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Wear particle dynamics drive the difference between repeated and non-repeated reciprocated sliding

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\textbf{A B S T R A C T}

The dependence of the sliding mode (repeated vs. non-repeated reciprocated sliding) on the friction and wear behavior of ball-on-flat, brittle non-metallic interfaces in ambient air conditions is evaluated. Repeated sliding promotes the formation of a third body (compressed wear particles) that stabilizes the friction. Non-repeated sliding shows reduced evidence of third body formation, and instead a steady increase in friction. The proposed mechanism driving the non-repeated friction behavior is attributed to a gradual reduction in the ball surface roughness, leading to an increased area of real contact and greater capillary bridge forming across non-contact regions of the interface.

1. Introduction

Friction and wear are of paramount importance to the performance and lifetime of applications with high economic and societal impact such as engines, wheels and industrial production machines [1]. Estimates show that tribological contacts consume 23% of the world energy budget: 20% is spent on overcoming friction and 3% on repairing or substituting components that have worn as a result of friction [2]. These staggering numbers and the applications that they represent have motivated numerous experiments that aim to understand, manipulate and reduce friction and wear under industrially relevant conditions, mimicking for instance the behavior of combustion engines [3], turbines [4], train wheels [5], robotic grasping [6] and micro- and nano-electromechanical systems [7]. While some of these applications involve repeated, reciprocated sliding of the same two surfaces against each other, others slide in a non-repeated fashion, meaning that the slider always contacts a fresh, unworn counter surface. Examples of the latter include rolling wheels [8], walking [9], grabbing [10] (grabbing or clamping typically involves slip due to misalignment and curvature of the touching surfaces), cutting [11], read/write cycles of hard drives [12] and even atomic force microscopy imaging [13,14]. Interestingly, experiments that are designed to reproduce the frictional behavior observed in various applications are almost exclusively performed in a repeated fashion [3,4,18] — also when the application involves non-repeated sliding [5,8,11,19]. Furthermore, the potential consequences this has for the tribological behavior is often overlooked. Here, we systematically study the friction and wear behavior of various types of nominally dry contacts between brittle non-metallic materials sliding in repeated or non-repeated fashion.

2. Experiment

In the ball-on-flat friction experiments (Fig. 1a inset), Al\textsubscript{2}O\textsubscript{3} (sapphire), silicon carbide (SiC) and glass balls were slid against Si wafers or glass flats in repeated and non-repeated fashion in ambient air (21–23 C, 20–60% relative humidity). These two different sliding modes are illustrated in the inset of Fig. 1a: in non-repeated sliding the ball was lifted after each forward and backward stroke (1 cycle), and placed back into contact with an untouched part of the flat. During repeated sliding the ball was reciprocated at a fixed location on the flat. In both the repeated and the non-repeated experiments, the stroke length (L), normal force (F\textsubscript{n}) and sliding speed (v) were kept constant at 20 mm, 0.1 N and 0.5 mm/s, respectively. These materials and experimental parameters were chosen to resemble the silicon wafer-on-support contacts that limit the positioning accuracy in nanolithography machines [20]. The used materials and their RMS roughness

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are further described in Table 1. The tribological experiments were carried out using a Universal Mechanical Tester (UMT TriboLab, Bruker) that was set to acquire force and position data at a rate of 5 Hz. Each friction experiment consisted of 150 cycles, totaling a sliding distance of 6 m. In the non-repeated experiments, the time required to move the ball between subsequent cycles was approximately 27 s.

To characterize the wear of the sliding bodies, optical focus variation profilometry measurements were performed ex-situ using a laser scanning confocal microscope (Keyence VK-X1000). Furthermore, the wear scars on the balls and the wear tracks on the Si flats were imaged and analyzed using scanning electron microscopy (SEM, FEI Verios 460) and SEM-integrated energy dispersive X-ray (EDX, Oxford) measurements. The surface topography of the sliding bodies was measured by tapping mode atomic force microscopy (AFM, Bruker Dimension Icon) where the nominal tip radius of AFM tips carried out in the measurements was 8 nm.

