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Resonance enhanced multiphoton ionization spectroscopy of carbonyl sulphide

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Rydberg excited states of the OCS molecule in the energy range 70500–86000 cm$^{-1}$ have been investigated via the two and three photon resonance enhancements they provide in the mass resolved multiphoton ionization (MPI) spectrum of a jet-cooled sample of the parent molecule. Spectral interpretation has been assisted by companion measurements of the kinetic energies of the photoelectrons that accompany the various MPI resonances. The present study supports the earlier conclusions of Weinkauf and Boesl [J. Chem. Phys. 98, 4459 (1993)] regarding five Rydberg origins in the 70500–73000 cm$^{-1}$ energy range, attributable to, respectively, states of $^2\Pi$, $^1\Pi$, $^3\Delta$, $^1\Delta$ and $^1\Sigma^+$ symmetry arising from the $4p\lambda\rightarrow 3\pi$ orbital promotion. We also identify a further 21 Rydberg origins at higher energies. These partition into clumps with quantum defects ca. 3.5 and 4.5, which we associate with the orbital promotions $np\lambda\rightarrow 3\pi$ ($n=5,6$), and others with near integer quantum defect which are interpretable in terms of excitation to $s,d$ and (possibly) $f$ Rydberg orbitals. We also identify MPI resonances attributable to CO($X\Sigma^+$) fragments and to S atoms in both their ground ($^3P$) and excited ($^1D$) electronic states. Analysis of the former resonances confirms that the CO($X\Sigma^+$) fragments resulting from one photon dissociation of OCS at excitation wavelengths ca. 230 nm are formed with a highly inverted, bimodal rotational state population distribution, whilst the latter are consistent with previous reports of the wavelength dependence for forming ground and excited state S atoms in the near uv photolysis of OCS. © 1996 American Institute of Physics.

INTRODUCTION

The spectroscopy$^{1–11}$ and photochemistry$^{7,12–20}$ of the excited electronic states of the 16 valence electron triatomic molecule carbonyl sulphide, OCS have both been the subject of much previous study. It has a linear ground state, with normal mode vibrational frequencies; $^3\nu_1$ ($\sigma^+$ symmetry, largely C–O stretch), 2062 cm$^{-1}$; $^3\nu_2$ ($\pi$, bend), 520 cm$^{-1}$ and $^3\nu_3$ ($\sigma^-$, C–S stretch), 859 cm$^{-1}$, and an electronic configuration

$$\cdots (8\sigma)^2(2\pi)^4(9\sigma)^2(3\pi)^4;X\Sigma^+.$$ (1)

The highest occupied $3\pi$ orbital is largely concentrated on the terminal sulphur atom, but the photoelectron spectrum$^{21–23}$ shows activity in $\nu_3$ and, to a lesser extent, in $\nu_1$, pointing to its having some antibonding character. Photoelectron spectroscopy (PES)$^{21–25}$ and photoionisation studies$^{26}$ give a value of 11.1736±0.0015 eV (90121±12 cm$^{-1}$)$^{26}$ for the ionisation threshold for forming the lower ($^2\Pi/3\Sigma$) spin–orbit state of the OCS$^+$ ion and a spin–orbit splitting (ca. 47 meV, 368 cm$^{-1}$) in good agreement with that deduced from analysis of the OCS$^+$(A–X) laser induced fluorescence (LIF)$^{27–29}$ and the laser induced fragmentation spectra of OCS$^+$. $^{30}$ Frequencies for the ground state ion stretching vibrations are well established: $\nu_3\sim2080$ cm$^{-1}$ and $\nu_3\sim698$ cm$^{-1}$. The frequency of the bending mode, $\nu_2$, has proven harder to define, with reported values ranging from 417 cm$^{-1}$ to 510 cm$^{-1}$. The present work fully supports the lower value.

The lower lying valence excited states of OCS arising from the $4\pi\rightarrow 3\pi$ and/or $10\sigma\rightarrow 3\pi$ electron promotion have been studied by conventional absorption techniques$^{1–3,5,6,8}$ and the broad continua centred around 222, 167, and 152 nm assigned in terms of $^1\Delta\rightarrow X\Sigma^+$, $^1\Pi\rightarrow X\Sigma^+$ and $^1\Sigma^+\rightarrow X\Sigma^+$ transitions, respectively. Studies of the temperature dependence of the longest wavelength absorption profile$^{31}$ led to the conclusion that the $^1\Delta$ excited state has a bent equilibrium geometry. Such a suggestion has been fully vindicated by more recent studies of the energy disposal and vector correlations amongst the primary photodissociation products—ground state CO molecules (carrying little vibrational excitation but with a highly inverted rotational state population distribution) together with excited S($^1D$) atoms.$^{16,17,19}$ The structured absorption centered around 152 nm shows a progression with a ca. 800 cm$^{-1}$ interval.$^{5,9}$ This structure has been attributed to an unstable periodic orbit on the excited $\Sigma^+$ surface;$^{20}$ the anisotropy parameter ($\beta\sim1.8$) measured for the dominant primary products resulting from the 157 nm photolysis of OCS–CO($X\Sigma^+$) molecules (in vibrational levels $v=0$–3, with comparatively modest levels of rotational excitation) and S($^1S$) atoms—serves to confirm

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the parallel character of the photoexcitation process and the very short excited state lifetime.\(^\text{18}\)

Absorption at shorter wavelengths is dominated by transitions to Rydberg states of OCS. The 70000–75000 cm\(^{-1}\) region has been studied extensively by one photon absorption spectroscopy,\(^\text{1–6}\) electron energy loss spectroscopy\(^\text{8}\) and, most recently, by resonance enhanced multiphoton ionisation (REMPI) spectroscopy.\(^\text{1,11}\) Yang \textit{et al.}\(^\text{10}\) identified two-photon resonances associated with three Rydberg states in this region: the \(\tilde{E}\) and \(\tilde{F}\) states, both of which, following Kopp,\(^\text{4}\) they labelled as having \(\Pi\) symmetry, and the so-called \(\tilde{P}\) state. This same region was investigated further by Weinkauf and Boesl,\(^\text{11}\) who combined 2+1 and 3+1 MPI with careful measurements of the kinetic energies of the accompanying photoelectrons. Five Rydberg states were identified, all attributable to the electron promotion \(4p\lambda = \pm 3\pi\), and a number of revisions of earlier analyses\(^\text{3,10}\) were suggested. To date there do not appear to have been any reported multiphoton spectra involving the higher energy Rydberg states of OCS, but previous one photon absorption studies,\(^\text{1–6}\) most notably that of Matsunaga and Watanabe,\(^\text{5}\) have identified higher members of two Rydberg series converging to the two spin–orbit components of the ground state ion.

