The radial supernova remnant distribution in the Galaxy

Verberne, S.; Vink, J.

DOI
10.1093/mnras/stab940

Publication date
2021

Document Version
Final published version

Published in
Monthly Notices of the Royal Astronomical Society

License
CC BY

Citation for published version (APA):
The radial supernova remnant distribution in the Galaxy

Sill Verberne and Jacco Vink

1Anton Pannekoek Institute for Astronomy, University of Amsterdam, NL-1090 Amsterdam, the Netherlands
2GRAPPA, University of Amsterdam, NL-1090 Amsterdam, the Netherlands

ABSTRACT

Supernovae (SNe) are the dominant source of chemical enrichment of galaxies, and they are a important source of energy to heat the interstellar medium and accelerate cosmic rays. Our knowledge of SNe in the Milky Way is based mostly on the study of Galactic supernova remnants (SNRs), providing an (incomplete) record to SN activity over the last \( \sim 100000 \) yr. Here, we report on an investigation of the spatial distribution of Galactic SNRs. Given the limited number of SNRs, it is common to assume a functional form for the Galactocentric distribution of SNRs. However, several functional forms have been used in the past, without much justification for the radial distribution. For example, one often used functional form implies that no SN activity is present in the Galactic Centre region. However, the presence of a magnetar and an SNR near the Galactic Centre suggests that a spatial distribution with zero SNRs at the Galactic Centre is not realistic. In light of these concerns, we re-evaluate the Galactic SNR distribution. We provide a brief outline of the main detection biases in finding SNRs and we investigate whether or not the use of the most common functional form is justified and how it compares to other models for the SNR distribution. We do this by analysing the longitudinal distribution of SNRs. We find that a simple exponential distribution is the most consistent and simplest model for describing the radial SNR distribution in the Galaxy and draw comparisons with the massive star formation and metallicity distributions.

Key words: supernovae: general – ISM: supernova remnants – Galaxy: structure.

1 INTRODUCTION

Supernova remnants (SNRs) are formed by both core-collapse and thermonuclear supernovae (SNe). Core-collapse SNe are associated with massive stars at the end of their life (e.g. Bethe 1990; Woosley & Weaver 1995) and thermonuclear SNe are associated with the explosions of CO white dwarfs (e.g. Hillebrandt & Niemeyer 2000; Maoz, Mannucci & Nelemans 2014). SNe occur at a rate of about two to three per century in the Galaxy (e.g. Tammann, Loeffler & Schroeder 1994; Li et al. 2011). The SNRs formed by these SNe are believed to remain visible for 20–80 kyr (Sarbadhicary et al. 2017). Given the rate at which they occur and their expected lifetime, only a fraction of the SNRs in the Galaxy have actually been observed. Furthermore, the sample of SNRs that have been discovered is heavily biased by selection effects such as distance and surface brightness.

SNRs are of interest to the study of cosmic rays since they are believed to be sites of cosmic ray acceleration up to PeV energies (e.g. Berezhko, Pühlhofer & Völk 2003; Blasi 2011; Bell et al. 2013). Furthermore, they can provide physical insight into the physics of SNe themselves. SNe are the most important source of alpha elements and Fe-group elements for enriching the interstellar medium, which can be studied in SNRs. Since \( \sim 80 \) per cent of SNe are of the core-collapse type (e.g. Li et al. 2011; Graur et al. 2017), which only occur for massive stars with relatively short lifetimes, the spatial SNR distribution is expected to follow massive star formation. A good model for the spatial distribution would allow more focused searching for new remnants by comparing the model with the observed SNRs and predicting the location of the highest number of undiscovered remnants. This could in turn result in a better understanding of SNe themselves, since SNRs contain information about the SNe that formed them.

Green (2015) relied on the 1D projection, Galactic longitude (\( l \)), of the spatial distribution, whereas Case & Bhattacharya (1998) attempted to reconstruct the full spatial distribution by using the \( \Sigma – D \) relation to infer distances to all SNRs. The aim of this work is to expand upon the study of Green (2015) by fitting and comparing multiple functional forms. In earlier work, the use of one of many possible functional forms seems to suggest there is only one possible description of the radial SNR distribution. The problem with the model usually adopted is that it approaches zero at the Galactic Centre, while SNRs are known to reside in this region as described by Maeda et al. (2002) and Kennea et al. (2013). Whereas Green (2015) fitted to minimize the residuals of the cumulative distribution function, we will fit directly to the observed number of SNRs per bin in \( l \).

