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A study of the composite supernova remnant MSH 15–56 with *Suzaku*

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Abstract

The Galactic supernova remnant (SNR) MSH 15–56 is a member of the class of composite SNRs that consists of the remnant shell and a displaced pulsar wind nebula (PWN). The earlier X-ray observations reported the comet-like morphology of the PWN and the ejecta distribution of the SNR. In this work, we present a study of MSH 15–56 using archival *Suzaku* data. We investigate the nature of the emission and spectral parameters of the remnant. The X-ray spectra are well fitted with a combination of a thermal and non-thermal model with temperature ~ 0.6 keV and photon index ~ 2.0 . The slightly enhanced abundances of Ne, Mg, S and enhanced abundance of Si confirm the presence of ejected material.

Keywords: ISM: individual objects: MSH 15–56 (G326.3–1.8) – ISM: supernova remnants – X-rays: ISM

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1. Introduction

Composite supernova remnants (SNRs) are characterized by a system containing an expanding shell into a surrounding medium and a pulsar wind nebula (PWN). The X-ray studies of composite SNRs provide detailed information about of these systems (e.g., MSH 11–62: [Slane et al. 2012](#); Kes 75: [Reynolds et al. 2018](#)).

MSH 15–56 (G326.3–1.8, Kes 25) is a member of the class of composite SNRs and has a PWN in a cometary morphology. The bow shock of the PWN runs supersonically through the interstellar medium (ISM). The SNR shell and the displaced plerionic component were found using high-resolution radio observations ([Green et al., 1997](#); [Dickel et al., 2000](#)). The compact hard X-ray nebula at the Southwest (SW) rim of the SNR was discovered by [Plucinsky et al. \(1998\)](#). The SNR has a circular shell structure in the radio band with a radius of $\sim 19'$ ([Green et al., 2014](#)). Its distance was assumed to be ~ 4.1 kpc by a study based on $H\alpha$ radial velocity measurements ([Rosado et al., 1996](#)). In the radio band, it has been revealed that there is a plerion-like feature closer to the center of the SNR ([Clark et al., 1975](#); [Milne et al., 1979](#)). In their *XMM-Newton* and *Chandra* study, [Yatsu et al. \(2013\)](#) detected a point source emitting a power-law spectrum with a photon index of ~ 1.3 , and they found spectral steepening between the spectral photon indices of ~ 1.7 and ~ 2.5 along the flow line from the apex toward the tail, and also an expanding shell at ~ 860 km s $^{-1}$ in the Sedov phase. Beside of the steepening photon index with distance from the putative pulsar, [Temim et al. \(2013\)](#) also detected a thermal component with a temperature of ~ 0.5 keV. Using the radio emission, [van der Swaluw et al. \(2001\)](#) estimated the ratio between the PWN radius and the SNR radius of MSH 15–56 ($R_{\text{pwn}}/R_{\text{snr}}=0.28$) and derived the initial spin period (P_0) of the pulsar residing in the SNR to be 37 ms. Using two different epochs of the archival *Chandra* data, [Temim et al. \(2017\)](#) found a pulsar velocity of ~ 720 km s $^{-1}$, and its direction of motion is $14^\circ \pm 22^\circ$ SW, and the ambient density increases from East to West.

In this work, we present the first results of the high spectral resolution *Suzaku* observation of the SW rim of MSH 15–56. Although the X-ray properties of the thermal and non-thermal emission were studied in detail by [Yatsu et al. \(2013\)](#) and [Temim et al. \(2013\)](#), we study this SNR using more recent atomic database (AtomDB)¹ version 3.0.9 and give a comparison of

¹<http://www.atomdb.org>

metal abundances with XSPEC² and SPEX³ software packages, which provide almost consistent results with each other except a few parameters. We also perform an *Athena* simulation to investigate the ejecta feature within this SNR. The paper is organized as follows. In Section 2, we describe *Suzaku* observation and data reduction. In Section 3, we present spectral analysis and results. In Section 4, we discuss the nature of the X-ray emission and the spectral properties of the SNR. Finally, we present our conclusions in Section 5.

