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Precisely Tailoring Upconversion Dynamics via Energy Migration in Core–Shell Nanostructures

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Abstract: Upconversion emission dynamics have long been believed to be determined by the activator and its interaction with neighboring sensitizers. Herein this assumption is, however, shown to be invalid for nanostructures. We demonstrate that excitation energy migration greatly affects upconversion emission dynamics. “Dopant ions’ spatial separation” nanostructures are designed as model systems and the intimate link between the random nature of energy migration and upconversion emission time behavior is unraveled by theoretical modelling and confirmed spectroscopically. Based on this new fundamental insight, we have successfully realized fine control of upconversion emission time behavior (either rise or decay process) by tuning the energy migration paths in various specifically designed nanostructures. This result is significant for applications of this type of materials in super resolution spectroscopy, high-density data storage, anti-counterfeiting, and biological imaging.

Lanthanide (Ln) ions doped upconversion (UC) materials have attracted extensive attention owing to their potential applications in a large variety of fields, for example, display,[1] biology,[2] anti-counterfeiting,[3] and solar cells.[4] Recent advances in upconversion emission applications rely more and more on precisely engineering the material structure on nanometer scale, in which energy migration behavior plays a key role. So far, most of the related investigations have focused on steady-state effects of energy migration.[5] That is, excited states could randomly “walk”, thus making the long-range energy transport (more than several nanometers) between Ln ions possible. However, the lack of a detailed microscopic picture of UC mechanism in such structures, especially the ambiguous recognition of the energy migration dynamics and its effect on the time behavior of UC emission, remains a major bottleneck for further development of these structures.

Based on traditional understanding, energy migration among sensitizers is always neglected in UC dynamics treatment. In that case, although UC emission involves complex interactions between multiple Ln ions, its dynamics are solely determined by the energy transfer processes within one activator–sensitizer pair (Figure 1a).[6] Despite some discussion of this assumption,[7] there is no direct evidence of the effect of the sensitizer-to-sensitizer energy migration on UC dynamics.

This unsatisfactory situation is primarily due to the lack of appropriate material systems. Indeed, even in the traditional sensitizer–activator co-doping materials, energy migration still plays an important role in their UC process.[7] However, its effect is difficult to evaluate as the energy transfer process is indistinguishable from energy transfer between sensitizers and activators. In our opinion, this dilemma can be solved by introducing a smart nanosystem termed “dopant ions’ spatial separation” (DISS) nanostructure.[8] By locating sensitizers and activators into different regions of a single nanoparticle, the three basic processes of UC (i.e. light absorption, energy migration, and UC emission) can be well separated (Figure 1b). Therefore, the role of energy migration in UC dynamics could be quantitatively determined by 1) tuning the migration layer thickness or 2) varying the dopant concentration of the ions as energy carriers in the migration layer. Secondly, traditional treatment of the UC mechanism based on simultaneous rate equations is not able to describe properly the actual energy migration processes, because it provides only an averaged macroscopic statistical result. Therefore, we set up a time-correlated Monte Carlo simulation model to fill the gap (Figure S1–S7, Table S1–S3, in the Supporting Information). Compared with previous approaches,[1,5,7,8] our model is capable of monitoring UC dynamics from the level of multiple ion-to-ion interactions,
yielding a more comprehensive vision without restriction of the distribution of Ln³⁺ ions in the nanostructure.

Combining the experimental and simulation results, we find that despite each step of the energy migration (e.g. \( \text{Yb}^{3+} \rightarrow \text{Yb}^{3+} \)) only costs a relatively short time (around microseconds), its randomly walking nature significantly prolongs the migration time up to several hundred microseconds. This result tells us that the conventional understanding of the UC mechanisms is incomplete and in some nanostructures it might be misleading by ignoring the actual energy-migration process. More importantly, by precisely tuning energy migration paths in a DISS nanostructure, either the rise or decay of UC emission dynamic traces can be quantitatively tailored in a wide range. Our study thus not only offers a new distinct microscopic picture on how energy migration affects UC and the temporal characteristics of UC emission, but is also an excellent starting point for the rational design of novel functional and efficient UC nanostructures.

