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Measurements of the Total and Differential Higgs Boson Production Cross Sections Combining the $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$ Decay Channels at $\sqrt{s} = 8$ TeV with the ATLAS Detector

G. Aad et al.*

(ATLAS Collaboration)

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Measurements of the total and differential cross sections of Higgs boson production are performed using 20.3 fb$^{-1}$ of $pp$ collisions produced by the Large Hadron Collider (LHC) [1] at a center-of-mass energy of $\sqrt{s} = 8$ TeV and recorded by the ATLAS detector. Cross sections are obtained from measured $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$ event yields, which are combined accounting for detector efficiencies, fiducial acceptances, and branching fractions. Differential cross sections are reported as a function of Higgs boson transverse momentum, Higgs boson rapidity, number of jets in the event, and transverse momentum of the leading jet. The total production cross section is determined to be $\sigma_{pp-H} = 33.0 \pm 5.3$ (stat) $\pm 1.6$ (syst) pb. The measurements are compared to state-of-the-art predictions.

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This Letter presents measurements of the total and differential cross sections of inclusive Higgs boson production using 20.3 fb$^{-1}$ of $pp$ collisions produced by the Large Hadron Collider (LHC) [1] at a center-of-mass energy of $\sqrt{s} = 8$ TeV and recorded by the ATLAS detector [2]. The measured cross sections probe the properties of the Higgs boson and can be directly compared to the theoretical modeling of different Higgs boson production mechanisms, such as the most recent gluon fusion (ggF) QCD calculations. They can also be used to constrain new physics scenarios, for example using the effective field theory framework as proposed in Refs. [3–7]. The analysis uses event yields measured in the $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$ decays and detector efficiencies, both determined as described in Refs. [8,9]. The statistical uncertainties on the Higgs boson signal yields in both channels are larger than the systematic uncertainties, while the total uncertainties in the two channels are similar. Combining the analyses improves the precision of the cross-section measurements by up to 40%, and by 25%—30% on average, with respect to the corresponding measurements in the most precise individual channel.

Distributions of the differential $pp \rightarrow H$ cross sections are reported as a function of the transverse momentum $p_T^H$ and the rapidity $y^H$ of the Higgs boson, the jet multiplicity $N_{jets}$, and the transverse momentum of the leading jet $p_T^{1\ell}$. The observables $p_T^H$ and $y^H$ describe the kinematics of the Higgs boson. They are sensitive to perturbative QCD modeling in ggF production, which is the dominant Higgs boson production mechanism in the Standard Model (SM). The $[y^H]$ distribution furthermore offers a clean probe of the gluon parton distribution function (PDF) and will play a role in future PDF fits. The $N_{jets}$ and $p_T^{1\ell}$ observables probe the theoretical modeling of partonic radiation in ggF production as well as the overall rate and modeling of jets in vector-boson fusion (VBF) and associated Higgs boson production ($VH$ and $t\bar{t}H$). Jets produced in VBF, $VH$, and $t\bar{t}H$ processes tend to have higher transverse momenta than those produced via ggF production; however, the sensitivity to measuring these contributions is weak with the current amount of data.

Cross sections are extracted using a combined likelihood built from the signal yields in the $H \rightarrow \gamma\gamma$ channel and the data and background yields in the $H \rightarrow ZZ^* \rightarrow 4\ell$ channel, as well as detector efficiencies, fiducial acceptances and SM branching fractions [10]. A complementary approach, using a separate likelihood, measures the shape of the differential distributions by imposing a unity normalization constraint, which removes the implicit SM assumption on the branching fractions. For the extraction of the signal yields and the corrections of detector efficiencies, it is assumed that the signal in both channels is due to a narrow resonance with a mass $m_H = 125.36 \pm 0.41$ GeV as measured by the ATLAS Collaboration [11]. The signal yield in the $H \rightarrow \gamma\gamma$ channel is obtained from fits to the diphoton mass spectra [8], and from the background subtracted data yield in a $m_{4\ell}$ mass window of 118 to 129 GeV for the $H \rightarrow ZZ^* \rightarrow 4\ell$ channel [9]. The fiducial acceptance in both channels [8,9] is derived using a set of Monte Carlo (MC) event generators. POWHEG-BOX [12–14], interfaced with PYTHIA8 [15] for showering, is used to generate ggF and VBF events, while PYTHIA8 is used to simulate $VH$ and

*Full author list given at the end of the article.

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associated production with top quarks ($t\bar{t}H$) and $b$-quarks ($b\bar{b}H$). The fiducial acceptance for events with $|y^{H}| < 1.2$ is approximately 72% for $H \to \gamma\gamma$, and 55%–59% for $H \to ZZ^* \to 4\ell$. For higher $|y^{H}|$, the acceptance decreases to 35%–38% in both channels. The fiducial acceptance is more constant as a function of the other variables and is in the range 56%–62% for the $H \to \gamma\gamma$ channel and 44%–53% for the $H \to ZZ^* \to 4\ell$ channel [16].

