Observation of electroweak production of a same-sign W boson pair in association with two jets in pp collisions at $\sqrt{s}=13$ TeV with the ATLAS detector

Aaboud, M.; ATLAS Collaboration

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Observation of Electroweak Production of a Same-Sign $W$ Boson Pair in Association with Two Jets in $pp$ Collisions at $\sqrt{s}=13$ TeV with the ATLAS Detector

M. Aaboud et al.
(ATLAS Collaboration)

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This Letter presents the observation and measurement of electroweak production of a same-sign $W$ boson pair in association with two jets using 36.1 fb$^{-1}$ of proton-proton collision data recorded at a center-of-mass energy of $\sqrt{s} = 13$ TeV by the ATLAS detector at the Large Hadron Collider. The analysis is performed in the detector fiducial phase-space region, defined by the presence of two same-sign leptons, electron or muon, and at least two jets with a large invariant mass and rapidity difference. A total of 122 candidate events are observed for a background expectation of 69 ± 7 events, corresponding to an observed signal significance of 6.5 standard deviations. The measured fiducial signal cross section is $\sigma^{\text{fid}} = 2.89^{+0.51}_{-0.48}$ (stat)$^{+0.26}_{-0.20}$ (syst) fb.

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The scattering of two massive vector bosons (VBS), $VV \rightarrow VV$ with $V = W$ or $Z$, is an important process for studying the mechanism of electroweak symmetry breaking [1–3]. VBS processes involve quartic gauge-boson self-interactions, and the $s$- and $t$-channel exchanges of a gauge or Higgs boson. The Higgs boson regularizes the VBS amplitude by canceling out the divergencies arising from longitudinally polarized vector bosons at high energy [4,5]. These cancellations depend on the gauge structure of the theory and are exact in the standard model (SM) [6,7]. The present measurement of $W$ boson scattering thus serves as a fundamental probe of the SM electroweak theory.

At the LHC, the VBS final state of two gauge bosons and two jets ($VVjj$) can be produced via two classes of mechanisms. The first class, referred to as strong production, involves both strong and electroweak interactions at Born level, and features diagrams where the incoming partons exchange color, as illustrated in Fig. 1(b). The second class, referred to as electroweak production, involves only weak interactions at Born level [8] and includes VBS diagrams. Figure 1(b) shows a typical VBS diagram where the gauge bosons are radiated off the incoming quarks and then scatter via the quartic self-interaction vertex. In VBS processes, the incoming partons do not exchange color and typically produce the two jets with a large invariant mass and with large rapidity difference [9].

The $W^{\pm}W^{\pm}jj$ final state has the largest ratio of electroweak to strong production cross sections compared to other VBS diboson processes [3]; this is because at leading-order (LO) accuracy in perturbative quantum chromodynamics (QCD) quark-gluon and gluon-gluon initiated diagrams are absent and contributions from quark and (anti-)quark annihilation diagrams are suppressed. This ratio is of order five in the fiducial phase-space region of this analysis. The $s$-channel VBS diagrams with trilinear self-interactions are absent in this final state. In addition, electroweak diagrams not involving self-interactions are suppressed [10], thus enhancing sensitivity of this final state to gauge-boson self couplings. Previously, an observation of $W^{\pm}W^{\pm}jj$ electroweak production was reported by the CMS Collaboration [11] and evidence was reported by the ATLAS Collaboration using a smaller dataset [12,13].

This Letter presents the observation and measurement of the electroweak production of $W^{\pm}W^{\pm}jj$ events in which both $W$ bosons decay into an electron or muon and a

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**FIG. 1.** Representative diagrams for $VVjj$ production where two electroweak gauge bosons are radiated off quarks. (a) Strong production diagrams involve both electroweak and strong interactions without gauge boson self-interactions. (b) In VBS diagrams, the gauge bosons can interact via gauge-boson self-interactions, e.g., a quartic $W$ boson vertex.
neutrino. This study uses 36.1 fb\(^{-1}\) of proton-proton (pp) collision data collected by the ATLAS detector at \(\sqrt{s} = 13\) TeV. The ATLAS detector [14] is a multipurpose particle detector with a forward-backward symmetric cylindrical geometry and almost 4\(\pi\) coverage in solid angle [15]. The inner tracking detector (ID) covers \(|\eta| < 2.5\) in pseudorapidity and consists of a silicon pixel detector, a silicon microstrip detector, and a transition radiation tracker. The ID is surrounded by a superconducting solenoid magnet and an almost hermetic calorimeter system, which provides three-dimensional reconstruction of particle showers up to \(|\eta| = 4.9\). The muon spectrometer (MS) has three air-core toroidal magnets: a barrel toroid and two end cap toroids. Three layers of precision tracking stations with drift tubes and cathode strip chambers allow precise muon momentum measurement up to \(|\eta| = 2.7\). Resistive-plate and thin-gap chambers provide muon triggering capability up to \(|\eta| = 2.4\).

Monte Carlo (MC) simulation is used to generate signal and background predictions. The generated events are processed through a detector simulation [16] based on GEANT4 [17] and are reconstructed using the same algorithms as used for data. The simulation includes the effect of multiple pp interactions per bunch crossing as detailed in Ref. [18].