In Section 2, the most important biases in the SNR catalogue will be outlined. Section 3 discusses the methods involved in the investigation. In Section 4, the results of the performed model fits will be given, and Sections 5 and 6 will give the discussion and conclusions, respectively.
2 Supernova Remnant Detection Biases

This work is based upon Green’s catalogue of Galactic SNRs (Green 2019) that contains 294 SNRs. In this section, an outline will be given of the most important selection effects in the discovery of these SNRs.

Green (2004, 2015) advocates for the use of a nominal surface brightness (Σ) completeness limit of $10^{-20}$ W m$^{-2}$ Hz$^{-1}$ sr$^{-1}$, which is a conserved quantity since

$$\Sigma = \frac{F}{\Omega}, \quad F = \frac{L}{4\pi R^2}, \quad \Omega = \frac{A}{R^2},$$

(1)

with $F$ the flux density, $L$ the luminosity, $\Omega$ the solid angle, $A$ the projected surface area, and $R$ the distance to the source. The reason for imposing a surface brightness threshold in the radio is that most remnants have been discovered in the radio and surface brightness is an important selection effect in the discovery of new remnants. Imposing the threshold thus decreases the bias of the sample, although, as mentioned by Green (2015), a higher surface brightness threshold might be required around the Galactic Centre due to higher Galactic background radio emission. Furthermore, the sample also suffers from a bias disfavouring small angular size remnants, since they might be too small to be properly identified. This problem with small angular size remnants is not resolved by imposing a surface brightness threshold. What we do know is that these small angular size remnants will be concentrated towards small $l$ (or large physical distances). Moreover, due to source confusion along the line of sight, even SNRs of sufficient size are more difficult to detect near the Galactic Centre. This means that the surface brightness, small angular size, and confusion biases are most important around the Galactic Centre.

Bias in the sample selection could also result from the correlation between the surface brightness of SNRs and the density of the medium in which they evolve. It is believed that for a remnant evolving in a relatively high density medium, its surface brightness will be high and its lifetime will be low compared to that evolving in a relatively low density environment (e.g. Blondin et al. 1998; Sarbadhikary et al. 2017). Fortunately for this study, these effects work to counteract each other; although an SNR evolving in a high-density environment is more likely to exceed a surface brightness completeness limit, it will be relatively short lived. We thus expect the net contribution of this bias to be relatively minor.

2.1 Unbiased sample selection

In order to investigate the bias introduced by the higher background around the Galactic Centre, we use the 408 MHz all-sky radio map from Haslam et al. (1982). In order for an SNR to be recognized as such, the surface brightness needs to exceed the noise level of the background sufficiently. Using only the background level for this purpose does not capture the full complexity by which SNRs are discovered (e.g. remnants are easier to discover from multifrequency observations by decomposing the thermal and non-thermal components). Also, it does not alleviate the bias induced by small angular size remnants or source confusion along the line of sight. However, it does give us the order of magnitude by which the SNR completeness limit of $10^{-20}$ W m$^{-2}$ Hz$^{-1}$ sr$^{-1}$ is off in the region around the Galactic Centre. By adjusting the surface brightness cutoff, we can thus limit the dominant sources of bias to small angular sizes and source confusion along the line of sight.

Since the radio sky is dominated by synchrotron emission at 408 MHz, we can use its spectral index to extrapolate the map to 1 GHz

$$T(\nu) \propto \nu^{-\beta},$$

(2)