3. Results and discussion

Fig. 1a shows the evolution of the measured friction force \( F \) as a function of the sliding distance for the repeated and non-repeated sapphire-on-Si wafer experiments. In the repeated sapphire-on-Si wafer experiments strong fluctuations in friction were initially observed, followed by a more stable friction force of roughly 50 mN after the first ~0.2 m of sliding (see Fig. S1a). The non-repeated version of the otherwise identical experiment resulted, in the majority of cases, in clearly different frictional behavior: the initial friction fluctuations observed during repeated sliding were not present and the friction force increased gradually with sliding distance. After 6 m of sliding the friction force had doubled, from ~55 mN to ~110 mN. It should be noted that this behavior was not observed in all experiments; in a few cases we observed run-in behavior followed by a stable friction force, see Fig. S2. This is discussed later in relation to third body formation. In the SiC-on-Si wafer experiments (Fig. 1b), the frictional behavior is qualitatively similar to that observed in the sapphire-on-Si wafer experiments. In the repeated SiC-on-Si wafer experiment, unstable friction run-in behavior during the first ~0.4 m of sliding is followed by a more stable evolution of the friction force. In the non-repeated experiment there is no run-in behavior and the friction force gradually increases with sliding distance. The relative increase in friction is, however, less pronounced in the SiC-on-Si wafer experiment than the sapphire-on-Si wafer experiment. Furthermore, in the repeated SiC experiment, the friction gradually decreases after the run-in phase (up to ~0.4 m sliding, Fig. S1b).

To characterize the wear behavior of the studied interfaces, optical profilometry imaging of the contact region of the balls was performed both before and after the friction experiments (Fig. 2, S4 and S5). The quantitative wear results are given in Table 2, and a summary of all predominant wear mechanisms can be found in Table S1. The profilometry images clearly demonstrate that a spherical cap is worn off from the sapphire ball during the non-repeated friction experiment, resulting in an average specific wear rate of \( K \approx 2.8 \times 10^{-10} \text{ mm}^3/	ext{Nm} \) (based on a minimum of three independent experiments); \( K = V/F_{\text{app}} \text{L} \) where \( V \) is the wear volume, determined based on the profilometry data. The wear scar on the sapphire ball that has undergone repeated sliding looks very different; a substantial amount of compressed debris — or third body [21] — has accumulated on the region of the ball that was in contact with the wafer. This was reproducible for three independent identical experiments. The third body was loosely attached and could easily be removed from the ball surface by cleaning for 60 min in an ultrasonic bath containing acetone (Fig. 2c). Removing the third body revealed that the ‘repeated’ ball had obviously worn less than in the non-repeated experiment: the average wear rate for the sapphire balls during repeated

<table>
<thead>
<tr>
<th>Ball</th>
<th>Sliding mode</th>
<th>Third body location</th>
<th>Third body origin</th>
</tr>
</thead>
<tbody>
<tr>
<td>Sapphire Si wafer</td>
<td>Repeated</td>
<td>1.0 0.4</td>
<td>Predominantly adhered to ball</td>
</tr>
<tr>
<td>SiC Si wafer</td>
<td>Non-repeated</td>
<td>6.5 2.8</td>
<td>Not observed^</td>
</tr>
<tr>
<td>S10</td>
<td>Non-repeated</td>
<td>9.7 0.5</td>
<td>Predominantly adhered to Si wafer</td>
</tr>
<tr>
<td>S11</td>
<td>Non-repeated</td>
<td>7.7 1.3</td>
<td>Small amount^ adhered to ball &amp; Si wafer</td>
</tr>
</tbody>
</table>

^ Note that, in a few exceptional cases where the relative humidity of the air environment was slightly higher, a small amount of third body formation was observed on the ball; this was much less than in the repeated case.