Processes producing four leptons were simulated using the event generator SHERPA v2.2.2 [19–22] with the NNPDF3.0NNLO [23] set of parton distribution functions (PDF). Electroweak diboson production in association with two jets was simulated with diagrams including exactly six orders of the EW coupling [24]. The simulation of the strong production processes includes diagrams with exactly four electroweak (EW) vertices [25]. The simulation of both the \(W^\pm W^\pm jj\) electroweak signal and \(W^\pm W^\pm jj\) strong background processes includes up to one additional parton at LO in QCD. For \(\ell\ell\ell\ell\) and \(4\ell\) final states, matrix elements include up to one parton at next-to-leading order (NLO) and three partons at LO in QCD. All multileg processes were simulated combining the various final-state topologies with the MEPS merging algorithm [22,26] and matched to the SHERPA internal parton shower, hadronization, and underlying-event modeling [21,27]. The renormalization and factorization scales were set to the invariant mass of the four-lepton system at the matrix-element level. A signal sample for alternative theoretical predictions was simulated with POWHEG-BOX [28–31] and PYTHIA8 [32] at NLO QCD accuracy and employing the VBS approximation [33], the NNPDF3.0NLO PDF set, and renormalization and factorization scales set to the W boson mass. The \(Y_T\) processes were simulated with exactly three EW vertices using SHERPA v2.1.1 at NLO (up to one parton) or LO accuracy (up to three partons) using the CT10NLO PDF set [34]. Electroweak \(Y_Tjj\) processes were simulated using SHERPA v2.2.4 with the NNPDF3.0NNLO PDF set, with exactly five orders of the EW coupling, and at LO QCD accuracy. The triboson production processes were simulated with SHERPA v2.1.1 and the CT10NLO PDF set. MADGRAPH5_AMC@NLO [35] with NNPDF3.0NLO PDF set and PYTHIA8.210 were used to simulate \(t\bar{t}V\) processes at NLO QCD accuracy.

The \(W^\pm W^\pm jj\) electroweak production cross section is measured in a fiducial region defined at the particle level in MC simulation by requiring exactly two same-sign leptons with a transverse momentum, \(p_T\), greater than 27 GeV and \(|\eta| < 2.5\); leptons are defined as electrons and muons produced in \(W\) boson decays and that do not originate from \(\tau\) decays. The lepton four-momentum includes the four-momenta of photons inside a cone of size \(\Delta R = 0.1\) around the lepton. The two leptons must be separated by a distance of \(\Delta R_{\ell\ell} > 0.3\) and must have an invariant mass, \(m_{\ell\ell}\), greater than 20 GeV. The magnitude of the vector sum of the transverse momenta of the two final-state neutrinos with highest \(p_T\) must be greater than 30 GeV. Jets are reconstructed using the anti-\(k_T\) algorithm [36] with radius parameter \(R = 0.4\) and using all final-state particles, except for neutrinos and charged leptons from \(W\) boson decays. Jets are required to have \(p_T > 35\) GeV and \(|\eta| < 4.5\). Events with a charged lepton that is within a cone of radius \(\Delta R_{jj} = 0.3\) around a jet are vetoed. The fiducial region requires at least two jets, including one with \(p_T > 65\) GeV and another with \(p_T > 35\) GeV. The two highest-\(p_T\) jets must have an invariant mass \(m_{jj} > 500\) GeV and a rapidity difference \(|\Delta y_{jj}| > 2\).

The fiducial cross section predicted by SHERPA for \(W^\pm W^\pm jj\) electroweak production is \(2.01^{+0.33}_{-0.23}\) fb. The uncertainty includes independent variations of the renormalization and factorization scales by factors of 0.5 and 2 with the constraint \(0.5 \leq \mu_F/\mu_R \leq 2\) which contribute \(+14\%\), \(-11\%\); it also includes uncertainties from the NNPDF3.0NNLO ensemble, as well as differences between the CT14 [37] and MMHT2014 [38] PDF sets \((+1.5\%)\) [39]. Uncertainties in the parton shower, hadronization, and underlying-event modeling are evaluated by varying the MEPS matching and resummation scales and amount to \(+8\%\), \(-8\%\).

POWHEG +PYTHIA8 predicts a signal fiducial cross section of \(3.08^{+0.45}_{-0.46}\) fb, with the uncertainties derived using the same procedures as for the SHERPA prediction, except for the uncertainty in the parton shower modeling, which is estimated as the difference relative to POWHEG +HERWIG7 [40,41]. The SHERPA electroweak samples suffer from a nonoptimal setting of the color flow, which leads to an excess of central jet emissions from the parton shower. Since up to one additional parton is included in the matrix element of the \(W^\pm W^\pm jj\) electroweak sample, the effect on its differential distributions is reduced but accompanied by a significant suppression of the predicted cross section [24].
Events are required to contain at least one reconstructed proton interaction vertex. The vertex with the highest $p_T^2$ sum of associated ID tracks is selected as the primary vertex. Electrons are reconstructed from energy clusters in the electromagnetic calorimeter that are matched to tracks reconstructed in the ID with the requirement of a hit in the innermost pixel layer [42]. Muon are reconstructed by combining ID and MS information [43]. Electron and muon candidates must satisfy loose identification criteria [42,43], have $p_T > 6$ GeV, and $|\eta| < 2.47$ and $|\eta| < 2.7$, respectively. The ID tracks associated to electron (muon) candidates are matched to the primary vertex by requiring their transverse impact parameter significance to satisfy $|d_0/\sigma_d| < 5(10)$; the longitudinal impact parameter multiplied by the sine of the polar angle of the lepton candidates must satisfy $|z_0 \sin \theta| < 0.5$ mm. Electrons and muons passing these selections are further referred to as baseline leptons.

