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Elemental chalcogens as a minimal model for combined charge and orbital order

Ana Silva, Jans Henke, and Jasper van Wezel

Institute for Theoretical Physics, Institute of Physics, University of Amsterdam, 1090 GL Amsterdam, The Netherlands

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I. INTRODUCTION

The bulk transition metal dichalcogenide $1T$-TiSe$_2$ has been shown, uniquely, to harbor a charge density wave transition that breaks inversion symmetry in a chiral way [1–5]. In contrast to the well-known chirality of spins in helical magnets, the formation of spirals within the scalar electronic density cannot occur by itself, and is necessarily accompanied by the onset of simultaneous orbital order [3,6,7]. The result is an example of a novel type of order, in which the orbital and charge degrees of freedom are combined into a single order parameter. Similar chiral order has been theoretically suggested to determine material properties of various transition metal dichalcogenides [8–10], and even cuprate high-temperature superconductors [11–13]. But the cooperation between charge and orbital degrees of freedom is not restricted to chiral phases. The two degrees of freedom have, for example, also already been proposed to combine into a polar order parameter [10], and there is no reason to believe this exhausts the list of possible novel phases.

Focusing first on the chiral combination of charge and orbitals, indirect evidence for the presence of spiral charge order was found in scanning-tunneling microscopy experiments on $1T$-TiSe$_2$ [1,5]. In addition, several predictions arising from a Ginzburg-Landau theory of the chiral phase transition were experimentally confirmed [3–5]. Nevertheless, it has proven difficult to obtain direct experimental confirmation of the broken inversion symmetry. The main reason for this is believed to be the presence of small, nanometer wide, domains of varying handedness [5], averaged over by almost all direct bulk probes. A microscopic understanding, going beyond the predictions of the macroscopic order parameter theory, is thus essential in the search for further experiments able to directly probe this novel type of combined charge and orbital order.

A microscopic theory for the chiral state of $1T$-TiSe$_2$ would necessarily involve all 22 orbitals in its unit cell. Even if such a model were constructed, its inherent complexity would obscure a general understanding of how charge and orbital degrees of freedom cooperate to form a single ordered phase, and would likely not be useful as a guide to identifying other possible materials that can harbor electronic spirals or other types of combined charge and orbital order. We therefore take an alternative approach, and formulate a minimal microscopic model for the appearance of spiral structures consisting of ordered occupied orbitals and increased charge density, both in terms of a microscopic Landau theory and a microscopic Hamiltonian. Both reproduce the known chiral crystal structure and are consistent with its observed thermal evolution and behavior under pressure. The combination of microscopic and macroscopic frameworks allows us to distill the essential ingredients in the emergence of combined charge and orbital order, and may serve as a general guide to predicting and understanding spontaneous chiral order as well as other, more general, types of combined charge and orbital order in other materials.

Helices of increased electron density emerging spontaneously in materials containing multiple, interacting density waves, are an example of how orbital and charge degrees of freedom may combine to form a single ordered phase. Although a macroscopic order parameter theory describing this behavior has been proposed and experimentally tested, a microscopic understanding of such simultaneous orbital and electronic order in specific materials is still lacking. Here we present the elemental chalcogens selenium and tellurium as models for the development of combined charge and orbital order. We formulate minimal models capturing the broken inversion symmetry. The main reason for this is proven difficult to obtain direct experimental confirmation of the chiral structure shown in Fig. 1(a) [14]. Although Se and Te do not exhibit a charge ordering transition at any temperature, the spiral bond order is understood as an instability of a simple cubic lattice structure [2,6,7]. The charge ordering transition from the hypothetical simple cubic to the actual chiral phase is known to be of the same type as the chiral transition in $1T$-TiSe$_2$ [3]. Owing to the simple lattice structure however, an explicit and easily accessible microscopic model can be formulated for Te and Se, elucidating how different types of electron-phonon coupling and Coulomb interactions conspire to form the spiral structure. This model is presented here as a minimal description for combined charge and orbital order in general.

II. INTUITIVE PICTURE

Before presenting both macroscopic and microscopic models of the chiral charge order, we first give an intuitive picture showcasing their basic ingredients. The starting point is the simple cubic lattice structure. Both Se and Te crystals possess the chiral structure shown in Fig. 1(a) for any temperature at ambient pressure. Upon melting Te however, the short-ranged...
The atomic displacement waves $u_j(x) = \tilde{u} \epsilon_j \sin(Q \cdot x)$, forming in response to the charge modulations, have polarizations $\epsilon_j$ whose direction is determined by the anisotropy of the local electron-phonon coupling matrix elements [3]. In a chain of $p_x$ orbitals with overlaps only along $x$, the electron-phonon coupling is maximally anisotropic, and the displacement direction $e$ will be purely along $x$. The three orthogonal chains running through each atom act independently, and the actual atomic displacement is the sum of the three contributions $u_j$.