Here, we report the first comprehensive study of the entire two and three photon resonance enhanced MPI spectrum of a jet-cooled sample of OCS up to ca. 86000 cm\(^{-1}\). Companion high resolution REMPI-PES measurements aid assignment of the numerous observed resonances, enabling confirmation of the conclusions reached by Weinkauf and Boesl\(^\text{11}\) and the identification of a further 21 origins (and much attendant vibronic structure), all of which we associate with Rydberg states built on the ground state ion core. The present study of the OCS molecule complements our recent REMPI and REMPI-PES studies of the isoelectronic molecules \(\text{CO}_2\)\(^\text{32}\) and \(\text{CS}_2\)\(^\text{33}\) but, because of its lower symmetry and the absence of the \(u/g\) parity restrictions, assignment of the numerous OCS resonances in terms of excitations to specific Rydberg orbitals is necessarily less definitive.

\textbf{EXPERIMENT}

The results reported here were obtained using two complementary experimental systems: In Bristol, a home-built time-of-flight (TOF) mass spectrometer was used to record mass resolved REMPI spectra of jet-cooled OCS, and in Amsterdam REMPI-PES studies were performed using a “magnetic bottle” spectrometer. The latter experiment has been described in detail previously,\(^\text{32–35}\) and only a cursory description is given here, but the Bristol TOF mass spectrometer is newly constructed and replaces our previous REMPI-TOF experiment, so is described in more detail.

The mass-resolved REMPI spectra [monitoring \(m/z\) 60 (OCS), 32 (S) and 28 (CO)] were recorded using a newly constructed (Wiley–McLaren configuration\(^\text{36}\)) linear TOF mass spectrometer. The spectrometer consists of two acceleration stages and a field-free flight tube, and is housed within a high vacuum system consisting of two differentially pumped chambers connected by a small hole. The first chamber is pumped by a 150 mm diameter oil diffusion pump (Edwards Diffstak; chamber base pressure ca. \(2 \times 10^{-7}\) Torr) and contains a repeller and an extractor plate (diameters 100 and 125 mm, respectively), separated by a distance of 26 mm and, under typical working conditions, held at +400 and \(-400\) V, respectively. A 25 mm diameter hole in the extractor plate, covered by 85% transmission nickel mesh (Buckbee–Mears) allows passage of any ions formed by REMPI between the extractor and repeller plates into the second acceleration region and the flight tube. The plate dividing the two vacuum chambers forms the front end of the field-free flight tube. It is 23 mm distant from the extractor plate, and has a 5 mm diameter hole covered with the 85% transmission mesh at its centre to pass the ions. This plate and the flight tube are normally held at \(-1400\) V. The second vacuum chamber houses the 600 mm long flight tube, which is constructed of perforated stainless steel sheet. This chamber is pumped by a turbomolecular pump (Edwards EXT 501), and reaches a base pressure of ca. \(7 \times 10^{-8}\) Torr, which rises to \(1 \times 10^{-6}\) Torr when the experiment is operational. Ions exiting the flight tube through a 25 mm diameter hole covered with the same 85% transmission nickel mesh are detected by a pair of 25 mm diameter, chevron-configuration microchannel plates (Galileo).

A pulsed nozzle (General Valve Series 9) located between the repeller and extractor plates injects the gas of interest (here OCS, seeded 10% in Ar) into the first chamber. A series of suitably biased thin steel strips located close beneath the nozzle shield the electric field within the ionization region from the perturbing effects of the metal nozzle, but are small enough not to affect the gas expansion. The voltages on the steel strips are determined by a chain of matched resistors and maintain the linearity of the electric field lines in the extraction region. The gas pulse is intersected by a focused laser beam which drives the MPI process. A Nd:YAG pumped dye laser system (Spectra Physics GCR-230-10 and Lumonics HD-500, operating at 10 Hz repetition rate) is used to generate the laser pulses, with the dye laser output frequency-doubled in BBO or KDP crystals when necessary, and focused into the vacuum chamber by a 200 mm focal length quartz lens.

The output signal from the microchannel plates is accumulated by a digital oscilloscope (LeCroy 9361i) and downloaded to a computer (486-PC, 50 MHz) via a GPIB interface for subsequent analysis. Mass resolved MPI spectra are obtained by scanning the dye laser wavelength continuously and monitoring just that fraction of the total ion signal with a TOF appropriate to the mass of interest, averaging typically 5–10 laser shots for each data point.

The REMPI-PES spectra were obtained in Amsterdam using an excimer pumped dye laser system with the laser output being frequency doubled when required, and focused (f.l. 25 mm) into the ionisation region of the magnetic bottle electron spectrometer. An effusive beam of pure OCS was used, and the resolution of the kinetic-energy resolved photoelectron spectra was typically 15 meV. Sulphur atoms were produced in the photolysis of OCS as a by product of the REMPI process. The measured electron kinetic energies

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from known S atom transitions provided a convenient cali-
bration that enabled the OCS photoelectron kinetic energies
to be placed on an absolute scale. In addition, xenon was
used as a calibration gas. Certain wavelength-resolved
REMPI spectra were also recorded in Amsterdam by mass-
resolved ion detection or by measuring the total photoelec-
tron yield as a function of excitation wavelength.