^ Much less than in the repeated case.
within sliding (a), Fig. 2.

Ex-situ experiment, calculated ball ball than to the third body, as compared to the background which consisted of Al$_2$O$_3$ (sapphire). (k) SEM image of the wear track and the Si wafer after the repeated SiC experiment (h-j). The orange dashed lines indicate the area within the wear track with increased silicon oxide content. The topography of the wear track matches that of the worn SiC sphere (see also Fig. 3). Scale bars, 50 μm.

sliding was $K = 1.0 \times 10^{-3}$ μm$^3$/Nm: an average of six times less than that measured for the non-repeated experiment.

To understand the nature of the third body formed on the sapphire ball during the repeated experiment, EDX analysis was performed on the ball from Fig. 2b (see Fig. 2g). The measurements indicate that the third body consists mainly of silicon and oxygen and no significant amounts of aluminium, strongly indicating that the third body is wear debris, primarily SiO$_x$ and originating from the Si wafer (Table 2). This is supported by SEM-EDX analysis of the Si flat performed after a repeated experiment, which clearly shows that SiO$_x$ debris particles remain next to the wear track but not inside the wear track (Fig. 5d-i). It is unlikely that the SiO$_x$ originates from accumulation of the native oxide, as the calculated volume of the SiO$_x$ in the contact region on the sapphire ball (Fig. 5f) was ~470 $10^3$ μm$^3$, which is much greater than the corresponding approximate volume of native oxide on the Si wafer wear track (20 $10^3$ 60 2 $10^3$ $2.4 \times 10^3$ μm$^3$). More SiO$_x$ debris was observed at the ends and either side the wear track (Fig. 6a, c), than in the central region of the flat. The large amount of debris suggests that an abrasive wear mechanism occurred during sliding. It is unclear whether the silicon wear debris found on the contacts becomes oxidized before, during or after attachment to the ball.

While in the SiC-on-Si wafer experiments the frictional behavior is qualitatively similar to that observed in the sapphire-on-Si wafer experiments, an important difference is that the third body is predominately located in the central region on the wear track of the Si wafer (Fig. 3b and c) rather than the SiC ball after the repeated sliding experiment (Fig. 2i and j). The third body in this case comprises a ridge of compressed SiO$_x$C$_y$ debris, likely originating from oxidized Si debris...
from the wafer and small fragments of SiC from the ball (Table 2). Furthermore, the profile of this third body, recorded perpendicular to the sliding direction, perfectly matched that observed on the SiC ball (along the same direction, Fig. 3d). This strongly suggests that in the repeated sliding SiC-on-Si wafer experiment, the sliding was accommodated at the ball-on-third body interface. This would then also explain why the difference in ball wear between repeated and non-repeated SiC-on-wafer experiments was much smaller than that in the sapphire-on-Si wafer experiments, in which the SiO$_2$ third body adhered to the ball (in repeated experiments). The SiO$_2$ third body may adhere more strongly to the sapphire ball than to the SiC ball because the sapphire ball was much smoother and therefore more susceptible to physical or capillary adhesion [23]. In all non-repeated experiments, we do not observe such dominant SiO$_2$ third bodies within the wear track on the wafer or on the ball, as seen for repeated experiments.

Our ex-situ observations thus suggest that the observed difference in friction between repeated and non-repeated sliding is related to the formation (or absence) of a third body at the sliding interface. To test this more directly, the friction experiments were repeated, but with different materials, in a microscopy setup (Fig. 4a) consisting of a rheometer mounted on top of an inverted microscope. In this setup, a glass ball is fixed off-center to the rheometer tool which can be rotated

Fig. 4. Imaging of the interface between a glass ball and a float glass coverslip and ex-situ optical images and height profiles recorded before and after the friction experiments.