The charge density wave in each orbital chain can be shifted along its propagation direction by the addition of a phase: $\rho_j(x) \propto \cos(Q \cdot x + \phi_j)$. A Coulomb interaction between charges in orthogonal orbitals on the same site, will cause the charge maxima along one chain to prefer to avoid those of other chains, effectively coupling the phases in different orbital sectors. The lowest energy configuration then produces precisely the charge redistribution and lattice deformations shown in Fig. 1(a), which agree with the experimentally observed crystal structure of Se and Te [2,6]. Each atom in this final structure has a single least occupied $p$ orbital. The chiral charge ordered structure is therefore also automatically an orbital ordered phase, as shown in Fig. 1(b).

III. MACROSCOPIC ORDER PARAMETER THEORY

A Landau free energy may be written in terms of the dimensionless order parameters $\alpha_j(x)$ representing the periodic modulation of the charge density within a given chain of head-to-toe orbitals: $\rho_j(x) = \rho_0 [1 + \alpha_j(x)]$. If the orbital sectors are noninteracting, their free energies are independent: $F_j = \int d^3x \, a(x)\alpha_j^2 + b(x)\alpha_j^2 + c(x)\alpha_j^2$. Notice that the presence of the lattice is taken into account by expanding the coefficients in terms of reciprocal lattice vectors, so that for example $a(x) = a_0 + a_1 \sum n \epsilon^n_0 x + \cdots$ [21]. Here $G_n$ denote the shortest reciprocal lattice vectors. Terms $a_j$ with $j > 0$ in this sum arise from the electron-phonon coupling in a more microscopic model.

The on-site Coulomb interaction provides the interaction terms $F_{\text{ Coulomb}} = \sum_j \int d^3x \, A_0 \alpha_j \alpha_{j+1}$. The periodic charge distributions can be written as $\alpha_j(x) = \psi_0 \cos(Q \cdot x + \varphi_j)$, with the amplitude $\psi_0$ equal for all three order parameters, and $\varphi_j$ is a spatial shift of the charge density wave along $j$. Performing the integration over $x$, the full free energy, per volume, becomes

$$F = \frac{3}{2} a_0 \psi_0^2 + \frac{9}{8} c_0 \psi_0^4 + \frac{1}{4} b_1 \psi_0^2 \sum_j \cos(3\varphi_j) + \frac{1}{2} A_0 \psi_0^2 \sum_{j} \cos(\varphi_j - \varphi_{j+1}).$$

As usual, the temperature dependence is considered by expanding $a_0$ as a function of $T - T_c$, near the critical temperature. In this way, $a_0$ determines when the free energy has a minimum at nonzero values of $\psi_0$, and charge order sets in. The final two terms, arising from the electron-phonon coupling and the Coulomb interaction, respectively, determine the values of the phases $\varphi_j$. They can be simultaneously minimized by taking $\varphi_1 = n\pi/3$, where $n$ is an odd or even integer depending on the sign of $b_1$. Physically, this corresponds to the charge order being either site or bond centered. Additionally, the relative
phase differences should be chosen as \( \varphi_j - \varphi_{j+1} = \pm 2\pi/3 \). These solutions are then precisely the left- and right-handed chiral configurations, one of which is shown in Fig. 1(a).

Comparing the free energy of Eq. (1) to the one given for 1T-TiSe\(_2\) in Ref. [3], it appears that in spite of the different underlying atomic and electronic configurations, the routes to chiral charge and orbital order are largely the same. The onset of charge order from a disordered state is determined by amplitude terms only. Electron-phonon coupling then favors locking of the different charge density wave contributions to the lattice. The on-site Coulomb interaction finally couples density waves in distinct orbital sectors, leading to relative phase shifts and the emergence of chiral charge and orbital order.

