

# INTEGRAL REPRESENTATIONS FOR HORN'S $H_2$ FUNCTION AND OLSSON'S $F_P$ FUNCTION

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**Abstract.** We derive some Euler type double integral representations for hypergeometric functions in two variables. In the first part of this paper we deal with Horn's  $H_2$  function, in the second part with Olsson's  $F_P$  function. Our double integral representing the  $F_P$  function is compared with the formula for the same integral representing an  $H_2$  function by M. Yoshida (Hiroshima Math. J. **10** (1980), 329–335) and M. Kita (Japan. J. Math. **18** (1992), 25–74). As specified by Kita, their integral is defined by a homological approach. We present a classical double integral version of Kita's integral, with outer integral over a Pochhammer double loop, which we can evaluate as  $H_2$  just as Kita did for his integral. Then we show that shrinking of the double loop yields a sum of two double integrals for  $F_P$ .

## 1. Introduction

This paper deals with integral representations for two bivariate hypergeometric functions: on the one hand the  $H_2$  Horn function [10, (5.7.14)] which occurs in Horn's list [14] of bivariate hypergeometric series of order two; on the other hand the  $F_P$  function which was introduced by Olsson [20, 21]. The integral representations under consideration will be integrals of classical type over two-dimensional real domains. Necessarily, because of convergence, this requires constraints on the allowed parameter values.

Horn [13] called a double power series  $\sum_{i,j=0}^{\infty} A(i, j)x^i y^j$  *hypergeometric* if the two quotients  $A(i+1, j)/A(i, j)$  and  $A(i, j+1)/A(i, j)$  are rational functions of  $i$  and  $j$ . Write the two rational functions as quotients of polynomials in  $i$  and  $j$  without common factors:

$$\frac{A(i+1, j)}{A(i, j)} = \frac{F(i, j)}{F'(i, j)} \quad \text{and} \quad \frac{A(i, j+1)}{A(i, j)} = \frac{G(i, j)}{G'(i, j)}.$$

In addition it is assumed that  $F'(i, j)$  contains the factor  $i+1$  and  $G'(i, j)$  the factor  $j+1$ . Then the highest degree in  $i, j$  of the four polynomials  $F, F', G, G'$  is called the *order* of the hypergeometric series.

Horn [14] listed all convergent bivariate hypergeometric series of order two. For the case that all polynomials  $F, F', G, G'$  have degree two he obtains 14 items, which include Appell's hypergeometric series  $F_1, F_2, F_3$  and  $F_4$ , dating back to 1880 [4–6], but also, among others, the  $H_2$  function. These bivariate hypergeometric series can be considered

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as the natural two-variable analogues of the Gauss hypergeometric series. Horn gives 20 further items which are confluent cases of the first 14 items. For the basics of hypergeometric functions of two variables see for instance [7, Ch. 9], [10, Sections 5.7–5.12], [11, Ch. 1] and [24, Ch. 8].

In [10, Section 5.9], for each function in Horn's list a system of two partial differential equations (PDEs) is given which has this function as a solution. This follows a method already indicated by Horn in [13] and [14]. Some of these systems can be transformed into each other. This is for instance the case with the systems for  $F_2$ ,  $F_3$  and  $H_2$ . Otherwise stated,  $F_3$  and  $H_2$ , when suitably adapted, occur as solutions of the system [10, 5.9(10)] for  $F_2$ . A comprehensive list of solutions of this  $F_2$  system was given by Olsson [22, p. 1289, Table I]. Each solution is in terms of the six bivariate functions  $F_2$ ,  $F_3$ ,  $H_2$ ,  $F_P$ ,  $F_Q$  and  $F_R$ . While  $F_R$  is not even a hypergeometric series in the sense of Horn,  $F_P$  and  $F_Q$  are hypergeometric series of order three.

Euler type double integral representations for the Appell hypergeometric functions  $F_1$ ,  $F_2$ ,  $F_3$  were already given by Appell [5, Ch. III] in 1882. Some such representations for other functions in Horn's list are scattered through the literature. Kita [17, pp. 56–58] gives a long list of double integral representations, however with integrals defined by a homological approach. Single integral representations with one or two hypergeometric functions in the integrand can be found in the literature in more cases. Usually double integral representations imply such single integral representations.

When we concentrate on solutions of the  $F_2$  system, it is natural to ask if, beside  $F_2$  and  $F_3$ , also other solutions have an Euler type double integral representation. In this paper we obtain these for  $H_2$  and  $F_P$ . The integral representation we derive for  $H_2$  turned out to have been given earlier in a forgotten paper by Tuan and Kalla [25, (77)]. We show that our method of derivation gives rise to several variants of this formula.

The integral representation we derive for  $F_P$  is puzzling when compared with Yoshida [26, (0.10)] and Kita [17, p. 57, item 9]. A straightforward rewriting of our integral as an integral over a triangular region is evaluated by these authors as an  $H_2$  function. However, as becomes clear from Kita's paper, their double integral is defined by a homological approach, involving twisted cycles. While Kita gives, for instance, an integral of this type for  $F_2$  which can be immediately matched with Appell's classical integral representation [10, 5.8(2)] (under suitable constraints on the parameter values), a similar relation apparently does not exist in the  $F_P$  case between Kita's and our formula. In this paper we will not go into the technicalities of the homological approach (see for instance [17, Section 3], [3, Section 2.3] and [28, Ch. IV]). Instead, we will show that the evaluation as  $H_2$  by Yoshida and Kita can also be achieved by a classical double integral, where the outer integral is taken over a Pochhammer double loop. Then it is clear that a branch point caused by an independent variable prevents the Pochhammer loop from being shrunk such that there results an integral over a segment. However, it is possible to do the shrinking in such a way that one arrives at a sum of two integrals over real intervals. Both integrals in this sum can be evaluated as an  $F_P$ . The resulting formula is a three-term relation [22, (53)] involving one  $H_2$  and two  $F_P$  terms. A similar phenomenon, but more simple, can already be illustrated for a variant of the Euler integral representation for the Gauss hypergeometric function. While it is quite standard to derive connection formulas by using the homological approach (see Remark 5.1), this does not seem to have been done before by using integrals over Pochhammer double loops.

The contents of this paper are as follows. In Section 2 we summarize the integral representations of the Gauss hypergeometric function. In Section 3 we derive some integral representations of the  $H_2$  function. We also indicate how to get many further integral representations of this function. In Section 4 we derive a double integral representation of the  $F_P$  function and we give some corollaries. In Section 5 we give another corollary of this double integral representation for  $F_P$ . We compare it with the Yoshida–Kita integral representation which uses the homological approach and yields  $H_2$ . Then we show that this  $H_2$  evaluation already occurs with a classical double integral, the outer one being over a Pochhammer double loop. Shrinking of the double loop leads to a known three-term relation. Finally, in Section 6 we briefly discuss how the obtained double integrals appear when written as integral representations of solutions of the system of PDEs for  $F_2$ .

*Notation.* For  $a \in \mathbb{C}$  and  $k \in \mathbb{Z}$  the *Pochhammer symbol* is defined by

$$(a)_k := \frac{\Gamma(a+k)}{\Gamma(a)} = \begin{cases} a(a+1) \cdots (a+k-1), & k > 0, \\ 1, & k = 0, \\ (-1)^k / (1-a)_{-k}, & k < 0. \end{cases}$$

## 2. Euler integral representations for the Gauss hypergeometric function

The *Gauss hypergeometric function* [2, Ch. 2] is defined as a power series by

$${}_2F_1 \left( \begin{matrix} a, b \\ c \end{matrix}; z \right) := \sum_{k=0}^{\infty} \frac{(a)_k (b)_k}{(c)_k k!} z^k \quad (|z| < 1). \quad (2.1)$$

It has, as a function of  $z$ , a one-valued analytic continuation to  $\mathbb{C} \setminus [1, \infty)$ , as is seen (at least under the given parameter restrictions) from the *Euler integral representation*

$${}_2F_1 \left( \begin{matrix} a, b \\ c \end{matrix}; z \right) = \frac{\Gamma(c)}{\Gamma(b)\Gamma(c-b)} \int_0^1 t^{b-1} (1-t)^{c-b-1} (1-zt)^{-a} dt \quad (2.2)$$

$$= \frac{\Gamma(c)}{\Gamma(b)\Gamma(c-b)} (1-z)^{c-a-b} \int_0^1 t^{c-b-1} (1-t)^{b-1} (1-zt)^{a-c} dt \quad (2.3)$$

$$= \frac{\Gamma(c)}{\Gamma(a)\Gamma(c-a)} \int_0^1 t^{a-1} (1-t)^{c-a-1} (1-zt)^{-b} dt \quad (2.4)$$

$$= \frac{\Gamma(c)}{\Gamma(a)\Gamma(c-a)} (1-z)^{c-a-b} \int_0^1 t^{c-a-1} (1-t)^{a-1} (1-zt)^{b-c} dt \quad (2.5)$$

with convergence conditions  $\operatorname{Re} c > \operatorname{Re} b > 0$  for (2.2) and (2.3), and  $\operatorname{Re} c > \operatorname{Re} a > 0$  for (2.4) and (2.5). Formula (2.2) is the usual formula for the Euler integral representation. Then (2.3) follows from (2.2) by the *Euler transformation formula*

$${}_2F_1 \left( \begin{matrix} a, b \\ c \end{matrix}; z \right) = (1-z)^{c-a-b} {}_2F_1 \left( \begin{matrix} c-a, c-b \\ c \end{matrix}; z \right). \quad (2.6)$$

Note its special case

$${}_2F_1 \left( \begin{matrix} a, b \\ a \end{matrix}; z \right) = (1-z)^{-b}. \quad (2.7)$$

Formulas (2.4) and (2.5) follow by the symmetry of (2.1) in  $a$  and  $b$ . The *Pfaff transformation formula*

$${}_2F_1\left(\begin{matrix} a, b \\ c \end{matrix}; z\right) = (1-z)^{-a} {}_2F_1\left(\begin{matrix} a, c-b \\ c \end{matrix}; \frac{z}{z-1}\right) \quad (2.8)$$

follows from (2.2) or (2.5) by the change of integration variable  $t \rightarrow 1-t$ . Similarly, its variant

$${}_2F_1\left(\begin{matrix} a, b \\ c \end{matrix}; z\right) = (1-z)^{-b} {}_2F_1\left(\begin{matrix} c-a, b \\ c \end{matrix}; \frac{z}{z-1}\right)$$

follows from (2.3) or (2.4). For generic values of the parameters, formulas (2.2)–(2.5) are the only integral representations of the Gauss hypergeometric function of the form

$$\text{const. } (1-z)^\lambda \int_0^1 t^\alpha (1-t)^\beta (1-zt)^\Gamma dt,$$

or else there would be other transformation formulas of the form (2.6).

