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Search for Large Extra Dimensions in Dielectron and Diphoton Production


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We report a search for effects of large extra spatial dimensions in $p\bar{p}$ collisions at a center-of-mass energy of 1.8 TeV with the D0 detector, using events containing a pair of electrons or photons. The data are in good agreement with the expected background and do not exhibit evidence for large extra dimensions. We set the most restrictive lower limits to date, at the 95% C.L. on the effective Planck scale between 1.0 and 1.4 TeV for several formalisms and numbers of extra dimensions.

The possibility that the universe has more than three spatial dimensions has long been discussed [1]. Recent developments in string theory suggest that there could be up to seven additional spatial dimensions, compactified at very small distances, on the order of $10^{-32}$ m. In a new model [2], possibly realizable within string theory [3], several of the compactified extra dimensions (ED) are suggested to be as large as 1 mm. These large ED are introduced to solve the hierarchy problem of the standard model (SM) by effectively lowering the Planck scale to the TeV energy range. The ED compactification radius, $R$, depends on the number of extra dimensions ($n$), on the fundamental higher dimensional (i.e., effective) Planck scale, $M_{S}$, and on the apparent four-dimensional Planck mass, $M_{Pl} = 1/\sqrt{G_{N}}$: $R \approx \frac{1}{n} \left( \frac{M_{Pl}}{M_{S}} \right)^{2/n}$ [2]. Since Newton’s law of gravity would be modified in the presence of compactified extra dimensions for interaction distances below the size of the largest extra dimension, current gravitational observations rule out the case of a single large extra dimension. Recent preliminary results from gravity experiments at submillimeter distances [4], as well as cosmological constraints from supernova cooling and distortion of cosmic diffuse gamma radiation [5], indicate that the case of $n = 2$ is likely ruled out as well. However, for $n \geq 3$, the size of the ED becomes microscopic and therefore eludes the reach of direct gravitational measurements. Cosmological constraints are also weak in this case. Therefore, high energy colliders, capable of probing very short distances, are crucial to test theories of large ED. In these theories, the effect of gravity is enhanced at high energies due to the accessibility of numerous excited states of the graviton (referred to as Kaluza-Klein [6] modes of graviton, $G_{KK}$), corresponding to extra degrees of freedom arising from the compactified dimensions. Since gravitons couple to the energy-momentum tensor, they can be produced in any SM process.

Large ED phenomenology at colliders has been studied in some detail [7–10]. One of the primary observable effects would be an apparent nonconservation of momentum caused by the direct emission of gravitons, which leave the three flat spatial dimensions. A typical signature would be the production of a single jet or vector boson at large transverse momentum. The other observable effect would be anomalous diferonn or diboson production at large invariant masses ($M$) via virtual graviton exchange. Direct graviton emission is expected to be suppressed by a factor $(M/M_{S^{n+2}}$, while virtual graviton effects depend only weakly on the number of extra dimensions [7,10]. Virtual graviton production therefore offers a potentially more sensitive way to search for manifestations of ED [11].

Both of the above effects have been sought at LEP [12], with lower limits on the effective Planck scale set on the order of 1 TeV. Virtual $G_{KK}$ exchange was also sought at HERA [13], but with less stringent limits. In this Letter, we report the results of the first specific search for ED at a hadron collider. The data are from the D0 experiment [14] at the Fermilab Tevatron, using proton-antiproton collisions at a center-of-mass energy of 1.8 TeV, and final states containing pairs of electrons or photons. We analyze the differential distribution [15] of dielectron or diphoton events in terms of their invariant mass and the scattering angle in their center-of-mass frame ($\theta^{*}$). These two variables completely define the $2 \rightarrow 2$ scattering processes, and their properties provide improved separation between the SM contributions and the effects of virtual graviton exchange when compared to analyses of invariant-mass distributions alone [15].