Upon applying uniform pressure, Se and Te undergo a series of structural transitions into nonchiral phases [22,23]. This can be understood within the free energy expansion as a pressure dependence of the critical temperature, which we write as \( a_0(T, P) = \frac{\beta T^2}{2}c_0 + \frac{3c_0}{2} + a_0(T) + b_1\psi_j^2 \cos(3\varphi_j) \)

\[ + \frac{1}{2} a_0\psi_j\psi_{j+1} \cos(\varphi_j - \varphi_{j+1}). \]  

Applying uniaxial pressure suppresses only one of the charge density wave components. As \( T \) is now decreased from high temperatures, one of the \( a_0(T, P) \) terms will remain positive while the others already cause order to set in their associated orbital sectors. There is thus a range of temperatures for which the charge order is confined to two orbital sectors only. This results in stacked planes, each containing zigzag charge order, as indicated in Fig. 2, which also shows the phase diagram resulting from this minimal model. The anisotropic structure agrees both with the predictions of an earlier semiclassical approach in terms of so-called vector charge density waves [6], and with the experimental observation of layered structures in Se under pressure [22,23].

**IV. MICROSCOPIC MODEL**

To see how the terms in the Landau free energy emerge from the interplay of microscopic degrees of freedom, we start again from a 2/3-filled \( p \) shell within the simple cubic lattice. Including hopping only between head-to-toe orbitals, the tight-binding Hamiltonian can be written as

\[ \hat{H} = \hat{H}_{\text{TB}} + \hat{H}_{\text{Coul}} + \hat{H}_{\text{e-ph}} + \hat{H}_{\text{boson}}, \]

\[ \hat{H}_{\text{TB}} = t \sum_{x,j} \hat{c}^\dagger_j(x)\hat{c}_{j+x}(x + a_j) + \text{H.c.}, \]

\[ \hat{H}_{\text{Coul}} = V \sum_{x,j} \hat{c}^\dagger_j(x)\hat{c}_{j}(x)^2\hat{c}_{j}(x), \]

\[ \hat{H}_{\text{boson}} = \hbar \omega \sum_{q,j} \hat{b}^\dagger_q(q)\hat{b}_q(q). \]

Here \( \hat{c}^\dagger_{j}(x) \) creates an electron in orbital \( j \) on position \( x \), and \( a_j \) is the simple cubic lattice vector in direction \( j \). The hopping amplitude \( t \) is positive, since overlapping orbital lobes on neighboring sites have opposite signs. The Coulomb interaction acts on-site only, and the displacement \( \hat{u}_j(x) \) of the atom on position \( x \) in the direction of \( j \) is written in terms of the phonon operator \( \hat{b}^\dagger_q(x) \), taken to be a dispersionless Einstein mode. The electron-phonon coupling \( \hat{H}_{\text{e-ph}} \) consists of two contributions:

\[ \hat{H}_{\text{e-ph}}^{(1)} = \alpha^{(1)} \sum_{x,j} (\hat{u}_j(x) - \hat{u}_j(x + a_j)), \]

\[ \times (\hat{c}^\dagger_j(x)\hat{c}_{j}(x + a_j) + \hat{c}^\dagger_j(x + a_j)\hat{c}_j(x)), \]

\[ \hat{H}_{\text{e-ph}}^{(2)} = \alpha^{(2)} \sum_{x,j} (\hat{u}_j(x + a_j) - \hat{u}_j(x - a_j))\hat{c}^\dagger_j(x)\hat{c}_j(x). \]

The first type of electron-phonon coupling affects the kinetic energy of electrons, by increasing the hopping amplitude if the interatomic distance decreases. The second process lowers the electronic potential energy in regions of increased ionic density.

The full Hamiltonian can be diagonalized in the mean field approximation, using ansatz averages that reflect the ordered states found in the Landau theory analysis:

\[ \langle \hat{c}^\dagger_j(x)\hat{c}_j(x) \rangle = \rho_0 + A \cos(Q \cdot x + \varphi_j), \]

\[ \langle \hat{c}^\dagger_j(x)\hat{c}_{j}(x + a_j) \rangle = \sigma_0 + B \cos(Q \cdot (x + a_j)/2 + \chi_j), \]

\[ \langle \hat{u}_j(x) \rangle = \tilde{u} \sin(Q \cdot x + \phi_j). \]
Course graining of the mean fields appearing in the Landau theory can in principle be obtained through diagonalization by introducing a new set of operators defined as \( \hat{b}_j \) and \( \hat{\gamma}_j \) for the density waves. The phases are \( \alpha \) and amplitudes are determined self-consistently, by requiring the expectation values of the displacement operators \( \hat{u}_j \) can then be computed in the diagonal basis, which yields the relation between atomic displacements and the electronic order parameters: \( \bar{\gamma}_j(q) = b_j(q) + v_{q,j} \delta(q - Q) + w_{q,j} \delta(q + Q) \) [25]. Requiring this transformation to bring the bosonic Hamiltonian into diagonal form determines the values of \( v_{q,j} \) and \( w_{q,j} \). The expectation values of the of the displacement operators \( \hat{u}_j \) can then be computed in the diagonal basis, which yields the relation between atomic displacements and the electronic order parameters: \( \bar{u} = 2\sqrt{3}/\hbar_0(2B\alpha^{(1)}_j \phi_j^2 - \phi_j^2) - \alpha^{(2)}_j \phi_j^2 \). Demanding the displacement to be real restricts the phase differences to be integer multiples of \( \pi \).