(a) Visualization/friction experiments were performed using an inverted confocal laser scanning microscope on top of which a rheometer was mounted. The glass ball is fixed to the rheometer tool at a distance of 12.98 mm from the rotation axis. By imposing an angular velocity of 3.52 rad/s, the ball is forced to slide with a velocity of 0.5 mm/s and makes 1 mm strokes. (b), (d) Initial contact at the interface and (b), (j) height profile of glass ball before sliding. Debris is collected at the interface (c) and is observed on the ball (i) and on the substrate (f) after repeated sliding. After non-repeated sliding, there is no visible debris at the interface (e) or on the ball (k), some debris is left on each of the sliding tracks (g). In both experiments the apparent contact area increases as a result of the wear. The contact force is 100 mN in all images. White and black scale bars are 50 μm and 100 μm, respectively. The color scale in (h–k) is identical to that used in Fig. 2.
and moved vertically, enabling measurement of and control over the normal and tangential force exerted at the ball-on-flat interface [24]. As with the UMT experiments, the ball diameter was 4 mm, the normal force was kept at 0.1 N and the sliding speed was 0.5 mm/s. To enable in situ visualization of the interface, the repeated and non-repeated experiments were performed with transparent materials: glass balls on glass flats. The microscope illuminates and images the ball-on-flat interface through the transparent flat. In a typical microscopy image of the interface, we observe a black central region enclosed by interference fringes, the first of which corresponds to a gap of 114.5 nm between the ball and the flat [25]. In the central black region the surfaces have approached to within 114.5 nm, this is the apparent area of contact.

Fig. 4b – e displays the contact images obtained before and after 40 cycles of repeated or non-repeated sliding. Even though the materials are different, it was observed that the glass-on-glass interfaces behave qualitatively similar to the sapphire or SiC-on-Si wafer interfaces studied above: wear debris is collected at the interface as a third body during repeated sliding (Fig. 4c), but not — or much less so — during non-repeated sliding (Fig. 4e). Furthermore, the friction measurements showed that non-repeated sliding results in a clear increase of friction with sliding distance, while repeated sliding does not, or much less so (Fig. 1b inset). Hence we observe that in three different materials systems the friction force measured during a repeated sliding experiment stabilizes after a run-in phase, a phenomenon that has also been observed elsewhere [26], while that measured during non-repeated sliding gradually increases with sliding distance. Based on the microscope and profilometry measurements (Fig. 4b – k) it can be concluded that, during the run-in phase observed in repeated experiments, debris particles were compressed into a third body that subsequently stabilizes the friction: supplementary Movies S1 – S3 show a typical example of the glass-on-glass interface recorded during repeated sliding (S1), non-repeated sliding (S2) and lifting of the sphere (S3) after repeated sliding. It is important to note that while non-repeated sliding suppresses the formation of a third body, we still observed a third body in some non-repeated sapphire-on-Si wafer experiments that were carried out in an air environment with a slightly higher average relative humidity of 45%, compared to other experiments at 34%. For these particular experiments, the friction measurements showed run-in behavior similar to that in the repeated experiments (Fig. S2). As mentioned previously, the exceptionally low roughness of the sapphire balls (Fig. 5 and S8) may enable wear debris to adhere to the ball surface through capillary or physical adhesion, even in some of the non-repeated experiments.

Supplementary video related to this article can be found at https://doi.org/10.1016/j.triboint.2019.105983.

Our results thus show that there is much less chance of third body formation at the interface during non-repeated sliding. In this case, it is important to consider the mechanism responsible for the gradual increase in friction observed in such experiments (e.g. Fig. 1). To address this question, the wear behavior for sapphire-on-Si wafer interfaces was analyzed in more detail.