with $T(\nu)$ the brightness temperature, $\nu$ the frequency, and $\beta$ the spectral index. Platania et al. (1998) determined that the spectral index for synchrotron radiation is 2.76 $\pm$ 0.11 in the range of 0.4–7.5 GHz. Using this spectral index, while adjusting for the cosmic microwave background, we created a radio map at 1 GHz. We then investigated the remaining bias around the Galactic Centre caused by using a surface brightness threshold of $10^{-20}$ W m$^{-2}$ Hz$^{-1}$ sr$^{-1}$ by making a contour plot. In Fig. 1, this plot is shown. It can be seen that the $10^{-20}$ W m$^{-2}$ Hz$^{-1}$ sr$^{-1}$ surface brightness threshold is not high enough to prevent bias for $|l| \approx 20^\circ$ due to the high background. By increasing the surface brightness threshold to $1.3 \times 10^{-20}$ W m$^{-2}$ Hz$^{-1}$ sr$^{-1}$, we can reduce this biased region to $|l| \approx 10^\circ$, as shown by the blue contours in Fig. 1. In this region, we know that there is bias anyway due to the small angular sizes and source confusion. We have thus effectively limited the biased region to this inner $10^\circ$ by increasing the surface brightness threshold. In Section 3, we will discuss how this remaining bias is addressed. Of the 294 SNRs in the catalogue, 67 pass the threshold of $10^{-20}$ W m$^{-2}$ Hz$^{-1}$ sr$^{-1}$ at 1 GHz and 57 that of $1.3 \times 10^{-20}$ W m$^{-2}$ Hz$^{-1}$ sr$^{-1}$ at 1 GHz. In Fig. 2, a histogram of the distribution in $l$ is shown of the remnants that pass the $1.3 \times 10^{-20}$ W m$^{-2}$ Hz$^{-1}$ sr$^{-1}$ surface brightness threshold at 1 GHz along with their cumulative distribution.
Figure 3. The SNR distribution as a function of the absolute value of $l$. Only SNRs satisfying the surface brightness threshold of $1.3 \times 10^{-20}$ W m$^{-2}$ Hz$^{-1}$ sr$^{-1}$ at 1 GHz have been included.

3 METHODS

3.1 Analysis

We follow Green (2015) in using the 1D projection of the SNR distribution in $l$ instead of the 3D $\Sigma$–$D$ relation (e.g. Case & Bhattacharya 1998). The reason for this is that the $\Sigma$–$D$ relation is an empirical relation between the surface brightness ($\Sigma$) and the diameter of the source ($D$) that suffers from scatter of about an order of magnitude in the derived distance to said source. The $l$ of SNRs is not subject to this problem and can still be compared with model distributions. Since we only consider cylindrically symmetric distributions, the absolute value of $l$ has been taken in order to improve the statistics per bin; 36 bins of 5° each between 0° and 180° were used. In Fig. 3, the resulting histogram of the surface brightness limited sample is shown.

Because the fitted distributions are functions of the Galactocentric radius ($R_{\text{gal}}$), which we do not directly observe, a conversion to a heliocentric radius ($R_{\text{helio}}$) needs to be made. This conversion can be made using the law of cosines as

$$R_{\text{gal}} = \left( R_{\odot}^2 + R_{\text{helio}}^2 - 2 R_{\odot} R_{\text{helio}} \cos l \right)^{1/2},$$

where $R_{\odot}$ is the distance to the centre of the Galaxy of about 8 kpc (Genzel, Eisenhauer & Gillessen 2010) and $l$ is the Galactic longitude. After this conversion, Simpson’s rule of numerical integration has been used in the range 0 kpc $\leq R_{\text{gal}} \leq 50$ kpc in order to obtain the number of SNRs along the line of sight in a given direction according to the model distribution. This model distribution is then fit to the observed distribution using the $\chi^2$ statistic

$$\chi^2 = \sum_{i=1}^{N} \frac{(n_i - e_i)^2}{\sigma_i^2},$$

where $n_i$ the observed counts per bin, $e_i$ the model counts per bin, and $\sigma_i$ the error on the observed counts per bin.

Since the SNR catalogue with the imposed surface brightness completeness limit only has 57 remnants, the catalogue is a sparse data set, which means that $\sigma_i$ cannot be approximated from the observed counts per bin. Instead, we need to use Poisson statistics to estimate $\sigma_i = \sqrt{n_i}$. Wheaton et al. (1995) note that in order to avoid biases, the weighting factors (i.e. $\sigma_i$’s) need to be taken as constants before the fitting process. This fitting process can then be iterated over while adjusting $\sigma_i$ from run to run. This results in a model fit with unbiased parameters even when $n_i \ll 1$. In order to convince the critical reader of its effectiveness, fitting has been repeated directly using the Poisson distribution (maximum likelihood). Parameter errors have been determined by grid searches, where (using $\chi^2$ minimization) the 1σ confidence limit coincides with where $\chi^2 = \chi^2_{\text{min}} + 1$. In these grid searches, the errors (i.e. $\sigma_i$) were based on the best-fitting model.