We prepared a DISS nanostructure: \( \text{YbEr}@\text{Yb}@\text{YbNd} \) (short for \( \text{NaYF}_4:20\%\text{Yb},2\%\text{Er} \oplus \text{NaYF}_4:20\%\text{Yb} \oplus \text{NaYF}_4:10\%\text{Nd},20\%\text{Yb} \)) to explore the temporal effect of energy migration. The pure hexagonal phase of the nanomaterials was confirmed by the X-ray powder diffraction (Figure S8). The uniformed sizes (or thickness) of the core and each shell were verified to be approximately 25.0 nm, 2.8 nm, 3.0 nm, respectively (Figure S9). To perform spectroscopic studies, a binary pulse excitation setup was designed, in which the interval \( \Delta t \) between a 800 nm and a 980 nm nanosecond pulse is tunable from ~200 µs to 1000 µs (Figure 2a). Typically, 980 nm laser excites \( \text{Yb}^{3+} \) ions all over the whole nanoparticle, whereas 800 nm only excites \( \text{Nd}^{3+} \) ions located in the outer layer (Figure 2b). The UC emission comes from 1) 980 nm excitation, 2) 800 nm excitation, and 3) 980 nm and 800 nm co-excitation. Clearly, the first two parts (1 and 2) are \( \Delta t \) independent, and part (3) depends on \( \Delta t \). Strikingly, the strongest UC emission occurs not at \( \Delta t = 0 \), but at the point when the 980 nm pulse is approximately 200 µs behind the 800 nm pulse (Figure 2c). Furthermore, since the time constant of \( \text{Nd}^{3+} \rightarrow \text{Yb}^{3+} \) energy transfer is very short (ca. 20 µs, Figure S10), the approximately 200 µs time delay should be mainly attributed to the additional \( \text{Yb}^{3+} \rightarrow \text{Yb}^{3+} \) energy migration time under the 800 nm excitation. Significance of this result is that: 1) it is the first acquisition in real time of energy migration contribution to UC emission, and 2) it confirms that the energy migration temporal effect is non-negligible. This result is consistent with the prediction of Monte Carlo simulation. In \( \text{NaYF}_4:20\%\text{Yb} \) sublattice, despite each step of \( \text{Yb}^{3+} \rightarrow \text{Yb}^{3+} \) energy migration only takes about 1.7 µs (Figure S5), and even if the migration distance is just a few nanometers, the time delay for UC will still be several hundred microseconds because 1) the random walk nature of energy migration, and 2) the non-linear nature of UC emission (Figure S11). To validate this assumption, two types of DISS nanostructure are designed to tune the rise and decay of UC dynamics, respectively.

To quantitatively tune the rise of the time trace of UC emission, the DISS nanostructure characterized with spatially separated absorption (sensitizer) and emission (activator) regions is suggested. As an example, the 800 nm excited \( \text{YbEr}@\text{Yb}@\text{Nd} \) core/shell/shell nanostructure (short for \( \text{NaYF}_4:20\%\text{Yb},2\%\text{Er} \oplus \text{NaLuF}_4:20\%\text{Yb} \oplus \text{NaYF}_4:20\%\text{Nd} \)) is employed. The UC emission of this structure relies fully on the \( \text{Yb}^{3+} \rightarrow \text{Yb}^{3+} \) energy transfer to transport the absorbed energy from the outer layer to the core area (Figure 3a). The pure hexagonal phase and narrow size distribution of nanoparticle are confirmed by Fourier-transform diffraction pattern and TEM images (Figures 3b and Figure S12). Besides, the core/shell/shell structure is directly confirmed by distinguishing the lighter lanthanide element (Y, brighter parts) and the heavier one (Lu, darker parts) in the TEM image (Figures 3b). Clearly, because of the non-negligible \( \text{Yb}^{3+} \rightarrow \text{Yb}^{3+} \) energy migration time (ca. 1.7 µs per step) in the \( \text{NaLuF}_4:20\%\text{Yb} \) middle layer, the required time of \( \text{Er}^{3+} \) receiving the migrated energy can be well controlled by the layer thickness, appearing as a tunable rise of UC luminescence time trace. Indeed, increasing the
next we turn to tune the decay of time-resolved UC emission in a wide range, which is not only required for practical applications, but also very important for an in-depth understanding of the UC mechanism. Up to now, modulation approach was almost exclusively addressed by changing the depopulation rate of the activator emitting energy level, such as surface modification, Ln\(^{3+}\) dopant concentration manipulation, plasmonic effect and introducing extra energy transfer channels. However, all these approaches usually bring along limitations including 1) a limited tuning range, and 2) UC efficiency reduction. Herein, we propose a novel strategy, that is, utilizing the energy migration process in DISS nanostructures characterized with partly overlapped sensitizer and activator regions.