RESULTS AND DISCUSSION

REMPI spectra

Given a value of 90121 cm\(^{-1}\) for the first ionization limit
of OCS\(^{26}\) we can calculate that the long wavelength thresh-
olds for four photon and three photon ionization of this mol-
ecule are, respectively, 444 and 333 nm. We have thus re-
corded REMPI spectra of OCS using linearly polarized laser
radiation throughout the entire wavelength range 444–230
nm. These are shown in Figs. 1–3. As detailed in the figure
captions, these spectra were recorded in one or other, or
both, of two ways—either using a jet-cooled sample, with
time-of-flight mass selection of the resulting ions (in Bristol),
or using a thermal effusive source with mass selected ion, or
kinetic energy selected photoelectron, detection (in Amster-
dam). Inevitably, a number of different laser dyes were used in
recording these spectra. The relative intensities of distant
peaks should therefore be viewed with some caution since
we have not taken any particular care to power normalize
either between dyes or within the tuning range of any one
dye. A few general points merit comment before embarking
on any detailed discussion of these spectra. As Figs. 1–3
show, most of the excited state levels that show strongly in
the 2+1 REMPI spectrum [panel (a) in Figs. 1–3] can be
discerned via 3+1 REMPI spectroscopy [panel (b)] also, and
vice versa, though on occasion their relative intensities show
quite dramatic variations. Figure 4, which shows MPI spec-
tra involving resonances in the 71000–72000 cm\(^{-1}\) region in
much greater detail, provides one rather striking illustration
of this point. The parallels between the 2+1 and 3+1
REMPI spectra of OCS reported here are in marked contrast
to our recent REMPI investigation of the related molecule,
CS\(_{2}\).\(^{33}\) This is simply a reflection of the fact that CS\(_{2}\)
is centrosymmetric, whilst OCS is not. As a result, symmetry
imposes far fewer restrictions on the allowed electronic ex-
citations in the latter molecule. Where comparisons are avail-
able, the wavenumbers of the resonances evident in Figs.
1–3 are in good agreement with those observed in previous
one photon absorption,\(^{1–6}\) electron energy loss\(^{8}\) and/or mul-
tiphonon studies.\(^{10,11}\) This is illustrated in Table I, which pro-
vides a complete listing of the wave numbers of all of the
significant resonances we have identified in the 2+1 and
3+1 REMPI spectra of OCS.

The effects of OCS fragmentation manifest themselves
in these spectra in several ways. Figure 3 highlights the pau-
city and weakness of features in the 2+1 REMPI spectrum for
forming parent ions (\(m/z 60\)) at excitation wavelengths
shorter than ca. 245 nm (\(\lambda > 81500 \text{ cm}^{-1}\)). This approxi-
mates to the long wavelength end of the region of continuous
uv absorption assigned in terms of excitation to the (bent)
\(^1A^\prime\) and (linear) \(^1A^\prime\) Renner–Teller components of a disso-
from the likely due to competitive population loss via predissociation energy. 6, 9, 12, 20

assigning ions with wavelengths. Figure 5 shows a part of the spectrum for formation of S(1 state, 19 provides confirmation of exclusive formation of S(1 1 REMPI spectra of OCS in the energy range 75000–80500 cm⁻¹ and monitoring the wavelength dependence of the m/z 60 signal. (b) was obtained using a thermal effusive source of OCS, again monitoring the parent ion yield save for the expanded view of the high frequency end of this spectrum which was obtained by monitoring the total photoelectron current. As with Fig. 1, the 3+1 REMPI spectrum is a composite obtained by splicing together spectra recorded using a number of different dyes, and the same comments apply. Vibronic features associated with the various Rydberg states identified in Tables I and II are indicated via the combs festooned above the spectra, whilst interfering resonances associated with neutral sulphur atoms in (a) are indicated by the symbol S.

FIG. 2. Linearly polarized (a) 2+1 and (b) 3+1 REMPI spectra of OCS in the energy range 75000–80500 cm⁻¹. (a) was recorded using a jet-cooled sample, the entire tuning range of a single dye (coumarin 500) and monitoring the wavelength dependence of the m/z 60 signal. (b) was obtained using a thermal effusive source of OCS, again monitoring the parent ion yield save for the expanded view of the high frequency end of this spectrum which was obtained by monitoring the total photoelectron current. As with Fig. 1, the 3+1 REMPI spectrum is a composite obtained by splicing together spectra recorded using a number of different dyes, and the same comments apply. Vibronic features associated with the various Rydberg states identified in Tables I and II are indicated via the combs festooned above the spectra, whilst interfering resonances associated with neutral sulphur atoms in (a) are indicated by the symbol S.

removing valence excited state which, in the linear limit, has \(^1\Delta\) symmetry.\(^6,7,31\) the weakness of the REMPI spectrum involving these shorter excitation wavelengths is attributable to this competing population loss channel at the energy of one absorbed photon. Several groups have studied the photochemistry of OCS following uv excitation in the wavelength range 220–250 nm and shown the primary products to be S(\(^1D\)) atoms and CO(X) fragments carrying substantial rotational excitation.\(^6,16,17,19\) Consistent with this, our simultaneous investigation of the REMPI of S atoms suggests exclusive formation of S(\(^1D\)) atoms in this range of excitation wavelengths. Figure 5 shows a part of the spectrum for forming ions with m/z 28 following one color excitation of OCS at wavelengths ca. 230 nm. This spectrum, which is readily assignable in terms of the well documented MPI spectrum of CO(X), resonance enhanced at the two photon energy by the zero-point level of the \(^B\) \(^1\Sigma^+\) state,\(^19\) provides confirmation of the importance of this fragmentation channel under the present experimental conditions and of the highly excited, bimodal rotational state population distribution in the resulting CO(X) photofragments. Note the appearance of the weak O and S branches in this spectrum, associated with the second rank \([T_0(A)]\) component of the two photon transition tensor. In the same vein, we suggest that the weakness of the m/z 60 signal in the 3+1 REMPI spectrum at excitation wavelengths less than ca. 352 nm (3 \(\tilde{\nu} > 85000\) cm⁻¹) is likely due to competitive population loss via predissociation from the \(^1\Pi\) excited state resonant at the two photon energy.\(^6,9,12,20\)

**REMPI-photoelectron spectra**

The available experimental resolution (ca. 15 meV, 120 cm⁻¹) is sufficient to allow us to identify not just the vibrational, but also the spin–orbit, structure of the ions formed in the various REMPI processes. As has been demonstrated many times previously, the analysis of such spectra can provide considerable insight into the electronic and vibronic characteristics of the intermediate state providing the resonance enhancement. Figures 6–8 show a number of representative REMPI-PE spectra. All those selected for display arise as a result of three photon resonant, four photon ionisation processes. Consider first the spectrum shown in Fig. 6(a), recorded at a laser wavelength of 395.3 nm (3 \(\tilde{\nu}=75897\) cm⁻¹). The kinetic energy of the larger, faster peak is entirely consistent with ionisation to the \(v^+=0\) level of the lower \((^2\Pi_{1/2})\) spin–orbit component of the ion, but the presence of a slower peak indicates that there is also some propensity for ionisation to the \(v^+=0\) level of the higher \((^2\Pi_{3/2})\) state of the ion. Figures 6(b) and 6(c), obtained following excitation at wavelengths of 391.4 and 412.2 nm (3 \(\tilde{\nu}=76650\) and 72777 cm⁻¹, respectively), show two further examples in which the final one photon ionization leads to formation of vibrationless \((v^+=0)\) ions in both spin–orbit states. In both these cases, however, the higher \((^2\Pi_{3/2})\) component is observed to make the dominant contribution. Franck–Condon considerations, together with the apparent simplicity of all three of these spectra, lead us to conclude that, in each case, REMPI proceeds via an electronic (Ryd-
Experimental observation: A very similar energy interval occurs, and we assign it as such. Franck–Condon consider energies lower than that predicted for formation of REMPI-PE spectra in which the peaks appear at kinetic energies, suggesting that the excited level of the neutral potassium resonance enhancement at 80508 cm⁻¹. These findings, and those derived from similar analyses of the REMPI-PE spectra accompanying virtually all other significant resonances in the 2+1 and 3+1 REMPI spectra of OCS [including those shown in Figs. 7 and 8 (see below)], are collected together in Table I.