The effect of the wear should be considered in terms of the continuous change in the real contact pressure and real area of contact occurring at the sliding interface during the wear process, which is difficult to determine experimentally. Ignoring roughness, the change in contact pressure between the start and end of the experiment can be estimated by calculating the Hertzian contact stress ($P_{\text{Hertz}}$) at the start and, assuming a fully conforming contact, the contact pressure at the end of the experiment ($P_{\text{final}}$). This gives a change in contact pressure from $P_{\text{Hertz}} = 428$ MPa to $P_{\text{final}} = 11.4$ MPa, where the latter value is based on the area of the flat cap worn off of the ball in Fig. 5e. Note that the real contact pressure at the start of the experiment is likely to be higher than $P_{\text{Hertz}}$ because the ball and wafer roughness is not taken into account. To estimate the effect of roughness, we carry out boundary element model (BEM) calculations [27], which in this case were performed using the Tribology Simulator that is publicly available at

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Fig. 5. The surface topography before and after a non-repeated sapphire-on-Si wafer experiment.
(a) SEM image of the last (150th) wear track on the wafer. (b) AFM line scan of the last wear track on the wafer recorded perpendicular to the sliding direction. The green arrow indicates the edge of the wear track at which debris particles were occasionally found. Inset figures show the surface roughness outside and inside the wear track. (c) Microscopy image of the sapphire ball. The width of the wear scar on the ball matches that of the track on the wafer (a). (d) AFM line profiles measured in- and outside the wear scar on the sapphire ball. The roughness of the wear scar ($R_{q} = 2.6$ nm) is significantly smaller than that measured on the unworn sapphire ($R_{q} = 9.0$ nm). Scale bars, 50 μm.
Tribonet [28] and using data from optical profilometry of the ball before and after the sliding experiment. These simulation results are limited by the resolution and quality of the optical profilometry measurements, which were carried out with a resolution of approximately 50 nm in lateral and 20 nm in height directions. The results of the calculated contact gap profiles are shown in the supplementary information Figs. S9a and S9b and yield real contact pressures of 2.62 GPa and 0.35 GPa for, respectively, the initial contact and after sliding. Although, as expected, these surface pressures are much higher than those calculated without roughness, it is evident that a large decrease in contact pressure occurs. Thus, the real contact pressure and real contact area rapidly change as a result of wear, and the wear rate is unlikely to be constant.

A different way to consider the wear behavior is to estimate the rate of material removal on the atomic scale. For rough surfaces, the maximum contact pressure exerted at the sapphire-on-Si-wafer interface can never exceed the hardness of the Si-wafer: this would cause the wafer to plastically deform such that the area of real contact increases and the contact pressure drops below the hardness again. This means that during the (non-repeated) sapphire-on-Si wafer friction experiments the minimum area of real contact \( A_{\text{min}} \) between the ball and the substrate is:

\[
A_{\text{min}} = \frac{F_n}{H}
\]

where \( F_n = 0.1 \text{ N} \) and \( H = 5.1 \text{ GPa} \) is the hardness of the boron doped single crystal silicon \(<100>\) wafer that we used [29], yielding:

\[
A_{\text{min}} = \frac{2 \times 10^{11}}{2 \times 10^{10}} \text{ m}^2 = 5 \times 10^8 \text{ atoms}
\]

Thus, over the whole area of the contact, at least \( 5 \times 10^8 \) atoms on the ball surface must touch and slide over the wafer surface at any time, where \( 2 \times 10^{10} \text{ m} \) is used as an approximate atom-to-atom distance within the ball (note that the goal of this calculation is to provide only order-of-magnitude estimates). Using as an example the total wear volume of the sapphire ball shown in Fig. 5c, measured after a sliding distance of 6 m, we can estimate how many atoms on the ball surface are worn off during a non-repeated experiment:

\[
\frac{3.9 \times 10^5}{2 \times 10^{10}} \text{ m}^3 = 5 \times 10^4 \text{ atoms}
\]

It follows that on average, atoms on the ball surface that make contact with the substrate slide at least 6 \( \mu \)m before they are worn off:

\[
\frac{6 \text{ m}}{5 \times 10^8} = 6 \mu \text{m}
\]

We therefore conclude that the ball wear is very mild and the above analysis suggests that this may even occur atom-by-atom [30,31]: when the ball slides over a distance that is the equivalent of more than tens of thousands of atomic spacings, only a single atomic layer is worn off from the sapphire ball.