Jets are reconstructed from calorimeter energy clusters [44,45] using the anti-$k_t$ algorithm with radius parameter $R = 0.4$. Jets are required to have $p_T > 30$ GeV in the forward region ($2.4 < |\eta| < 4.5$) and $p_T > 25$ GeV in the central region ($|\eta| < 2.4$). Central jets with $p_T < 60$ GeV must be matched to the primary vertex [46]. Jets containing $b$ hadrons ($b$ jets) are identified in the range of $|\eta| < 2.5$ with an efficiency of 85% using techniques described in Ref. [47]. Selected electron, muon, and jet candidates are required to be nonoverlapping using the procedures described in Ref. [18]. The missing transverse momentum, $E_T^{miss}$, is computed using selected electrons, muons, and jets, and the track-based soft term defined in Ref. [48].

Events are selected online by single-electron or single-muon triggers [49]. Candidate events are selected by requiring exactly two same-sign baseline leptons, electron or muon, with $m_{ee} > 20$ GeV and by requiring $E_T^{miss} > 30$ GeV. They are required to contain at least two jets, including one with $p_T > 65$ GeV and another with $p_T > 35$ GeV. Events with at least one identified $b$ jet are rejected in order to reduce background contributions from top-quark pair production ($t\bar{t}$).

The two highest-$p_T$ jets are required to have $m_{jj} > 200$ GeV and $|\Delta y_{jj}| > 2$. These jet selection criteria were optimized to separate the $W^\pm W^\pm j j$ electroweak process from the strong production and other background processes.

After these selections, the dominant source of background events is due to leptons originating from decays of heavy-flavor hadrons and jets misidentified as electrons, collectively referred to as nonprompt leptons. Additional selection criteria are applied to reduce their contributions. Signal electrons are required to satisfy tight identification criteria [42], to have $p_T > 27$ GeV, and to be outside the calorimeter transition region (1.37 < |\eta| < 1.52). Signal muons are required to satisfy medium identification criteria [43], and to have $p_T > 27$ GeV and $|d_0/\sigma_d| < 3$. Signal electrons and muons are further required to be isolated from nearby particles, with isolation criteria defined using calorimeter clusters and ID tracks. These isolation criteria are optimized to have an efficiency of at least 90% for $p_T > 25$ GeV and at least 99% for $p_T > 60$ GeV [42,43]. For dielectron events, the electron pseudorapidity is restricted to $|\eta| < 1.37$ and events with $m_{ee} - 91.2$ GeV < 15 GeV are discarded. These criteria reduce the background from electron charge misreconstruction described later. Candidate events with exactly two signal leptons are said to pass the full event selection.

The contributions from the $WZ$, $V\gamma$, $ZZ$, $t\bar{t}V$, and triboson production are estimated using simulation. The predicted event yields of the $WZ$ and $V\gamma$ processes are normalized to data in dedicated control regions. The normalization of the $WZ$ background is determined using events with exactly three baseline leptons, two of which are required to pass the signal lepton selection, and that satisfy the dijet and $E_T^{miss}$ selection criteria. Events from $V\gamma$ production enter the signal region when a photon is misidentified as an electron. The modeling of this misidentification process in simulation is corrected using $Z \rightarrow \mu^+\mu^-\gamma$ events where a photon is emitted by a muon and then misidentified as an electron. These events are selected by requiring exactly two opposite-sign signal muons, one signal electron, $E_T^{miss} < 30$ GeV and a trilepton invariant mass satisfying 75 GeV < $m_{\mu\mu e} < 100$ GeV. A normalization factor of 1.8 is derived from this control region and used to correct the simulated $V\gamma$ events. To account for the differences between the $Z\gamma$ and $W\gamma$ processes, the full effect of this correction factor is assigned as a systematic uncertainty, corresponding to 44% of the estimated $V\gamma$ yield. The relative contributions from electroweak and strong production of $WZ$ and $V\gamma$ processes are estimated from simulation since this analysis is not sensitive to their different admixtures. Theoretical uncertainties in the predictions of the $ZZ$, $V\gamma$, triboson, and $t\bar{t}V$ backgrounds vary from 20% to 30% [25,50,51].