After correcting the differential cross sections and normalized shapes for fiducial acceptance and branching fractions, the corresponding measurements in both channels are found to be in good agreement with each other; $p$ values obtained from $\chi^2$ compatibility tests are in the range 56%–99% [16].

In the binned maximum-likelihood fit, the statistical uncertainty of the $H \to \gamma\gamma$ event yield is modeled using a Gaussian distribution, while the event yield in the $H \to ZZ^* \to 4\ell$ channel follows a Poisson distribution due to the small sample size. Experimental and theoretical systematic uncertainties affecting the signal yields, detector efficiencies, branching fractions, and fiducial acceptance corrections are taken into account in the likelihood as constrained nuisance parameters. Nuisance parameters describing the same uncertainty sources are treated as fully correlated between bins and channels. Systematic uncertainties on the $H \to \gamma\gamma$ and $H \to ZZ^* \to 4\ell$ background estimates and efficiency correction factors, as well as the uncertainty on the integrated luminosity, are described in detail in Refs. [8,9]. The branching fraction uncertainty due to the assumed quark masses and other theoretical uncertainties are evaluated following the recommendations of Ref. [17], considering uncertainty correlations between the $H \to \gamma\gamma$ and $H \to ZZ^* \to 4\ell$ decay channels. Uncertainties on the acceptance correction related to the choice of PDF set are evaluated by taking the envelope of the sum in quadratures of eigenvector variations of the baseline (CT10 [18]) and the central values of alternative (MSTW2008NLO [19] and NNPDF2.3 [20]) PDF sets. Uncertainties on the acceptance correction associated with missing higher-order corrections are evaluated by varying the renormalization and factorization scales coherently and individually by factors of 0.5 and 2 from their nominal values, and by reweighting the $p_T^H$ distribution from POWHEG-BOX to the prediction of the HRes 2.2 calculation [21,22].

The envelope of the maximum deviation of the combined scale variations and the $p_T^H$ reweighting is used as the systematic variation. To account for the uncertainty in the mass measurement, the Higgs boson mass is varied by $\pm 0.4$ GeV. To assess the systematic uncertainty due to the assumption of SM cross-section fractions of the Higgs boson production modes, the VBF and $VH$ fractions are varied by factors of 0.5 and 2 from the SM prediction and the fraction of $t\bar{t}H$ is varied by factors of 0 and 5. These factors are based on current experimental bounds [23–27].

The total uncertainties on the acceptance correction range from 1% to 6%, depending on the channel, distribution and bin.

The total systematic uncertainties on the combined differential cross sections range from 4% to 12%, depending on the distribution and bin. For the kinematic variables $p_T^H$ and $|y^{H}|$, the largest systematic uncertainties on the differential cross sections are due to the luminosity and the background estimates in both channels. For the jet variables $N_{jets}$ and $p_T^j$, the largest systematic uncertainties on the differential cross sections are due to the jet energy scale and resolution. In the shape combination, the normalization uncertainties including luminosity, branching fractions, and efficiency uncertainties do not apply. Statistical uncertainties dominate all resulting distributions, ranging from 23% to 75%. The total $pp \to H$ cross section is determined in the $H \to \gamma\gamma$ channel to be $31.4 \pm 7.2$ (stat) $\pm 1.6$ (syst) pb and in the $H \to ZZ^* \to 4\ell$ channel to be $35.0 \pm 8.4$ (stat) $\pm 1.8$ (syst) pb. Combining the analyses yields $\sigma_{pp \to H}$ = $33.0 \pm 5.3$ (stat) $\pm 1.6$ (syst) pb. Figure 1 presents a comparison of these measurements with two ggF predictions to which contributions from other relevant Higgs boson production modes (VBF, $VH$, $t\bar{t}H$, $b\bar{b}H$) are added using cross sections and uncertainties from Ref. [10]. The LHC-XS ggF prediction, recommended in Ref. [10], is accurate to next-to-next-to-leading order (NNLO) in QCD and utilizes threshold resummation accurate to next-to-next-to-leading logarithms (NNLL). A significant effort has been undertaken by the theory community to provide ggF cross sections beyond this precision through various improvements in the perturbative calculations [28–33]. Recently, the ADDFGHLM group has provided a fixed-order calculation accurate to next-to-next-to-leading order ($N^3$LO) [34–37]. A PDF uncertainty

![FIG. 1 (color online). Measured total cross section of Higgs boson production compared to two calculations of the ggF cross section. Contributions from other relevant Higgs boson production modes (VBF, $VH$, $t\bar{t}H$, $b\bar{b}H$) are added using cross sections and uncertainties from Ref. [10]. Details of the predictions are presented in Table I.](091801-2)
FIG. 2 (color online). Measured Higgs boson production cross sections in inclusive and exclusive jet multiplicity bins compared to different theoretical predictions (see Table I for details and references).