On the right-hand side of (2.2) we can apply a one-parameter group of transformations of integration variable

$$\phi_p: t \mapsto \frac{t}{p + (1-p)t} \quad (p > 0), \quad (2.9)$$

which map the integration interval  $[0, 1]$  onto itself. Note that  $\phi_p \circ \phi_q = \phi_{pq}$  and  $\phi_1 = \text{id}$ . The integral representation resulting from applying the transformation of integration variable  $\phi_p$  to (2.2) is

$$\begin{aligned} {}_2F_1\left(\begin{matrix} a, b \\ c \end{matrix}; z\right) &= \frac{\Gamma(c)}{\Gamma(b)\Gamma(c-b)} p^{-b} \\ &\times \int_0^1 t^{b-1} (1-t)^{c-b-1} \left(1 - \frac{p-1}{p}t\right)^{a-c} \left(1 - \frac{z+p-1}{p}t\right)^{-a} dt \end{aligned} \quad (2.10)$$

$$= p^{-b} F_1\left(b, c-a, a, c; \frac{p-1}{p}, \frac{z+p-1}{p}\right), \quad (2.11)$$

where the constraints in (2.10) are the same as in (2.2), and where the second equality follows from the integral representation [10, 5.8(5)] for the Appell hypergeometric function  $F_1$ . Then (2.11) gives a reduction formula for  $F_1(a, b, c, b+c; x, y)$ , which is [10, 5.10(1)] combined with (2.8).

Similarly, application of the transformation of integration variable

$$\psi_p(t) := \phi_p(1-t) = \frac{1-t}{1-(1-p)t} \quad (p > 0) \quad (2.12)$$

to the right-hand side of (2.2) gives

$$\begin{aligned} {}_2F_1\left(\begin{matrix} a, b \\ c \end{matrix}; z\right) &= \frac{\Gamma(c)}{\Gamma(b)\Gamma(c-b)} p^{c-b} (1-z)^{-a} \\ &\times \int_0^1 t^{c-b-1} (1-t)^{b-1} (1-(1-p)t)^{a-c} \left(1 - \frac{1-p-z}{1-z}t\right)^{-a} dt \end{aligned} \quad (2.13)$$

$$= p^{c-b} (1-z)^{-a} F_1\left(c-b, c-a, a, c; 1-p, \frac{1-p-z}{1-z}\right), \quad (2.14)$$

again with the same constraints in (2.13) as in (2.2), and with (2.14) being the reduction formula [10, 5.10(1)] combined with (2.6). Note that the special case  $p = 1$  of (2.13) or (2.14) gives (2.8).

*Remark 2.1.* The same method of applying  $\phi_p$  or  $\psi_p$  to the integration variable works for the single integral representation [10, 5.8(5)] for the Appell  $F_1$  function or, more generally, for the single integral representation [6, p. 116, formula (8)] for the Lauricella  $F_D$  function in  $n$  variables. For instance, in the Lauricella case the result is an integral representation for Lauricella  $F_D$  in  $n + 1$  variables, by which we also obtain a reduction formula for Lauricella  $F_D$  with a dependence between the parameters.

### 3. Integral representations for the $H_2$ function

We start with the double power series of the  $H_2$  function [10, 5.7(14) and (49)]

$$H_2(a, b, c, d, e; x, y) = H_2(x, y) := \sum_{i,j=0}^{\infty} \frac{(a)_{i-j}(b)_i(c)_j(d)_j}{(e)_i i! j!} x^i y^j \quad (3.1)$$

with convergence region (see Figure 1)

$$\Omega_1 := \{(x, y) \in \mathbb{C}^2 \mid |x| < 1, |y| < (|x| + 1)^{-1}\}.$$

By inserting

$$(a)_{i-j} = (-1)^j \frac{(a-j)_i}{(1-a)_j}$$

in (3.1) we get

$$\begin{aligned} H_2(a, b, c, d, e; x, y) &= \sum_{j=0}^{\infty} \frac{(c)_j(d)_j}{(1-a)_j j!} (-y)^j \sum_{i=0}^{\infty} \frac{(a-j)_i(b)_i}{(e)_i i!} x^i \\ &= \sum_{j=0}^{\infty} \frac{(c)_j(d)_j}{(1-a)_j j!} (-y)^j {}_2F_1 \left( \begin{matrix} a-j, b \\ e \end{matrix}; x \right) \end{aligned} \quad (3.2)$$

for  $(x, y) \in \Omega_1$ . Specialization of (3.2) gives

$$H_2(a, b, c, d, e; 0, y) = {}_2F_1 \left( \begin{matrix} c, d \\ 1-a \end{matrix}; -y \right), \quad (3.3)$$

initially for  $|y| < 1$ , and after analytic continuation for complex  $y$  outside  $(-\infty, -1]$ .

Substitution of Euler's integral representation (2.2) in (3.2) gives

$$\begin{aligned} &H_2(a, b, c, d, e; x, y) \\ &= \sum_{j=0}^{\infty} \frac{(c)_j(d)_j}{(1-a)_j j!} (-y)^j \frac{\Gamma(e)}{\Gamma(b)\Gamma(e-b)} \int_0^1 u^{b-1} (1-u)^{e-b-1} (1-xu)^{-a+j} du \end{aligned}$$

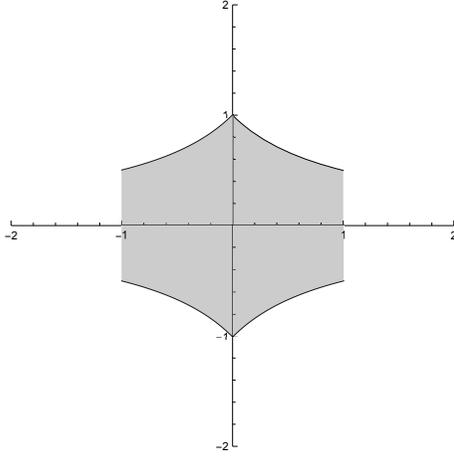


FIGURE 1. Convergence region  $\Omega_1 \cap \mathbb{R}^2$  of the  $H_2$  function.

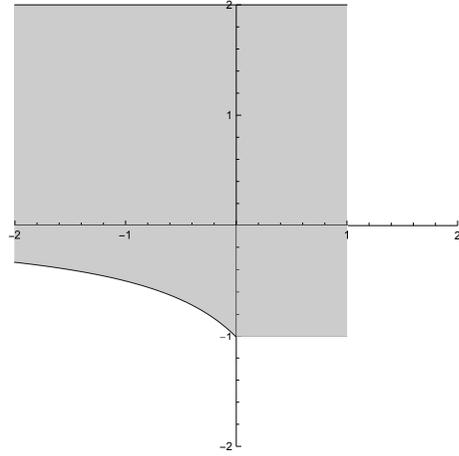


FIGURE 2. Region of analytic continuation  $\Omega_2 \cap \mathbb{R}^2$  of  $H_2$ .

with condition  $0 < \operatorname{Re} b < \operatorname{Re} e$ . Because, for  $(x, y) \in \Omega_1$ , we have

$$\begin{aligned} & \sum_{j=0}^{\infty} \left| \frac{(c)_j (d)_j}{(1-a)_j j!} (-y)^j \right| \int_0^1 |u^{b-1} (1-u)^{e-b-1} (1-xu)^{-a+j}| du \\ & \leq \text{const.} \sum_{j=0}^{\infty} \left| \frac{(c)_j (d)_j}{(1-a)_j j!} \right| |y|^j (|x|+1)^j \int_0^1 |u^{b-1} (1-u)^{e-b-1} (1-xu)^{-a}| du < \infty, \end{aligned}$$

the sum and integral can be interchanged by dominated convergence. So

$$\begin{aligned} & H_2(a, b, c, d, e; x, y) \\ & = \frac{\Gamma(e)}{\Gamma(b)\Gamma(e-b)} \int_0^1 u^{b-1} (1-u)^{e-b-1} (1-xu)^{-a} \sum_{j=0}^{\infty} \frac{(c)_j (d)_j}{(1-a)_j j!} (-y)^j (1-xu)^j du \\ & = \frac{\Gamma(e)}{\Gamma(b)\Gamma(e-b)} \int_0^1 u^{b-1} (1-u)^{e-b-1} (1-xu)^{-a} {}_2F_1 \left( \begin{matrix} d, c \\ 1-a \end{matrix}; -y(1-xu) \right) du. \end{aligned} \tag{3.4}$$

The absolute value of the integrand of (3.4) is dominated by  $\text{const.} |u^{b-1} (1-u)^{e-b-1}|$ , uniformly for  $(x, y)$  in compact subsets of the region

$$\Omega_2 := \{(x, y) \in \mathbb{C}^2 \mid x \notin [1, \infty) \text{ and } \forall u \in [0, 1] \ y(xu - 1) \notin [1, \infty)\}.$$

Therefore, under the constraints  $0 < \operatorname{Re} b < \operatorname{Re} e$ , (3.4) provides an analytic continuation of  $H_2(x, y)$  for  $(x, y) \in \Omega_2$ . The intersection of  $\Omega_2$  with  $\mathbb{R}^2$  is

$$\{(x, y) \in \mathbb{R}^2 \mid x \leq 0, (x-1)y < 1 \vee 0 < x < 1, y > -1\}$$

(see Figure 2), much larger than the convergence region for the power series (3.1) intersected with  $\mathbb{R}^2$  (see Figure 1).

Substitution of Euler's integral representation (2.2) in (3.4) gives

$$\begin{aligned} H_2(a, b, c, d, e; x, y) &= \frac{\Gamma(e)}{\Gamma(b)\Gamma(e-b)} \frac{\Gamma(1-a)}{\Gamma(c)\Gamma(1-a-c)} \\ &\quad \times \int_0^1 u^{b-1} (1-u)^{e-b-1} (1-xu)^{-a} \int_0^1 v^{c-1} (1-v)^{-a-c} (1+yv-xyuv)^{-d} dv du, \end{aligned} \quad (3.5)$$

where we require moreover that  $0 < \operatorname{Re} c < \operatorname{Re}(1-a)$ . By Fubini's theorem formula (3.5) can be rewritten as

$$\begin{aligned} H_2(a, b, c, d, e; x, y) &= \frac{\Gamma(e)}{\Gamma(b)\Gamma(e-b)} \frac{\Gamma(1-a)}{\Gamma(c)\Gamma(1-a-c)} \\ &\quad \times \int_0^1 \int_0^1 u^{b-1} v^{c-1} (1-u)^{e-b-1} (1-v)^{-a-c} (1-xu)^{-a} (1+yv-xyuv)^{-d} du dv, \end{aligned} \quad (3.6)$$

because the absolute value of the integrand is dominated by  $\operatorname{const} \cdot |u^{b-1} (1-u)^{e-b-1} v^{c-1} \times (1-v)^{-a-c}|$ , for which the integral over  $(u, v) \in [0, 1] \times [0, 1]$  is finite. So formula (3.6) is valid for  $x \in \Omega_2$  and  $0 < \operatorname{Re} b < \operatorname{Re} e$ ,  $0 < \operatorname{Re} c < \operatorname{Re}(1-a)$ .

