H.E.S.S. observations of the flaring gravitationally lensed galaxy PKS 1830–211


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1 THE PKS 1830–211 GRAVITATIONALLY LENSED QUasar

PKS 1830–211 is a high-redshift (z = 2.5; Lidman et al. 1999) flat-spectrum radio quasar (FSRQ) that has been detected in all wavelengths from radio to high-energy gamma-rays. It is a known gravitationally lensed object with two compact images of the quasar nucleus visible in the radio (Jauncey et al. 1991) and optical (Meylan et al. 2005) passbands. The Einstein ring, well visible at radio frequencies, comes from the imaging of the quasar jet (Kochanek & Narayan 1992). The quasar source is lensed by a foreground galaxy at z = 0.89 (Wiklind & Combes 1996). The angular size of the Einstein ring and separation of compact images is roughly 1 arcsec so that it cannot be resolved with high-energy instruments such as High Energy Stereoscopic System (H.E.S.S.; 50 GeV–50 TeV range) or Fermi-Large Area Telescope (LAT; 100 MeV–100 GeV range). PKS 1830–211 is seen as a bright, high-energy source by the Fermi-LAT instrument and had several flaring periods during the decade of Fermi-LAT observations. PKS 1830–211 is listed in the 1FHL (Ackermann et al. 2013) and the 3FHL (Ajello et al. 2017) catalogues with a photon index above 10 GeV of 3.55 ± 0.34 that corresponds to the average 'low-state' spectrum. No significant curvature in the spectrum was detected. Photons up to 35 GeV, potentially detectable by H.E.S.S., have been observed by Fermi-LAT (Ajello et al. 2017). Observations of these very high energy photons and the measurement of the very high energy tail of the spectrum would give useful constraints on extragalactic background light (EBL) at redshift z = 2.5.

Since the components of the lens cannot be resolved at high or very high energy, the evidence for lensing was searched indirectly on the observed light curve. Because of the different travel paths, the light curves of the two compact components of the lens have a relative time delay, measured in the radio (Lovell et al. 1998) and microwave (Wiklind & Combes 2001) passbands, of 26 ± 5 d. Barnacka, Glicenstein & Moudden (2011) have studied the first 3 yr of the Fermi-LAT light curve with cepstral and autocorrelation
methods. Evidence for a delay of $27.5 \pm 1.3$ d was found with a 3σ significance. The time delay between the compact images of PKS 1830–211 was also studied by the Fermi-LAT Collaboration (Abdo et al. 2015). They selected several flaring periods and calculated the autocorrelation function of the light curve. No significant peak was found. A possible peak of $\sim 20$ d was found with a 1-d binning of the data, which could be attributed to the $\sim 20$ d separation between two flaring events and perhaps to gravitational lensing. Barnacka et al. (2015) have argued that the time delay measured by high-energy instruments could be very different than the value measured by radio telescopes. The delay measured by Lovell et al. (1998) is obtained from the emission of the compact images. Since the jet of the PKS 1830–211 source is imaged close to the Einstein ring, the time difference between the initial burst and its lensed image can be much smaller if the source of high-energy emission is located inside the jet.

PKS 1830–211 is monitored by Fermi-LAT and its light curve is posted on the internet1 on a daily basis. H.E.S.S. observations of PKS 1830–211 were triggered by an alert posted by the Fermi-LAT team on 2014 August 2 (Krauss et al. 2014). The flare seen by the Fermi-LAT instrument started on July 27 and lasted $\sim 4$ d. The H.E.S.S. observations are described in Section 2 and data analysis in Section 3. The H.E.S.S. limits are compared to the Fermi-LAT signal in Section 4 and discussed in Section 5.