Background contributions with nonprompt leptons are estimated by weighting data events from dedicated control regions by scale factors. These scale factors are measured in dijet events containing exactly one lepton that is $p_T$ balanced by a $b$ jet. The $b$-jet requirement enhances nonprompt lepton contributions and suppresses contributions from $W/Z$ bosons, which are subtracted from data using simulation. The scale factor is defined as the ratio of the number of signal leptons to the number of leptons passing a dedicated background selection. The background leptons are required to pass the baseline lepton selection and fail the signal lepton selection, where background electrons are in addition required to satisfy medium...
The detector systematic uncertainties arising from the mismodeling of the reconstructed objects are estimated primarily from data and their impact on the analysis is assessed using simulated events. The dominant source is the uncertainty of the jet energy scale, which amounts to 2% for the signal and 10% for the WZ background. The uncertainty in the measurement of the integrated luminosity is 2.1%.

The theory modeling uncertainties of the \( m_{jj} \) distributions predicted by SHERPA for \( W^\pm W^{\pm} jj \) and WZ processes are evaluated using the procedures described above. They account for uncertainties in the total cross section, the acceptance of the fiducial selection, the modeling of the event selection efficiency and the shape of the \( m_{jj} \) distribution. Only the latter two affect the measured fiducial cross section of the \( W^\pm W^{\pm} jj \) signal, since absolute normalization uncertainties cancel in this measurement. The uncertainty in the modeling of the event selection efficiency also accounts for extrapolations from the fiducial phase space to the detector level, in particular for the \( \eta \) acceptance in dielectron events. Effects of the NLO electroweak corrections [53] and of the interference between electroweak and strong \( W^{\pm}W^{\pm} jj \) production [9] are assigned as an uncertainty in the \( m_{jj} \) shape of the \( W^\pm W^{\pm} jj \) signal, amounting to 6% and 4%, respectively. This approach is used because no event generator implemented the complete NLO calculation until recently [54] and because the interference contribution is defined only at the leading order [8]. The overlap of the photon radiation in the SHERPA parton shower model with the NLO EW corrections is found to be negligible.

Signal events are categorized by their lepton flavor and charge into six mutually exclusive channels: \( e^\pm e^\pm \), \( e^\pm \mu^\pm \), and \( \mu^\pm \mu^\pm \), in order to exploit their different signal and background compositions. The signal region is defined as \( m_{jj} > 500 \) GeV and further split into four \( m_{jj} \) bins, optimized to increase the expected signal sensitivity. Events with 200 GeV < \( m_{jj} < 500 \) GeV serve as additional control regions, dominated by contributions from nonprompt leptons and WZ backgrounds. The resulting 30 bins of the \( m_{jj} \) distributions in the signal and control regions are combined in a profile likelihood fit [55] to extract the fiducial cross section.

The signal strength, a free parameter in the fit, multiplies the expected fiducial \( W^\pm W^{\pm} jj \) electroweak production cross section used to produce the signal template. The signal template of reconstructed \( W^\pm W^{\pm} jj \) electroweak events also includes candidate events with electrons and muons produced in \( W \) decays into \( \tau \) lepton. Since the fiducial cross section prediction does not include such events, their fractional contribution predicted by the simulation is removed from the fiducial cross section measurement. Systematic uncertainties are included in the fit as nuisance parameters constrained by Gaussian functions. The WZ control region is also included in the fit as a single
ATLAS

We separate the nonprompt events into categories and the constrained mainly by the observed number of data events.

The signal region contains 122 data events, compared with POWHEG-BOX. The hatched band represents the statistical and systematic uncertainties of the background predictions added in quadrature. The backgrounds from $V\gamma$ production and electron charge misreconstruction are combined in the $e/\gamma$ conversions category. The other prompt category combines $ZZ$, $VVV$, and $t\bar{t}V$ background contributions.

Table I compares the numbers of data events in the signal region with the background and signal event yields after the fit; the signal region contains 122 data events, compared with a best-fit yield of 69 $\pm$ 7 background events. By fitting the data and background events in the signal and control regions, the background-only hypothesis is rejected with a significance of 6.5$\sigma$. Figure 2 shows the control region events separated into categories and the $m_{jj}$ distribution in the signal region after the fit. All nuisance parameters remain within their 1 standard deviation uncertainty after the fit. The normalization of the WZ background is scaled by a factor of $0.86^{+0.07}_{-0.08} (\text{stat})^{+0.18}_{-0.10} (\text{exp syst})^{+0.31}_{-0.23} (\text{mod syst})$, constrained mainly by the observed number of data events in the WZ control region. Figure 3 shows the $m_{ee}$ distribution in the signal region after the fit.

A signal strength of $1.44^{+0.26}_{-0.24}(\text{stat})^{+0.28}_{-0.22}(\text{syst})$ is measured with respect to the POWHEG fiducial cross section prediction for $W^{\pm}W^{\pm}jj$ electroweak production, where the systematic uncertainty also includes the absolute normalization uncertainty of this prediction. This corresponds to a fiducial signal cross section of

$$\sigma_{\text{fid}} = 2.89^{+0.51}_{-0.48}(\text{stat})^{+0.21}_{-0.22}(\text{exp syst}) \times ^{+0.14}_{-0.16}(\text{mod syst})^{+0.08}_{-0.06}(\text{lumi}) \text{ fb},$$

where the uncertainties correspond to the statistical, experimental systematic, theory modeling systematic, and luminosity uncertainties, respectively. The experimental systematic uncertainty includes the detector systematic uncertainties and the uncertainties in estimating all background processes except for the WZ and $W^{\pm}W^{\pm}jj$ strong