Formula (3.3) was earlier given by Olsson [22, (44)]. The integral representation (3.6) was earlier given by Tuan and Kalla [25, (77)]. In both references the same constraints on the parameters are given as we have, but the region on which the integral representation is valid is not specified in these two references. For the proof of [25, (77)] Tuan and Kalla just mention that the formula follows by termwise integration of the double power series of the integrand. This is essentially the proof as given above.

For  $H_2$  we have the transformation formula

$$H_2(a, b, c, d, e; x, y) = (1-x)^{-a} H_2\left(a, e-b, c, d, e; \frac{x}{x-1}, y(1-x)\right), \quad (3.7)$$

which is a rewriting of [22, (25)]. Formula (3.7) also follows from (3.2) together with (2.8) or from (3.4) or (3.6) by the change of integration variable  $u \rightarrow 1-u$ .

In the proof of (3.6) we have twice substituted Euler's integral representation (2.2) for a  ${}_2F_1$ : after (3.2) and after (3.4). In both places we might have substituted one of the alternatives (2.3), (2.4) or (2.5) for (2.2). After (3.2) only the alternative (2.3) will be an option, because  $\operatorname{Re}(a-j)$  would not be positive for  $j$  large enough and therefore the conditions for (2.4) and (2.5) would be violated.

After substitution in (3.2) of (2.3) we can proceed in the same way as above for the derivation of (2.5). We obtain

$$\begin{aligned} H_2(a, b, c, d, e; x, y) &= \frac{\Gamma(e)}{\Gamma(b)\Gamma(e-b)} (1-x)^{e-a-b} \\ &\quad \times \int_0^1 u^{e-b-1} (1-u)^{b-1} (1-xu)^{a-e} {}_2F_1\left(\begin{matrix} d, c \\ 1-a \end{matrix}; \frac{y(x-1)}{1-xu}\right) du \end{aligned} \quad (3.8)$$

$$\begin{aligned}
&= \frac{\Gamma(e)}{\Gamma(b)\Gamma(e-b)} \frac{\Gamma(1-a)}{\Gamma(c)\Gamma(1-a-c)} \\
&\quad \times \int_0^1 \int_0^1 u^{e-b-1} (1-u)^{b-1} v^{c-1} (1-v)^{-a-c} (1-xu)^{a+d-e} \\
&\quad \times (1-xu+y(1-x)v)^{-d} du dv. \tag{3.9}
\end{aligned}$$

Formulas (3.7) and (3.8) give an analytic continuation of  $H_2(x, y)$  to the region

$$\Omega_3 := \left\{ (x, y) \in \mathbb{C}^2 \mid x \notin [1, \infty) \text{ and } \forall u \in [0, 1] \frac{y(x-1)}{1-xu} \notin [1, \infty) \right\}$$

(note that  $\Omega_3$  has the same intersection with  $\mathbb{R}^2$  as  $\Omega_2$ ). The parameter constraints for (3.8) and (3.9) are the same as those for (3.4) and (3.6), respectively.

We can obtain two other variants of the double integral representations (3.6) and (3.9) by substituting (2.3) in (3.4) and (3.8), respectively, but we omit the explicit formulas. On the other hand, nothing of interest comes out when we substitute (2.4) or (3.6) in (3.4) or (3.8). The resulting formulas only differ from the four earlier integral representations by interchange of  $c$  and  $d$ . This is the trivial symmetry which is obvious from (3.1).

We may also vary the four double integral representations for  $H_2$  by applying the change of integration variable  $\phi_p$  or  $\psi_p$  (see (2.9), (2.12)) to  $u$  or  $v$ . As already remarked,  $\psi_1$  applied to  $u$  corresponds to the transformation formula (3.7). Other cases possibly give specializations of double integral representations for hypergeometric functions of three variables (see also Remark 2.1). Indeed,  $\phi_p$  or  $\psi_p$  applied to the  $u$  variable in (3.6) would yield a double integral with a relation between the parameters which corresponds to the result of applying (2.11) or (2.14) to the  ${}_2F_1$  in (3.2). If we were next to substitute for the resulting  $F_1$  in (3.2) its double power series expansion then a hypergeometric triple series with a relation between the parameters would arise which is equal to  $H_2$ . The identities below, where we have omitted the constraints, give the explicit result for the case that  $\phi_p$  is applied:

$$\begin{aligned}
&H_2(a, b, c, d, e; x, y) \\
&= \frac{\Gamma(e)}{\Gamma(b)\Gamma(e-b)} \frac{\Gamma(1-a)}{\Gamma(c)\Gamma(1-a-c)} p^{-b} \int_0^1 \int_0^1 u^{b-1} v^{c-1} (1-u)^{e-b-1} (1-v)^{-a-c} \\
&\quad \times (1-(1-p^{-1})u)^{a+d-e} (1-p^{-1}(x+p-1)u)^{-a} \\
&\quad \times (1-(1-p^{-1})u+yv-p^{-1}(x+p-1)yuv)^{-d} du dv \\
&= p^{-b} \sum_{i,j,k=0}^{\infty} \frac{(a)_{k-i} (b)_{j+k} (e-a)_{i+j} (c)_i (d)_i}{(e)_{j+k} (e-a)_i i! j! k!} y^i \left(\frac{p-1}{p}\right)^j \left(\frac{x+p-1}{p}\right)^k.
\end{aligned}$$

Quite possibly the second identity can be generalized to a case with less dependence between the parameters in the double integral and in the triple series.

#### 4. Integral representations for the $F_p$ function

The function  $F_p$  was first introduced by Olsson [20, 21] as a certain solution of the system of PDEs associated with Appell's hypergeometric function  $F_2$  given in [10, 5.9(10)]. In [20]

the notation  $Z_1$  instead of  $F_P$  is used. The definitions [20, (4)] and [21, (1a)] of  $F_P$  are also different. But we can put together formulas for  $F_P$  in [20] and [21] as the following string of equalities:

$$F_P(a, b_1, b_2, c_1, c_2; x, y) = \sum_{i,j=0}^{\infty} \frac{(a)_{i+j}(a-c_2+1)_i(b_1)_i(b_2)_j}{(a+b_2-c_2+1)_{i+j}(c_1)_i i! j!} x^i (1-y)^j \quad (4.1)$$

$$= \sum_{i=0}^{\infty} \frac{(a)_i(a-c_2+1)_i(b_1)_i}{(c_1)_i(a+b_2-c_2+1)_i i!} x^i {}_2F_1\left(\begin{matrix} a+i, b_2 \\ a+b_2-c_2+1+i \end{matrix}; 1-y\right) \quad (4.2)$$

$$= y^{-a} \sum_{i=0}^{\infty} \frac{(a)_i(a-c_2+1)_i(b_1)_i}{(c_1)_i(a+b_2-c_2+1)_i i!} \left(\frac{x}{y}\right)^i {}_2F_1\left(\begin{matrix} a+i, a-c_2+1+i \\ a+b_2-c_2+1+i \end{matrix}; \frac{y-1}{y}\right) \quad (4.3)$$

$$= y^{-a} \sum_{i,j=0}^{\infty} \frac{(a)_{i+j}(a-c_2+1)_{i+j}(b_1)_i}{(a+b_2-c_2+1)_{i+j}(c_1)_i i! j!} \left(\frac{x}{y}\right)^i \left(\frac{y-1}{y}\right)^j. \quad (4.4)$$

The double power series (4.1) and (4.4) are both of order three. By Horn's rule for determination of the convergence region (see [10, Section 5.7.2]) the double power series (4.1) is absolutely convergent on the region  $|x|, |y-1| < 1$  in  $\mathbb{C}^2$ , while (4.4) is absolutely convergent on the region  $|xy^{-1}| + |1-y^{-1}| < 1$  in  $\mathbb{C}^2$ , which has intersection  $|x| < 1, y > \frac{1}{2}(1+|x|)$  with  $\mathbb{R}^2$ . Both convergence regions are neighborhoods of  $(0, 1)$ . The equalities (4.1)=(4.2) and (4.3)=(4.4) on the respective convergence regions of (4.1) and (4.4) follow by rewriting the double series with an inner  $j$ -sum and an outer  $i$ -sum. The equality (4.2)=(4.3), which follows by Pfaff's transformation (2.8), is initially valid on the intersection of the two convergence regions, but next yields an analytic continuation of  $F_P$  to the union of the two convergence regions. The intersection of this union with  $\mathbb{R}^2$  is  $|x| < 1 \wedge y > 0$ .

By using the double power series (4.1) it can be seen that  $F_P(a, b_1, b_2, c_1, c_2; x, y)$  is the solution of the system of PDEs for  $F_2(a; b_1, b_2; c_1, c_2; x, y)$  which is regular and equal to 1 at the point  $(0, 1)$ .

In [21] and [22] Olsson takes (4.4) as a definition of  $F_P$ , while he starts in [20] with

$$F_P(a; b_1, b_2; c_1, c_2; x, y) = \sum_{i=0}^{\infty} \frac{(a)_i(b_2)_i}{(a+b_2-c_2+1)_i i!} (1-y)^i {}_3F_2\left(\begin{matrix} a+i, b_1, a-c_2+1 \\ c_1, a+b_2-c_2+1+i \end{matrix}; x\right) \quad (4.5)$$

and then gives (4.1) and (4.2). Note that (4.5) follows on the convergence region of (4.1) by rewriting (4.1) with an inner  $j$ -sum and an outer  $i$ -sum.

The expression of  $F_P$  by the double power series (4.4) implies the symmetry for  $F_P$  which is given by the first equality in [22, (26)]. However, no immediate symmetry can be derived from (4.1) because the five parameters occur there in six shifted factorials.

We now state our main result on  $F_P$ : a double integral representation.

**THEOREM 4.1.** For  $\operatorname{Re} a, \operatorname{Re}(a - b_2), \operatorname{Re}(b_2 + c_1 - a), \operatorname{Re}(b_2 - c_2 + 1), \operatorname{Re}(c_1 + c_2 - a - 1) > 0$  and  $(x, y) \in \mathbb{C}^2$  such that  $x \notin [1, \infty)$  and  $y \notin (-\infty, 0]$ , we have

$$\begin{aligned} & F_P(a, b_1, b_2, c_1, c_2; x, y) \\ &= \frac{\Gamma(a + b_2 - c_2 + 1)\Gamma(c_1)}{\Gamma(a)\Gamma(a - c_2 + 1)\Gamma(b_2 - c_2 + 1)\Gamma(c_1 + c_2 - a - 1)} \\ & \times y^{-b_2} \int_0^1 \int_0^1 u^{a-1}(1-u)^{b_2+c_1-a-1} v^{b_2-c_2}(1-v)^{c_1+c_2-a-2} (1-xu)^{-b_1} \\ & \times (u + (1-u)v y^{-1})^{-b_2} du dv \end{aligned} \quad (4.6)$$

with absolutely convergent double integral.