2 H.E.S.S. OBSERVATIONS

The very high energy (50 GeV–50 TeV range) gamma-ray observatory of the H.E.S.S. Collaboration consists of five Imaging Atmospheric Cherenkov Telescopes (IACTs) located in the Khomas Highland of Namibia (23°16′18″ S, 16°30′11″ E), 1800 m above sea level. From 2004 January to 2012 October, the array was a four-telescope instrument, with telescopes labelled CT1–4. Each of the telescopes, located at the corners of a square with a side length of 120 m, has an effective mirror surface area of 107 m², and is able to detect cosmic gamma-rays in the energy range 0.1–50 TeV. In 2012, the telescope instrument, with telescopes labelled CT1–4. Each of the telescopes, located at the corners of a square with a side length of 120 m, has an effective mirror surface area of 107 m², and is able to detect cosmic gamma-rays in the energy range 0.1–50 TeV. In 2012, a fifth telescope CT5, with an effective mirror surface area of 600 m² and an improved camera (Bolmont et al. 2014), was installed at the centre of the original square, giving access to energies below 100 GeV (H.E.S.S. Collaboration et al. 2017).

PKS 1830–211 was observed by the five telescopes of the H.E.S.S. IACT array between 2014 August 12 and August 26, to allow for the detection of delayed flares with time delays ranging from 20 to 27 d. The observations were taken at an average zenith angle of $12°$.

3 DATA ANALYSES

This paper is based on a sample of 12.4 h of high-quality data. Data selection cuts have been described in H.E.S.S. Collaboration et al. (2017). Data were next analysed with the Model analysis (de Naurois & Rolland 2009) and cross-checked with the ImpACT analysis (Parsons & Hinton 2014), the two methods giving compatible results. The two analyses use different calibration chains. With both reconstruction chains, data of CT5 were analysed either alone (Mono reconstruction) or combined with the CT1–4 data (Combined reconstruction). The Mono reconstruction has an energy threshold of 67 GeV. The Combined reconstruction has a higher threshold of 144 GeV, but a larger effective area.

A point source is searched at the location of PKS 1830–211. Fig. 1 shows the distribution of the squared angular distance $\theta^2$ of candidate photons from the target position. This distribution, obtained in the Mono analysis, is compared to the background from hadrons misidentified as photons. The background is calculated with the ring background method (Berge, Funk & Hinton 2007), other methods giving similar results.

Table 1 summarizes the number of candidate photons in the signal region, the expected background, and the significance of the excess, calculated with Li and Ma formula 17 (Li & Ma 1983). No significant excess of photons over background is seen by H.E.S.S. at the position of PKS 1830–211. A similar search using the Combined analysis also gives a negative result.

Because of the very soft spectrum measured by Fermi-LAT in the low state, PKS 1830–211 has a chance of being detectable by H.E.S.S. only during flares. The delayed flare lasts only less than about 4 d, however, due to the uncertainties on the date of the flare, it could have happened at any time between August 17 = MJD 56886 (time delay of 20 d) and August 24 = MJD 56893 (radio time delay of 27 d) as explained in Section 1. Fig. 2 shows the evolution over time of significance, binned by 28-min runs. No significant daily photon excess was detected during the H.E.S.S. observation period.

4 FLUX UPPER LIMITS AND COMPARISON TO THE FERMI-LAT SPECTRA

The non-detection by H.E.S.S. translates into 99 per cent confidence level (C.L.) upper limits on the average very high energy flux of PKS 1830–211 during H.E.S.S. observations. These upper limits are shown in Fig. 3. Red (respectively blue) arrows show the limits obtained from the Mono (respectively Combined) analysis and the corresponding solid lines show the effect of deabsorption using the EBL model of Gilmore et al. (2012).

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1 https://fermi.gsfc.nasa.gov/ssc/data/access/lat/mls1jc/source/PKS_1830–211

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**Figure 1.** $\theta^2$ plot of PKS 1830–211 obtained with the Mono reconstruction. The background, shown by crosses, is estimated with the ring background method.

**Table 1.** Analysis results of observations of PKS 1830–211 by H.E.S.S.

<table>
<thead>
<tr>
<th>Reconstruction</th>
<th>$N_{\text{ON}}$</th>
<th>$N_{\text{background}}$</th>
<th>Significance ($\sigma$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Mono</td>
<td>1641</td>
<td>1649.2</td>
<td>−0.2</td>
</tr>
<tr>
<td>Combined</td>
<td>935</td>
<td>954.4</td>
<td>−0.6</td>
</tr>
</tbody>
</table>
H.E.S.S. observations of PKS 1830–211

Figure 2. Significance of the H.E.S.S. signal versus date, obtained with the Mono analysis. The red arrow shows the expected date of the delayed flare for a lensing time delay of 27 d.