<table>
<thead>
<tr>
<th></th>
<th>$e^+e^+$</th>
<th>$e^-e^-$</th>
<th>$e^+\mu^+$</th>
<th>$e^-\mu^-$</th>
<th>$\mu^+\mu^-$</th>
<th>$\mu^-\mu^-$</th>
<th>Combined</th>
</tr>
</thead>
<tbody>
<tr>
<td>WZ</td>
<td>1.48 $\pm$ 0.32</td>
<td>1.09 $\pm$ 0.27</td>
<td>11.6 $\pm$ 1.9</td>
<td>7.9 $\pm$ 1.4</td>
<td>5.0 $\pm$ 0.7</td>
<td>3.4 $\pm$ 0.6</td>
<td>30 $\pm$ 4</td>
</tr>
<tr>
<td>Nonprompt</td>
<td>2.2 $\pm$ 1.1</td>
<td>1.2 $\pm$ 0.6</td>
<td>5.9 $\pm$ 2.5</td>
<td>4.7 $\pm$ 1.6</td>
<td>0.56 $\pm$ 0.05</td>
<td>0.68 $\pm$ 0.13</td>
<td>15 $\pm$ 5</td>
</tr>
<tr>
<td>$e/\gamma$ conversions</td>
<td>1.6 $\pm$ 0.4</td>
<td>1.6 $\pm$ 0.4</td>
<td>6.3 $\pm$ 1.6</td>
<td>4.3 $\pm$ 1.1</td>
<td>...</td>
<td>...</td>
<td>13.9 $\pm$ 2.9</td>
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<tr>
<td>Other prompt</td>
<td>0.16 $\pm$ 0.04</td>
<td>0.14 $\pm$ 0.04</td>
<td>0.90 $\pm$ 0.20</td>
<td>0.63 $\pm$ 0.14</td>
<td>0.39 $\pm$ 0.09</td>
<td>0.22 $\pm$ 0.05</td>
<td>2.4 $\pm$ 0.5</td>
</tr>
<tr>
<td>$W^+W^-jj$ strong</td>
<td>0.35 $\pm$ 0.13</td>
<td>0.15 $\pm$ 0.05</td>
<td>2.9 $\pm$ 1.0</td>
<td>1.2 $\pm$ 0.4</td>
<td>1.8 $\pm$ 0.6</td>
<td>0.76 $\pm$ 0.25</td>
<td>7.2 $\pm$ 2.3</td>
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<td>Expected background</td>
<td>5.8 $\pm$ 1.4</td>
<td>4.1 $\pm$ 1.1</td>
<td>28 $\pm$ 4</td>
<td>18.8 $\pm$ 2.6</td>
<td>7.7 $\pm$ 0.9</td>
<td>5.1 $\pm$ 0.6</td>
<td>69 $\pm$ 7</td>
</tr>
<tr>
<td>$W^\pm W^{\mp}jj$ electroweak</td>
<td>5.6 $\pm$ 1.0</td>
<td>2.2 $\pm$ 0.4</td>
<td>24 $\pm$ 5</td>
<td>9.4 $\pm$ 1.8</td>
<td>13.4 $\pm$ 2.5</td>
<td>5.1 $\pm$ 1.0</td>
<td>60 $\pm$ 11</td>
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<tr>
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<td>4</td>
<td>44</td>
<td>28</td>
<td>25</td>
<td>11</td>
<td>122</td>
</tr>
</tbody>
</table>

FIG. 2. Event yields for data, signal, and background in the WZ and $200 < m_{jj} < 500$ GeV control regions (left) and the $m_{jj}$ distribution for events meeting all selection criteria for the signal region (right). Signal and background distributions are shown as predicted after the fit. The hatched band represents the statistical and systematic uncertainties of the background predictions added in quadrature. The backgrounds from $V\gamma$ production and electron charge misreconstruction are combined into the $e/\gamma$ conversions category. The other prompt category combines $ZZ$, $VVV$, and $t\bar{t}V$ background contributions. The last bin on the right figure includes the overflow.

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FIG. 3. The $m_{jj}$ distribution for events meeting all selection criteria for the signal region is shown as predicted after the fit. The hatched band represents the statistical and systematic uncertainties of the background prediction added in quadrature. The fitted signal strength and nuisance parameters have been propagated, with the exception of the uncertainties due to the interference and electroweak corrections for which a flat uncertainty is assigned. The backgrounds from $V\gamma$ production and electron charge misreconstruction are combined in the $e/\gamma$ conversions category. The other prompt category combines ZZ, VV, and $t\bar{t}V$ background contributions. The last bin of the distribution includes the overflow.

TABLE II. Impact of different components of systematic uncertainty on the measured fiducial cross section, without taking into account correlations. The impact of one source of systematic uncertainty is computed by first performing the fit with the corresponding nuisance parameter fixed to 1 standard deviation up or down from the value obtained in the nominal fit, then these up and down variations are symmetrized. The impacts of several sources of systematic uncertainty are added in quadrature for each component. The categorization of sources of systematic uncertainties into experimental and theory modeling correspond to those used for the measured fiducial cross section.