As a goodness-of-fit test, the $Y^2$ statistic (Lucy 2000) has been used. This is, like Pearson’s statistic, a $\chi^2$ statistic with the added benefit of being more accurate when the amount of counts per bin is low. In fact, the statistic remains reasonably accurate at the 2σ level when the mean counts per bin are $\ll 1$ as long as the total number of events is $\gtrsim 30$ (Lucy 2000). The $Y^2$ statistic, in the multinomial case, is given by

$$Y^2 = v + \xi^{-1}(X^2 - v), \quad \xi^{-1} = \sqrt{\frac{2v}{V(X^2)}},$$

where $v$ is the number of degrees of freedom and

$$V(X^2) = 2v + \sum_i \frac{1}{e_i} - \frac{1}{N} (I^2 + 2I - 2), \quad X^2 = \sum_{i=1}^{N} \frac{(n_i - e_i)^2}{e_i},$$

with $I$ being the number of bins used and $N$ the total number of SNRs in the sample. The multinomial case is used since the model only describes the distribution of SNRs and not the number of them, meaning that the number of predicted SNRs is simply the number of observed ones ($\tilde{N} = N$).

Given the relatively large bias in the data set around the Galactic Centre, the model fits have also been performed when omitting the inner $\pm 10°$ (2 bins) of SNRs. What this allows us to do is fit the distributions to a data set with less observational bias. However, note that the fits performed on this modified data set should not be used to infer properties of the inner $\pm 10°$ of the Galactic SNR distribution.

3.2 Functional forms

A number of distributions have been fitted using equation (4) to see which one best describes the data. Since the distributions are normalized to the data, no normalization parameters have been included.

In earlier works (see Stecker & Jones 1977; Case & Bhattacharya 1998; Green 2015), there has been an insistence on the use of a distribution of the form

$$f(r) = \left( \frac{r}{r_\odot} \right)^\alpha \exp \left( - \beta \frac{r-r_\odot}{r_\odot} \right),$$

which is a modified version of a gamma distribution (henceforth referred to as MGD), with $\alpha$ and $\beta$ being free parameters and $r_\odot$ the distance to the centre of the Galaxy. No attempts have so far been made to see whether this distribution is appropriate in this case and no real justification has been given for its use. Therefore, in order to see whether this distribution is truly fit for describing the SNR distribution, model fits have also been performed with simple power-law (PL) and exponential distributions of the forms:

$$f(r) = \left( \frac{r}{r_\odot} \right)^\alpha$$

and

$$f(r) = \exp \left( - \beta \frac{r-r_\odot}{r_\odot} \right).$$
where both functions have a single free parameter. The benefit of looking at these two distributions is that we can easily compare them to equation (7) since they are nested models.

A Sérsic profile (Sérsic 1963), which is a generalization of a de Vaucouleurs profile (de Vaucouleurs 1948), has also been included in the analysis. The Sérsic profile can be used to describe the stellar or intensity distribution in a galaxy and is given by

$$f(r) = \exp \left( -b_n \left( \left( \frac{r}{r_e} \right)^{1/n} - 1 \right) \right),$$

where $n$ is a free parameter, $r_e$ is the half-light radius of the galaxy, and $b_n$ is a constant that can be approximated by

$$b_n \simeq 1.999n - 0.3271$$

for $0.5 < n < 10$. An extensive discussion on this profile can be found in Graham & Driver (2005). This profile has an additional constraint in comparison to the exponential, since we need the half-light radius that is not treated as a free parameter and assumed to be 3 kpc.

Finally, a sum between an exponential and PL has been included in the form of

$$f(r) = \left( \frac{r}{r_\odot} \right)^{a} + \exp \left( -\beta \left( \frac{r - r_\odot}{r_\odot} \right) \right).$$

This gives us a set of five models that are combinations between PLs and exponentials that will be used to determine which one best describes the radial SNR distribution in the Galaxy.

Of course one could imagine many more possible models to fit to the data, but in doing so the chances of finding a model that purely by coincidence fits very well would also increase.