*Remark 4.2.* In (4.6) the function of  $(x, y)$  defined by the right-hand side is analytic on the given region and provides an analytic continuation of the left-hand side. A similar remark applies to the subsequent corollaries.

Before proving Theorem 4.1 we mention three immediate consequences, which are, in a sense, equivalent to (4.6). The last two of these give formulas occurring without proof in papers by Olsson. We postpone a further corollary, a rewritten form of (4.6) which is close to the Yoshida–Kita integral, to Section 5.

**COROLLARY 4.3.** For  $\operatorname{Re} a, \operatorname{Re}(a - c_2 + 1), \operatorname{Re}(b_2 + c_1 - a) > 0$  and  $(x, y) \in \mathbb{C}^2$  such that  $x \notin [1, \infty)$  and  $y \notin (-\infty, 0]$ , we have

$$\begin{aligned} F_P(a, b_1, b_2, c_1, c_2; x, y) &= \frac{\Gamma(c_1)\Gamma(a + b_2 - c_2 + 1)}{\Gamma(a)\Gamma(a - c_2 + 1)\Gamma(b_2 + c_1 - a)} \\ & \times y^{-b_2} \int_0^1 u^{a-b_2-1}(1-u)^{b_2+c_1-a-1} (1-xu)^{-b_1} \\ & \times {}_2F_1\left(\begin{matrix} b_2, b_2 - c_2 + 1 \\ -a + b_2 + c_1 \end{matrix}; y^{-1}(1-u^{-1})\right) du. \end{aligned} \quad (4.7)$$

The integral representation (4.7) for  $F_P$  is different from the two integral representations [22, (43), (48)] for  $F_P$  in terms of  ${}_2F_1$ . By (2.7) and (2.2) we obtain from (4.7) the specialization formula

$$F_P(a, b_1, b_2, a, c_2; 0, y) = y^{-c_2+1} {}_2F_1\left(\begin{matrix} b_2 - c_2 + 1, a - c_2 + 1 \\ a + b_2 - c_2 + 1 \end{matrix}; 1 - y\right) \quad (y \notin (-\infty, 0]). \quad (4.8)$$

**COROLLARY 4.4.** With the assumptions of Theorem 4.1 we have

$$\begin{aligned} & F_P(a, b_1, b_2, c_1, c_2; x, y) \\ &= \frac{\Gamma(a + b_2 - c_2 + 1)\Gamma(c_1)\Gamma(b_2 + c_1 - a)}{\Gamma(b_2 + c_1)\Gamma(a - c_2 + 1)\Gamma(b_2 - c_2 + 1)\Gamma(c_1 + c_2 - a - 1)} \\ & \times \int_0^1 v^{-c_2}(1-v)^{c_1+c_2-a-2} F_1\left(a, b_1, b_2, b_2 + c_1; x, \frac{v-y}{v}\right) dv. \end{aligned} \quad (4.9)$$

Formula (4.9) was earlier stated without proof by Olsson<sup>†</sup> [22, (65)].

<sup>†</sup>In [22, (65)] the factor  $(c_1 + c_2 - a + 1)$  should be  $\Gamma(c_1 + c_2 - a + 1)$ .

COROLLARY 4.5. *With the assumptions of Theorem 4.1 and moreover  $\operatorname{Re} y > \frac{1}{2}$  we have*

$$\begin{aligned} & F_P(a, b_1, b_2, c_1, c_2; x, y) \\ &= \frac{\Gamma(c_1)\Gamma(a+b_2-c_2+1)}{\Gamma(a-c_2+1)\Gamma(b_2+c_1)} y^{-b_2} \\ & \quad \times \sum_{j=0}^{\infty} \frac{(b_2)_j (b_2+c_1-a)_j}{(b_2+c_1)_j j!} {}_2F_1\left(\begin{matrix} a, b_1 \\ b_2+c_1+j \end{matrix}; x\right) {}_2F_1\left(\begin{matrix} -j, b_2-c_2+1 \\ b_2+c_1-a \end{matrix}; y^{-1}\right). \end{aligned} \quad (4.10)$$

Formula (4.10) was earlier stated without proof by Almström and Olsson [1, (22)]. See also [22, (55)].

*Proof of Theorem 4.1.* Denote the right-hand side of (4.6) by  $I(a, b_1, b_2, c_1, c_2; x, y)$  and consider this under the given constraints for  $x, y$  and the parameters. Then all occurring Gamma factors have argument with positive real part. Because

$$\text{const. } u \leq |u + t(1-u)vy^{-1}| \leq \max(|y|^{-1}, 1),$$

uniformly for  $y$  in compact subsets of  $\mathbb{C}$  outside  $(-\infty, 0]$ , we see that the double integral converges absolutely, and that  $I(x, y)$  is analytic for  $(x, y)$  in the given region. Now restrict  $(x, y)$  to  $-1 < x < 1$  and  $y > 0$ . We will show that  $I(x, y)$  is then equal to the series (4.2). Thus the functions given by (4.2) and by (4.1) on their initial domains will have analytic continuation to the domain of  $I(x, y)$ .

Expand the factor  $(1-xu)^{-b_1}$  in the integrand as binomial series and interchange the sum and the integrals by dominated convergence. We obtain

$$I(a, b_1, b_2, c_1, c_2; x, y) = \sum_{j=0}^{\infty} \frac{(b_1)_j}{j!} x^j I_j(a, b_2, c_1, c_2; y) \quad (4.11)$$

with

$$\begin{aligned} & I_j(a, b_2, c_1, c_2; y) = I_j(y) \\ &= \frac{\Gamma(a+b_2-c_2+1)\Gamma(c_1)}{\Gamma(a)\Gamma(a-c_2+1)\Gamma(b_2-c_2+1)\Gamma(c_1+c_2-a-1)} \\ & \quad \times \int_0^1 v^{-c_2}(1-v)^{c_1+c_2-a-2} \int_0^1 u^{a+j-1}(1-u)^{b_2+c_1-a-1} \left(1 - \frac{v-y}{v}u\right)^{-b_2} du dv. \end{aligned}$$

The inner integral can be evaluated by (2.2). Then

$$\begin{aligned} I_j(y) &= \frac{\Gamma(a+b_2-c_2+1)\Gamma(c_1)\Gamma(b_2+c_1-a)}{\Gamma(b_2+c_1)\Gamma(a-c_2+1)\Gamma(b_2-c_2+1)\Gamma(c_1+c_2-a-1)} \frac{(a)_j}{(b_2+c_1)_j} \\ & \quad \times \int_0^1 v^{-c_2}(1-v)^{c_1+c_2-a-2} {}_2F_1\left(\begin{matrix} b_2, a+j \\ b_2+c_1+j \end{matrix}; \frac{v-y}{v}\right) dv \end{aligned}$$

$$\begin{aligned}
&= \frac{\Gamma(a + b_2 - c_2 + 1)\Gamma(c_1)\Gamma(b_2 + c_1 - a)}{\Gamma(b_2 + c_1)\Gamma(a - c_2 + 1)\Gamma(b_2 - c_2 + 1)\Gamma(c_1 + c_2 - a - 1)} \frac{(a)_j}{(b_2 + c_1)_j} y^{1-c_2} \\
&\quad \times \int_{-\infty}^{1-y} (1-w)^{a-c_1} (1-y-w)^{c_1+c_2-a-2} {}_2F_1\left(\begin{matrix} b_2, a+j \\ b_2 + c_1 + j \end{matrix}; w\right) dw.
\end{aligned}$$

Now use [18, (4.11)]

$$\begin{aligned}
&\int_{-\infty}^x (1-y)^{a+b-c} {}_2F_1\left(\begin{matrix} a, b \\ c \end{matrix}; y\right) \frac{(x-y)^{\mu-1}}{\Gamma(\mu)} dy \\
&= \frac{\Gamma(c-a-\mu)}{\Gamma(c-a)} \frac{\Gamma(c-b-\mu)}{\Gamma(c-b)} \frac{\Gamma(c)}{\Gamma(c-\mu)} (1-x)^{a+b-c+\mu} {}_2F_1\left(\begin{matrix} a, b \\ c-\mu \end{matrix}; x\right),
\end{aligned}$$

with convergence conditions  $x < 1$ ,  $\operatorname{Re}(c-a)$ ,  $\operatorname{Re}(c-b) > \operatorname{Re}(\mu) > 0$ , in order to obtain

$$I_j(y) = \frac{(a)_j(a-c_2+1)_j}{(c_1)_j(a+b_2-c_2+1)_j} {}_2F_1\left(\begin{matrix} a+j, b_2 \\ a+b_2-c_2+1+j \end{matrix}; 1-y\right),$$

with  $y > 0$  and  $\operatorname{Re}(c_1)$ ,  $\operatorname{Re}(a-c_2+1)$ ,  $\operatorname{Re}(c_1+c_2-a-1) > 0$ . These constraints are implied by what we had already assumed. Substitution in (4.11) and combination with (4.2) gives the desired result (4.6).  $\square$

*Proof of Corollary 4.3.* First assume the conditions of Theorem 4.1. In the integrand of (4.6) extract the factor  $u^{-b_2}$  from  $(u + (1-u)vy^{-1})^{-b_2}$  and then substitute Euler's integral (2.1) for the  $v$ -integral. This yields (4.7). The conditions on the parameters in Corollary 4.3 are implied by those in Theorem 4.1 and they are sufficient for absolute convergence of the integral (use [10, 2.10(2)]). Hence, by analytic continuation in the parameters one can show that (4.7) holds under the given constraints.  $\square$

*Proof of Corollary 4.4.* Rewriting of (4.6) gives

$$\begin{aligned}
&F_P(a, b_1, b_2, c_1, c_2; x, y) \\
&= \frac{\Gamma(a + b_2 - c_2 + 1)\Gamma(c_1)}{\Gamma(a)\Gamma(a - c_2 + 1)\Gamma(b_2 - c_2 + 1)\Gamma(c_1 + c_2 - a - 1)} \\
&\quad \times \int_0^1 v^{-c_2} (1-v)^{c_1+c_2-a-2} \\
&\quad \times \int_0^1 u^{a-1} (1-u)^{b_2+c_1-a-1} (1-xu)^{-b_1} \left(1 - \frac{v-y}{v}u\right)^{-b_2} du dv.
\end{aligned}$$