Figures 7 and 8 provide further illustrations of the way in which REMPI-PES can be used to aid in the vibronic assignment of the resonance enhancing states. Figure 7(a), obtained monitoring photoelectrons at an excitation wavelength of 375.9 nm (3τ = 79812 cm⁻¹), provides yet another example of ionisation occurring to the τ⁺ = 0 level of the ground state ion, predominantly the 2Π₁/₂ spin–orbit component in this case. Figures 7(b) and 7(c), in contrast, show REMPI-PE spectra in which the peaks appear at kinetic energies lower than that predicted for formation of τ⁺ = 0 ions, indicating that the accompanying ions are formed with non-zero internal energy. The PE spectrum shown in Fig. 7(b), recorded at 372.6 nm (3τ = 80508 cm⁻¹), exhibits two peaks with kinetic energies of 2.005 and 2.052 eV. These energies imply that the accompanying ions are formed in both the 2Π₃/₂ and 2Π₁/₂ spin–orbit states with, in each case, ca. 690 cm⁻¹ of internal energy. This energy accords well with previous measurements of the ν₁ C–S stretching mode in the ion,²⁴,²⁵ and we assign it as such. Franck–Condon considerations then suggest that the excited level of the neutral providing the resonance enhancement at 80508 cm⁻¹ likely also carries one quantum of ν₁. Such is entirely consistent with experimental observation: A very similar energy interval (696 cm⁻¹) separates the 80508 cm⁻¹ resonance from the 79812 cm⁻¹ band which Fig. 7(a) showed unequivocally to be an electronic origin. Similar arguments operate in the case of the PE spectrum shown in Fig. 7(c), recorded at 366.4 nm (3τ = 81882 cm⁻¹). The dominant feature here is consistent with the formation of ground state OCS⁺ ions with a 2Π₁/₂ core and ca. 2090 cm⁻¹ of vibrational energy—an energy gap that accords well with the interval between the MPI resonance in question and the 79812 cm⁻¹ origin. We can envisage two possible explanations for an energy separation of this size—either one quantum of the excited state τ₁ (C–O stretch) vibration or three quanta of ν₁. Both interpretations have some merit; it may well be that both contribute, and that their energy difference falls inside our available resolution.

Our final three examples (Fig. 8) involve the resonances observed at laser excitation wavelengths of 423.2, 417.9, and 412.5 nm (3τ = 70881, 71786, and 72726 cm⁻¹, respectively). The REMPI-PE spectrum associated with the first of these [Fig. 8(a)] is dominated by a peak whose kinetic energy (0.545 eV) is consistent with formation of ground state ions in the lower 2Π₃/₂ spin–orbit state, suggesting that the 70881 cm⁻¹ resonance involves another electronic origin. Such a result accords with the conclusions of most recent analyses.⁴–⁷,¹⁰,¹¹ Three other peaks are discernible: The next fastest is consistent with formation of 2Π₁/₂, ν⁺ = 0 ions. More interestingly, the two slower peaks (which indicate formation of ions with ca. 810 and 1640 cm⁻¹ of internal energy, respectively) are most readily assigned in terms of 2Π₃/₂ ions carrying, respectively, two and four quanta of ν₂.
bending excitation. Figures 8(b) and 8(c) serve to reinforce this conclusion. The former is dominated by a peak corresponding to formation of OCS$^+$ ions in their $^2\Pi_{3/2}$ spin–orbit state with ca. 810 cm$^{-1}$ of internal energy. This internal energy, and the 905 cm$^{-1}$ interval between the 70881 and 71786 cm$^{-1}$ resonances, both encourage the interpretation that the resonance enhancement is provided by one or more of the Renner–Teller components associated with the $^2\Pi$ vibronic level of the state whose origin appears at 70881 cm$^{-1}$.

The weaker features evident in Fig. 8(b) with kinetic energies of 0.695 and 0.54 eV are most sensibly assigned to ion formation in their $^2\Pi_{3/2}$, $v_2^+ = 0$ and $^2\Pi_{1/2}$, $v_2^+ = 2$ levels, respectively. In the same vein, the structure apparent in the REMPI-PE spectrum shown in Fig. 8(c) is most plausibly interpreted in terms of photoionisation to (in order of decreasing kinetic energy) the $^3\Pi_{3/2}$ and $^3\Pi_{1/2}$, $v_3^+ = 2$ levels and the $^3\Pi_{3/2}$ and $^3\Pi_{1/2}$, $v_3^+ = 4$ levels. Franck–Condon considerations suggest that the 72726 cm$^{-1}$ resonance should be associated with one (or more) of the levels associated with the $^2\delta$ vibronic level of the state originating at 70881 cm$^{-1}$.

All of these assignments lend strong support to the lower ($\sim 417$ cm$^{-1}$) of the values suggested for the $v_2$ bending vibrational frequency in the ground state OCS$^+$ ion.

**Assigning the multiphoton resonances**

Wave numbers ($\tilde{\nu}$) for the maxima of the various origin bands identified in this work are listed in Table II together with their effective quantum number \([n^* = n - \delta]\), where \(n\) is the principal quantum number and \(\delta\) the quantum defect) calculated using the relationship

\[
\tilde{\nu} = E_i - R(n^*)^2,
\]

where \(R\) is the Rydberg constant (109737 cm$^{-1}$). The ionization limit, \(E_i\), in this expression is taken as 90121 or 90500 cm$^{-1}$, according to whether the companion REMPI-PE spectrum identifies the Rydberg state as being associated with a series converging to the $^2\Pi_{3/2}$ or $^2\Pi_{1/2}$ spin–orbit component of the ion. For completeness, we also
The kinetic energy of the faster peak is consistent with ionization to the dominant peak indicates a propensity for forming ions in their ground state, whilst the quantum states of the ion align vertically.