### 4 Results

The fitting results of the models to the surface brightness limited sample are shown in Table 1. We list here a complete sample in the range of $0^\circ \leq |l| \leq 180^\circ$ and a centre-excluded sample in the range of $10^\circ \leq |l| \leq 180^\circ$. The separate model fits are all shown in Appendix (FigsA1 and A2). Of the 57 remnants with a surface brightness at 1 GHz of at least $1.3 \times 10^{-20}$ W m$^{-2}$ Hz$^{-1}$ sr$^{-1}$, 47 are outside of the inner $\pm 10^\circ$ of the Galaxy. As mentioned in Section 3.1, the fits have also been performed directly using the Poisson distribution. The results (i.e. the fitted parameters) from using the Poisson distribution were exactly equal in the reported significance to those obtained using the $\chi^2$ statistic with iterations over the errors. Note that for models that include a PL component, this component is not allowed to become smaller than $-1$, since this would result in an infinite number of predicted SNRs at the Galactic Centre. Table 2 gives the nested model significances between the MGD and PL, and exponential models. A small significance ($\sim 0.05$) indicates that the MGD fits significantly better, while a large significance ($> 0.05$) indicates that it does not.

#### 4.1 Modified gamma distribution

The model fits of the MGD from equation (7) show good fits for both the SNR samples with $p$-values of 0.44 and 0.62, respectively. For the complete sample fit, the distribution goes to zero at the Galactic Centre and has its surface density peak at about 1.7 kpc from the Galactic Centre. A more meaningful number, however, is the peak of the number of SNRs at a given radius. In this way, we need to also consider the area $2\pi r_\Delta r$ of a ring at radius $r$. The highest number of SNRs is expected at a Galactocentric radius of 3.9 kpc.

#### 4.2 Exponential + power-law distribution

The distribution made up of the sum of an exponential and a PL gives similar results between the two samples with $p$-values of 0.67 and 0.73, indicating that this model again fits the data well. Contrary to the MGD, this distribution does not go to zero near the Galactic Centre. Instead it goes to $\infty$ since $f(r) \propto r^{-1.00}$ in the PL component. As mentioned before, for a PL component $<-1$ the number of expected SNRs would go to infinity at the Galactic Centre. To avoid this, we forced $\alpha > -1$. This is also the reason that there is no negative error on this value. For the complete surface brightness limited sample, the best-fitting PL component is $-1$. This means that when integrating over the surface the PL falls out and, since we are then left with an exponential, the maximum number of SNRs is expected at $R_{gal} = 0$.

#### 4.3 Exponential distribution

The exponential distribution again fits both samples well with $p$-values of 0.57 and 0.51 for the complete and centre-excluded

---

**Table 1.** Parameters and goodness-of-fit values for the fits performed with the models described in Section 3.2. All SNRs with a surface brightness above $1.3 \times 10^{-20}$ W m$^{-2}$ Hz$^{-1}$ sr$^{-1}$ at 1 GHz have been used in this analysis. The peak refers to the radius at which the highest number of SNRs is predicted.

<table>
<thead>
<tr>
<th>Model</th>
<th>Parameters</th>
<th>Complete sample Peak (kpc)</th>
<th>$\chi^2$</th>
<th>$P$-value</th>
<th>Sample excluding inner $\pm 10^\circ$ Parameters</th>
<th>$\chi^2$</th>
<th>$P$-value</th>
</tr>
</thead>
<tbody>
<tr>
<td>MGD (equation 7)</td>
<td>$\alpha = 0.74^{+0.38}_{-0.91}$</td>
<td>3.9</td>
<td>33.65</td>
<td>0.44</td>
<td>$\alpha = -0.96^{+0.38}_{-0.04}$</td>
<td>28.03</td>
<td>0.62</td>
</tr>
<tr>
<td></td>
<td>$\beta = 3.52^{+0.91}_{-2.43}$</td>
<td></td>
<td></td>
<td></td>
<td>$\beta = 1.71^{+1.51}_{-1.58}$</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Exponential + PL (equation 12)</td>
<td>$\alpha = -1.00^{+0.41}_{-0.00}$</td>
<td>0</td>
<td>28.89</td>
<td>0.67</td>
<td>$\alpha = -0.91^{+0.61}_{-0.09}$</td>
<td>25.80</td>
<td>0.73</td>
</tr>
<tr>
<td></td>
<td>$\beta = 2.90^{+0.54}_{-0.06}$</td>
<td></td>
<td></td>
<td></td>
<td>$\beta = 3.70^{+0.39}_{-0.91}$</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Exponential (equation 9)</td>
<td>$\beta = 2.46^{+0.39}_{-0.33}$</td>
<td>3.3</td>
<td>31.82</td>
<td>0.57</td>
<td>$\beta = 2.88^{+0.56}_{-0.46}$</td>
<td>31.13</td>
<td>0.51</td>
</tr>
<tr>
<td>PL (equation 8)</td>
<td>$\beta = -1.00^{+0.03}_{-0.00}$</td>
<td>–</td>
<td>58.02</td>
<td>0.006</td>
<td>$\beta = -1.00^{+0.04}_{-0.00}$</td>
<td>59.27</td>
<td>0.002</td>
</tr>
<tr>
<td>Sérsic profile (equation 10)</td>
<td>$n = 1.95^{+0.84}_{-0.67}$</td>
<td>0.9</td>
<td>51.25</td>
<td>0.03</td>
<td>$n = 2.56^{+1.49}_{-0.74}$</td>
<td>27.39</td>
<td>0.70</td>
</tr>
</tbody>
</table>