The inner integral is an  $F_1$  Appell function, see [10, 5.8(5)].  $\square$

*Proof of Corollary 4.5.* Assume moreover  $\operatorname{Re} y > \frac{1}{2}$ . Then  $|1 - vy^{-1}| \leq v$ . So we can expand  $(u + (1-u)vy^{-1})^{-b_2} = (1 - (1-y^{-1}v)(1-u))^{-b_2}$  in (4.6) as binomial series. By dominated convergence we obtain

$$\begin{aligned}
&F_P(a, b_1, b_2, c_1, c_2; x, y) \\
&= \frac{\Gamma(a + b_2 - c_2 + 1)\Gamma(c_1)}{\Gamma(a)\Gamma(a - c_2 + 1)\Gamma(b_2 - c_2 + 1)\Gamma(c_1 + c_2 - a - 1)} y^{-b_2}
\end{aligned}$$

$$\begin{aligned} & \times \sum_{j=0}^{\infty} \frac{(b_2)_j}{j!} \int_0^1 u^{a-1} (1-u)^{b_2+c_1-a-1+j} (1-xu)^{-b_1} \\ & \times \int_0^1 v^{b_2-c_2} (1-v)^{c_1+c_2-a-2} (1-y^{-1}v)^j du dv. \end{aligned}$$

Then substitution of (2.2) or (2.4) gives the result.  $\square$

## 5. The Yoshida–Kita integral for $H_2$

In this section we compare a rewritten version of the double integral representation (4.6) for  $F_P$  with Yoshida [26, (0.10)] and Kita [17, p. 57, item 9], who, on first superficial view, seem to give an evaluation of this integral as an  $H_2$  function. The difference is caused because Kita's integral is defined quite differently, by using homological methods involving twisted cycles, while we are integrating classically over a domain in  $\mathbb{R}^2$ . We will consider a variant of our double integral (4.6) in which the outer integral is taken over a Pochhammer double loop. This turns out to have the same  $H_2$  evaluation as Kita's integral. There is a branch point inside the double loop caused by an independent variable. This prevents shrinking of the double loop outer integral such that one arrives at the double integral (4.6).

For a better understanding we first discuss a comparable case of a double loop integral version of Euler's integral representation for  ${}_2F_1$  which does not allow immediate shrinking. It turns out that there shrinking of the contour is possible to a sum of two Euler type integrals, by which we arrive at a standard three-term identity for  ${}_2F_1$ .

Then we do a similar, but more complicated, exercise for the double integral with double loop outer integral which has an evaluation as an  $H_2$  function. We will keep the inner integral as an ordinary integral over  $[0, 1]$ , and we will evaluate it by Euler's integral (2.2). Then we are dealing with a double loop integral having a Gauss hypergeometric function in the integrand. After shrinking of the double loop this integral will split as a sum of two ordinary integrals over intervals. We will treat only one of the four occurring cases in more detail. It will turn out that both terms can be evaluated as an  $F_P$  function by means of (4.7). The resulting three-term identity will now be an identity given by Olsson [22], involving two  $F_P$  terms and one  $H_2$  term.

### 5.1. Pochhammer's double loop integral for the Gauss hypergeometric function

The notion of a *double loop* integral was introduced by Jordan [15, pp. 243–244] (1887) and Pochhammer [23] (1890). This means an integral over a contour in the complex  $t$ -plane starting at some point  $t_0$  on the interval  $(0, 1)$ , for which we assume that  $\arg t_0 = 0$  and  $\arg(1-t_0) = 0$ , then making successively a loop around 1 in the positive sense, around 0 in the positive sense, around 1 in the negative sense, and around 0 in the negative sense, and finally returning at  $t_0$  (see picture, for instance, in [10, p. 272]). The prototypical evaluation of such an integral was done by Pochhammer [23] for the generalized beta integral, see also [10, 1.6(7)]:

$$\frac{1}{(1-e^{2\pi ia})(1-e^{2\pi ib})} \int^{(1+,0+,1-,0-)} t^{a-1} (1-t)^{b-1} dt = \frac{\Gamma(a)\Gamma(b)}{\Gamma(a+b)} \quad (a, b \in \mathbb{C} \setminus \mathbb{Z}). \quad (5.1)$$

For  $\operatorname{Re} a, \operatorname{Re} b > 0$  it can be reduced to the classical beta integral by shrinking the loops around 0 and 1.

Since the specifications for the loop are only of topological nature, the actual choice of the loop allows much freedom. However, a convenient standard form, in particular if we want to shrink the loop, is to take  $\varepsilon \in (0, \frac{1}{2})$ , define  $C_\varepsilon(z)$  as the circle of radius  $\varepsilon$  around  $z \in \mathbb{C}$ , and then let the loop be given as follows. Start at the point  $\varepsilon$  on the interval  $[0, 1]$ . First go to  $1 - \varepsilon$  along the interval, then turn around 1 along  $C_\varepsilon(1)$  in the positive sense. Then go from  $1 - \varepsilon$  to  $\varepsilon$  and turn around 0 along  $C_\varepsilon(0)$  in the positive sense. Then go from  $\varepsilon$  to  $1 - \varepsilon$  and turn around 1 along  $C_\varepsilon(1)$  in the negative sense. Finally go from  $1 - \varepsilon$  to  $\varepsilon$  and turn around 0 along  $C_\varepsilon(0)$  in the negative sense.

As already observed by Pochhammer [23], not just the beta integral can be extended to a wider parameter range by integrating over a double loop, see (5.1), but this similarly works for Euler's integral representation (2.2) for the Gauss hypergeometric function, see [10, 2.1(13)]:

$$\begin{aligned} {}_2F_1 \left( \begin{matrix} a, b \\ c \end{matrix} ; z \right) &= \frac{\Gamma(c)}{\Gamma(b)\Gamma(c-b)} \frac{1}{(1 - e^{2\pi i b})(1 - e^{2\pi i(c-b)})} \\ &\times \int^{(1+, 0+, 1-, 0-)} t^{b-1} (1-t)^{c-b-1} (1-tz)^{-a} dt \quad (5.2) \\ &(b, 1-c, c-b \notin \mathbb{Z}_{>0}, z^{-1} \notin [0, 1] \text{ and } z^{-1} \text{ outside the double loop}). \end{aligned}$$

The proof is by reduction to (5.1), where we now assume that  $b$  and  $c - b$  avoid the integers completely and that  $|z| < 1$ . Indeed, deform the double loop such that  $|tz| < 1$  for  $t$  on the loop (allowed because  $z^{-1}$  is outside the loop and  $|z^{-1}| > 1$ ) and expand  $(1 - tz)^{-a}$  as a binomial series. Note also that, for  $\operatorname{Re} b, \operatorname{Re}(c - b) > 0$  and with  $z^{-1} \notin [0, 1]$  and fixed outside the double loop, we can shrink the loops around 0 and 1 in order to arrive at (2.2).

Now consider the double loop integral in (5.2) under the assumption that  $z \notin [0, \infty)$ , that  $z^{-1}$  is not on the loop, and that there is a path connecting  $z^{-1}$  with 0 which does not cross the double loop. Thus we consider

$$\begin{aligned} &\int^{(1+, \{z^{-1}, 0\}+, 1-, \{z^{-1}, 0\}-)} t^{b-1} (1-t)^{c-b-1} (1-tz)^{-a} dt \\ &= (-z)^{-a} \int^{(1+, \{z^{-1}, 0\}+, 1-, \{z^{-1}, 0\}-)} t^{b-a-1} (1-t)^{c-b-1} (1-t^{-1}z^{-1})^{-a} dt =: I(z). \end{aligned} \quad (5.3)$$

Assume that moreover  $|z^{-1}| < 1$  and that  $b - a, c - b \notin \mathbb{Z}$ . Then we can deform the double loop in  $I(z)$  such that  $|t^{-1}z^{-1}| < 1$  for  $t$  on the loop. Binomial expansion, interchange of sum and integral by dominated convergence, and application of (5.1) give that

$$I(z) = \frac{\Gamma(b-a)\Gamma(c-b)}{\Gamma(c-a)} (1 - e^{2\pi i(b-a)})(1 - e^{2\pi i(c-b)})(-z)^{-a} \sum_{k=0}^{\infty} \frac{(a)_k (a-c+1)_k}{(a-b+1)_k k!} z^{-k}.$$

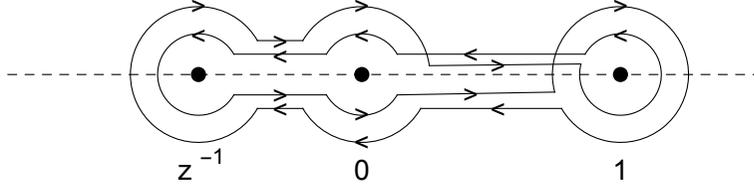


FIGURE 3. Double loop prepared for shrinking in the integral in (5.4).

 TABLE 1. Essential data while  $t$  goes through the double loop.

|              | $(0,1)$<br>$\rightarrow$ | $(0,1)$<br>$\leftarrow$ | $(z^{-1},0)$<br>$\leftarrow$ | $(z^{-1},0)$<br>$\rightarrow$ | $(0,1)$<br>$\rightarrow$ | $(0,1)$<br>$\leftarrow$ | $(z^{-1},0)$<br>$\leftarrow$ | $(z^{-1},0)$<br>$\rightarrow$ |
|--------------|--------------------------|-------------------------|------------------------------|-------------------------------|--------------------------|-------------------------|------------------------------|-------------------------------|
| $\arg(t)$    | 0                        | 0                       | $\pi$                        | $\pi$                         | $2\pi$                   | $2\pi$                  | $\pi$                        | $\pi$                         |
| $\arg(1-t)$  | 0                        | $2\pi$                  | $2\pi$                       | $2\pi$                        | $2\pi$                   | 0                       | 0                            | 0                             |
| $\arg(1-tz)$ | 0                        | 0                       | 0                            | $2\pi$                        | $2\pi$                   | $2\pi$                  | $2\pi$                       | 0                             |

Thus we have shown that

$$\begin{aligned}
 & \frac{\Gamma(c)\Gamma(b-a)}{\Gamma(b)\Gamma(c-a)} (-z)^{-a} {}_2F_1 \left( \begin{matrix} a, a-c+1 \\ a-b+1 \end{matrix}; z^{-1} \right) \\
 &= \frac{\Gamma(c)}{\Gamma(b)\Gamma(c-b)} \frac{1}{(1-e^{2\pi i(b-a)})(1-e^{2\pi i(c-b)})} \\
 & \quad \times \int^{(1+, \{z^{-1}, 0\}+, 1-, \{z^{-1}, 0\}-)} t^{b-1} (1-t)^{c-b-1} (1-tz)^{-a} dt \tag{5.4}
 \end{aligned}$$

( $b-a, c-b \notin \mathbb{Z}$ ,  $c \notin \mathbb{Z}_{\leq 0}$ ,  $z \notin [0, \infty)$ ,  $z^{-1}$  connected with 0 inside the double loop).