<table>
<thead>
<tr>
<th>Source</th>
<th>Impact [%]</th>
</tr>
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<tbody>
<tr>
<td>Experimental</td>
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</tr>
<tr>
<td>Electron energy scale and resolution, efficiency</td>
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</tr>
<tr>
<td>Muon momentum scale and resolution, efficiency</td>
<td>1.3</td>
</tr>
<tr>
<td>Jet energy and $E_T^{miss}$ scale and resolution</td>
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</tr>
<tr>
<td>$b$-tagging inefficiency</td>
<td>2.1</td>
</tr>
<tr>
<td>Pileup modeling</td>
<td>1.6</td>
</tr>
<tr>
<td>Background, statistical</td>
<td>3.2</td>
</tr>
<tr>
<td>Background, misid. leptons</td>
<td>3.3</td>
</tr>
<tr>
<td>Background, charge misrec.</td>
<td>0.3</td>
</tr>
<tr>
<td>Background, other</td>
<td>1.8</td>
</tr>
<tr>
<td>Theory modeling</td>
<td></td>
</tr>
<tr>
<td>$W^\pm W^\pm jj$ electroweak-strong interference</td>
<td>1.0</td>
</tr>
<tr>
<td>$W^\pm W^\pm jj$ electroweak, EW corrections</td>
<td>1.4</td>
</tr>
<tr>
<td>$W^\pm W^\pm jj$ electroweak, shower, scale, PDF &amp; $\alpha_s$</td>
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<td>Luminosity</td>
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</table>

production processes that are accounted for in the modeling systematic uncertainty. Table II summarizes the impacts of different components of systematic uncertainty.

The measured fiducial cross section includes contributions from both the $W^\pm W^\pm jj$ electroweak production and its interference with the $W^\pm W^\pm jj$ strong production, estimated to be approximately 6% of the predicted fiducial cross section for $W^\pm W^\pm jj$ electroweak production. The fiducial cross section for the $W^\pm W^\pm jj$ electroweak production, without the interference effect, is predicted by SHERPA and POWHEG+PYTHIA8 to be $2.01^{+0.33}_{-0.23}$ fb and $3.08^{+0.45}_{-0.46}$ fb, respectively. The impact on the measured fiducial cross section of using POWHEG+PYTHIA8 instead of SHERPA to generate the $m_{jj}$ signal template was tested and found to be smaller than the 3.6% signal modeling uncertainty.

In summary, the electroweak $VVjj$ production process was studied in the $W^\pm W^\pm jj$ final state using 36.1 fb$^{-1}$ of $pp$ collision data recorded at $\sqrt{s} = 13$ TeV by the ATLAS detector at the LHC. This process includes VBS diagrams that probe the mechanism of electroweak symmetry breaking. An excess of events is observed and the background-only hypothesis is rejected with a significance of 6.5$\sigma$. The fiducial cross section for $W^\pm W^\pm jj$ electroweak production is measured to be $\sigma^{fid}_{VWjj} = 2.89^{+0.51}_{-0.48}$ (stat)$^{+0.24}_{-0.22}$ (exp syst)$^{+0.14}_{-0.16}$ (mod syst)$^{+0.08}_{-0.06}$ (lumi) fb.