2. Observation and data reduction

We use 76.3 ks archival *Suzaku* data of the SW region of the SNR MSH 15–56 (Obs.ID: 507039010, PI: Y. Yatsu) that was observed on 2013 February 3. The observation was performed by the X-ray Imaging Spectrometer (XIS; Koyama et al., 2007). The XIS instrument consists of four X-ray CCD cameras on the focal planes of the X-Ray Telescope (XRT; Serlemitsos et al., 2007). The *Suzaku* CCDs have a high spectral resolution and low background. The XIS0, 2 and 3 cameras have front-illuminated (FI) CCDs, where the XIS1 is back-illuminated (BI). XIS0, XIS1 and XIS3 are available in this observation. In order to show the location of the regions on the structure of the SNR, we also use a *ROSAT* PSPC image that was obtained with an exposure time of 3975 s on 1991 March 14.

For data reduction and analysis, we used HEADAS software version 6.20 and XSPEC version 12.9.1 (Arnaud, 1996) with AtomDB 3.0.9 (Foster et al., 2012). Furthermore, spectral analysis was done using the SPEX package (Kaastra et al., 1996, 2018) version 3.04.00 with SPEXACT version 3.04.00. We used the xisrmfgen and xissimarfgen tools to create the redistribution matrix file and the ancillary response file, respectively (Ishisaki et al., 2007). All spectra were grouped with a minimum of 30 counts bin⁻¹.

3. Analysis and result

3.1. X-ray image

Figure 1 (top) shows the *ROSAT* image of the SNR MSH 15–56 in the 0.1–2.4 keV energy band. In this figure, the blue elliptical area represents the