Several modifications to the method of Ref. [15] are introduced to optimize the sensitivity of the search. First, because the efficiency and resolution for high-energy electromagnetic (EM) objects at D0 are superior to those for muons, we use only the dielectron and diphoton channels, with the sensitivity to ED coming primarily from the diphoton events [15]. Second, to optimize the efficiency for diphoton and dielectron selection, we eliminate tracking requirements from electron and photon identification, thereby effectively combining them. Ignoring tracking information maximizes both the dielectron and diphoton efficiencies since neither electrons with unreconstructed tracks nor photons with matching tracks from conversion or random overlap are lost. In what follows, we refer to electron or photon objects that do not use tracking information as EM objects. The dominant background at high mass, where ED effects are enhanced, is the irreducible SM background from direct diphoton production, rather than instrumental background from misidentification of jets as EM objects (see, e.g., Ref. [16]). Increasing the di-EM identification...
checks on possible variation of selection ef-
We choose this EM-object-based vertex-
we choose the one that best matches these directions.
most probable direction of each EM object using the trans-
energy, pseudorapidity, and invariant mass indicate that the
data. The MC includes SM contributions (19) distributions to esti-
ation, and the effect of different parton distributions. We
resolution for the EM objects, smearing and misidenti-
same way.
presence of two EM objects and corresponds to an inte-
to, (i) deposit more than 95% of its energy in
the EM calorimeter; (ii) have an energy deposition pattern
consistent with that expected for an EM object; and
(iii) be isolated [17]. The ef-
ficiency for these selections, based on $Z \rightarrow ee$ events, is $87 \pm 2\%$ per EM object.
We also require (i) two EM objects with $E_T > 45$ GeV
in good fiducial regions of the detector ($|\eta_d| < 1.1$ or
$1.5 < |\eta_d| < 2.5$, where $\eta_d$ is the pseudorapidity relative
to the center of the detector [18]); (ii) missing transverse
energy $E_T < 25$ GeV; and (iii) no additional EM objects
with $E_T > 5$ GeV. The trigger is fully efficient for this set
of offline selections. The above selections, with the overall
efficiency of $79 \pm 2\%$ per event, define our base sample,
which contains 1282 candidate events. The error on the
efficiency includes uncertainties due to the statistics and
background parametrization in the $Z \rightarrow ee$ sample. Other
checks on possible variation of selection efficiency with
energy, pseudorapidity, and invariant mass indicate that the
efficiency is constant within uncertainties.

To determine the hard-scattering vertex, we calculate the
most probable direction of each EM object using the trans-
verse and longitudinal segmentation of the EM calorimeter
[17]. Among all the reconstructed vertices in the event,
we choose the one that best matches these directions.
We choose this EM-object-based vertex-finding technique
since it treats the diphoton and dielectron events in the
same way.

We model the effects of ED via the parton-level leading-
order (LO) Monte Carlo (MC) generator of Ref. [15], aug-
mented with a parametric simulation of the D0 detector.
The simulation takes into account detector acceptance and
resolution for the EM objects, smearing and misidentifi-
cation of the primary interaction vertex, initial state radia-
tion, and the effect of different parton distributions. We
used leading order CTEQ4LO [19] distributions to es-
imate the nominal prediction. The parameters of the detector
model are tuned using independent samples of collider
data. The MC includes SM contributions ($Z/\gamma'$ and direct
diphoton production), Kaluza-Klein graviton exchange di-
agrams, and their interference in di-EM object production.