**Table 2.** The nested-model significances of the MGD and exponential, and PL distributions for both the complete and centre-excluded samples.

<table>
<thead>
<tr>
<th></th>
<th>Complete sample significance</th>
<th>Centre-excluded sample significance</th>
</tr>
</thead>
<tbody>
<tr>
<td>MGD and exponential</td>
<td>1.0</td>
<td>0.078</td>
</tr>
<tr>
<td>MGD and PL</td>
<td>$8 \times 10^{-7}$</td>
<td>$2 \times 10^{-8}$</td>
</tr>
</tbody>
</table>
samples, respectively. The peak of the distribution is at the Galactic Centre, but one again needs to realize that the surface area at a given radius plays a role as well. This means that the maximum number of SNRs at a given radius is at about 3.3 kpc.

4.4 Power-law distribution

The PL model fits both samples poorly (p-values <0.05). The problem with this model is that, just like with the sum of an exponential and PL model, the distribution goes to infinity at the Galactic Centre. In this case, however, this constraint on the PL is more limiting to the model due to it being the only free parameter.

4.5 Sérsic profile

In the complete SNR sample set, the Sérsic profile provides a good fit with a p-value of 0.03 and a peak in the expected number of SNRs at 0.9 kpc. However, a large improvement can be seen when moving to the centre-excluded model, where we find a p-value of 0.70. Looking at Fig. A1, we see that the poor fit for the complete sample is likely due to the model distribution becoming very small at large angles.

4.6 Nested models

The values reported in Table 2 show what the significances are when comparing how well the shown distributions fit. The difference in their $\chi^2$ values can be compared using that it is distributed as a $\chi^2$ distribution with $\nu$ the difference in the number of degrees of freedom between the models, since by going from the MGD to either the PL or exponential model a degree of freedom effectively gets added. The value of $\nu$ implies that the model with the larger amount of degrees of freedom (i.e. the exponential) fits better than the one with fewer (i.e. the MGD). In Section 5, this will be discussed.

5 DISCUSSION

One of the caveats of this investigation is that it is not clear from a theoretical point of view what the SNR distribution should look like. As discussed in the introduction, we in general expect it to follow the massive star formation. But a complicating factor are type Ia SNe that cannot be expected to follow massive star formation in any trivial sense. For this reason, the models investigated here were chosen so to have a mix of distributions and investigate whether the use of the MGD from equation (7) is really appropriate.

Instead of the $10^{−20}$ W m$^2$ Hz$^{-1}$ sr$^{-1}$ surface brightness completeness limit that Green (2004, 2015) use, we opted for a slightly higher limit of $1.3 \times 10^{−20}$ W m$^2$ Hz$^{-1}$ sr$^{-1}$ based upon the Galactic background near the Galactic Centre. Using this surface brightness limited sample, we effectively limited the bias in the SNR sample to the inner $\pm10^\circ$ of the Galaxy.

Green (2015) used the least sum of squares of the differences between the observed and model cumulative distributions. Since in the cumulative distribution the bins are not independent, we instead use the $\chi^2$ statistic from equation (4) to compare the observed and model number of SNRs per bin in $l$. This method was chosen in order to avoid biases from using unweighted fitting to data bins that are not independent. The fact that the minimization of the $\chi^2$ statistic with iterations over $\sigma_n$ provided equal results to fits performed with the use of the Poisson distribution shows that the parameter estimates are reliable.