The quantum defects $\delta \sim 1.5 - 1.6$, if we assume that $n=4$ suggest that all are first members of series arising from excitation to Rydberg orbitals with dominant $p$-character. The resonances at 70881 and 71385 cm$^{-1}$, evident in both 2+1 and 3+1 MPI spectroscopy and which the companion PE spectra show to involve, respectively, a $^2\Pi_{3/2}$ and a $^2\Pi_{1/2}$ ion core, have no obvious counterparts in the $^2\Pi_{1/2}$ ion core. As in Fig. 6, the kinetic energy scales have been offset so that the internal energy states of the ion align vertically.

[FIG. 6. Illustrative REMPI-PE spectra of OCS obtained following excitation at (a) $\lambda = 395.3$ nm, where we excite a three photon resonance (3$\nu$=75897 cm$^{-1}$). The kinetic energy of the faster peak is consistent with ionization to the $v^+ = 0$ level of the lower spin–orbit component of the ion, whilst the slower peak indicates some propensity for ionisation to the $v^+ = 0$ level of the higher $^2\Pi_{1/2}$ state also. (b) $\lambda = 391.4$ nm (3$\nu$=76650 cm$^{-1}$); here the dominant peak indicates a propensity for forming ions in their $^2\Pi_{1/2}$ $v^+ = 0$ level. (c) $\lambda = 412.2$ nm (3$\nu$=72777 cm$^{-1}$). Our assignment of this three photon resonance in terms of the $^2\Pi_{3/2}4p\pi;^3\Sigma^+ ← \tilde{X}^1\Sigma^+$ origin transition accords with the appearance of peaks indicating ionization to the $v^+ = 0$ level of both spin–orbit states of the ion. The kinetic energy scales have been offset so that the quantum states of the ion align vertically.]

indicate those resonances that have been identified in previous one photon absorption$^{1-6}$ and electron impact studies.$^8$

The 70000–75000 cm$^{-1}$ region shown in Fig. 1 has been the subject of many previous studies, including one high resolution absorption study$^4$ and two previous multiphoton investigations.$^{10,11}$ The five origins we identify in this region are consistent with the findings of Weinkauf and Boesl.$^{11}$

The companion REMPI-PE spectra allow us to deduce the dominant core configuration of each of these origins. The associated quantum defects ($\delta \sim 1.5 - 1.6$, if we assume that $n=4$), suggest that all are first members of series arising from excitation to Rydberg orbitals with dominant $p$-character. The resonances at 70881 and 71385 cm$^{-1}$, evident in both 2+1 and 3+1 MPI spectroscopy and which the companion PE spectra show to involve, respectively, a $^2\Pi_{3/2}$ and a $^2\Pi_{1/2}$ ion core, have no obvious counterparts in the one photon absorption spectrum.$^{1-6}$

Following Weinkauf and Boesl,$^{11}$ it thus seems plausible to assign these two resonances in terms of the respective excitations $^2\Pi_{3/2}4p\pi;^3\Delta ← \tilde{X}^1\Sigma^+$ and $^2\Pi_{1/2}4p\pi;^1\Delta ← \tilde{X}^1\Sigma^+$. Here, as in our recent REMPI study of the Rydberg states of CS$_2$, we choose to indicate the upper state both in terms of a Hunds case (c) coupling scheme in which we specify the angular momentum of the ion core, $\Omega_c$ (3/2 or 1/2), along with the $n\lambda\sigma$ character of the Rydberg electron, and in terms of the most appropriate $\Delta, S$ label, e.g., $^3\Delta$.

These two origins are in close coincidence with two other resonances, at 70911 and 71432 cm$^{-1}$. These latter resonances appear in both the 2+1 and 3+1 MPI spectra (Fig. 1) and in one photon absorption.$^{1-6}$ The associated REMPI-PE spectra indicate that these are origins, built on $^2\Pi_{3/2}$ and $^2\Pi_{1/2}$ ion cores, respectively. The quantum defects of these features suggest assignment in terms of a $4p\pi$ orbital excitation, whilst their one photon band contours$^8$ and their presence both in one photon absorption and in the multiphoton excitation spectra can be accommodated if we adopt the respective assignments $^2\Pi_{3/2}4p\sigma;^3\Pi ← \tilde{X}^1\Sigma^+$ and $^2\Pi_{1/2}4p\sigma;^1\Pi ← \tilde{X}^1\Sigma^+$. Such an assignment also accords with the conclusions of Weinkauf and Boesl.$^{11}$ It is compatible with the earlier absorption work of Kopp$^4$ (who labelled...
unable to identify any obvious candidate for the core mixing in the et al. 7, the kinetic energy scales in these PE spectra have been offset so that the ~423.2 nm ~20 origin. The kinetic energies of the dominant peaks in (a) indicates the formation of e− = 0 ions in their lower spin–orbit state, consistent with assignment of this three photon resonance as the 4p− π orbital promotion called these respective origins in terms of excitations to the so-called REMPI-PE spectrum feature4. The term values of these 3 ~20 resonances involve the 2v2 and 2v4 vibronic bands of this same Rydberg transition. As in Figs. 6 and 7, the kinetic energy scales in these PE spectra have been offset so that the internal energy states of the ion align vertically.

these respective origins in terms of excitations to the so-called E and F states), but contradicts the earlier multiphoton work of Yang et al.10 Our final origin in this region, at 72777 cm−1, appears in one photon absorption (but was not recognized as an origin in the only high resolution study of this feature6) and is assigned in terms of the only remaining dipole allowed transition associated with the orbital promotion 4p− π − π, viz. the 1Σ+ − X 1Σ+ transition involving an excited state traditionally labelled as the F state.10,11 The associated REMPI-PE spectrum [Fig. 6(c)] shows “core mixing” to be more developed in this particular excited state: it may be relevant to note that our recent REMPI-PES study of the Rydberg states of CS2 revealed a similar high degree of core mixing in the Σ members of the 4p− π Rydberg complex of this molecule also.33 Again, as with CS2, we are unable to identify any obvious candidate for the 4p− π, Σ − X 1Σ+ counterpart of the 72777 cm−1 resonance. Weinkauf and Boesl11 tentatively attributed a very weak feature at 70810 cm−1 in terms of this excitation but, unfortunately, its weakness precluded our being able to confirm or disprove such an origin assignment via REMPI-PES.