The fermionic part of the Hamiltonian can be diagonalized numerically. The values of the order parameter phases and amplitudes are determined self-consistently, by requiring that the expectation values computed using particular ansatz averages match the ansatz values. With zero on-site Coulomb interaction, the orbital sectors independently develop charge density waves. The phases are \( \phi_j = n_j \pi/3 \), with \( n_j \) integer, which includes both nonchiral solutions in which the \( n_j \) are all equal, and chiral ones. For any nonzero value of the Coulomb interaction, this degeneracy is lifted, and the left- and right-handed chiral charge ordered configurations with \( n_j - n_{j+1} = \pm 2\pi/3 \) become the lowest energy states.

For each handedness, the \( n_j \) may be odd or even multiples of \( \pi/3 \). These solutions correspond to either bond-centered or site-centered charge density waves, as indicated in the insets of Fig. 3. The bond-centered solution dominates for large \( \alpha^{(1)}/\alpha^{(2)} \), while the site-centered one is consistent with the opposite regime. The atomic structure observed in elemental Se and Te corresponds to the bond-centered charge order [14], with \( \alpha^{(1)} \) prevailing.

Within the chiral phase, the short bonds in the three orbital chain directions connect to form a spiral, in agreement with the experimentally observed structure shown schematically in Fig. 1(a). The displacements in the \( x, y, \) and \( z \) directions arise from charge order in chains of \( px, py, \) and \( pz \) orbitals, respectively. The modulation of charge density can thus also be seen as a spatial modulation of orbital occupation, as shown explicitly in Fig. 1(b). The emergence of orbital order in conjunction with chiral charge order is inevitable, since both arise from the same relative phase shifts between charge density waves in distinct orbital sectors. The presence of an interaction between charge density wave components leading to relative phase shifts, can thus be interpreted as a generic route to combined charge and orbital order, which should be applicable to a wide range of materials hosting multicomponent charge order.

V. DISCUSSION

Although a phase of combined charge and orbital order has been proposed to exist in the low-temperature phase of \( IT\text{-TiSe}_2 \) [1,3,4], the broken inversion symmetry is yet to be observed directly. In addition, the interplay between the great number of atoms within the unit cell of \( \text{TiSe}_2 \) complicates the extraction of physical insight from microscopic models [26,27]. Having a model material, which harbors a similar charge and orbital ordered state but is structurally simple and well understood, is therefore crucial to aid in building a general understanding of this novel state of matter, and in allowing the identification of related novel types of order in other materials. We argue that the elemental chalcogens tellurium and selenium constitute precisely such model materials.

In both the elemental chalcogens and \( IT\text{-TiSe}_2 \) multiple density wave instabilities occur in distinct orbital sectors. These give rise to multiple, differently polarized, displacement waves in both materials. The on-site Coulomb repulsion then causes maxima of different density waves to repel each other. This results in shifting of the density waves, thus breaking inversion symmetry and yielding a chiral crystal structure (as long as no mirror symmetries in the parent lattice reduce the chiral order to polar [8–10]). The density waves originating in distinct orbital sectors, necessarily implies that orbital order accompanies the charge modulations, creating a combined charge and orbital ordered phase.

In contrast to the chalcogens, the propagation vectors for different density waves in \( \text{TiSe}_2 \) are all distinct. The order is also site centered in Se and Te, but bond centered in \( \text{TiSe}_2 \). Finally, the on-site Coulomb repulsion coupling different density waves, yields only an indirect interaction between bond-centered charges in the case of \( \text{TiSe}_2 \). Despite these and other differences, including for example the different driving mechanisms underlying the density wave formation [25,28], a common general mechanism for combining charge and orbital order is identified: as long as multiple density wave instabilities occur in distinct sections of the Fermi surface, which correspond to distinct orbital textures, any local interaction between orbital sections will cause the combination of charge and orbital order into a single ordered phase. We thus predict that a type of combined order can be found in any charge...
ordered material involving multiple orbital sectors, including (families of) materials with much more complicated structures than the elemental chalcogens. The theoretical understanding developed here for Se and Te can be used as a guiding principle in looking for such novel states of matter.

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