²<https://heasarc.gsfc.nasa.gov/xanadu/xspec/>

³<https://www.sron.nl/astrophysics-splex>

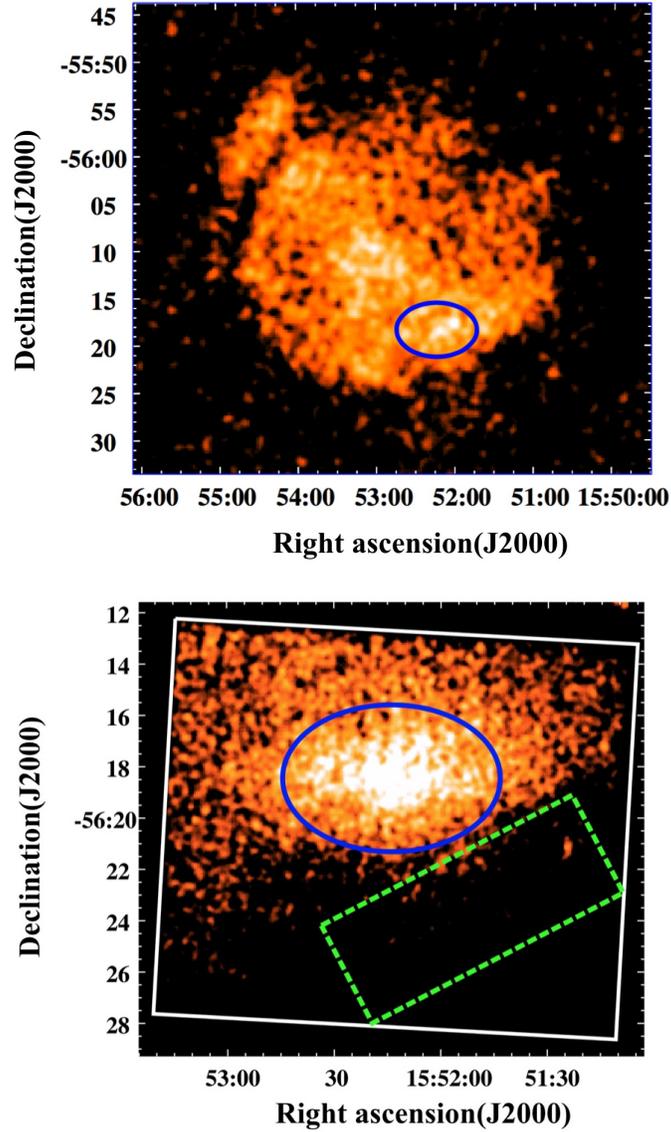


Figure 1: Top: *ROSAT* PSPC image of MSH 15–56 in the 0.1–2.4 keV energy band. The source region that contains SW region is shown by the blue elliptical area and utilized for the *Suzaku* spectral analysis. Bottom: *Suzaku* XIS1 image of the SW rim of MSH 15–56 in the 0.3–10 keV energy band. The SW region is shown by the blue elliptical area. The background region is indicated by the dashed box, and the white box represents the FoV of XIS1. In both images, North is up and East is to left.

SW region that was selected for the *Suzaku* spectral analysis. Subsequently, Figure 1 (bottom) shows the XIS1 image of this region in the 0.3–10 keV energy band. In order to characterize the emission, we extract the spectrum from an elliptical area with 4.2 arcmin \times 2.9 arcmin centered at R.A.(J2000) = 15^h52^m13^s.89, Dec.(J2000) = $-56^{\circ}18'28''.23$, which is shown by blue ellipse on the figure. The background region is shown by the dashed box.

3.2. Background estimation

The background spectrum was extracted from the region near the edge of the rim for the purposes of the spectral modeling (see Figure 1, bottom). The average surface brightness of the background region is up to 23 per cent of those of the source region in the 0.6–7 keV energy band. The Non-X-ray Background (NXB) spectrum for the source and background was estimated from the night-earth data using `xisnxbgen` (Tawa et al., 2008). From the source and background spectrum, we subtracted the NXB spectrum using the `MATHPHA` tool before the spectral fit.

3.3. Spectral fitting

We carried out the spectral fitting in the source region both with `XSPEC` and `SPEX`. We first tried to fit the spectra with a single-component non-equilibrium ionization (NEI) model. The free parameters are the normalization, the absorption (N_{H}), the electron temperature (kT_e), the ionization time-scale ($\tau = n_e t$, where n_e and t represent the mean electron density and the elapsed time after the shock heating of the plasma to a constant temperature, kT_e), and the abundances of Ne, Mg, Si and S, respectively. Because of the statistically unacceptable fit ($\chi^2/\text{d.o.f} > 1.2$), we added a power-law model to search for the non-thermal emission component and eventually obtained statistical improvement on the fits between the 0.6–7 keV energy range ($\chi^2/\text{d.o.f} \sim 1$).

Fitting the spectrum with this model, we were able to satisfactorily characterize the SW region with the following components: power-law plus VNEI model, both coupled with the same absorption model, which was set to vary. We used the absorption models of Tuebingen-Boulder ISM absorption model (TBABS) (Wilms et al., 2000) (in `XSPEC`) and `HOT` (in `SPEX`), (see `SPEX` manual⁴ and Kaastra et al. 2009) to account for the interstellar absorption