Since the parton-level generator involves only the
$2 \rightarrow 2$ hard-scattering process, we model next-to-leading
order (NLO) effects by adding a transverse momentum
to the di-EM system, based on the transverse momentum
spectrum of di-EM events observed in the data. In the
presence of the NLO corrections, the scattering angle $\theta^*$
is defined in the di-EM helicity frame, i.e., relative to the
direction of the boost of the di-EM system. Since the
parton-level cross section is calculated at LO, we account
for NLO effects in the SM background by scaling the
cross sections by a constant $K$ factor of 1.3 [20]. We
assign a 10\% systematic uncertainty on the value of the $K$
factor to account for a possible growth of the $K$ factor at
high mass. Because NLO corrections to the Kaluza-Klein
diagrams have not yet been calculated, we use the same
constant $K$-factor for the signal. The $K$ factor for graviton
exchange is expected to grow with invariant mass, similar
to that for $Z/\gamma'$ exchange; consequently, our assump-
tion tends to underestimate the ED contribution at high
invariant mass and is conservative.

The main SM sources of di-EM production, Drell-Yan
and direct diphoton production, are included in the MC. Other SM
sources, such as $W_Z$, $Z_W$, $WW$, $Z \rightarrow \tau \tau$, and $t\bar{t}$
production are negligible, as most have small cross section
compared to that for Drell-Yan and direct diphoton pro-
pesses, and are reduced further by the requirements of low
$E_T$ and exactly two EM objects in the event. The only non-
negligible sources of background arise from single-photon and
dijet events, in which jet(s) are misidentified as EM ob-
jects. We estimate this background using a separate data
sample collected with a trigger requiring a single EM ob-
ject. We require that a combination of a jet and the EM
object satisfy all the requirements for signal, except that the
EM ID requirements are applied only to the EM object. We
obtain the instrumental background by scaling the number
of jet-EM combinations passing our selection criteria by
the probability of a jet to be mistaken as an EM object,
measured to be $0.18 \pm 0.04\%$, and, within the uncertain-
ties, being independent of $E_T$ or $\eta$ [21]. The other source
of instrumental background, $W +$ jets production with a jet
misidentified as an EM object, is negligible due to the re-
quirement of low $E_T$. The total instrumental background
in the base sample is $87 \pm 22$ events, which is less than
$7\%$ of the dominant SM background. The sum of the SM
background and the instrumental background reproduces
the main kinematic characteristics of the base sample, as
illustrated in Fig. 1.

The cross section in the presence of large ED is given by
[7,8,10]:

$$\frac{d^2\sigma}{dM d \cos \theta^*} = f_{SM} + f_{int} \eta_G + f_{KK} \eta^2_G, \quad (1)$$

where $f_{SM}$, $f_{int}$, and $f_{KK}$ are functions of ($M, \cos \theta^*$)
and denote the SM, interference, and $G_{KK}$ terms. The
effects of ED are parametrized via a single variable $\eta_G = \mathcal{F} / M^4_t$,
where $\mathcal{F}$ is a dimensionless parameter of order
unity, reflecting the dependence of virtual $G_{KK}$ exchange
on the number of extra dimensions. Different formalisms
use different definitions for $\mathcal{F}$:

$$\mathcal{F} = 1, \quad (GRW \ [7]); \quad (2)$$

$$\mathcal{F} = \begin{cases} 
\log \left( \frac{M^2_t}{M^2} \right), & n = 2, \quad (HLZ \ [8]); \\
\frac{2}{n-2}, & n > 2.
\end{cases} \quad (3)$$
\[ F = \frac{2\lambda}{\pi} = \pm \frac{2}{\pi}, \quad \text{(Hewett [10])}. \tag{4} \]

Note that only within the HLZ formalism \( F \) depends explicitly on \( n \).

Figure 2 shows the two-dimensional distribution in \( M \) and \( |\cos\theta^*| \) for data in (a), total background in (b), and the sum of the background and the ED signal for \( \eta_G = 1 \text{ TeV}^{-4} \) in (c). The data agree well with the background prediction and do not exhibit evidence for large \( M \), which would produce an excess of events at large mass and small values of \( |\cos\theta^*| \) [see Fig. 2(c)]. A comparison of Figs. 2(b) and 2(c) shows that use of both \( M \) and \( \cos\theta^* \) can improve the sensitivity over just one variable. The projection on \( \cos\theta^* \) for high mass values in Fig. 2(d) shows the relative increase of signal over SM processes at small \( |\cos\theta^*| \). The use of both variables allows about 10% improvement in the sensitivity to \( \eta_G \).