Several aspects of this spectral region merit further comment. The singlet–triplet splitting of both the Δ and Π Rydberg states we associate with the 4p− π orbital promotion (504 and 521 cm−1, respectively) are both substantially greater than the spin–orbit splitting of the ground state ion (ca. 368 cm−1). This fact, and the observed core mixing revealed by the REMPI-PE spectra involving these states, can both be rationalized by appreciating that, as in the corresponding transitions of CS2,33 exchange interactions make a significant contribution to the observed splittings in these low n Rydberg states. For example, in the case of the two Π states, we can use the expression

\[ K = \sqrt{\frac{\Delta E}{2}} - A^2, \]

where \( \Delta E \) and \( A \) are, respectively, the observed singlet–triplet splitting (521 cm−1) and the spin–orbit coupling constant appropriate to the ion core (2Δ ~ 368 cm−1), to estimate a value of 185 cm−1 for \( K \), the exchange contribution to the observed splitting. As Table I showed, given the associated REMPI-PES data, it is possible to assign all remaining features in this region in terms of transitions involving excited vibronic levels built upon the various recognized origins. All three excited state vibrational modes are seen to be active in each of the more intense (singlet–singlet) transitions. The Δ − X and, to a lesser extent, the Δ − X transitions are noteworthy in as much that they exhibit more activity in the bending mode, and quite prominent resonances attributable in terms of various of the components of their respective 2[1 hot band transitions. This bending activity, and the change in bending frequency upon electronic excitation implied by the observed red shift of the 2[1 hot bands, suggests that the bending potentials for these excited states have a smaller curvature (around the linear equilibrium configuration) than that of the ground state neutral and the ion. Hopefully future ab initio calculations might be able to shed light on this observation, and also provide a rationale for the surprisingly low (cf. CS233) term values of these 3 Δ and 1 Δ Rydberg states viz a viz the Π states we attribute to the 4pσ− π orbital promotion. We reserve discussion of the quite striking intensity variations shown by some of the resonances upon changing from 2+1 REMPI to 3+1 REMPI (as, for example, is illustrated in Fig. 4) until later.

Likely higher members of the various np− π Rydberg series are indicated in Figs. 2 and 3; the term values of the various origins and proposed assignments are listed in Table II. The companion REMPI-PE spectra show the 80850 and 81204 cm−1 features to involve Rydberg states with, respectively, 2Π1/2 and 2Π1/2 ion cores. Their quantum defects (δ ~ 1.56 assuming n = 5) plus the fact that they are clearly apparent in the one photon absorption spectrum encourages the respective assignments: [2Π1/2]5pσ; Π − X 1Σ+ and [2Π1/2]5pσ; Π − X 1Σ+. The reduced “singlet–triplet
TABLE I. Ion internal energies and vibronic assignments of multiphoton resonances observed in the two and three photon REMPI spectra (0=origin). Dominant photoelectron peaks are indicated in bold and the values in brackets represent the relative intensities of the peaks. Ion internal energies are deduced from the energy difference between the peak of interest in the REMPI-PE spectrum and that for formation of vibrationless \( ^2\Pi_{1/2} \) and \( ^2\Pi_{3/2} \) ions.

<table>
<thead>
<tr>
<th>Wave number of resonance (cm(^{-1}))</th>
<th>Ion internal energy (in cm(^{-1})) deduced from REMPI-PES</th>
<th>Proposed assignment of strongest vibronic feature</th>
</tr>
</thead>
<tbody>
<tr>
<td>70741 0 (10), 810 (1), 1650 (1)</td>
<td>( ^2\Pi_{1/2} ) core 0 (4) ( ^2\Pi_{3/2} ) core 2 (1)</td>
<td></td>
</tr>
<tr>
<td>70805 0 (1), 810 (10)</td>
<td>840 (1) ( ^2\Pi_{1/2} ) core 2 (1) ( ^2\Pi_{3/2} ) core 2 (1)</td>
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<tr>
<td>72534 850 (1), 1600 (10)</td>
<td>1620 (3) ( ^2\Pi_{1/2} ) core 2 (1) ( ^2\Pi_{3/2} ) core 2 (1)</td>
<td></td>
</tr>
<tr>
<td>72726 820 (4), 1630 (10)</td>
<td>810 (2), 1620 (2) ( ^2\Pi_{1/2} ) core 2 (1) ( ^2\Pi_{3/2} ) core 2 (1)</td>
<td></td>
</tr>
<tr>
<td>73175 0 (10)</td>
<td>0 (2) ( ^2\Pi_{1/2} ) core 0 (2) ( ^2\Pi_{3/2} ) core 0 (2)</td>
<td></td>
</tr>
<tr>
<td>73257 70911 ( \text{from REMPI-PES} )</td>
<td>2 (1) ( ^2\Pi_{1/2} ) core 2 (1) ( ^2\Pi_{3/2} ) core 2 (1)</td>
<td></td>
</tr>
<tr>
<td>73341 0 (10) 700 (10)</td>
<td>0 (10) ( ^2\Pi_{1/2} ) core 0 (10) ( ^2\Pi_{3/2} ) core 0 (10)</td>
<td></td>
</tr>
<tr>
<td>73577 0 (10) ( \text{from REMPI-PES} )</td>
<td>2 (1) ( ^2\Pi_{1/2} ) core 2 (1) ( ^2\Pi_{3/2} ) core 2 (1)</td>
<td></td>
</tr>
<tr>
<td>73809 2070 (5), 2900 (10)</td>
<td>2090 (2), 2930 (2) ( ^2\Pi_{1/2} ) core 2 (1) ( ^2\Pi_{3/2} ) core 2 (1)</td>
<td></td>
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</tbody>
</table>
splitting” (ca. 354 cm⁻¹, comparable with the spin–orbit splitting of the ion) and the much “cleaner” REMPI-PE spectrum (see Table I) both serve to illustrate the rapid evolution towards the expected Hund’s case (c) limit with increasing n. The two other origins identified in this region, at 81072 and 81438 cm⁻¹, show a very similar splitting. Analogy with the 4p→π Rydberg complex and the fact that these resonances have not been identified in one photon absorption leads us to propose the respective assignments [2Π₃/2]5p π⁺Δ→X⁺Σ⁺ and [2Π₃/2]5p π⁺Δ→X⁺Σ⁺. We note that such an assignment implies a reversal of the energetic ordering of the Π(npnσ→π) and Δ(npπ→π) states on going from n=4 to n=5. Two resonances attributable to members of the 6p→π complex are discernible, most notably in the 3+1 REMPI spectrum, with respective wave numbers 84627 and 84996 cm⁻¹. Given that these appear in one photon absorption also, it seems reasonable to suggest the respective assignments [2Π₃/2]6p π⁺Π→X⁺Σ⁺ and [2Π₁/2]6p σ⁺Π→X⁺Σ⁺.