⁴<https://www.sron.nl/astrophysics-spex/manual>

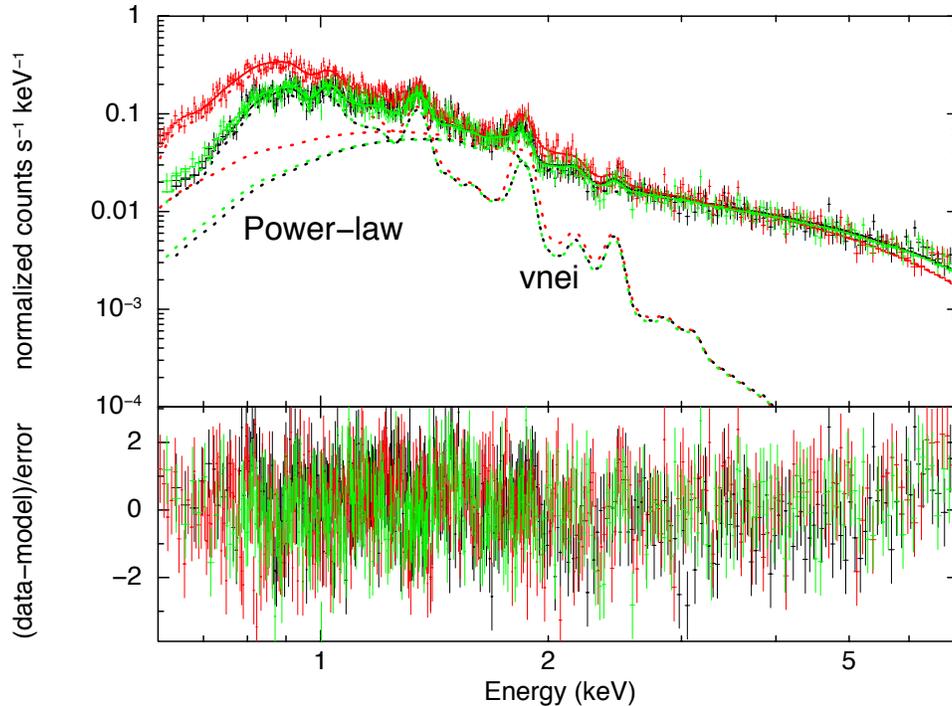


Figure 2: The XIS spectra of the SW region of MSH 15–56 in the 0.6–7 keV energy band fitted with the $\text{tbabs}\times(\text{vnei}+\text{power-law})$. The XIS0, 1 and 3 data are shown by crosses in black, red and green, respectively. The dotted lines represent vnei and power-law components. Residuals are plotted in the bottom panel.

from our Galaxy. To reckon with the neutral plasma gas limit, we fixed the temperature of the `HOT` model to 0.5 eV (Kaastra et al., 2009). According to the residuals and goodness-of-fit, we thawed the elemental abundances of NEI, both in XSPEC (with NEI version 3.0) and SPEX. Correlative to the best-fit parameters, the elemental abundances of Ne, Mg, Si and S are free parameters, while the other metal abundances were fixed to the solar values relative to the abundance tables of Anders & Grevesse (1989) to compare between the XSPEC and SPEX fit directly, which is available in both spectral codes.

The best-fit spectral results are listed in Table 1, and the XIS spectra are displayed in Figure 2. We calculated the uncertainties at the 90 per cent confidence level with the `error` command of XSPEC and SPEX.

Table 1: *Suzaku* spectral fitting results with XSPEC and SPEX.

Component	Parameter	Value	
		XSPEC	SPEX
Absorption	N_{H} (10^{21} cm $^{-2}$)	3.2±0.3	2.2±0.2
NEI	kT_e (keV)	0.64±0.03	0.63 ±0.01
	Ne (solar)	1.6±0.2	1.3±0.1
	Mg (solar)	1.7±0.2	1.8±0.1
	Si (solar)	2.2±0.4	2.3±0.2
	S (solar)	1.7±0.4	1.6±0.5
	$n_e t$ (10^{11} cm $^{-3}$ s)	1.4±0.2	2.0±0.1
	norm † (10^{-3})	1.0±0.2	...
	norm ‡	...	1.1±0.1
Power-law	Γ	2.04±0.04	2.02±0.02
	norm $^{\lrcorner}$ (10^{-4})	7.8±0.4	...
	norm $^{\llcorner}$...	9.9±0.3
	χ^2 (d.o.f.)	1.06 (1119)	1.03 (409)

Notes. Abundances are given relative to the solar values in the abundance table of [Anders & Grevesse \(1989\)](#).

† The normalization (norm) of the thermal model (VNEI in XSPEC) is derived by $10^{-14} \int n_e n_{\text{H}} dV / (4\pi d^2)$ in units of 10^{-3} cm $^{-5}$, where n_e and n_{H} are the electron and Hydrogen densities and V is the volume of the source.

‡ The norm of the thermal model (NEI in SPEX) is the emission measure ($Y \equiv n_e n_{\text{H}} V$) in units of 10^{66} m $^{-3}$.