As shown in Figure 4a,b, the UC processes in this DISS nanostructure can be roughly divided into three parts. In part I, Yb\(^{3+}\) and Er\(^{3+}\) are co-doped, therefore its UC emission dynamics presents as a relatively sharp rise due to the minimal energy migration time (black curve in Figure 4b). On the contrary, for parts II and III, the spatially separated Yb\(^{3+}\) and Er\(^{3+}\) require the excitation energy to migrate with more time to achieve UC emission, which leads to 1) slower rises, and 2) relatively low UC efficiency due to the energy loss in the migration process (the green line and blue line in Figure 4b). Notably, the UC emission time trace is a profile of contributions of all the three parts, thus it makes the decay longer than any of the individual parts (Figure 4b). As indicated in Figure 4c, the UC emission (ca. 540 nm) decay lifetime (the required time for transient emission intensity to decay from its maximum value to its 1/e value) could be tuned over a wide range from 139 to 648 μs by varying the active shell thickness from 0 to 15 nm. Meanwhile, the rises change little (from 36 to 144 μs) (Figure S17) because they are all dominated by the common part I. The significance of the results is that 1) it breaks through the limitation of the traditional co-doping system.
Without the assistance of energy migration in the shell, YbEr@Y core/inert shell structures (short for NaYF₄: 20% Yb, 2% Er@NaYF₄), as a control, offer a relatively narrow tuning range (from 139 to 330 μs) via eliminating the surface quenching effect (Figure 4d). 2) It guides us to reconsider the relationship between UC efficiency and decay lifetime. Based on the empirical views developed from sensitizer-activator codoping systems, the longer UC emission decay lifetime is often related to a higher UC efficiency due to the reduction of nonradiative relaxation rate. This view is, however, not always valid to DISS nanostructures. As indicated in Figure 4e, compared with YbEr@Y nanostructure, the YbEr@Yb DISS nanostructure exhibits a relatively long decay lifetime but a relatively low UC efficiency when the shell thickness is over 6 nm. This “abnormal” phenomenon can be explained by 1) the energy migration effect of the active shell (prolonging the decay lifetime) and 2) the efficient energy back-transfer from core to active shell (decreasing the UC efficiency) (Figure S18,19 and Table S5). It is quite gratifying to notice that the “abnormal” phenomenon is well predicted by our Monte Carlo simulations (Figure S20 and Table S6).

In summary, long-standing puzzle of the intimate link between excitation energy migration and UC emission dynamics is unraveled through theoretical modeling and spectroscopic experiments of specifically designed “dopant ions’ spatial separation” nanostructure. The contribution of excitation migration to the UC emission dynamics is demonstrated to be significant as a result of the accumulative effect of multiple-step random walks of the excitation energy. Based on this improved understanding, we propose and demonstrate a convenient and effective approach for tailoring UC dynamics (either the rise or decay) through tuning the excitation energy migration paths in “dopant ions’ spatial separation” nanostructures. Our study paves the way for application-optimized design of novel functional UC nanostructures and even to improve the UC emission efficiency.

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Conflict of interest

The authors declare no conflict of interest.

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