The SNR catalogue has been updated since Green (2015)’s investigation and along with the increased surface brightness threshold this has resulted in 12 remnants fewer in the sample used here compared to that used in Green (2015). In combination with the improved fitting presented here, this causes the distribution to be more centrally concentrated than the one found by Green (2015). We find that the MGD fit of the surface density peaks at a radius of 1.68 kpc, while Green (2015) and Case & Bhattacharya (1998) find values of 2.25 and 4.53 kpc, respectively. Like Green (2015) noted, the relatively large difference with the distribution found by Case & Bhattacharya (1998) can be explained by their use of the $\Sigma$–$D$ relation that causes bias in their sample.

5.1 Model discussion

The results shown in Table 1 indicate that neither in the complete nor centre-excluded samples, fitting an MGD provides significantly better fits than a simple exponential distribution. Although the MGD does come close to being significantly better in the centre-excluded one, it is in the complete sample where the exponential actually performed better while having less free parameters. The reason this can happen is that the iterations over the errors that were fitted make the parameters deviate from the global minimum of the $\chi^2$ statistic. What we also see in the MGD fit for the complete sample is that the distribution in $l$ does not have its peak at 0°. In Section 2, it was explained that it is likely that not all SNRs with a surface brightness above our threshold of $1.3 \times 10^{−20}$ W m$^2$ Hz$^{-1}$ sr$^{-1}$ at 1 GHz have been found near the Galactic Centre. The drop in the number of SNRs predicted by the MGD below about $\pm10^\circ$ could thus be due to bias in the data. However, this bias is also present for the other fitted models that do not show this feature. This makes it interesting to investigate whether this feature remains present when a more complete SNR sample is analysed.

Apart from the statistical significance argument, there is another important observational argument to be made against the use of the MGD. While the MGD does not predict SNRs near the Galactic Centre, observations (e.g. Maeda et al. 2002; Kennea et al. 2013) have shown SNe to take place near the centre of the Galaxy thus causing a mismatch between the model and observations.

The reason that we have focused mainly on the MGD and exponential profiles is that the PL provides poor fits to both samples. Furthermore, the exponential + PL model cannot be compared in the same way since the PL and exponential distributions are not nested models of it. Moreover, the predicted maximum in the number of SNRs from this model is at the Galactic Centre itself, which we know to be untrue. The Sérsic profile also does not fit the complete sample well and we find that the Sérsic index $n$ for the complete sample is higher than the $1.3 \pm 0.3$ found by Widrow, Pym & Dubinski (2008). When going to the sample with the inner $\pm10^\circ$ excluded, we find an even higher number for $n$ that would correspond to a bulge-dominated system (Driver et al. 2006). In combination with the poor fit to the complete sample, we can quite safely say that the Galactic SNR distribution does not follow a Sérsic profile.

From the above discussion, we conclude that the simple exponential distribution provides the best fit to the data. In Figs 4 and 5, we therefore show a normalized surface density and number density comparison, respectively, between this distribution and the best-fitting MGDs found by Case & Bhattacharya (1998) and Green (2015).

5.2 Galactic Centre

As mentioned before, our SNR sample cannot be assumed to be unbiased near the Galactic Centre ($|l| \leq 10^\circ$). The analysis of the
Even if these regions would not follow the same SNR distribution. The disc and bar regions based on a purely radial SNR distribution, we would not expect a sharp cut-off between the two. The study done by Urquhart et al. (2014) also found other peaks in the distribution that they attribute to local structure in the disc (e.g. spiral arms). Although Urquhart et al. (2014) were not sensitive to Galacticcentric radii ≤3 kpc, their distribution does seem to plateau if not decrease for 3 kpc ≤ R_{gal} ≤ 5 kpc. This seems to correspond to a local minimum in the H_2 surface density at ~2 kpc (Nakanishi & Sofue 2006). Although we see no clear evidence for these structures in our data, this could be the consequence of our use of the 1D projection in l.

Having discussed the Galactic Centre and massive star formation, it is also interesting to discuss the shapes of the distributions shown in Figs 4 and 5. Given the interplay between surface area and surface density, a peak in the number of SNRs at R_{gal} ≠ 0 is not unexpected. The peak in the surface density of SNRs found by Green (2015) and Case & Bhattacharya (1998), on the other hand would imply that the SNR formation mechanism would be most efficient at that particular radius. Given that the SNR sample around the Galactic Centre is likely incomplete, we expect that if such a peak were to be present it would be located at R_{gal} ≲ 1.7 kpc, which is the peak radius we find using an MGD. This would place it within the Galactic bar region (e.g. Babusiaux & Gilmore 2005). We expect the presence of this peak to be unlikely, since there is a lack of observational evidence supporting it and no clear interpretation for a peak in this region. It could therefore be an artefact of the use of an MGD to describe the radial SNR distribution.