It is also possible to shrink the double loop in (5.4) and thus arrive at a sum of two classical Euler type integrals. For convenience assume that  $z < 0$  and deform the double loop such that it consists of segments  $[z^{-1} + \varepsilon, -\varepsilon]$ ,  $[\varepsilon, 1 - \varepsilon]$ , circles of radius  $\varepsilon$  around  $z^{-1}$  and 1, and two half circles of radius  $\varepsilon$  around 0, where each part is traversed various times and in different directions (see Figure 3). In Table 1 we list the arguments of the factors in the integrand for  $t$  on the various parts of the double loop.

Also assume that  $\operatorname{Re} a < 1$  and  $\operatorname{Re} c > \operatorname{Re} b > 0$  with  $b-a, c-b \notin \mathbb{Z}$ ,  $c \notin \mathbb{Z}_{\leq 0}$ . Then, as  $\varepsilon \downarrow 0$ , the right-hand side of (5.4) tends to

$$\begin{aligned}
 & \frac{\Gamma(c)}{\Gamma(b)\Gamma(c-b)} \left( \int_0^1 t^{b-1} (1-t)^{c-b-1} (1-tz)^{-a} dt \right. \\
 & \quad \left. - \frac{\sin(\pi a)}{\sin(\pi(a-b))} \int_{z^{-1}}^0 (-t)^{b-1} (1-t)^{c-b-1} (1-tz)^{-a} dt \right)
 \end{aligned}$$

$$\begin{aligned}
&= {}_2F_1\left(\begin{matrix} a, b \\ c \end{matrix}; z\right) - \frac{\Gamma(c)\Gamma(a-b)\Gamma(b-a+1)}{\Gamma(b)\Gamma(c-b)\Gamma(a)\Gamma(1-a)} (-z)^{-b} \\
&\quad \times \int_0^1 s^{b-1}(1-s)^{-a}(1-z^{-1}s)^{c-b-1} ds \\
&= {}_2F_1\left(\begin{matrix} a, b \\ c \end{matrix}; z\right) - \frac{\Gamma(c)\Gamma(a-b)}{\Gamma(a)\Gamma(c-b)} (-z)^{-b} {}_2F_1\left(\begin{matrix} b, b-c+1 \\ b-a+1 \end{matrix}; z^{-1}\right), \quad (5.5)
\end{aligned}$$

where we used (2.2) and [10, 1.2(6)]. It follows that (5.5) is equal to (5.4), by which we have recovered [10, 2.10(2)].

*Remark 5.1.* In the homological approach it is quite standard to prove connection formulas as above. See for instance Mimachi [19] for the more general case dealing with  ${}_3F_2$ .

## 5.2. A classical analysis version of Kita's integral for $H_2$

We start with another corollary to Theorem 4.1.

**COROLLARY 5.2.** *For  $\operatorname{Re} a, \operatorname{Re} d, \operatorname{Re}(a+c), \operatorname{Re}(a+d), \operatorname{Re}(e-a-d) > 0$  and  $(x, y) \in \mathbb{C}^2$  such that  $x \notin [1, \infty)$  and  $y \notin [0, \infty)$ , we have*

$$\begin{aligned}
&\frac{\Gamma(a+c)\Gamma(a+d)}{\Gamma(a+c+d)\Gamma(a)} (-y)^{-c} F_P(a+c, b, c, e, c-d+1; x, -y^{-1}) \\
&= \frac{\Gamma(e)}{\Gamma(e-a-d)\Gamma(a)\Gamma(d)} \\
&\quad \times \int_0^1 \int_0^1 u^{a-1}(1-u)^{e-a-1} v^{d-1}(1-v)^{e-a-d-1}(1-ux)^{-b} \\
&\quad \times (1-(1-u)u^{-1}vy)^{-c} du dv \quad (5.6)
\end{aligned}$$

$$\begin{aligned}
&= \frac{\Gamma(e)}{\Gamma(e-a-d)\Gamma(a)\Gamma(d)} \int_{u=0}^1 \int_{v=0}^{u^{-1}-1} u^{a+d-1}(1-ux)^{-b} v^{d-1}(1-vy)^{-c} \\
&\quad \times (1-u-uv)^{e-a-d-1} dv du \quad (5.7)
\end{aligned}$$

$$\begin{aligned}
&= \frac{\Gamma(e)}{\Gamma(e-a-d)\Gamma(a)\Gamma(d)} \int_{u=0}^1 \int_{v=0}^{1-u} u^{a-1}(1-ux)^{-b} v^{d-1}(1-u^{-1}vy)^{-c} \\
&\quad \times (1-u-v)^{e-a-d-1} dv du. \quad (5.8)
\end{aligned}$$

*Proof.* Formula (5.6) is a trivial rewriting of (4.6). We went from (5.6) to (5.7) by transforming  $(u, v) \mapsto (u, uv/(1-u))$  and from (5.7) to (5.8) by  $(u, v) \mapsto (u, u^{-1}v)$ .  $\square$

Expression (5.7) was evaluated as  $H_2(a, b, c, d, e; x, y)$  by Yoshida<sup>†</sup> [26, (0.10)]. Kita [17, p. 57, item 9] gave the same evaluation<sup>‡</sup> for (5.8) as Yoshida gave for (5.7), but he specified that integration should be understood by the homological approach, which changes things drastically. Dwork and Loeser [9, p. 105] also reproduce Yoshida's formula, but without specification of the type of integral.

<sup>†</sup>In [26, (0.10)] a power of  $-(uv+u-1)$  should be taken, because the integral is over a region with  $uv+u-1 \leq 0$ .

<sup>‡</sup>In [17, p. 57, item 9] the second occurrence of  $z_1$  should be  $z_2$ . Also, on [17, p. 27, line 3] the denominator  $\Gamma(c)$  should be  $\Gamma(\gamma)$ .

We will now show that we can reproduce Kita's evaluation of (5.8) as  $H_2$  by working with classical integrals. For our purpose it is easier to work with a variant of the right-hand side of (5.6).

**THEOREM 5.3.** *Let  $\operatorname{Re} d, \operatorname{Re}(e - a - d) > 0$  and  $a, e - a \notin \mathbb{Z}$ . Then*

$$\begin{aligned} H_2(a, b, c, d; e, x, y) &= \frac{\Gamma(e)}{\Gamma(e - a - d)\Gamma(a)\Gamma(d)} \frac{1}{(1 - e^{2\pi i a})(1 - e^{2\pi i(e-a)})} \\ &\quad \times \int^{(1+, 0+, 1-, 0-)} u^{a-1}(1-u)^{e-a-1}(1-ux)^{-b} \\ &\quad \times \left( \int_0^1 v^{d-1}(1-v)^{e-a-d-1}(1-(u^{-1}-1)vy)^{-c} dv \right) du. \end{aligned} \quad (5.9)$$

Here  $(x, y) \in \mathbb{C}^2$  and the double loop are such that  $ux \notin [1, \infty)$  and  $(u^{-1} - 1)y \notin [1, \infty)$  for  $u$  on the double loop.

*Proof.* First choose  $x \in \mathbb{C}$  with  $|x| < 1$ . Then choose a double loop  $(1+, 0+, 1-, 0-)$  (as defined in the beginning of Section 5.1) such that  $x^{-1}$  is not on the double loop and can be connected with  $\infty$  without crossing the double loop. Then  $u^{a-1}(1-u)^{e-a-1}(1-ux)^{-b}$  is well defined on the loop. Since the set  $\{(u^{-1} - 1)v \mid u \in (1+, 0+, 1-, 0-), v \in [0, 1]\}$  is compact, we will have  $|(u^{-1} - 1)vy| < 1$  in the double integral (5.9) for  $|y|$  small enough. Thus, for such  $x, y$ , the whole integrand is well defined, and so is the double integral.

Write the right-hand side of (5.9) as

$$\frac{\Gamma(e)}{\Gamma(e - a - d)\Gamma(a)\Gamma(d)} J.$$

In  $J$  we can expand  $(1-ux)^{-b}$  and  $(1-(1-u)u^{-1}vy)^{-c}$  as binomial series and interchange the double sum with the double integral by dominated convergence. We obtain that

$$\begin{aligned} J &= \sum_{k,l=0}^{\infty} \frac{(b)_k (c)_l}{k! l!} x^k y^l \int_0^1 v^{d+l-1} (1-v)^{e-a-d-1} dv \\ &\quad \times \int^{(1+, 0+, 1-, 0-)} u^{a+k-l-1} (1-u)^{e-a+l-1} du \\ &= \sum_{k,l=0}^{\infty} \frac{(a)_{k-l} (b)_k (c)_l (d)_l}{(e)_k k! l!} x^k y^l = H_2(a, b, c, d, e; x, y). \end{aligned}$$

Analytic continuation in  $x$  and  $y$  with simultaneous deformation of the double loop is possible as long as  $ux \notin [1, \infty)$  and  $(u^{-1} - 1)y \notin [1, \infty)$  for  $u$  on the double loop.  $\square$

### 5.3. The integral (5.9) as a solution of the $F_2$ system and analytic continuation

We can rewrite (5.9) as

$$\begin{aligned}
& y^{-b_2} H_2(a - b_2, b_1, b_2, b_2 - c_2 + 1, c_1; x, -y^{-1}) \\
&= \frac{\Gamma(c_1)}{\Gamma(-a + c_1 + c_2 - 1)\Gamma(a - b_2)\Gamma(b_2 - c_2 + 1)} \\
&\quad \times \frac{y^{-b_2}}{(1 - e^{2\pi i(a-b_2)})(1 - e^{2\pi i(-a+b_2+c_1)})} \\
&\quad \times \int^{(1+,0+,1-,0-)} u^{a-b_2-1}(1-u)^{-a+b_2+c_1-1}(1-xu)^{-b_1} \\
&\quad \times \left( \int_0^1 v^{b_2-c_2}(1-v)^{-a+c_1+c_2-2}(1-y^{-1}(1-u^{-1})v)^{-b_2} dv \right) du =: I(x, y),
\end{aligned} \tag{5.10}$$

where  $\operatorname{Re}(b_2 - c_1 + 1), \operatorname{Re}(-a + c_1 + c_2 - 1) > 0$  and  $a - b_2, -a + b_2 + c_1 \notin \mathbb{Z}$ , while  $|x| < 1$  and  $|y|$  is large enough, with the same branch choice for the two occurrences of  $y^{-b_2}$ . By [22, (14)] the left-hand side of (5.10) is a solution of the system [10, 5.9(10)] of PDEs for  $F_2$ . Actually, see also (3.1), it is the unique solution of this system of the form  $y^{-b_2} \sum_{i,j=0}^{\infty} c_{i,j} x^i y^{-j}$  with  $c_{0,0} = 1$  which is analytic in a neighborhood of  $(x, y) = (0, \infty)$ .