We identify a further fifteen electronic origins in the 75000–85000 cm⁻¹ region (Figs. 2 and 3), most of which have no clear counterparts in the one photon absorption spectrum of OCS. We start by considering the pair of features at 78144 and 78486 cm⁻¹ and their attendant vibronic structure (Table I), which dominate the 2+1 REMPI spectrum in this region and have been identified in one photon absorption also.⁵ The companion REMPI-PE spectra prove that these are origins involving excited states with, respectively, a 2Π₃/2 and a 2Π₁/2 ion core; the peak separation (342 cm⁻¹) is actually slightly smaller than the spin–orbit splitting of the ground state ion. Their effective principal quantum numbers (both n²=3.02, calculated relative to the appropriate ionisation limits) are compatible with assignment in terms of excitation to either a 3d or the 5sσ Rydberg orbital. It is likely that the pair of features at 83514 and 83874 cm⁻¹ (which the REMPI-PE spectra confirm to be built on 2Π₃/2 and 2Π₁/2 ion cores, respectively) are the next spin–orbit split member of this Rydberg series, though we recognize that their quantum defect would also be compatible with a 4f→π Rydberg assignment.

We identify a further six origins in the vicinity of the 78144 cm⁻¹ feature, only one of which appears to have been identified in one photon absorption.⁶ REMPI-PES shows the lowest energy resonance (at 75897 cm⁻¹) to involve a 2Π₃/2 core, whereas the next three (at 76650, 76833, and 77292 cm⁻¹) all involve the excited spin–orbit core. Assignment of any one of these three as the “singlet” partner of the 75897 cm⁻¹ feature would imply an abnormally large exchange interaction. The 76650 cm⁻¹ transition, and its associated vibronic features (Table I), shows particularly strongly in the 3+1 REMPI spectrum—behavior very reminiscent of that exhibited by the previously discussed [2Π₁/2]4pσ⁺Π→X⁺Σ⁺ resonance at longer wavelength. REMPI-PE spectroscopy shows the two origins to higher energy (at 79356 and 79812 cm⁻¹) to involve 2Π₃/2 and 2Π₁/2 ion cores, respectively; both show n* values ca. 3.2,

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TABLE I. (Continued).

<table>
<thead>
<tr>
<th>Wave number of resonance (cm⁻¹)</th>
<th>Ion internal energy (in cm⁻¹) deduced from REMPI-PES</th>
<th>Proposed assignment of strongest vibronic feature</th>
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<tr>
<td></td>
<td>2Π₃/2 core</td>
<td>2Π₁/2 core</td>
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<td>810 (4)</td>
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<td>0 (2)</td>
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</tr>
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<td>690 (10)</td>
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<td>2050 (2)</td>
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</tr>
<tr>
<td>84996</td>
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<td>0 (10)</td>
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TABLE II. Wave numbers (ν), effective quantum defects (n*) and proposed assignments for the various origin bands assigned in the 2+1 and 3+1 REMPI spectra. Where no definite assignments are possible we list the spin–orbit component of the ion core and state number so as to allow identification in Figs. 1–3. Bands identified in previous one photon absorption (Refs. 2–5) and electron impact studies (in brackets) (Ref. 8) are also indicated.

\[
\begin{array}{cccc}
\nu (\text{cm}^{-1}) & n^* & \text{One photon Assignment} & \text{Assignment} \\
70881 & 2.39 & [\Pi_{1/2}]4p^2 \pi^1 \Delta \leftrightarrow X^1 \Sigma^+ & \\
70911 & 2.39 & 70913^a & [\Pi_{1/2}]4p^2 \pi^1 \Pi \leftrightarrow X^1 \Sigma^+ & \\
71938 & 2.40 & (71935) & [\Pi_{1/2}]4p^2 \pi^1 \Delta \leftrightarrow X^1 \Sigma^+ & \\
71432 & 2.40 & 71430^b & [\Pi_{1/2}]4p^2 \pi^1 \Pi \leftrightarrow X^1 \Sigma^+ & \\
72277 & 2.49 & 72276^c (72275) & [\Pi_{1/2}]4p^2 \pi^1 \Sigma^* \leftrightarrow X^1 \Sigma^+ & \\
75897 & 2.78 & 75899^d & [\Pi_{1/2}]; \text{State I} & \\
75670 & 2.81 & [\Pi_{1/2}]; \text{State II} & \\
78663 & 2.83 & [\Pi_{1/2}]; \text{State III} & \\
77292 & 2.88 & (77299) & [\Pi_{1/2}]; \text{State IV} & \\
78144 & 3.02 & 78156^e & [\Pi_{1/2}]; \text{State V} & \\
78486 & 3.02 & 78497^f & [\Pi_{1/2}]; \text{State VI} & \\
79356 & 3.19 & [\Pi_{1/2}]; \text{State VII} & \\
79812 & 3.20 & [\Pi_{1/2}]; \text{State VII} & \\
80850 & 3.44 & 80840^g (80849) & [\Pi_{1/2}]5p^2 \pi^1 \Pi \leftrightarrow X^1 \Sigma^+ & \\
81072 & 3.48 & [\Pi_{1/2}]5p^2 \pi^1 \Delta \leftrightarrow X^1 \Sigma^+ & \\
81204 & 3.44 & 81200^h (81203) & [\Pi_{1/2}]5p^2 \pi^1 \Sigma^* \leftrightarrow X^1 \Sigma^+ & \\
81438 & 3.48 & [\Pi_{1/2}]5p^2 \pi^1 \Delta \leftrightarrow X^1 \Sigma^+ & \\
82278 & 3.74 & [\Pi_{1/2}]; \text{State IX} & \\
82775 & 3.77 & [\Pi_{1/2}]; \text{State X} & \\
83148 & 3.86 & [\Pi_{1/2}]; \text{State XI} & \\
83514 & 4.07 & [\Pi_{1/2}]; \text{State XII} & \\
83874 & 4.07 & [\Pi_{1/2}]; \text{State XIII} & \\
84030 & 4.25 & [\Pi_{1/2}]; \text{State XIV} & \\
84411 & 4.25 & [\Pi_{1/2}]; \text{State XV} & \\
84627 & 4.47 & 84620^i (84607) & [\Pi_{1/2}]6p^2 \pi^1 \Pi \leftrightarrow X^1 \Sigma^+ & \\
84996 & 4.47 & 84999^j & [\Pi_{1/2}]6p^2 \pi^1 \Pi \leftrightarrow X^1 \Sigma^+ & \\
\end{array}
\]