As we noted earlier, SNe are the dominant source of chemical enrichment of galaxies with heavy elements. For this reason, a comparison with the radial metallicity distribution in the Galaxy is quite natural. However, given the complex formational history of the Galaxy, one of course needs to be careful in direct comparisons. By studying Cepheids with R_{gal} ≲ 15 kpc, Luck & Lambert (2011) find an exponential relation between metallicity and Galactocentric radius. Assuming that this enrichment happens on a time-scale smaller than or similar to that of the radial diffusion of matter in the Galaxy in combination with the assumption that the current SNR distribution is representative for that in the past, this result seems to support an exponential SNR distribution for R_{gal} ≳ 3–5 kpc.

5.4 Future work

It would be interesting to see how well the SNR distribution in other spiral galaxies is described by the exponential profile in comparison with the other models. In this way it could be investigated whether, and how far, the distribution found here holds a universal power in describing SNR distributions. What would also make this advantageous is that the full spatial Σ–D relation that causes a large distance determination scatter in the Galaxy is not needed for extra-galactic sources. Especially in face-on spiral galaxies, all remnants are at approximately the same distance in a 2D configuration. The
distance to the centre of their host galaxy can then be determined using geometry. Unfortunately it would be difficult to compare the results with what has been done here. Extragalactic SNRs suffer from different detection biases that would change per host galaxy. Quantifying these biases would therefore be another study in itself. Since comparing the distributions directly as if no biases are present would be of questionable value, we have opted to focus solely on the Galactic distribution instead.

With regards to discovering new SNRs, the region surrounding the Galactic Centre (|l| ≲ 10°) is the most promising. Although the observational challenges are greatest in this region, it is also expected to harbour the greatest undiscovered Galactic SNR density per square unit area on the sky.

6 CONCLUSIONS

By imposing a surface brightness completeness limit of 1.3 × 10⁻²⁰ W m⁻² Hz⁻¹ sr⁻¹ at 1 GHz, we have analysed the radial distribution of SNRs in the Galaxy. We fitted multiple model distributions to determine whether the use of an MGD model that has seen persistent use in the literature is really justified. The analysis has been done for two samples. The first sample included every SNR above the surface brightness threshold and the second sample only included, from these remaining SNRs, those with l > ±10°, since biases are mainly present near the Galactic Centre. In neither of these samples does the MGD provide the best-fitting results. Moreover, the MGD does not provide significantly better fits than a simple exponential profile in either of these samples. We find that a simple exponential profile of the form given in equation (9), with β = 2.46±0.39 and an integrated SNR density peak at 3.3 kpc, is the most consistent and, with one free parameter, least complicated model.

Given the relatively high bias in the data set around the Galactic Centre, no strict inferences have been made about the central region of the Galaxy. By comparing our findings to those of studies investigating the radial distribution of massive star formation and metallicity, we find indirect evidence supporting an exponential SNR surface density for Rgal ≳ 3–5 kpc. Our study therefore suggests that preference should be given to a simple exponential distribution over the MGD given in equation (7) when describing the radial SNR distribution in the Galaxy.

ACKNOWLEDGEMENT

We would like to give special thanks to Ciaran Rogers and Stefanie Fijma for their valuable feedback on an earlier version of this paper.

DATA AVAILABILITY

The Galactic SNR data underlying this article are available at https://doi.org/10.1007/s12036-019-9601-6. The 408 MHz all-sky radio map data underlying this article are available in MPIfR’s Survey Sampler at http://www3.mpifr-bonn.mpg.de/survey.html.

REFERENCES

de Vaucouleurs G., 1979, Ann. Astrophys., 11, 247
Hillebrandt W., Niemeyer J. C., 2000, ARA&A, 38, 191
Sérsic J. L., 1963, BAAA, 6, 41

APPENDIX A: MODEL FITS

Here, we provide all the model fits that have been performed with the different distributions. Both the model fits to the complete and centre-excluded data sets are shown.
Figure A1. The best model fits to the complete surface brightness limited data set.
Figure A2. The best model fits to the centre-excluded sample.

This paper has been typeset from a TeX/LaTeX file prepared by the author.