The *singular loci* of the  $F_2$  system are the points around which the local solution space has dimension less than the generic dimension four. It is well known (and easily derived) that the singular loci are the lines  $x = 0, x = 1, y = 0, y = 1, x + y = 1$  and (after one-point compactification of the two  $\mathbb{C}$ -factors of  $\mathbb{C}^2$ )  $x = \infty, y = \infty$ . We conclude that the left-hand side of (5.10) can be certainly analytically extended to a one-valued analytic function on the regions in  $\mathbb{R}^2$  given by  $y > 1, 1 - y < x < 1$  and by  $x < 1, y < 0$ . We will show that this can also be read off from the right-hand side  $I(x, y)$  of (5.10).

First we evaluate the inner integral in  $I(x, y)$  as an Euler integral (2.2):

$$\begin{aligned}
I(x, y) &= \frac{\Gamma(c_1)}{\Gamma(a - b_2)\Gamma(-a + b_2 + c_1)} \frac{y^{-b_2}}{(1 - e^{2\pi i(a-b_2)})(1 - e^{2\pi i(-a+b_2+c_1)})} \\
&\quad \times \int^{(1+,0+,1-,0-)} u^{a-b_2-1}(1-u)^{-a+b_2+c_1-1}(1-xu)^{-b_1} \\
&\quad \times {}_2F_1 \left( \begin{matrix} b_2, b_2 - c_2 + 1 \\ -a + b_2 + c_1 \end{matrix}; y^{-1}(1-u^{-1}) \right) du.
\end{aligned} \tag{5.11}$$

Then the constraints  $\operatorname{Re}(b_2 - c_1 + 1), \operatorname{Re}(-a + c_1 + c_2 - 1) > 0$  for (5.10) can be dropped, while the other constraints are still kept. For analytic continuation in (5.11) we can simultaneously bring  $(x, y)$  to a larger region and deform the double loop. We will certainly have a one-valued analytic continuation as long as, for  $u$  in the double loop,  $xu$  and  $y^{-1}(1-u^{-1})$  do not cross the cut  $[1, \infty)$ . For  $(x, y) \in \mathbb{R}^2$  and in the mentioned domains  $y > 1, 1 - y < x < 1$  or  $x < 1, y < 0$  of unique analytic continuation for  $y^{-b_2} H_2(x, -y^{-1})$  we distinguish four cases. We will see that in each case the two cuts are no obstruction for the double loop.

- (1)  $x \leq 0$  and  $x + y > 1$ . Then  $u$  has to remain outside  $(-\infty, x^{-1}]$  and  $[(1-y)^{-1}, 0]$ , so  $u$  can pass  $\mathbb{R}$  through the holes  $(1, \infty)$  and  $(x^{-1}, (1-y)^{-1})$  (left from 0).

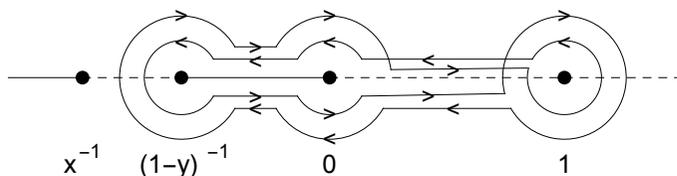
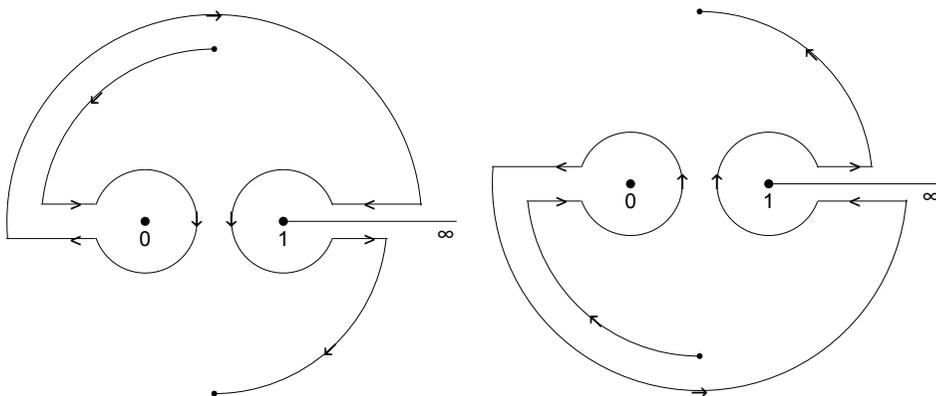


FIGURE 4. Double loop prepared for shrinking in the integral in (5.11).


 FIGURE 5. Orbits of  $(1 - u^{-1})y^{-1}$  during the first and second parts of the double loop.

- (2)  $0 \leq x < 1$  and  $y > 1$ . Then  $u$  has to remain outside  $[(1 - y)^{-1}, 0]$  and  $[x^{-1}, \infty)$ , so  $u$  can pass  $\mathbb{R}$  through the holes  $(1, x^{-1})$  and  $(-\infty, (1 - y)^{-1})$  (left from 0).
- (3)  $x \leq 0$  and  $y < 0$ . Then  $u$  has to remain outside  $(-\infty, x^{-1}]$  and  $[0, (1 - y)^{-1}]$ , so  $u$  can pass  $\mathbb{R}$  through the holes  $(1, \infty)$  and  $(x^{-1}, 0)$ .
- (4)  $0 \leq x < 1$  and  $y < 0$ . Then  $u$  has to remain outside  $[0, (1 - y)^{-1}]$  and  $[x^{-1}, \infty)$ , so  $u$  can pass  $\mathbb{R}$  through the holes  $(1, x^{-1})$  and  $(-\infty, 0)$ .

In a way depending on the case the double loop can be next shrunk, analogous to our treatment of (5.4). As an illustration we will only treat the first case.

#### 5.4. Shrinking the double loop in (5.11)

We will continue with case (1) above. Assume that  $x < 0$  and  $x + y > 1$ . Deform the double loop in (5.11) such that it consists of segments  $[(1 - y)^{-1} + \varepsilon, -\varepsilon]$ ,  $[\varepsilon, 1 - \varepsilon]$ , circles of radius  $\varepsilon$  around  $(1 - y)^{-1}$  and 1, and half circles of radius  $\varepsilon$  around 0, where each part is traversed various times and in different directions (see Figure 4). One has to do good bookkeeping of the arguments of  $u$ ,  $1 - u$  and  $y^{-1}(1 - u^{-1})$  while  $u$  goes through the double loop, since these will affect the integrand. See Figure 5 for the orbit of  $(1 - u^{-1})y^{-1}$  while  $u$  goes through the double loop. See Table 2 for the essential data associated with the double loop.

We already had the constraints  $a - b_2$ ,  $-a + b_2 + c_1 \notin \mathbb{Z}$ . In order to allow the shrinking as  $\varepsilon \downarrow 0$  we also need the constraints

- at  $u = 1$ :  $\text{Re}(b_2 + c_1 - a) > 0$ ;

TABLE 2. Essential data while  $u$  goes through the double loop.

|                        | (0,1)<br>→     | (0,1)<br>←     | $((1-y)^{-1},0)+i0$<br>← | $((1-y)^{-1},0)-i0$<br>→ |
|------------------------|----------------|----------------|--------------------------|--------------------------|
| $\arg(u) =$            | 0              | 0              | $\pi$                    | $\pi$                    |
| $\arg(1-u) =$          | 0              | $2\pi$         | $2\pi$                   | $2\pi$                   |
| $(1-u^{-1})y^{-1} \in$ | $(-\infty, 0)$ | $(-\infty, 0)$ | $(1, \infty) + i0$       | $(1, \infty) - i0$       |
|                        | (0,1)<br>→     | (0,1)<br>←     | $((1-y)^{-1},0)-i0$<br>← | $((1-y)^{-1},0)+i0$<br>→ |
| $\arg(u) =$            | $2\pi$         | $2\pi$         | $\pi$                    | $\pi$                    |
| $\arg(1-u) =$          | $2\pi$         | 0              | 0                        | 0                        |
| $(1-u^{-1})y^{-1} \in$ | $(-\infty, 0)$ | $(-\infty, 0)$ | $(1, \infty) - i0$       | $(1, \infty) + i0$       |

- at  $u = 0$ :  $\operatorname{Re} a, \operatorname{Re}(a - c_2 + 1) > 0$  (use [10, 2.10(2)]);
- at  $u = (1 - y)^{-1}$ :  $\operatorname{Re}(-a - b_1 + c_1 + c_2) > 0$  (use [10, 2.10(1)]).

In the limit for  $\varepsilon \downarrow 0$  there will result the sum of an integral over  $u \in [0, 1]$  and an integral over  $u \in [(1 - y)^{-1}, 0]$ . For both integrals we collect four instances of the integrand in (5.11) with branch choices following Table 2. These four instances are added, where a term has positive sign if the arrow in Table 2 is in forward direction, and negative sign for an arrow in negative direction. We obtain (also use [10, 1.2(6)]):

$$I(x, y) = I_1(x, y) + I_2(x, y) \quad (5.12)$$

with

$$I_1(x, y) = \frac{\Gamma(c_1)}{\Gamma(a - b_2)\Gamma(-a + b_2 + c_1)} y^{-b_2} \int_0^1 u^{a-b_2-1} (1-u)^{-a+b_2+c_1-1} (1-xu)^{-b_1} \times {}_2F_1\left(\begin{matrix} b_2, b_2 - c_2 + 1 \\ -a + b_2 + c_1 \end{matrix}; y^{-1}(1-u^{-1})\right) du \quad (5.13)$$

and

$$I_2(x, y) = \frac{\Gamma(c_1)\Gamma(-a + b_2 + 1)}{\Gamma(-a + b_2 + c_1)} \frac{y^{-b_2}}{2\pi i} \int_{(1-y)^{-1}}^0 (-u)^{a-b_2-1} (1-u)^{-a+b_2+c_1-1} (1-xu)^{-b_1} \times \left( {}_2F_1\left(\begin{matrix} b_2, b_2 - c_2 + 1 \\ -a + b_2 + c_1 \end{matrix}; y^{-1}(1-u^{-1}) + i0\right) - {}_2F_1\left(\begin{matrix} b_2, b_2 - c_2 + 1 \\ -a + b_2 + c_1 \end{matrix}; y^{-1}(1-u^{-1}) - i0\right) \right) du. \quad (5.14)$$

In (5.13) we can relax the parameter constraints to those given above at  $u = 0$  and  $u = 1$ . By (4.7) we see that

$$I_1(x, y) = \frac{\Gamma(a)\Gamma(a - c_2 + 1)}{\Gamma(a + b_2 - c_2 + 1)\Gamma(a - b_2)} F_P(a, b_1, b_2, c_1, c_2; x, y). \quad (5.15)$$