$^{\lrcorner}$ The norm of the power-law (XSPEC) is in units of photons cm $^{-2}$ s $^{-1}$ keV $^{-1}$ at 1 keV.

$^{\llcorner}$ The norm of the power-law (SPEX) is in units of 10^{45} ph s $^{-1}$ keV $^{-1}$.

4. Discussion

4.1. Thermal emission

We detected thermal and non-thermal emission from the SW rim with *Suzaku* data, which indicates the interactive relation between the PWN and

the thermal emission, pointing out the material in the remnant interior and an interaction of the SNR reverse shock with the PWN (Gaensler & Wallace, 2003; Borkowski et al., 2016). It is therefore likely that the thermal component integrated with the PWN spectrum may be explained by the morphological relation between the PWN and the shell region. This component dominates nearly half of the X-ray spectrum, accounting for ~ 54 per cent of the total unabsorbed flux [$F_{\text{thermal}} = (3.85 \pm 0.74) \times 10^{-12}$ ergs cm^{-2} s^{-1}]. Besides, we found slightly enhanced abundances of Ne, Mg, S and enhanced abundance of Si, which presence of an ejecta feature and confirm the earlier studies. In addition, the abundance table of Anders & Grevesse (1989) which we used during spectral fitting, is an earlier table, but the difference with the recent solar abundances (Lodders & Palme, 2009) is at maximum about 10 per cent for our fitted abundances and is comparable or smaller to the size of the error bars. Thus, impact of using this table on our fitting is not that high.

The *Suzaku* X-ray spectrum was well described by the electron temperature of $kT_e \sim 0.64$ keV, and with an ionization time-scale of $\tau \sim 1.4 \times 10^{11}$ cm^{-3} s, which indicates that the plasma has not yet reached ionization equilibrium. The previous study of Yatsu et al. (2013) has shown that the thermal component in the SW rim has an electron temperature of $kT_e = 0.61$ (0.58–0.64) keV, which is consistent with our result [$kT_e = 0.64$ (0.61–0.67) keV]. Furthermore, compared with the results those of Yatsu et al. (2013), we found higher elemental abundances of Ne, Mg and Si, and the ionization parameter, $n_e t$, is smaller than their result. A note of caution is due here since they extracted spectra from two regions in the same SW rim area, we only selected a single region which is covering them all. Therefore, the difference between the results might be explained by this factor. Temim et al. (2013) also found thermal emission parameters in their spectral regions [$kT_e = 0.62$ (0.42–0.92) keV, for *region b (trail)*, *XMM-Newton*; $kT_e = 0.56$ (0.54–0.57) keV, for *region c*, *Chandra*]. In our fittings, the abundance of Si and S are higher than the solar abundances and these results reflect those of Temim et al. (2013). The comparison of our findings with those of other studies confirms the interrelation between PWN and ejecta material.

Based on the normalization of the vnei model, we estimated the SNR parameters considering the distance of 4.1 kpc. We then calculated the volume of the SNR to be $V \sim 10.7 \times 10^{57} f d_{4.1}^3$ cm^3 , where f is the volume filling factor ($0 < f < 1$) and $d_{4.1}$ is the distance in units of 4.1 kpc. Assuming

$n_e = 1.2n_H$, we found an electron density of $n_e \sim 0.15f^{-1/2}d_{4.1}^{-1/2} \text{ cm}^{-3}$ and age of $\sim 30f^{1/2}d_{4.1}^{1/2} \text{ kyr}$, suggesting that it is a middle-aged SNR. Derived age and electron density results corroborate the finding of [Yatsu et al. \(2013\)](#) and [Temim et al. \(2013\)](#). We then calculated the total mass of the emitting gas, $M_X \sim 2f^{1/2}d_{4.1}^{5/2} M_\odot$ using $M_X = 1.4m_H n_e V$. The small X-ray-emitting mass indicates that the X-ray emission in the SW region arises from the ejecta.