Figure 3. 99 per cent C.L. upper limits (arrows) on the PKS 1830–211 flux between 67 GeV and 1 TeV for the 2014 August H.E.S.S. observations. A constant photon index of $-3$ was assumed. The solid lines show the effect of EBL deabsorption, assuming the EBL model of Gilmore et al. (2012).

H.E.S.S. upper limits are compared to Fermi-LAT GeV spectra in Fig. 4. The Fermi-LAT observations have been analysed with Fermi Science Tools v10-r0p5 and Pass8 data, in the Enrico framework (Sanchez & Deil 2013). The spectral data from 2014 July 26 to 30 (flare) are well described by a power-law spectrum with an index of $n_{\text{flare}} = -2.36 \pm 0.17$ for photon energies $>1$ GeV. The relatively high low-energy cut was used to avoid contamination from the Galactic plane. The flare spectrum is much harder than the spectrum measured in the low state of PKS 1830–211, but $n_{\text{flare}}$ is compatible with the photon indices of previous flare spectra, as measured by Fermi-LAT (Abdo et al. 2015). The spectrum of PKS 1830–211 obtained from the Fermi-LAT observations within the H.E.S.S. observation window is well described by a power law with an index of $n_{\text{low}} = -2.97 \pm 0.44$ above 1 GeV. The value of $n_{\text{low}}$ is compatible with the value published in the 3FHL catalogue.

A proper comparison between H.E.S.S. upper limits and the Fermi signal has to take into account the effect of the absorption of the flux of PKS 1830–211 by the EBL and the difference between the flare duration and H.E.S.S. exposure. Since no significant curvature of the spectrum was measured by the Fermi-LAT Collaboration, the unabsorbed spectrum was modelled by a power law. The effect of light absorption by EBL from PKS 1830–211 has been estimated with the models of Gilmore et al. (2012) (black dash–dotted line), Finke et al. (2010) (dotted line), and Franceschini et al. (2008) (black solid line). Fig. 4 shows that there is a substantial difference between the predictions of these models for a source at redshift 2.5 such as PKS 1830–211. Note that EBL absorption could also be affected by the lens environment. Light from lensed active galactic nucleus (AGN) is expected to be more absorbed than average, due to the presence of galaxies along the line of sight. Indeed, absorption from the intervening galaxy has been detected by de Rosa et al. (2005) in the X-ray spectrum of PKS 1830–211. However, Barnacka, Böttcher & Sushch (2014) and Böttcher & Thiersen (2016) have argued that gravitational lensing could help gamma-rays from a distant source avoiding excess absorption. The Fermi-LAT flare spectrum from Fig. 4 is a 4-night average while the H.E.S.S. exposure amounts to 10 nights of data taking. The steady source upper limits from Fig. 3 are thus a factor of $\sim \sqrt{10}/4$ too constraining, which is corrected for in Fig. 4.

5 CONCLUSION

No significant delayed flare from PKS 1830–211 was detected by either H.E.S.S. or Fermi-LAT. The flare did not repeat or was too faint to be detected. Fig. 4 shows, however, that the detection of a strong flare would have been possible close to the Mono analysis energy threshold if the level of EBL absorption was at or below the absorption predicted by the model of Franceschini et al. (2008). Because of its lensed nature, observation of flaring event of PKS 1830–211 in the TeV passband could be useful to constrain EBL models at redshift as large as 2.5. The detection of the lensing time delay in future very high energy observations would help pinpoint the spatial origin of the high-energy emission (Barnacka et al. 2015). It would also permit more exotic applications such as constraining photon mass (Glicenstein 2017) or testing Lorentz invariance violation (Biesiada & Piorkowska 2009).
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