In (5.14) we can write the difference of the two  ${}_2F_1$  as one  ${}_2F_1$  in view of [10, 2.10(1)], and next apply Euler's transformation (2.6):

$$\begin{aligned} & \frac{1}{2\pi i} \left( {}_2F_1 \left( \begin{matrix} a, b \\ c \end{matrix}; x + i0 \right) - {}_2F_1 \left( \begin{matrix} a, b \\ c \end{matrix}; x - i0 \right) \right) \\ &= \frac{\Gamma(c)}{\Gamma(a)\Gamma(b)\Gamma(c-a-b+1)} (x-1)^{c-a-b} {}_2F_1 \left( \begin{matrix} c-a, c-b \\ c-a-b+1 \end{matrix}; 1-x \right) \\ &= \frac{\Gamma(c)}{\Gamma(a)\Gamma(b)\Gamma(c-a-b+1)} x^{1-c} (x-1)^{c-a-b} {}_2F_1 \left( \begin{matrix} 1-a, 1-b \\ c-a-b+1 \end{matrix}; 1-x \right) \\ & \quad (x \in (1, \infty)). \end{aligned}$$

Then (5.14) takes the form

$$\begin{aligned} I_2(x, y) &= \frac{\Gamma(c_1)\Gamma(b_2-a+1)}{\Gamma(b_2)\Gamma(b_2-c_2+1)\Gamma(c_1+c_2-a-b_2)} y^{b_2-c_2} \\ & \quad \times \int_{(1-y)^{-1}}^0 (-u)^{a+b_2-c_2-1} (1+u(y-1))^{c_1+c_2-a-b_2-1} (1-xu)^{-b_1} \\ & \quad \times {}_2F_1 \left( \begin{matrix} c_2-b_2, 1-b_2 \\ c_1+c_2-a-b_2 \end{matrix}; 1-y^{-1}(1-u^{-1}) \right) du \\ &= \frac{\Gamma(c_1)\Gamma(b_2-a+1)}{\Gamma(b_2)\Gamma(b_2-c_2+1)\Gamma(c_1+c_2-a-b_2)} (y-1)^{-a} \left( \frac{y}{y-1} \right)^{b_2-c_2} \\ & \quad \times \int_0^1 w^{a+b_2-c_2-1} (1-w)^{c_1+c_2-a-b_2-1} \left( 1 - \frac{xw}{1-y} \right)^{-b_1} \\ & \quad \times {}_2F_1 \left( \begin{matrix} c_2-b_2, 1-b_2 \\ c_1+c_2-a-b_2 \end{matrix}; \frac{y-1}{y}(1-w^{-1}) \right) dw \\ &= \frac{\Gamma(b_2-a+1)\Gamma(a)\Gamma(a-c_2+1)}{\Gamma(b_2)\Gamma(a-b_2+1)\Gamma(b_2-c_2+1)} (y-1)^{-a} \\ & \quad \times F_P \left( a, b_1, c_2-b_2, c_1, c_2; \frac{x}{1-y}, \frac{y}{y-1} \right). \end{aligned} \quad (5.16)$$

Here we have made the change of integration variable  $u = w(1-y)^{-1}$  in the second equality, and we have used (4.7) in the third equality. Furthermore, we can relax in (5.16) the parameter constraints to those given above at  $u = 0$  and  $u = (1-y)^{-1}$ . These are in agreement with the parameter constraints in (4.7) (with  $b_2$  replaced by  $c_2 - b_2$ ).

*Remark 5.4.* The identity (5.12) with (5.10), (5.13) and (5.16) substituted agrees with Olsson's identity [22, (53)].

*Remark 5.5.* If in the present section from (5.11) until here we make the specializations  $x = 0$ ,  $c_1 = a$  then the  ${}_2F_1$  in (5.11) specializes to  $(1-y^{-1}(1-u^{-1}))^{c_2-b_2-1}$  by (2.7). The subsequent decomposition (5.12) (i.e.,  $I(0, y) = I_1(0, y) + I_2(0, y)$ ) gives then another example of what we wrote in Section 5.1 starting with (5.3). There results a formula connecting three Gauss hypergeometric functions (also seen by applying (3.3) and (4.8) to (5.10) and (5.15), (5.16), respectively) which is essentially the connection formula [10, 2.10(3)].

*Remark 5.6.* Kato [16] also discusses the solutions of the  $F_2$  system at the various singular points, but he does not refer to Olsson [22], and he gives, apart from  $F_2$ , few explicit solutions. However, the first double power series in his Remark on p. 330 can be recognized as  $F_P(a, b_2, b_1, c_2, c_1; y, x)$  expanded as (4.4). The three solutions  $I(x, y)$ ,  $I_1(x, y)$  and  $I_2(x, y)$  of the  $F_2$  system occurring above in (5.10), (5.15) and (5.16), respectively, can be matched up to a constant factor with the following solutions in [16]:  $I_1(x, y)$  with the first solution in §3.2.1,  $I_2(x, y)$  with the first solution in §3.3.5, and  $I(x, y)$  with the third solution in §3.3.5. Kato gives many connection formulas, but apparently not the one connecting  $I$ ,  $I_1$  and  $I_2$ .

*Remark 5.7.* Kita's integral representations for Appell  $F_2$  and  $F_3$  [17, p. 56, items 2.2 and 3], which go back to Hattori and Kimura [12], and where the double integral is defined by the homological approach, remain valid if the double integral is defined classically, but then further constraints on the parameters are needed. It would be interesting to consider also the other new double integral representations given by Yoshida [26, (0.6)–(0.9)], which come back in Kita [17] (for Yoshida's alternative  $F_2$ -integral [26, (0.6)] see also Yoshida [27, p. 71]). Are these easy cases like the double integrals for  $F_2$  and  $F_3$  or would their classical counterparts be a sum of two integral representations?

## 6. Double integrals for solutions of the $F_2$ system

Olsson's list [22, p. 1289, Table I] of solutions of the system of PDEs [10, 5.9(10)] for  $F_2$  includes, besides  $F_2$  itself, for instance functions expressible in terms of  $F_3$ , of  $H_2$ , and of  $F_P$ . It is interesting to rewrite the well-known double integral representations† [10, 5.8(2),(3)] for  $F_2$  and  $F_3$  and our double integral representations (3.6) for  $H_2$  and (4.6) for  $F_P$  as integral representations for solutions of the  $F_2$  system. We obtain

$$\begin{aligned}
 & F_2(a, b_1, b_2, c_1, c_2; x, y) \\
 &= \frac{\Gamma(c_1)\Gamma(c_2)}{\Gamma(b_1)\Gamma(b_2)\Gamma(c_1 - b_1)\Gamma(c_2 - b_2)} x^{1-c_1} y^{1-c_2} \\
 &\quad \times \int_{u=0}^x \int_{v=0}^y u^{b_1-1} v^{b_2-1} (1-u-v)^{-a} (x-u)^{c_1-b_1-1} (y-v)^{c_2-b_2-1} dv du \\
 &\quad (0 < x < 1 - y < 1, \operatorname{Re} b_1, \operatorname{Re} b_2, \operatorname{Re}(c_1 - b_1), \operatorname{Re}(c_2 - b_2) > 0), \tag{6.1}
 \end{aligned}$$

$$\begin{aligned}
 & x^{-b_1} y^{-b_2} F_3(b_1, b_2, 1 + b_1 - c_1, 1 + b_2 - c_2, b_1 + b_2 - a + 1; x^{-1}, y^{-1}) \\
 &= \frac{\Gamma(b_1 + b_2 - a + 1)}{\Gamma(b_1)\Gamma(b_2)\Gamma(1 - a)} x^{1-c_1} y^{1-c_2} \\
 &\quad \times \int_{v=0}^1 \int_{u=0}^{1-v} u^{b_1-1} v^{b_2-1} (1-u-v)^{-a} (x-u)^{c_1-b_1-1} (y-v)^{c_2-b_2-1} du dv \\
 &\quad (x, y > 1, \operatorname{Re} b_1, \operatorname{Re} b_2, \operatorname{Re}(1 - a) > 0), \tag{6.2}
 \end{aligned}$$

†In [10, 5.8(3)], in the exponent of  $(1 - u - v)$ , one should replace  $-\gamma$  by  $\gamma$ .

$$\begin{aligned}
& y^{-b_2} H_2(a - b_2, b_1, b_2 - c_2 + 1, b_2, c_1; x, -y^{-1}) \\
&= \frac{\Gamma(b_2 - a + 1)\Gamma(c_1)}{\Gamma(1 - a)\Gamma(b_2)\Gamma(b_1)\Gamma(c_1 - b_1)} x^{1-c_1} y^{1-c_2} \\
&\quad \times \int_{u=0}^x \int_{v=0}^{1-u} u^{b_1-1} v^{b_2-1} (1 - u - v)^{-a} (x - u)^{c_1-b_1-1} (y - v)^{c_2-b_2-1} dv du \\
&\quad (0 < x < 1, y > 1, \operatorname{Re}(1 - a), \operatorname{Re} b_2, \operatorname{Re} b_1, \operatorname{Re}(c_1 - b_1) > 0), \tag{6.3}
\end{aligned}$$

$$\begin{aligned}
& |x|^{1-c_1} y^{1-c_2} F_P(a - c_1 - c_2 + 2, b_1 - c_1 + 1, b_2 - c_2 + 1, 2 - c_1, 2 - c_2; x, y) \\
&= \frac{\Gamma(a + b_2 - c_1 - c_2 + 2)\Gamma(2 - c_1)}{\Gamma(a - c_1 - c_2 + 2)\Gamma(a - c_1 + 1)\Gamma(b_2)\Gamma(1 - a)} |x|^{1-c_1} y^{1-c_2} \\
&\quad \times \int_{u=1}^{\infty} \int_{v=1-u}^0 u^{b_1-1} (-v)^{b_2-1} (u + v - 1)^{-a} (u - x)^{c_1-b_1-1} (y - v)^{c_2-b_2-1} dv du \\
&\quad (x \in (-\infty, 0) \cup (0, 1), y > 0; \\
&\quad \operatorname{Re} a < 1, \operatorname{Re} b_2 > 0, \operatorname{Re}(a - c_1 - c_2 + 2), \operatorname{Re}(a - c_1 + 1), \operatorname{Re} b_2, \operatorname{Re}(1 - a) > 0). \tag{6.4}
\end{aligned}$$

Here we have restricted to domains of  $(x, y)$  in  $\mathbb{R}^2$ . For going from (4.6) to (6.4) make the change of integration variables  $(u, v) \rightarrow (u^{-1}, v/(1 - u))$ .

Note that the right-hand sides of (6.1)–(6.4) all have the form

$$\text{const. } |x|^{1-c_1} |y|^{1-c_2} \iint |u|^{b_1-1} |v|^{b_2-1} |1 - u - v|^{-a} |x - u|^{c_1-b_1-1} |y - v|^{c_2-b_2-1} dv du,$$

where the integral is over a suitable  $(u, v)$  region. This will enable uniform proofs of such identities by showing that the right-hand sides are solutions of the system of PDEs for  $F_2$ , quite similar as was done in [18] in the one-variable case. In Ch. 6 of the PhD Thesis [8] by the first author, the double integrals (6.1)–(6.3) appear as integral kernels in certain approximations of two-dimensional fractional integrals.

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