The wear tracks that are left on the silicon wafer after the non-repeated experiment were also analyzed using AFM and SEM imaging. Although the very first strokes on the wafer may involve some abrasive wear resulting in scratches on the wafer, no, or very little, evidence of wear was found during the subsequent strokes (Figs. 5a and 5b); the wafer roughness is greater than the height difference measured inside and outside the tracks (Fig. 5b).

These observations and calculations suggest that adhesive friction, controlled by the area of real contact, is the main friction mechanism in the non-repeated experiment. The wear may however be governed by the tribocorrosional behavior at the interface [32]. An adhesion friction mechanism is also supported by the observation that for pristine sapphire-on-Si wafer interfaces the friction force — measured at varying normal forces during 100 \( \mu \)m strokes — is strictly proportional to the normal force (Fig. 6a), despite the fact that the apparent area of contact does not increase linearly with the normal force for ball-on-flat interfaces [33]. To further study the interplay between this coefficient of friction and the wear that occurs in non-repeated experiments, a sapphire ball was subjected to a milling procedure. The ball was worn by sliding it along a 6 m spiral-shaped track over the wafer at a fixed normal force of 50 mN. This spiral-shaped sliding track ensured that the ball always encountered fresh wafer surface during the whole experiment. Based on the observations above, it is reasonable to assume that this type of sliding avoids the build-up of a third body on the sphere; indeed AFM PeakForce Tapping measurements [34] performed after the milling procedure confirm this (Fig. S10). Repetition of the non-repeated friction experiment after the ball had been worn in this way showed that the coefficient of friction had substantially increased (Fig. 6a), similar to the behavior observed in the non-repeated sapphire-on-Si wafer experiments (Fig. 1a). Interestingly, if the surrounding air is replaced with dry nitrogen (N\(_2\)), a lower coefficient of friction (\( \mu = 0.57 \) was measured, a value between that measured before wearing the ball (\( \mu = 0.50 \)) and after milling (\( \mu = 0.65 \)) in ambient conditions.

These experiments can be interpreted as follows: During the milling procedure, and also during non-repeated sliding, the surface roughness on the balls is reduced as a consequence of mild wear (Fig. 5). This reduced roughness leads to a larger area of real contact and with that a greater adhesive friction and friction coefficient [35]. Additionally, as Fig. 6. The friction at a sapphire-on-Si wafer interface.

(a) Friction measurements at varying normal forces were performed in a non-repeated fashion using 100 mN (unidirectional) strokes and a sliding speed of 50 \( \mu \)m/s (black data points). The coefficient of friction increases substantially (red data points) after the ball has been worn (see main text). The coefficient of friction is lower when measured in dry N\(_2\) (blue data points, measured after milling) compared to the ambient measurement. The inset figures show the height profile of the sapphire ball before and after the friction experiments. The color scale is identical to that used in Fig. 2. Scale bar, 50 \( \mu \)m. (b) The area of apparent contact \( A_p \) between a sapphire ball and a Si wafer is calculated using the width \( w \) of the wear tracks on the wafer:

\[ A_p = \frac{w \times w}{2}\pi /4. \]
the ball becomes smoother, the total area, within which the gap between the ball and the flat is only a few nanometers, becomes larger (Fig. 5d). Across such nanometric gaps, the water layers that cover most surfaces in ambient conditions can form capillary bridges [36–38], thereby pulling the surfaces into closer contact and increasing the friction [39]. This capillary effect can however be reversed by changing the atmosphere at the interface to dry N₂, thereby removing the influence of the capillary bridges (Fig. 6a). Furthermore, since the capillary effect is active at locations across the interface with a finite gap, the friction can be expected to correlate with the apparent area of contact: the larger the apparent area of contact, the larger the area within which the contacting surfaces experience adhesion: adhesion becomes significant and dominates the friction force when there is high surface conformity and the average gap between the contacting surfaces, defined as the sum of the RMS roughness of the two surfaces, is less than 10 nm [40]. This is indeed observed; by analyzing the wear tracks on the Si wafer, the apparent area of contact was measured at various stages during the non-repeated sphare-on-Si wafer experiment. The result is plotted in Fig. 6b alongside the friction measurements in ambient conditions to show that the gradual growth of the apparent area of contact closely tracks that of the friction force.

