fysiska institutionen, Lunds universitet, Lund, Sweden
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81 Departamento de Fisica Teorica C-15, Universidad Autonoma de Madrid, Madrid, Spain
82 Institut für Physik, Universität Mainz, Mainz, Germany
83 School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
84 CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
85 Department of Physics, University of Massachusetts, Amherst, Massachusetts, USA
86 Department of Physics, McGill University, Montreal, QC, Canada
87 School of Physics, University of Melbourne, Victoria, Australia
88 Department of Physics, The University of Michigan, Ann Arbor, Michigan, USA
89 Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan, USA
90a INFN Sezione di Milano, Italy
90b Dipartimento di Fisica, Università di Milano, Milano, Italy
91 B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus
92 National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Republic of Belarus
93 Department of Physics, Massachusetts Institute of Technology, Cambridge, Massachusetts, USA
94 Group of Particle Physics, University of Montreal, Montreal, QC, Canada
95 P.N. Lebedev Institute of Physics, Academy of Sciences, Moscow, Russia
96 Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia
97 Moscow Engineering and Physics Institute (MEPhI), Moscow, Russia
98 D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
99 Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
100 Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
101 Nagasaki Institute of Applied Science, Nagasaki, Japan
102 Graduate School of Science and Kobe University Institute, Nagoya University, Nagoya, Japan
103 INFN Sezione di Napoli, Italy
104 Dipartimento di Fisica, Università di Napoli, Napoli, Italy
105 Department of Physics and Astronomy, University of New Mexico, Albuquerque, New Mexico, USA
106 Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, The Netherlands
107 Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, The Netherlands
108 Department of Physics, Northern Illinois University, DeKalb, Illinois, USA
109 Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia
110 Department of Physics, New York University, New York, New York, USA
111 Department of Physics, Ohio State University, Columbus, Ohio, USA
112 Faculty of Science, Okayama University, Okayama, Japan
113 Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, Oklahoma, USA
114 Department of Physics, Oklahoma State University, Stillwater, Oklahoma, USA
115 Palacký University, RCPTM, Olomouc, Czech Republic
116 Center for High Energy Physics, University of Oregon, Eugene, Oregon, USA
117 LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France
118 Graduate School of Science, Osaka University, Osaka, Japan
119 Department of Physics, University of Oslo, Oslo, Norway
120 Department of Physics, Oxford University, Oxford, United Kingdom
121 INFN Sezione di Pavia, Italy
122 Dipartimento di Fisica, Università di Pavia, Pavia, Italy
123 Department of Physics, University of Pennsylvania, Philadelphia, Pennsylvania, USA
124 Petersburg Nuclear Physics Institute, Gatchina, Russia
125 INFN Sezione di Pisa, Italy
126 Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
127 Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, Pennsylvania, USA
128 Laboratorio de Instrumentacion y Fisica Experimental de Particulas-LIP, Lisboa, Portugal
129 Faculdade de Ciências, Universidade de Lisboa, Lisboa, Portugal
130 Department of Physics, University of Coimbra, Coimbra, Portugal
131 Centro de Física Nuclear da Universidade de Lisboa, Lisboa, Portugal
132 Departamento de Física, Universidade do Minho, Braga, Portugal
133 Departamento de Física Teorica y del Cosmos and CAFPE, Universidad de Granada, Granada, Spain
134 Dep Fisica and CEFITEC of Faculdade de Ciencias e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal

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166 Department of Physics, University of Illinois, Urbana, Illinois, USA
167 Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden
168 Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNMT), University of Valencia and CSIC, Valencia, Spain
169 Department of Physics, University of British Columbia, Vancouver, BC, Canada
170 Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada
171 Department of Physics, University of Warwick, Coventry, United Kingdom
172 Waseda University, Tokyo, Japan
173 Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel
174 Department of Physics, University of Wisconsin, Madison, Wisconsin, USA
175 Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany
176 Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany
177 Department of Physics, Yale University, New Haven, Connecticut, USA
178 Yerevan Physics Institute, Yerevan, Armenia
179 Centre de Calcul de l’Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France

Deceased.
Also at Department of Physics, King’s College London, London, United Kingdom.
Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.
Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom.
Also at TRIUMF, Vancouver, BC, Canada.
Also at Department of Physics, California State University, Fresno, California, USA.
Also at Tomsk State University, Tomsk, Russia.
Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France.
Also at Università di Napoli Parthenope, Napoli, Italy.
Also at Institute of Particle Physics (IPP), Canada.
Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia.
Also at Chinese University of Hong Kong, China.
Also at Department of Financial and Management Engineering, University of the Aegean, Chios, Greece.
Also at Louisiana Tech University, Ruston, Louisiana, USA.
Also at Instituto Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain.
Also at Department of Physics, The University of Texas at Austin, Austin, Texas, USA.
Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia.
Also at CERN, Geneva, Switzerland.
Also at Ochadai Academic Production, Ochanomizu University, Tokyo, Japan.
Also at Manhattan College, New York, New York, USA.
Also at Novosibirsk State University, Novosibirsk, Russia.
Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.
Also at LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France.
Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan.
Also at Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France.
Also at School of Physical Sciences, National Institute of Science Education and Research, Bhubaneswar, India.
Also at Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy.
Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia.
Also at Section de Physique, Université de Genève, Geneva, Switzerland.
Also at International School for Advanced Studies (SISSA), Trieste, Italy.
Also at Department of Physics and Astronomy, University of South Carolina, Columbia, South Carolina, USA.
Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China.
Also at Faculty of Physics, M.V.Lomonosov Moscow State University, Moscow, Russia.
Also at Moscow Engineering and Physics Institute (MEPhI), Moscow, Russia.
Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.
Also at Department of Physics, Oxford University, Oxford, United Kingdom.
Also at Department of Physics, Nanjing University, Jiangsu, China.
Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany.
Also at Department of Physics, The University of Michigan, Ann Arbor, Michigan, USA.
Also at Discipline of Physics, University of KwaZulu-Natal, Durban, South Africa.
Also at University of Malaya, Department of Physics, Kuala Lumpur, Malaysia.