4.2. Non-thermal emission

The non-thermal component of the XSPEC and SPEX models yielded a photon index (Γ) values of 2.04 ± 0.04 and 2.02 ± 0.02 , respectively. These findings are consistent with that of [Temim et al. \(2013\)](#) [$\Gamma = 2.01$ (1.90–2.13), for *region b (trail)*, *XMM-Newton*; $\Gamma = 2.05$ (1.97–2.16), for *region c*, *Chandra*].

The photon index values we found are also slightly higher than those of [Yatsu et al. \(2013\)](#) [$\Gamma = 1.87$ (1.75–1.99)]. On the other hand, the estimated non-thermal X-ray spectrum of a typical PWN has a spectral photon index value in the range of 1.3–2.3 ([Gotthelf et al., 2002](#)), and our result is consistent with this value. From the X-ray emission in [Figure 2](#), it is apparent that there is a hard X-ray continuum above 2 keV, and in the study of [Yatsu et al. \(2013\)](#), it is attributed to the presence of the arc and PWN components.

4.3. Athena X-IFU simulation

The future X-ray mission *Athena* ([Nandra et al., 2013](#)) may be able to reveal the nature of the ejecta of MSH 15–56. With this aim, in this subsection we simulated the *Athena* X-ray Integral Field Unit (X-IFU) ([Barret et al., 2016](#); [Pajot et al., 2018](#)) 50-ks observation using XSPEC and the latest X-IFU response files. [Figure 3](#) shows *Athena* X-IFU spectrum fitted with an absorbed VNEI+power-law model. Our simulation demonstrates the potential of X-IFU for revealing the thermal emission from MSH 15–56. Spectroscopy with X-IFU ($\sim 2.5 \text{ eV}$ energy resolution) would help us separate the thermal and non-thermal components, and constrain the properties of SNR and progenitor.

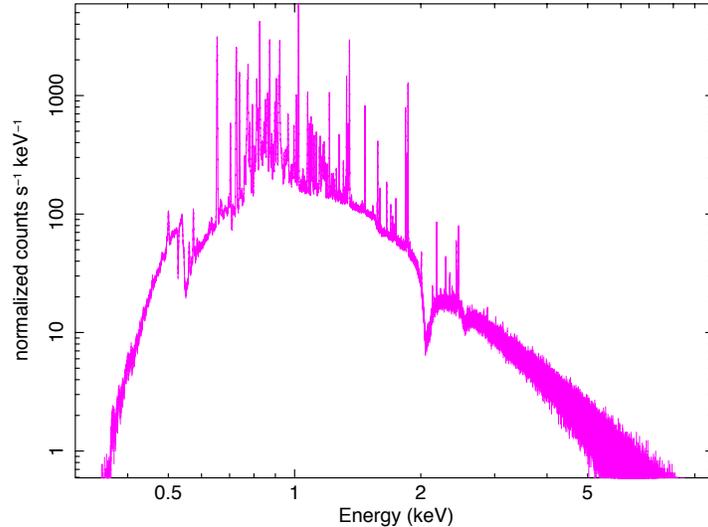


Figure 3: 50-ks *Athena* X-IFU spectrum (an absorbed vnei+power-law model) of the SW region of MSH 15–56 in the 0.3–10 keV energy band.

5. Conclusions

We presented a study of the SW region of the composite SNR MSH 15–56, utilizing the archival *Suzaku* data. We described our main findings as follows.

1. The X-ray spectrum is well described by a two-component model with an absorption of $N_{\text{H}} \sim 3.2 \times 10^{21} \text{ cm}^{-2}$: A non-thermal power-law model with a photon index of ~ 2.04 , and a thermal NEI model with an electron temperature, $kT_e \sim 0.64 \text{ keV}$, an ionization time-scale, $\tau \sim 1.4 \times 10^{11} \text{ cm}^{-3} \text{ s}$.
2. The slightly enhanced abundances of Ne, Mg, S and enhanced abundance of Si support evidence of ejecta heated by the reverse shock. This result together with the small X-ray-emitting mass suggests that its emission arises from the shock heated ejecta.
3. Future X-ray observations with a high-resolution spectrometer such as *Athena* will allow us to examine the properties of thermal emission and the ejecta nature of MSH 15–56.

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