Photons are excluded from jet-finding if they lie inside a cone of radius $\Delta R < 0.1$ of an electron or muon, and neither the photon nor lepton originate from a hadron decay. To allow comparisons with the unfolded measurements, the analytical calculations are corrected for effects of hadronization and multiple particle interactions. These correction factors and their associated uncertainties are obtained using the PYTHIA8 and HERWIG [41] MC event generators with different tunes [42]. The total cross sections from the ML merged predictions are lower than from fully inclusive NNLO to NLO. The central value of the measured total cross section is larger than the SM predictions presented in Fig. 1. A likelihood-ratio test statistic is used to quantify the agreement, using a bifurcated Gaussian to model the asymmetric theory uncertainties. The resulting $p$ values are 5.5% and 9.0% for the agreement between data and the predictions from LHC-XS and ADDFGHLM, respectively. The ratio of the measured cross section to the LHC-XS prediction is larger than the SM predictions presented in Fig. 1.

The larger Higgs event yield observed in data motivates the search for deviations from the SM. The accuracy of the ggF predictions is increased by a factor of 10.5, which is larger than the SM predictions presented in Fig. 1. A likelihood-ratio test statistic is used to quantify the agreement between data and the predictions from the ML merged predictions.

The PDF uncertainties further include resummation scale uncertainties due to the choice of merging scale. The PDF variations. The ML-merged ggF predictions also have been resummed to NNLL order, which is the same as the analytical calculations. Contributions from other relevant Higgs boson production modes are generated using POWHEG for VBF and PYTHIA8 for VH, $t\bar{t}H$, and $bbH$, and are scaled to the cross sections in Ref. [10]. Uncertainties are assigned to all MC predictions from QCD scale and PDF variations. The total cross sections from the ML merged predictions are lower than from fully inclusive NNLO + NLL calculations. However, for $N_{\text{jets}} \geq 1$, the MC predictions formally have NLO accuracy, which is the same as the analytical calculations. Contributions from other relevant Higgs boson production modes are generated using POWHEG for VBF and PYTHIA8 for VH, $t\bar{t}H$, and $bbH$, and are scaled to the cross sections in Ref. [10]. Uncertainties are assigned to all MC predictions from QCD scale and PDF variations. The total cross sections from the ML merged predictions are lower than from fully inclusive NNLO + NLL calculations. However, for $N_{\text{jets}} \geq 1$, the MC predictions formally have NLO accuracy, which is the same as the analytical calculations. Contributions from other relevant Higgs boson production modes are generated using POWHEG for VBF and PYTHIA8 for VH, $t\bar{t}H$, and $bbH$, and are scaled to the cross sections in Ref. [10]. Uncertainties are assigned to all MC predictions from QCD scale and PDF variations. The total cross sections from the ML merged predictions are lower than from fully inclusive NNLO + NLL calculations.
FIG. 3 (color online). Differential cross sections (left) and normalized cross-section shapes (right) for inclusive Higgs boson production measured by combining the $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ \rightarrow 4\ell$ channels. The measured variables are the Higgs boson transverse momentum $p_T^H$ (top) and its rapidity $|y|^{H}$ (middle), and the transverse momentum of the leading jet $p_T^{j_1}$ (bottom). The 0–30 GeV bin of the $p_T^{j_1}$ distributions corresponds to events without jets above 30 GeV. Various theoretical predictions are presented, using the same bin widths as the measurement.
agreement between data and predictions can be found in the inclusive and exclusive 1-jet bins, with local $p$ values ranging between 0.1% and 3.6%. Normalizing the total expected cross section to the data results in an improved agreement for these bins, with local $p$ values ranging from 4%–29%.

The combined differential cross sections as a function of $p_T^H$, $|y^H|$, and $p_T^{ll}$ are shown in Fig. 3 (left). The measured $p_T^H$ and $|y^H|$ distributions are compared to the HREM calculation and the $p_T^{ll}$ measurement is compared to STWZ and JetVHeto predictions. Figure 3 (right) shows the comparisons of the normalized shapes to predictions from the MC event generators NNLOPS, SHERPA 2.1.1, and MG5_aMC@NLO, as well as the HREM calculation. The uncertainties on the predicted shapes are evaluated following the same approach as for the differential cross-section predictions. They are derived from the impact of QCD scale, merging scale, and PDF variations. The mean of the measured $p_T^H$ distribution is $40.1 \pm 3.0$ GeV, while the means of the MC predictions range from 34 to 37 GeV.