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[15] ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the center of the detector and the z axis along the beam direction. The x axis points from the IP to the center of the LHC ring, and the y axis points upward. Cylindrical coordinates (r, φ) are used in the (x, y) plane, φ being the azimuthal angle around the beam direction. The rapidity is defined as y = 1/2 ln[(E + p_φ)/|E − p_φ|], where E is the energy of the particle and p_φ is the projection of the momentum along the z axis. The pseudorapidity is defined in terms of the polar angle θ as η = − ln tan(θ/2). The distance ΔR is defined as ΔR = √((Δη)^2 + (Δφ)^2).
Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey
LAPP, Université Grenoble Alpes, Université Savoie Mont Blanc, CNRS/IN2P3, Annecy, France
High Energy Physics Division, Argonne National Laboratory, Argonne, Illinois, USA
Department of Physics, University of Arizona, Tucson, Arizona, USA
Department of Physics, University of Texas at Arlington, Arlington, Texas, USA
Physics Department, National and Kapodistrian University of Athens, Athens, Greece
Physics Department, National Technical University of Athens, Zografou, Greece
Department of Physics, University of Texas at Austin, Austin, Texas, USA
Bahcesehir University, Faculty of Engineering and Natural Sciences, Istanbul, Turkey
Istanbul Bilgi University, Faculty of Engineering and Natural Sciences, Istanbul, Turkey
Department of Physics, Bogaziçi University, Istanbul, Turkey
Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey
Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
Institut de Física d’Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona, Spain
Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China
Physics Department, Tsinghua University, Beijing, China
Department of Physics, Nanjing University, Nanjing, China
University of Chinese Academy of Science (UCAS), Beijing, China
Institute of Physics, University of Belgrade, Belgrade, Serbia
Department for Physics and Technology, University of Bergen, Bergen, Norway
Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, California, USA
Institut für Physik, Humboldt Universität zu Berlin, Berlin, Germany
Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
Facultad de Ciencias y Centro de Investigaciones, Universidad Antonio Nariño, Bogota, Colombia
INFN Bologna and Universita’ di Bologna, Dipartimento di Fisica, Italy
INFN Sezione di Bologna, Italy
Physikalisches Institut, Universität Bonn, Bonn, Germany
Department of Physics, Boston University, Boston, Massachusetts, USA
Department of Physics, Brandeis University, Waltham, Massachusetts, USA
Transilvania University of Brasov, Brasov, Romania
Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania
Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi, Romania
National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj-Napoca, Romania
University Politehnica Bucharest, Bucharest, Romania
West University in Timisoara, Timisoara, Romania
Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava, Slovak Republic
Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
Physics Department, Brookhaven National Laboratory, Upton, New York, USA
Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
California State University, California, USA
Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
Department of Physics, University of Cape Town, Cape Town, South Africa
Department of Mechanical Engineering Science, University of Johannesburg, Johannesburg, South Africa
School of Physics, University of the Witwatersrand, Johannesburg, South Africa
Department of Physics, Carleton University, Ottawa, Ontario, Canada
Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Énergies-Université Hassan II, Casablanca, Morocco
Faculté des Sciences, Université Ibn-Tofail, Kénitra, Morocco
Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech, Morocco
Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda, Morocco
Faculté des sciences, Université Mohammed V, Rabat, Morocco
CERN, Geneva, Switzerland
Enrico Fermi Institute, University of Chicago, Chicago, Illinois, USA
LPC, Université Clermont Auvergne, CNRS/IN2P3, Clermont-Ferrand, France
Nevis Laboratory, Columbia University, Irvington, New York, USA
Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
Dipartimento di Fisica, Università della Calabria, Rende, Italy
INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Italy
Physics Department, Southern Methodist University, Dallas, Texas, USA
Universidade Federal de São João del Rei (UFSJ), São João del Rei, Brazil
Instituto de Física, Universidade de São Paulo, São Paulo, Brazil
KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
Graduate School of Science, Kobe University, Kobe, Japan
AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland
Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland
Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland
Faculty of Science, Kyoto University, Kyoto, Japan
Kyoto University of Education, Kyoto, Japan
Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka, Japan
Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
Physics Department, Lancaster University, Lancaster, United Kingdom
Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom
Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana, Slovenia
School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom
Department of Physics, Royal Holloway University of London, Egham, United Kingdom
Department of Physics and Astronomy, University College London, London, United Kingdom
Louisiana Tech University, Ruston, Louisiana, USA
Fysiska institutionen, Lunds universitet, Lund, Sweden
Centre de Calcul de l’Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France
Departamento de Física Teorica C-15 and CIAFF, Universidad Autónoma de Madrid, Madrid, Spain
Institut für Physik, Universität Mainz, Mainz, Germany
School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France
Department of Physics, University of Massachusetts, Amherst, Massachusetts, USA
Department of Physics, McGill University, Montreal, Quebec, Canada
School of Physics, University of Melbourne, Victoria, Australia
Department of Physics, University of Michigan, Ann Arbor, Michigan, USA
Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan, USA
B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Belarus
Research Institute for Nuclear Problems of Byelorussian State University, Minsk, Belarus
Group of Particle Physics, University of Montreal, Montreal, Quebec, Canada
P.N. Lebedev Physical Institute of the Russian Academy of Sciences, Moscow, Russia
Institute for Theoretical and Experimental Physics of the National Research Centre Kurchatov Institute, Moscow, Russia
National Research Nuclear University MEPhI, Moscow, Russia
D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
Nagasaki Institute of Applied Science, Nagasaki, Japan
Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan
Department of Physics and Astronomy, University of New Mexico, Albuquerque, New Mexico, USA
Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
Department of Physics, Northern Illinois University, DeKalb, Illinois, USA
Budker Institute of Nuclear Physics and NSU, Novosibirsk, Russia
Novosibirsk State University Novosibirsk, Russia
Institute for High Energy Physics of the National Research Centre Kurchatov Institute, Protvino, Russia
Department of Physics, New York University, New York, New York, USA
Ochanomizu University, Otsuka, Bunkyo-ku, Tokyo, Japan
Ohio State University, Columbus, Ohio, USA
Faculty of Science, Okayama University, Okayama, Japan
Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, Oklahoma, USA
Department of Physics, Oklahoma State University, Stillwater, Oklahoma, USA
Palacký University, RCPTM, Joint Laboratory of Optics, Olomouc, Czech Republic
Center for High Energy Physics, University of Oregon, Eugene, Oregon, USA
LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay, France
Graduate School of Science, Osaka University, Osaka, Japan
Department of Physics, University of Oslo, Oslo, Norway
Department of Physics, Oxford University, Oxford, United Kingdom
Konstantinov Nuclear Physics Institute of National Research Centre “Kurchatov Institute”, PNPI, St. Petersburg, Russia

LPNHE, Sorbonne Université, Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, Paris, France