\(a\)Reference 4.
\(b\)Reference 5.
\(c\)Reference 2.

in Fig. 4 were obtained by monitoring just the parent ion yield. We concur with previous workers\(^{11}\) that none of these factors can explain the very different relative strengths observed for the \([\Pi_{1/2}]4p^2 \pi^1 \Delta \leftrightarrow X^1 \Sigma^+\) and \([\Pi_{1/2}]4p^2 \pi^1 \Delta \leftrightarrow X^1 \Sigma^+\) resonances (Fig. 4). Weinkauf and Boesl\(^{11}\) suggested an extension of the parity selection rule to argue that the \(\Delta \leftrightarrow X^1 \Sigma^+\) transition should show strongly in two photon excitation, whilst \(\Pi \leftrightarrow X^1 \Sigma^+\) transitions should be favored by one and three photon excitations (consistent with the observations in Fig. 4). Whilst such a propensity is frequently observed to hold for changes in the \(l\) quantum number in atoms and hydride molecules in which the HOMO is well approximated as an atomic orbital,\(^{37}\) it is far from clear why it should apply to changes in the \(\Lambda\) quantum number in a molecule like OCS. Clearly angular momentum constraints exclude one photon \(\Delta \leftrightarrow X^1 \Sigma^+\) transitions in a linear molecule, but we suspect that additional information (e.g., some knowledge of the composition of the various \(l\) functions contributing to the various low lying Rydberg orbitals) will be needed before one can offer a wholly convincing explanation as to why this particular \(\Delta \leftrightarrow X^1 \Sigma^+\) resonance is favored by two photon excitation whilst the neighboring \(\Pi \leftrightarrow X^1 \Sigma^+\) transition appears more strongly in 3+1 REMPI.

### Fragmentation

The mass resolved REMPI spectra showed numerous resonances in the mass/charge (m/z) 32 (S) mass channel and, as Fig. 5 showed, a few notable features in the mass/charge 28 (CO) mass channel also. The structured mass/charge 28 resonance observed at excitation wavelengths ca. 230 nm is readily attributable to MPI of ground state CO molecules, resonance enhanced at the two photon energy by the zero-point level of the \(B^1 \Sigma^+\) state.\(^{19}\) The highly excited, bimodal rotational state population distribution is consistent with previous measurements of the energy disposal in the CO(X) fragments resulting from one uv photon dissociation of OCS at 230 nm\(^{19}\) and 222 nm,\(^{16,17}\) i.e., following excitation within its \(\Delta \leftrightarrow X^1 \Sigma^+\) absorption system. Some signal attributable to S\(^+\) ions was observed at all wavelengths where OCS undergoes 2+1 REMPI; this has been commented on previously,\(^{11}\) and is attributable to one photon excitation and subsequent fragmentation of the ground state OCS\(^+\) ions formed in the initial REMPI process. Of greater interest are the strong resonances appearing at excitation wavelengths less than ca. 260 nm, attributable to REMPI of neutral S atom fragments. We have been able to assign many, though by no means all, of these lines by referring to previous listings of S atom excited state term values.\(^{38,40}\) It is inappropriate to list the wave numbers of all of the observed resonances here, but a complete listing is available from the authors upon request. We see a clear dividing line at 251.5 nm. All assigned resonances at longer wavelength are attributable to two photon transitions originating from the ground \((^3P)\) state, whilst at all shorter wavelengths we only identify two photon resonances associated with the excited \((^1D)\) state. The latter observation is particularly significant, since the region down to
ca. 241 nm should be rich with two photon resonances originating from the ground state and involving triplet members of Rydberg series converging to the ground \(^3S^0\) state of the ion. All of these observations are explicable in terms of known near uv photochemistry of OCS. The ground state atoms observed at the longest wavelengths (\(>251.5\) nm) probably arise as a result of one photon absorption to (and dissociation from) triplet valence states which must contribute in this region of the near uv absorption spectrum of the parent, whilst the observed onset of the \(S(1D)\) resonances marks the long wavelength limit of the much stronger parent \(1\Delta ← X^3\Sigma^+\) absorption. The lack of \((1P)\) resonances at these shorter excitation wavelengths is consistent with previous reports\(^{17,19}\) that the \(1\Delta\) state of the parent fragments to singlet products \([\text{CO}(X)+S(1D)]\) with a (near) unit quantum yield.

**CONCLUSION**

The present work reports the use of two and three photon resonant MPI spectroscopy to investigate Rydberg excited states of the OCS molecule in the energy range 70500–86000 cm\(^{-1}\). Our interpretation of the various observed resonances has been greatly assisted by companion measurements of the kinetic energies of the accompanying photoelectrons. The present study serves to reinforce earlier conclusions\(^{11}\) regarding five Rydberg origins in the 70500–73000 cm\(^{-1}\) energy range, which have been attributed to, respectively, states of \(^3\Pi\), \(^1\Pi\), \(^3\Delta\), \(^1\Delta\) and \(^1\Sigma^+\) symmetry arising from the \(4p\lambda ← 3\pi\) orbital promotion. At higher energies we identify no fewer than 21 more Rydberg origins. These appear in a number of clumps. Those with quantum defects ca. 3.5 and 4.5 we associate with the orbital promotions \(np\lambda ← 3\pi\) \((n=5,6)\), whilst the others with near integer quantum defect we suggest must be associated with excitation to \(s,d\) and (possibly) \(f\) Rydberg orbitals. We also identify a number of MPI resonances with \(m/z\) 28 and 32 which we attribute to, respectively, \([\text{CO}(X)+3\Sigma^+]\) fragments and S atoms [in both their ground \(^3P\) and excited \(^3D\) electronic states]. Analysis of the former feature confirms that the \([\text{CO}(X)+3\Sigma^+]\) fragments resulting from one photon dissociation of OCS at excitation wavelengths ca. 230 nm are formed with a highly inverted, bimodal rotational state population distribution, whilst the latter are consistent with previous reports of the wavelength dependence for forming ground and excited state S atoms in the near uv photolysis of OCS.

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Once again the Bristol group is happy to acknowledge financial support from the EPSRC (previously SERC) in the form of equipment grants and a studentship (to R.A.M.), and the unstinting practical help and encouragement of Mr K. N. Rosser. A.J.O.-E. is grateful to the Royal Society for the award of the Eliz. Challener Research Fellowship, whilst M.N.R.A. is grateful to the Ciba Fellowship Trust for their support of the Bristol-Amsterdam collaboration. D.A. acknowledges the financial support of the European Union Erasmus scheme. The Amsterdam group gratefully acknowl-