In the absence of evidence for ED, we set limits on the effective Planck scale. We perform a Bayesian fit of the sum of the cross section given by Eq. (1) and the instrumental background, to the data in the entire \((M, |\cos\theta^*|)\) plane shown in Fig. 2, with \( \eta_G \) as free parameter, with an assumed uniform prior distribution. The systematic uncertainties on signal and background in the fit include systematics of the \( K \) factor (10%), choice of parton distribution functions (5%), integrated luminosity (4%), EM ID efficiency (3%), and the uncertainty on the instrumental background (25%).

The best estimate of the parameter \( \eta_G \) is consistent with the SM value of \( \eta_G = 0 \), and the one-sided 95% C.L. limits on \( \eta_G \) are

\[ \eta_G < 0.46 \text{ TeV}^{-4} \quad (\eta_G \geq 0), \tag{5} \]

\[ \eta_G > -0.60 \text{ TeV}^{-4} \quad (\eta_G \leq 0), \tag{6} \]

in good agreement with the expected sensitivity to \( \eta_G \), as obtained in an ensemble of MC trials (0.44 TeV\(^{-4}\) for \( \eta_G > 0 \)). We can express these results in terms of limits on the effective Planck scale for the three formalisms of Eqs. (2)–(4). In the formalism of Ref. [10], both signs of \( \eta_G \) are possible and therefore both limits (5) and (6) are relevant. In the other two formalisms, \( \eta_G \) is always positive, and only the first limit is relevant. For the HLZ formalism, the case of \( n = 2 \) is special since \( F \), and therefore \( \eta_G \), depend on \( M \). To relate \( \eta_G \) to \( M_S \) for \( n = 2 \), we use an average \( M \) for the \( G_{KK} \) term at the Tevatron of \((0.6 \text{ TeV})^2 \) [15]. Limits for different formalisms and for different numbers of extra dimensions are given in Table I. They correspond to the ED compactification radius ranging from \( R < 0.3 \text{ mm} \) (\( n = 2 \)) to \( R < 2 \text{ fm} \) (\( n = 7 \)).

In summary, we have performed the first search for large extra spatial dimensions at hadron colliders by looking for effects of virtual Kaluza-Klein gravitons in the production of dielectrons and diphotons at high energies. No evidence is found for large ED, and we set a model-independent 95% C.L. upper limit of 0.46 TeV\(^{-4}\) on the parameter \( \eta_G \).
that describes the size of the ED contribution. This result corresponds to limits on the effective Planck scale ranging between 1.0 and 1.4 TeV for several formalisms and numbers of ED. These are the most restrictive limits on large ED to date, and are complementary to analogous limits from LEP that probe a different range of invariant masses.

Note that our results can also be treated as model-independent limits on the general class of dimension eight $T_{\mu\nu}T^{\mu\nu}/M_p^3$ operators, where $T_{\mu\nu}$ is the conventional four-dimensional energy-momentum tensor. Operators of this kind are found in a variety of models based on conformal field theory and strings (see, e.g., [22]).

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[11] Strictly speaking, virtual graviton effects are sensitive to the ultraviolet cutoff required to keep the divergent sum of Kaluza-Klein states finite [7,8,10]. This cutoff is expected to be of the order of the effective Planck scale. Dependence on the value of the cutoff is discussed, e.g., in Refs. [15,23].
[18] Pseudorapidity $\eta_\ell$ is defined as $-\ln(\tan\theta/2)$, where $\theta_\ell$ is the polar angle with respect to the proton beam, as measured relative to the center of the detector.

Table I. Lower limits at 95% C.L. on the effective Planck scale, $M_5$, in TeV.

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