Laboratório de Instrumentação e Física Experimental de Partículas-LIP, Portugal

Depto de Física, Pontificia Universidad Católica de Chile, Santiago, Chile

Department of Physics and Astronomy, University of California Santa Cruz, Santa Cruz, California, USA

Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, California, USA

IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France

High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia

Institute of Physics of the Czech Academy of Sciences, Prague, Czech Republic

Czech Technical University in Prague, Prague, Czech Republic

Department of Animal Sciences, Ohio State University, Columbus, Ohio, USA

Charles University, Faculty of Mathematics and Physics, Prague, Czech Republic

Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom

SLAC National Accelerator Laboratory, Stanford, California, USA

Physics Department, Royal Institute of Technology, Stockholm, Sweden

Department of Physics and Astronomy, Stony Brook University, Stony Brook, New York, USA

Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom

School of Physics, University of Sydney, Sydney, Australia

Department of Physics, Shinshu University, Nagano, Japan

Department Physik, Universität Siegen, Siegen, Germany

Department of Physics, Simon Fraser University, Burnaby, British Columbia, Canada

SLAC National Accelerator Laboratory, Stanford, California, USA

Physics Department, Royal Institute of Technology, Stockholm, Sweden

Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom

Department of Physics, University of Sheffield, Sheffield, United Kingdom

Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel

Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece

International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo, Japan

Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan

Department of Physics, Tokyo Institute of Technology, Tokyo, Japan

Tomsk State University, Tomsk, Russia

Department of Physics, University of Toronto, Toronto, Ontario, Canada

TRIUMF, Vancouver, British Columbia, Canada

Department of Physics and Astronomy, York University, Toronto, Ontario, Canada

Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan

Department of Physics and Astronomy, Tufts University, Medford, Massachusetts, USA

Department of Physics and Astronomy, University of California Irvine, Irvine, California, USA

Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden

Department of Physics, University of Illinois Urbana Champaign, Urbana, Illinois, USA

Instituto de Física Corpuscular (IFIC), Centro Mixto Universidad de Valencia-CSIC, Valencia, Spain

Department of Physics, University of British Columbia, Vancouver, British Columbia, Canada

Department of Physics and Astronomy, University of Victoria, Victoria, British Columbia, Canada

Fakultät für Physik und Astronomie, Julius-Maximilians-Universität Würzburg, Würzburg, Germany

Department of Physics, University of Warwick, Coventry, United Kingdom

Waseda University, Tokyo, Japan

Department of Particle Physics, Weizmann Institute of Science, Rehovot, Israel

Department of Physics, University of Wisconsin, Madison, Wisconsin, USA

Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany

Department of Physics, Yale University, New Haven, Connecticut, USA

Yerevan Physics Institute, Yerevan, Armenia
Deceased.

Also at Department of Physics, King’s College London, London, United Kingdom.

Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.

Also at TRIUMF, Vancouver, British Columbia, Canada.

Also at Department of Physics and Astronomy, University of Louisville, Louisville, Kentucky, USA.

Also at Physics Department, An-Najah National University, Nablus, Palestine.

Also at Department of Physics, California State University, Fresno, USA.

Also at Department of Physics, University of Fribourg, Fribourg, Switzerland.

Also at Departament de Fisica de la Universitat Autonoma de Barcelona, Barcelona, Spain.

Also at Tomsk State University, Tomsk, and Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia.

Also at The Collaborative Innovation Center of Quantum Matter (CICQM), Beijing, China.

Also at Universita di Napoli Parthenope, Napoli, Italy.

Also at Institute of Particle Physics (IPP), Canada.

Also at II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen, Germany.

Also at Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania.

Also at CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France.

Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia.

Also at Borough of Manhattan Community College, City University of New York, New York, New York, USA.

Also at Department of Financial and Management Engineering, University of the Aegean, Chios, Greece.

Also at Centre for High Performance Computing, CSIR Campus, Rosebank, Cape Town, South Africa.

Also at Department of Physics, California State University, East Bay, USA.

Also at Institucio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain.

Also at Department of Physics, University of Michigan, Ann Arbor, Michigan, USA.

Also at LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay, France.

Also at Graduate School of Science, Osaka University, Osaka, Japan.

Also at Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg, Germany.

Also at Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands.

Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia.

Also at CERN, Geneva, Switzerland.

Also at Manhattan College, New York, New York, USA.

Also at Hellenic Open University, Patras, Greece.

Also at LPNHE, Sorbonne Université, Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, Paris, France.

Also at The City College of New York, New York, New York, USA.

Also at Universidad de Granada, Granada (Spain), Spain.

Also at Department of Physics, California State University, Sacramento, USA.

Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia.

Also at Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève, Switzerland.

Also at Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom.

Also at Louisiana Tech University, Ruston, Louisiana, USA.

Also at School of Physics, Sun Yat-sen University, Guangzhou, China.

Also at Faculty of Physics, M.V. Lomonosov Moscow State University, Moscow, Russia.

Also at Department of Applied Physics and Astronomy, University of Sharjah, Sharjah, United Arab Emirates.

Also at National Research Nuclear University MEPhI, Moscow, Russia.

Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.

Also at Giresun University, Faculty of Engineering, Giresun, Turkey.

Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.