The $p$ values quantifying the compatibility of the measured cross sections and predictions range from 2% to 26%, and for the shapes from 8% to 88%. For the calculation of these values, the theory uncertainties are assumed to be Gaussian distributed and fully correlated between bins [16].

In conclusion, this Letter presents the first measurements of total and differential cross sections and shapes for inclusive $pp \rightarrow H$ production. The measurements were performed in the $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ' \rightarrow 4\ell^\pm$ channels using the full 2012 data set, which consists of 20.3 fb$^{-1}$ of $pp$ collisions produced by the LHC at a center-of-mass energy of $\sqrt{s} = 8$ TeV and recorded by the ATLAS detector. The results of the two channels are compatible and have similar precision. The measurements indicate that the total production cross section of the Higgs boson is larger, and that it is produced with larger transverse momentum and more associated jets than predicted by the current most advanced SM calculations; however, more data is needed to confirm these observations.

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Laboratoire de Physique Nucléaire et de Hautes Énergies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France

Department of Physics, The Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China

Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands

Department of Physics, Hampton University, Hampton, Virginia, USA

Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, Massachusetts, USA

Kirchhoff-Institut für Physik, Rupprecht-Karls-Universität Heidelberg, Heidelberg, Germany

Physikalisches Institut, Rupprecht-Karls-Universität Heidelberg, Heidelberg, Germany

ZITI Institut für technische Informatik, Rupprecht-Karls-Universität Heidelberg, Mannheim, Germany

Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan

Department of Physics, The Chinese University of Hong Kong, Shatin, N.T., Hong Kong, China

Department of Physics, The University of Hong Kong, Hong Kong, China

Department of Physics, Indiana University, Bloomington, Indiana, USA

Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria

University of Iowa, Iowa City, Iowa, USA

Department of Physics and Astronomy, Iowa State University, Ames, Iowa, USA

Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia

KEK, High Energy Accelerator Research Organization, Tsukuba, Japan

Graduate School of Science, Kobe University, Kobe, Japan

Faculty of Science, Kyoto University, Kyoto, Japan

Kyoto University of Education, Kyoto, Japan

Department of Physics, Kyushu University, Fukuoka, Japan

Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina

Physics Department, Lancaster University, Lancaster, United Kingdom

INFN Sezione di Lecce, Italy

Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom

Department of Physics, Jožef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia

School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom

Department of Physics, Royal Holloway University of London, Surrey, United Kingdom

Department of Physics and Astronomy, University College London, London, United Kingdom

Louisiana Tech University, Ruston, Louisiana, USA

Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France

Fysiska institutionen, Lunds universitet, Lund, Sweden

Departamento de Física Teórica C-15, Universidad Autonoma de Madrid, Madrid, Spain

Institut für Physik, Universität Mainz, Mainz, Germany

School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom

CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France

Department of Physics, University of Massachusetts, Amherst, Massachusetts, USA

Department of Physics, McGill University, Montreal QC, Canada

School of Physics, University of Melbourne, Victoria, Australia

Department of Physics, The University of Michigan, Ann Arbor, Michigan, USA

Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan, USA

INFN Sezione di Milano, Italy

Dipartimento di Fisica, Università di Milano, Milano, Italy

B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus

National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Republic of Belarus

Department of Physics, Massachusetts Institute of Technology, Cambridge, Massachusetts, USA

Group of Particle Physics, University of Montreal, Montreal, QC, Canada

P.N. Lebedev Institute of Physics, Academy of Sciences, Moscow, Russia

Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia

National Research Nuclear University MEPhI, Moscow, Russia

D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia

Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany

Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany

Nagasaki Institute of Applied Science, Nagasaki, Japan

Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan

INFN Sezione di Napoli, Italy

Dipartimento di Fisica, Università di Napoli, Napoli, Italy

Department of Physics and Astronomy, University of New Mexico, Albuquerque, New Mexico, USA

Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands

Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands

Department of Physics, Northern Illinois University, DeKalb, Illinois, USA
5 Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.
7 Also at LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France.
6 Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan.
8 Also at School of Physics, Shandong University, Shandong, China.
9 Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia.
10 Also at Section de Physique, Université de Genève, Geneva, Switzerland.
11 Also at International School for Advanced Studies (SISSA), Trieste, Italy.
12 Also at Department of Physics and Astronomy, University of South Carolina, Columbia, SC, USA.
13 Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China.
14 Also at Faculty of Physics, M.V.Lomonosov Moscow State University, Moscow, Russia.
15 Also at National Research Nuclear University MEPhI, Moscow, Russia.
16 Also at Department of Physics, Stanford University, Stanford, CA, USA.
17 Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.
18 Also at Department of Physics, The University of Michigan, Ann Arbor, MI, USA.
19 Also at Discipline of Physics, University of KwaZulu-Natal, Durban, South Africa.
20 Also at University of Malaya, Department of Physics, Kuala Lumpur, Malaysia.