4. Summary and closing remarks

Summarizing, the difference between repeated and non-repeated sliding has been studied systematically for various interfaces between non-metallic brittle materials in ambient air conditions. It has been demonstrated that the sliding mode is a very important parameter: Idential systems were shown to give widely varying friction and wear behavior depending on whether the sliding is repeated or non-repeated; for sapphire-on-Si wafer contacts the Sapphire wear changed by a factor 6 and the friction force varied by a factor 2. In the case of SiC-on-Si wafer, the increase in friction observed during a non-repeated experiment was smaller. It is hypothesized that this difference between sapphire and SiC may be explained by the grain structure of the sintered SiC balls, which is not present in the single crystal sapphire balls; such grain structure may enable the SiC balls to maintain a minimum level of roughness that is larger than that of the wafer. Indeed, the worn SiC surfaces are rougher than the worn sapphire surfaces (Fig. 8B). An additional important difference between the sapphire and SiC experiments is that, in repeated experiments, the SiO₂ third body forms on the ball when a sapphire ball is used but within the Si-wafer wear track when a SiC ball is used, indicating that the wear debris adheres more readily to sapphire than SiC. This difference between sapphire and SiC may again be caused by the surface roughness; since the sapphire surface is smoother than the SiC surface, debris may adhere more readily to the sapphire through physical or capillary adhesion. In fact, we observe that even in some of the non-repeated sapphire-on-Si wafer experiments wear debris can adhere to the sapphire surface (Fig. S2). Indeed those experiments in which we observed this third body formation and the associated run-in behavior, were performed at a slightly higher relative humidity compared to the otherwise identical experiments in which we did not observe run-in. If capillary adhesion is responsible for the sticking of the wear debris, this correlation could be expected. However, further research is clearly needed in order to investigate and fully explain this phenomenon. The sliding distance required to achieve stable friction in the repeated experiments was also longer for SiC than for sapphire. This difference may be a consequence of the fact that in the SiC experiment, the SiO₂ third body needs to form over the entire 20 mm sliding track on the wafer, while in the sapphire experiment it is sufficient to form this SiO₂ body only on top of the ball. Furthermore, because in this SiC experiment the third body is immobilized on the wafer rather than on the ball, the ball can be expected to wear more, as observed (Fig. 2B). It has thus been shown that while repeated sliding steers all systems toward the formation of a third body that stabilizes the friction — which was observed in situ for glass-on-glass interfaces — in non-repeated sliding no evidence for such third body formation is found. In all the studied systems, non-repeated sliding leads to a gradual increase of the friction coefficient with sliding distance. Based on an in-depth analysis of the sapphire-on-Si wafer system, including wear calculations, AFM, contact pressure calculations and dedicated wear experiments, it is proposed here that the mechanism behind the increase in friction with sliding distance is a gradual loss of slider surface roughness that not only increases the (nominally dry) area of real contact, but also leads to more capillary bridges across the interface and potentially allows van der Waals forces to become significant by decreasing the average gap between the surfaces; all of these effects result in higher friction. This interplay between surface roughness, capillarity and the area of real contact is complex, but likely universal since virtually all surfaces are rough and covered by water layers.

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Appendix A. Supplementary data

Supplementary data to this article can be found online at https://doi.org/10.1016/j.triboint.2019.105983.

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